

$$1. \quad u(-x,t) = \int_{-\infty}^{\infty} S(-x-y,t) \phi(y) dy$$

$$= \int_{-\infty}^{\infty} S(x+y,t) \phi(y) dy$$

(Here we used $S(x,t) = \frac{1}{\sqrt{4\pi k t}} e^{-\frac{x^2}{4kt}} = S(-x,t)$)

$$= - \int_{-\infty}^{\infty} S(x+y,t) \phi(-y) dy \quad (\phi \text{ odd})$$

$$= - \int_{-\infty}^{\infty} S(x-y,t) \phi(y) dy \quad (\text{change of variable})$$

$$= -u(x,t).$$

$$2. \quad u(-x,t) = \int_{-\infty}^{\infty} S(-x-y,t) \phi_{\text{even}}(y) dy$$

$$= \int_{-\infty}^{\infty} S(x+y,t) \phi_{\text{even}}(y) dy \quad (S(x,t) = S(-x,t))$$

$$= \int_{-\infty}^{\infty} S(x+y,t) \phi_{\text{even}}(-y) dy \quad (\phi_{\text{even even}})$$

$$= \int_{-\infty}^{\infty} S(x-y,t) \phi_{\text{even}}(y) dy \quad (\text{change of variable})$$

$$= u(x,t)$$

Now $u(x,t) = \int_{-\infty}^{\infty} S(x-y,t) \phi_{\text{even}}(y) dy$ clearly

satisfies $ut - k u_{xx} = 0$ for $-\infty < x < \infty$, $t > 0$

and $u(x,0) = \phi_{\text{even}}(x) = \phi(x)$ for $0 < x < \infty$.

$$\text{Then } u_x(x,t) = \int_{-\infty}^{\infty} \frac{\partial}{\partial x} S(x-y,t) \phi_{\text{even}}(y) dy$$

$$\Rightarrow u_x(0,t) = \int_{-\infty}^{\infty} \frac{\partial}{\partial x} S(-y,t) \phi_{\text{even}}(y) dy$$

$$\begin{aligned}
 \left(\frac{\partial}{\partial x} S(x,t) \right) &= \frac{\partial}{\partial x} \left(\frac{1}{\sqrt{4\pi k t}} e^{-\frac{x^2}{4kt}} \right) \\
 &= \frac{1}{\sqrt{4\pi k t}} \cdot \left(-\frac{x}{2kt} \right) \cdot e^{-\frac{x^2}{4kt}} \\
 &= \int_{-\infty}^{\infty} \frac{1}{\sqrt{4\pi k t}} \cdot \frac{y}{2kt} \cdot e^{-\frac{y^2}{4kt}} \cdot \underbrace{\phi_{\text{even}}(y)}_{\text{even in } y} dy \\
 &\quad \underbrace{\text{odd in } y}_{\text{odd in } y} \\
 &= 0. \quad \left(\int_{-\infty}^{\infty} \text{an odd function} = 0 \right)
 \end{aligned}$$

We have checked that $u(x,t)$ is a sol. to (*).

$$\begin{aligned}
 (2) \quad u(x,t) &= \int_{-\infty}^{\infty} S(x-y,t) \phi_{\text{even}}(y) dy \\
 &= \int_0^{\infty} S(x-y,t) \phi(y) dy + \int_{-\infty}^0 S(x-y,t) \phi(-y) dy \\
 &= \int_0^{\infty} S(x-y,t) \phi(y) dy + \int_0^{\infty} S(x+y,t) \phi(y) dy \\
 &\quad (\text{change of variable}) \\
 &= \int_0^{\infty} [S(x-y,t) + S(x+y,t)] \phi(y) dy.
 \end{aligned}$$

3. $u(-x,t) = -u(x,t)$:

$$\begin{aligned}
 u(-x,t) &= \frac{1}{2} [\phi_{\text{ext}}(-x-ct) + \phi_{\text{ext}}(-x+ct)] + \\
 &\quad \frac{1}{2c} \int_{-x-ct}^{-x+ct} \chi_{\text{ext}}(y) dy
 \end{aligned}$$

$$= \frac{1}{2} [-\phi_{\text{ext}}(x+ct) - \phi_{\text{ext}}(x-ct)]$$

$$- \frac{1}{2c} \int_{-x-ct}^{-x+ct} \gamma_{\text{ext}}(-y) dy$$

(Using the oddness of ϕ_{ext} & γ_{ext} at $x=0$)

$$= -\frac{1}{2} [\phi_{\text{ext}}(x-ct) + \phi_{\text{ext}}(x+ct)]$$

$$+ \frac{1}{2c} \int_{x+ct}^{x-ct} \gamma_{\text{ext}}(y) dy \quad (\text{change of variable})$$

$$= -\frac{1}{2} [\phi_{\text{ext}}(x-ct) + \phi_{\text{ext}}(x+ct)]$$

$$- \frac{1}{2c} \int_{x-ct}^{x+ct} \gamma_{\text{ext}}(y) dy$$

$$= -u(x,t)$$

$$u(zL-x,t) = -u(x,t) :$$

$$u(zL-x,t) = \frac{1}{2} [\phi_{\text{ext}}(zL-x-ct) + \phi_{\text{ext}}(zL-x+ct)]$$

$$+ \frac{1}{2c} \int_{zL-x-ct}^{zL-x+ct} \gamma_{\text{ext}}(y) dy$$

$$= \frac{1}{2} [-\phi_{\text{ext}}(x+ct) - \phi_{\text{ext}}(x-ct)]$$

$$- \frac{1}{2c} \int_{zL-x-ct}^{zL-x+ct} \gamma_{\text{ext}}(zL-y) dy$$

(Using the oddness of ϕ_{ext} & γ_{ext} at $x=L$)

$$= -\frac{1}{2} [\phi_{\text{ext}}(x-ct) + \phi_{\text{ext}}(x+ct)]$$

(3)

$$\begin{aligned}
& + \frac{1}{2c} \int_{x+ct}^{x-ct} \gamma_{\text{ext}}(y) dy \quad (\text{change of variable})^{\textcircled{4}} \\
& = - \frac{1}{2} [\phi_{\text{ext}}(x-ct) + \phi_{\text{ext}}(x+ct)] \\
& \quad - \frac{1}{2c} \int_{x-ct}^{x+ct} \gamma_{\text{ext}}(y) dy \\
& = -u(x, t).
\end{aligned}$$

Finally, $u(x, t)$ as given by the d'Alembert's formula clearly satisfies $u_{tt} - c^2 u_{xx} = 0$ for all x and t ,

$$\text{and } u(x, 0) = \phi_{\text{ext}}(x) = \phi(x) \text{ for } 0 < x < L$$

$$\text{and } u_t(x, 0) = \gamma_{\text{ext}}(x) = \gamma(x) \text{ for } 0 < x < L.$$

Lastly, $u(x, t) = -u(-x, t) \Rightarrow u(0, t) = 0$

and $u(x, t) = -u(2L-x, t) \Rightarrow u(L, t) = 0$.

Thus $u(x, t)$ is a sol to (*).

$$4. \quad \phi'_{\text{ext}}(x) + a\phi_{\text{ext}}(x) = \begin{cases} \phi'(x) + a\phi(x), & x > 0 \\ \phi'_o(x) + a\phi_o(x), & x \leq 0. \end{cases}$$

$$\text{Oddness} \Rightarrow \phi'_o(x) + a\phi_o(x) = -\phi'(-x) - a\phi(-x), \quad x \leq 0$$

Solving it for $\phi_o(x)$ by integrating factor :

$$(e^{ax}\phi_o(x))' = -e^{ax}\phi'(-x) - e^{ax}a\phi(-x)$$

$$\phi_o(x) = -e^{-ax} \left(\int_0^x e^{ay} \phi'(-y) + e^{ay} a\phi(-y) dy + C \right)$$

$$\phi_o(0) = \phi(0) \Rightarrow C = -\phi(0).$$

$$\Rightarrow \phi_0(x) = -e^{-ax} \left[\int_0^x e^{ay} \phi'(-y) + e^{ay} a \phi(-y) dy - \phi(0) \right] \quad (5)$$

(2) Since $u(x,t)$ is a sol. to the heat equation,

$u_x(x,t)$ is also a sol. to the heat equation

$$((u_x)_t - k(u_x)_{xx} = (u_t - k u_{xx})_x = 0)$$

$\Rightarrow w = u_x + a u$ is also a sol. to the heat eq.

$$((u_x + a u)_t - k(u_x + a u)_{xx} = (u_x)_t - k(u_x)_{xx} \\ + a(u_t - k u_{xx}) = 0)$$

$$\text{Now } u(x,t) = \int_{-\infty}^{\infty} S(x-y,t) \phi_{\text{ext}}(y) dy$$

$$\text{satisfies } u(x,0) = \phi_{\text{ext}}(x).$$

$$\begin{aligned} \text{Then } w(x,0) &= u_x(x,0) + a u(x,0) \\ &= \phi'_{\text{ext}}(x) + a \phi_{\text{ext}}(x), \end{aligned}$$

which is the initial condition for $w(x,t)$.

There is a mathematical subtlety in the derivation of $u_x(x,0) = \phi'_{\text{ext}}(x)$ from

$u(x,0) = \phi_{\text{ext}}(x)$. The continuity of $\phi_{\text{ext}}(x)$ becomes important here. This is the reason that we require $\phi_0(0) = \phi(0)$ in the definition of $\phi_{\text{ext}}(x)$ in order to make $\phi_{\text{ext}}(x)$ continuous. We will not explain this subtlety in detail here.

(3) Now since $w(x,0) = \phi'_{\text{ext}}(x) + a \phi_{\text{ext}}(x)$ is an

(6)

odd function, by Problem 1 in HW7. $w(x,t)$

is odd w.r.t. x . Thus in particular $w(0,t) = 0$.

That is. $U_x(0,t) + aU(0,t) = 0$, for all $0 < t < \infty$. Now clearly $U(x,t) = \int_{-\infty}^{\infty} S(x-y,t) \phi_{ext}(y) dy$ satisfies the heat equation and that

$$U(x,0) = \phi_{ext}(x) = \phi(x) \text{ for } x \geq 0.$$

so in conclusion, $U(x,t)$ is a sol. to (2).

3.2.6. We look for solutions given by the d'Alembert's

$$U(x,t) = \frac{1}{2c} \int_{x-a}^{x+ct} \gamma_{ext}(y) dy,$$

where γ_{ext} is the extension of the constant function V defined for $x > 0$ to the whole line. Then $U(x,t)$ already satisfies $\begin{cases} Ut - c^2 U_{xx} = 0 \\ U(x,0) = 0 \end{cases}$,

$$\begin{cases} Ut - c^2 U_{xx} = 0 \\ U(x,0) = 0 \\ U_t(x,0) = \gamma_{ext}'(x) = V, \text{ when } x > 0. \end{cases}$$

It suffices to check for the boundary condition

$$Ut(0,t) + aU_x(0,t) = 0$$

$$\text{Computation: } U_t(x,t) = \frac{1}{2c} [\gamma_{ext}'(x+ct) \cdot c - \gamma_{ext}'(x-ct) \cdot (-c)]$$

$$U_x(x,t) = \frac{1}{2c} [\gamma_{ext}'(x+ct) - \gamma_{ext}'(x-ct)]$$

$$\Rightarrow Ut(0,t) + aU_x(0,t) = \frac{1}{2c} [(c+a)\gamma_{ext}'(ct) + (c-a)\gamma_{ext}'(-ct)] \\ = 0$$

$$\Rightarrow \gamma_{ext}'(-ct) = \frac{a+c}{a-c} \gamma_{ext}'(ct) = \frac{a-c}{a+c}. V, t > 0.$$

$$\text{Thus } \gamma_{\text{ext}}(x) = \begin{cases} V, & x > 0 \\ \frac{a+c}{a-c}V, & x < 0 \end{cases}$$

(7)

$$\text{Then } U(x,t) = \frac{1}{2c} \int_{x-ct}^{x+ct} \gamma_{\text{ext}}(y) dy$$

$$= \frac{1}{2c} \int_{x-ct}^{x+ct} V dy = V \cdot t, \text{ if } x-ct > 0$$

$$= \frac{1}{2c} \left(\int_{x-ct}^0 \frac{a+c}{a-c}V dy + \int_0^{x+ct} V dy \right) \text{ if } x-ct < 0.$$

$$= \frac{1}{2c} \left(\frac{a+c}{a-c} \cdot V \cdot (ct-x) + V \cdot (x+ct) \right).$$

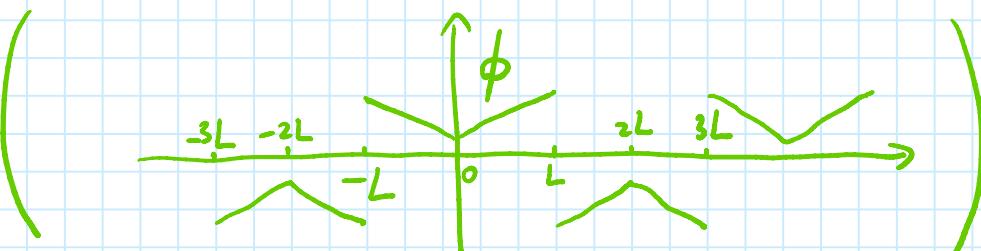
$$(x > 0, t \geq 0)$$

6. To make sure $U_x(0,t) = 0$, we require that ϕ_{ext} be even at $x=0$, i.e.

$$\phi_{\text{ext}}(x) = \phi_{\text{ext}}(-x);$$

To make sure $U(L,t) = 0$, we require that ϕ_{ext} be odd at $x=L$, i.e.

$$\phi_{\text{ext}}(x) = -\phi_{\text{ext}}(2L-x).$$



Thus ϕ_{ext} is $4L$ -periodic. and

(8)

$$\phi_{\text{ext}}(x) = \begin{cases} \phi(x - 4Ln), & 4Ln \leq x < 4Ln + L \\ -\phi(2L - (x - 4Ln)), & 4Ln + L \leq x < 4Ln + 2L \\ -\phi(x - 4Ln - 2L), & 4Ln + 2L \leq x < 4Ln + 3L \\ \phi(4L - (x - 4Ln)) & 4Ln + 3L \leq x < 4Ln + 4L \end{cases}$$

$(n \in \mathbb{Z})$

Then (let $U(x, t) = \int_{-\infty}^{\infty} S(x-y, t) \phi_{\text{ext}}(y) dy$,

then $U(x, t)$ is a sol. to the heat equation

satisfying $U(x, 0) = \phi_{\text{ext}}(x) = \phi(x)$, $0 < x < L$

and the evenness of ϕ_{ext} at $x=0$

would imply that $U_x(0, t) = 0$ and

the oddness of ϕ_{ext} at $x=L$

would imply that $U(L, t) = 0$.

(Verification details omitted. Compare with
Problem 1 & 2 in HW7.)

thus $U(x, t)$ is a sol. to (X).

Finally we write $U(x, t)$ as an infinite sum of
integrals :

$$U(x,t) = \sum_{n \in \mathbb{Z}} \int_{4Ln}^{4Ln+4L} S(x-y,t) \phi_{0xt}(y) dy$$

$$= \sum_{n \in \mathbb{Z}} \left(\int_{4Ln}^{4Ln+L} + \int_{4Ln+L}^{4Ln+2L} + \int_{4Ln+2L}^{4Ln+3L} + \int_{4Ln+3L}^{4Ln+4L} \right)$$

(9)

where

$$\int_{4Ln}^{4Ln+L} = \int_{4Ln}^{4Ln+L} S(x-y,t) \phi(y - 4Ln) dy$$

$$= \int_0^L S(x - 4Ln - y, t) \phi(y) dy$$

$$\int_{4Ln+L}^{4Ln+2L} = - \int_{4Ln+L}^{4Ln+2L} S(x-y,t) \phi(2L - (y - 4Ln)) dy$$

$$= - \int_0^L S(x - 4Ln - 2L + y, t) \phi(y) dy$$

$$\int_{4Ln+2L}^{4Ln+3L} = - \int_{4Ln+2L}^{4Ln+3L} S(x-y,t) \phi(y - 4Ln - 2L) dy$$

$$= - \int_0^L S(x - 4Ln - 2L - y, t) \phi(y) dy$$

$$\int_{4Ln+3L}^{4Ln+4L} = \int_{4Ln+3L}^{4Ln+4L} S(x-y,t) \phi(4L - (x - 4Ln)) dy$$

$$= \int_0^L S(x - 4Ln - 4L + y, t) \phi(y) dy.$$

Thus

$$U(x,t) = \sum_{n \in \mathbb{Z}} \int_0^L \left(S(x - 4Ln - y, t) - S(x - 4Ln - 2L + y, t) - S(x - 4Ln - 2L - y, t) + S(x - 4Ln - 4L + y, t) \right) \phi(y) dy$$

7. By Exercise 2 in HW7. a sol. to

$$\begin{cases} u_t - k u_{xx} = 0, & x > 0, \quad t > 0 \\ u(x, 0) = \phi(x), & x > 0 \\ u_x(0, t) = 0 & t > 0 \end{cases}$$

is $u(x, t) = \int_0^\infty (S(x-y, t) + S(x+y, t)) \phi(y) dy.$

Denote this formula by $\mathcal{G}(t)\phi$.

Then

$$\begin{aligned} & \int_0^t \mathcal{G}(t-s)f(\cdot, s) ds \\ &= \int_0^t \int_0^\infty (S(x-y, t-s) + S(x+y, t+s)) f(y, s) dy ds. \end{aligned}$$

Thus by the Duhamel's principle,

$$\begin{aligned} u(x, t) &= \mathcal{G}(t)\phi + \int_0^t \mathcal{G}(t-s)f(\cdot, s) ds \\ &= \int_0^\infty (S(x-y, t) + S(x+y, t)) \phi(y) dy + \\ & \quad \int_0^t \int_0^\infty (S(x-y, t-s) + S(x+y, t-s)) f(y, s) dy ds \end{aligned}$$

is a sol. to the inhomogeneous problem (*).

8. Provided u is a sol. to (*),

$$w(x, t) = u(x, t) - x g(t) \text{ is a sol. to}$$

(10)

(11)

$$(\ast\ast) \quad \left\{ \begin{array}{l} w_t - k w_{xx} = f(x,t) - x g'(t) \\ w(x,0) = \phi(x) - x g(0) \\ w_x(0,t) = g(t) - \bar{g}(t) = 0 \end{array} \right.$$

and vice versa (that is, provided $w(x,t) = u(x,t) - x g(t)$
 is a sol. to the above problem $(\ast\ast)$,
 $u(x,t)$ is a sol. to (\ast) .)

By the previous exercise,

$$\begin{aligned} w(x,t) &= \int_0^\infty (S(x-y,t) + S(x+y,t)) (\phi(y) - y g(0)) dy \\ &+ \int_0^t \int_0^\infty (S(x-y,t) + S(x+y,t)) (f(y,s) - y g'(s)) dy ds \end{aligned}$$

is a sol. to $(\ast\ast)$.

$$\begin{aligned} \text{Then } u(x,t) &= w(x,t) + x g(t) \\ &= \int_0^\infty (S(x-y,t) + S(x+y,t)) (\phi(y) - y g(0)) dy + \\ &\int_0^t \int_0^\infty (S(x-y,t) + S(x+y,t)) (f(y,s) - y g'(s)) dy ds + x g(t) \end{aligned}$$

is a sol. to (\ast) .