Simultaneous determination of the cosmic birefringence and miscalibrated polarisation angles from CMB experiments arXiv:1904.12440 [astro-ph.CO]

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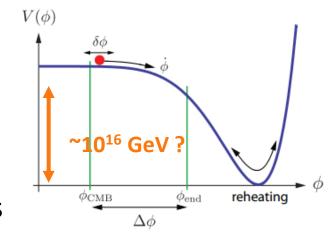
Before we dive into the content of the paper

Introduction to cosmic microwave background (CMB) measurement

Inflation: Physics of high energy scale

- > Exponential expansion of space in the early universe
- ➤ Hope to explore GUT-scale physics
- > The potential of single field slow-roll model is

$$V^{1/4} \sim 1.04 \cdot \left(\frac{r}{0.01}\right)^{\frac{1}{4}} \times 10^{16} \text{ GeV}$$



r: tensor-to-scalar ratio
Power-spectrum ratio of the tensor type
perturbation to scalar type perturbation

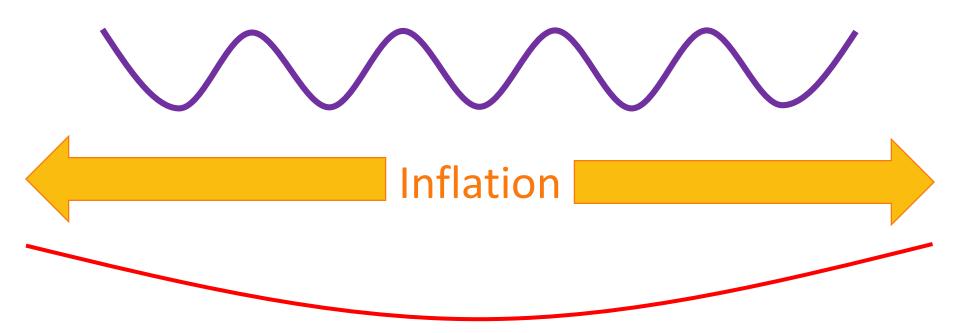
$$r \equiv \frac{\langle h_{ij} h^{ij} \rangle}{\langle \zeta^2 \rangle}$$

- > Two types of fluctuation in inflation
 - \triangleright scalar $\langle \zeta^2 \rangle$: density
 - \succ tensor $\langle |h|^2 \rangle$: primordial gravitational wave (PGW)

$$ds^{2} = a^{2}(t) (1 + 2\zeta(x,t)) (\delta_{ij} + h_{ij}(x,t)) dx^{i} dx^{j}$$

Inflation and Gravitational Waves (GWs)

Inflation expands the space by $\sim 10^{26}$ times in $\sim 10^{-36}$ sec

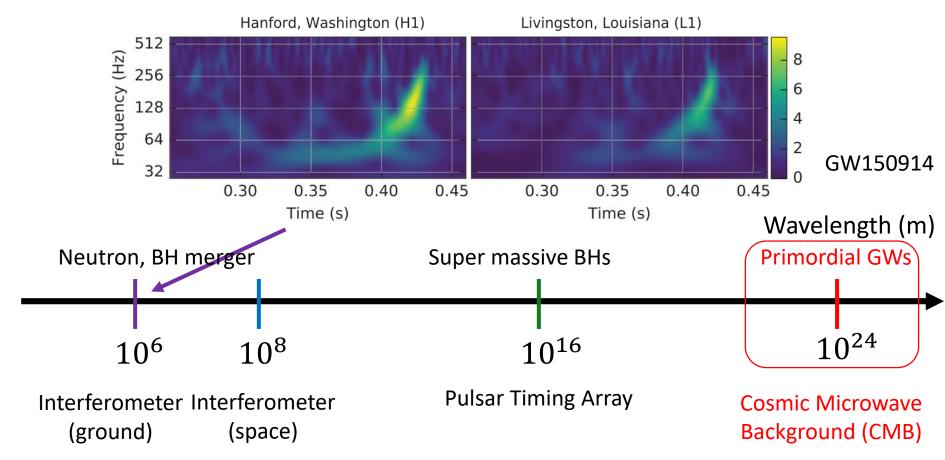


Quantum fluctuation becomes large scale fluctuation

Primordial gravitational waves

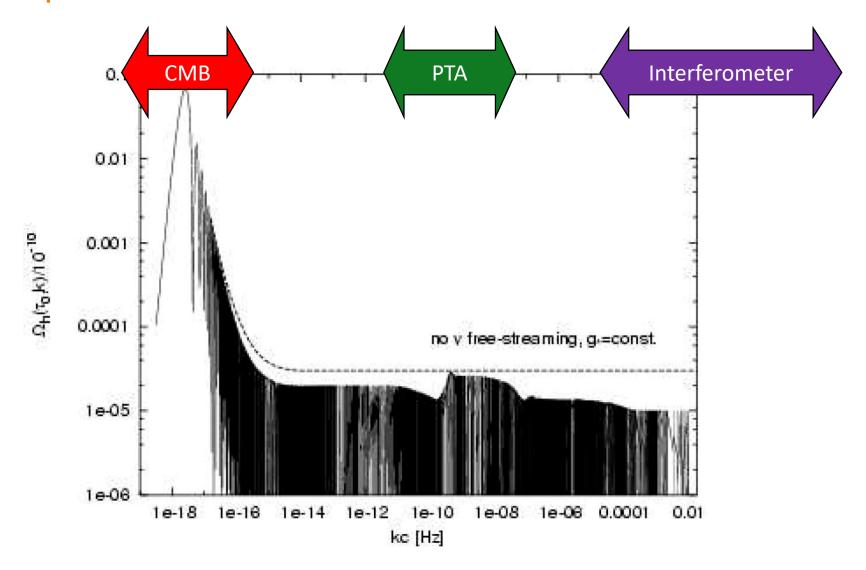
Primordial gravitational wave (PGW)

> LIGO made a first observation of gravitational wave > Wavelength is $\sim 10^6$ m



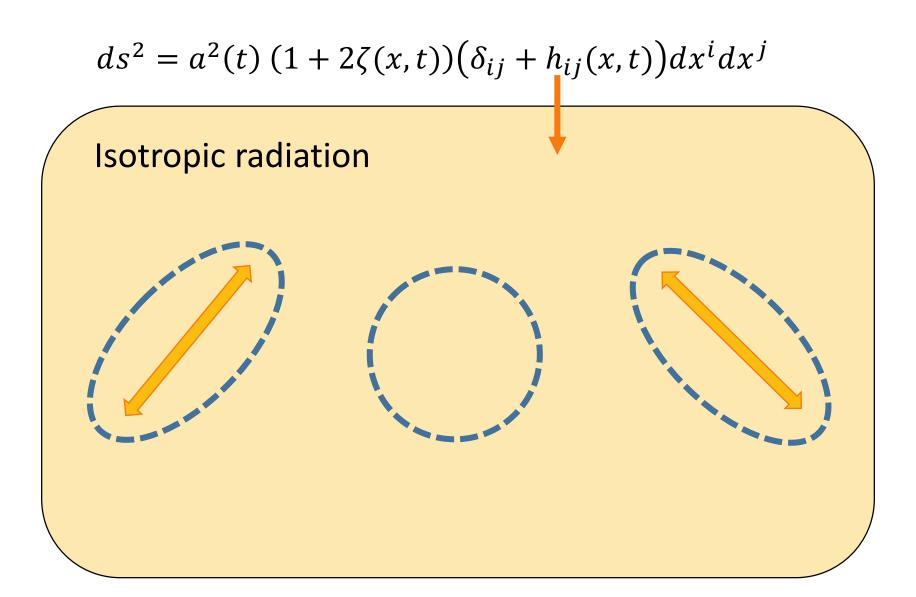
We target the 10 billion light year scale of gravitational waves

Power spectrum of PGWs



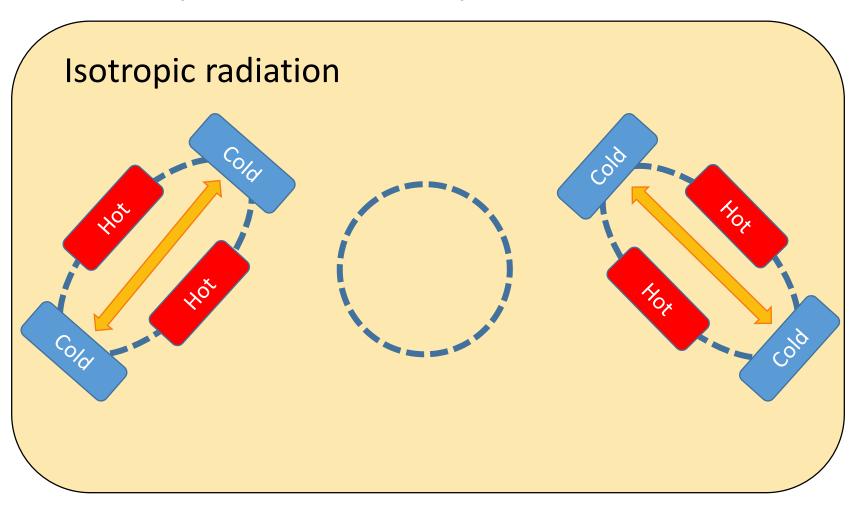
Watanabe and Komatsu (2006)

How to measure PGWs with CMB?



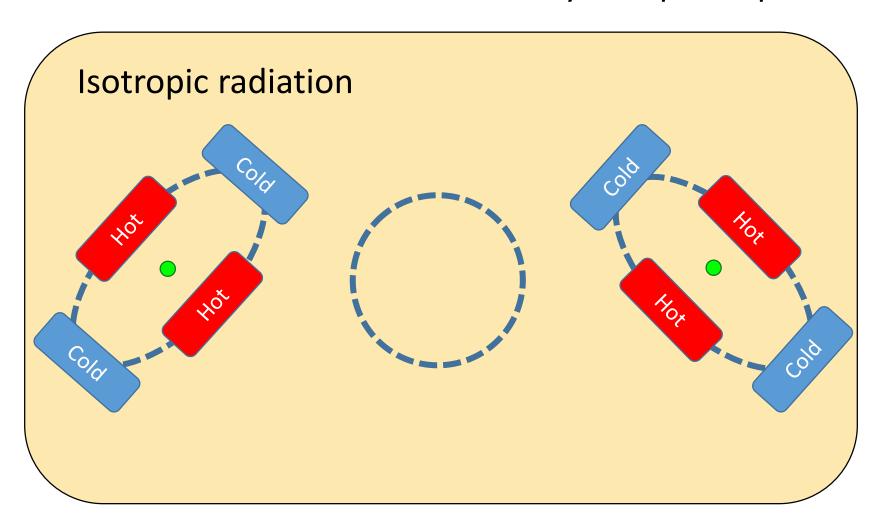
How to measure PGWs with CMB?

Space is distorted by fluctuation

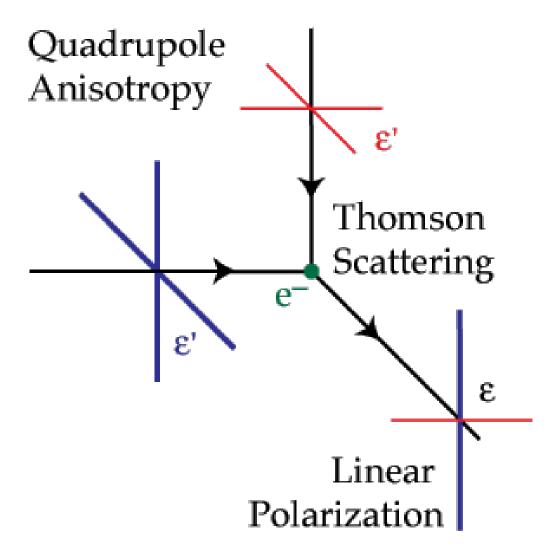


How to measure PGWs with CMB?

What if electrons are surrounded by the quadrupole?



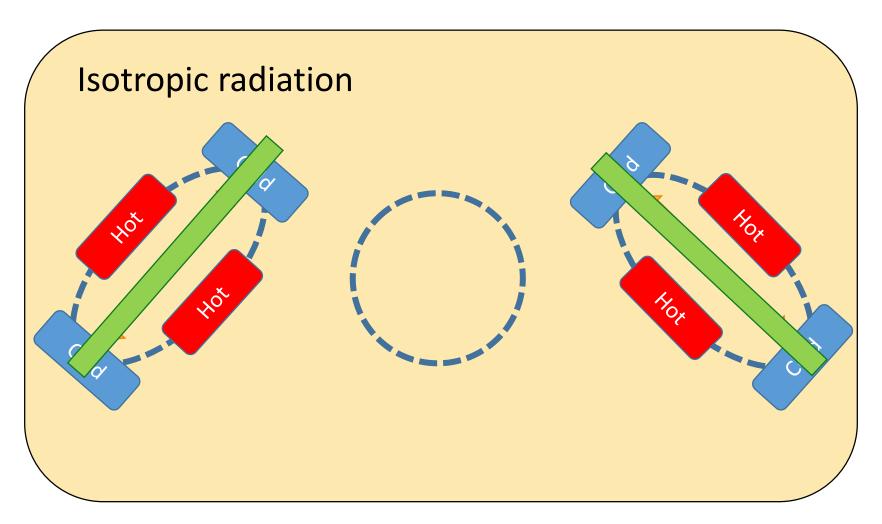
Linear polarization from quadrupole anisotropy



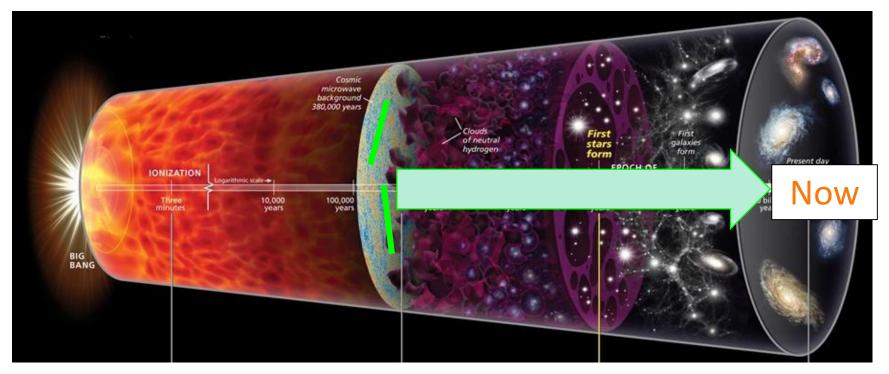
by Wayne Hu

Linear polarization from gravitational waves

Quadrupole creates linear polarization



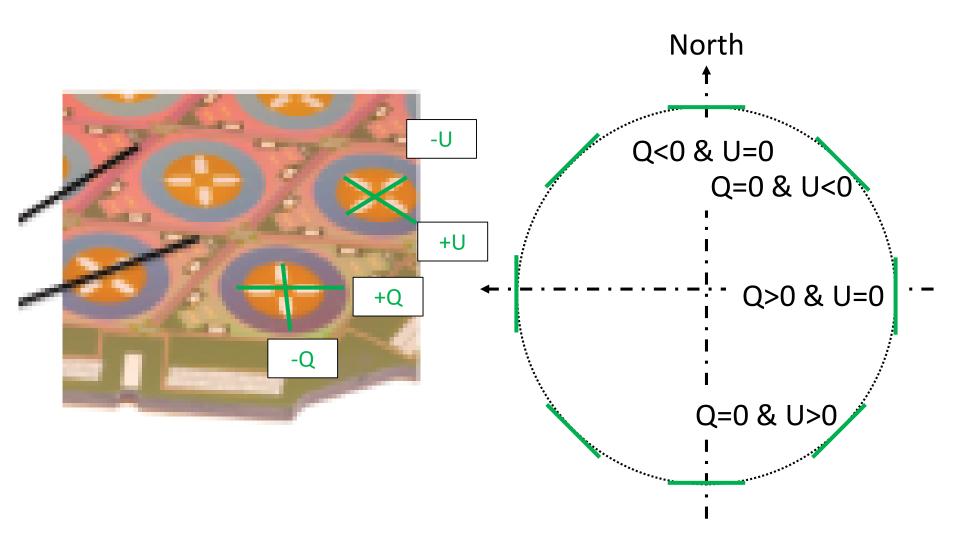
Cosmic microwave background



Credit: Roen Kelly, Discovermagazine

Last scatterings at the recombination era imprint the quadrupole anisotropy on CMB

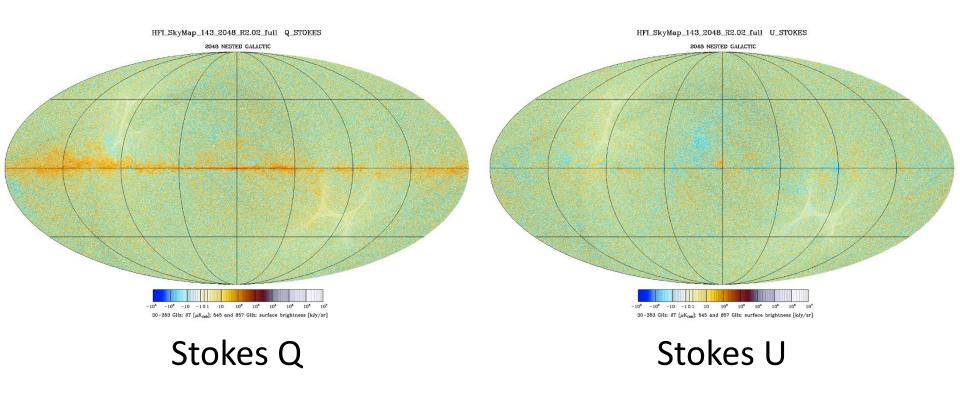
Stokes parameter



Linear polarization is measured as Stokes parameters

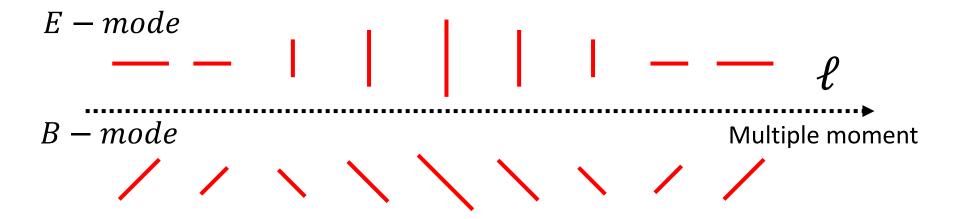


Planck Map of 143 GHz channel



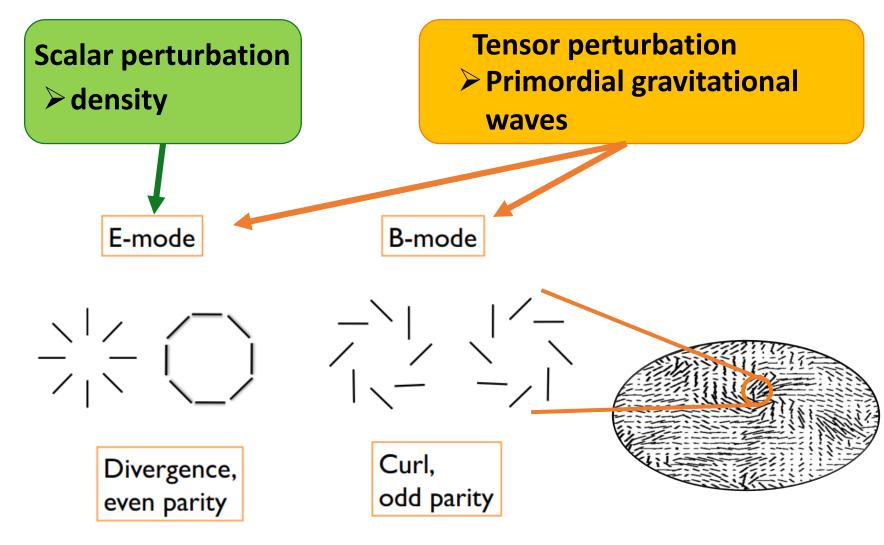
E-mode and B-mode

$$E(\ell) \pm iB(\ell) = e^{\mp 2i\phi_{\ell}} \int d\widehat{\boldsymbol{n}} \left[Q(\widehat{\boldsymbol{n}}) + iU(\widehat{\boldsymbol{n}}) \right] e^{-i\ell \cdot \widehat{\boldsymbol{n}}}$$



Rotation invariant *E*- and *B*- modes are used for the analysis of CMB polarization

E-mode and B-mode polarization



r = (tensor perturbation)/(scalar perturbation)

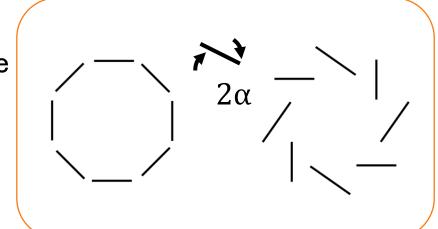
B-mode search is needed to determine r!

Next: Contents of the paper

Introduction:

 \blacktriangleright Miscalibration of detector rotation angle (α) creates spurious B mode from E mode

$$C_\ell^{BB,\mathrm{o}} = C_\ell^{EE} \sin^2(2\alpha) + C_\ell^{BB} \cos^2(2\alpha)$$



observed

We need to estimate α and calibrate rotation angle

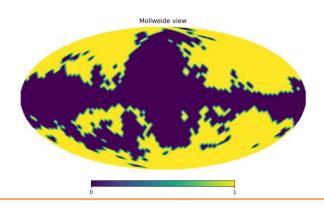
 \blacktriangleright In past experiments, this lpha was calculated assuming EB correlation of CMB is zero

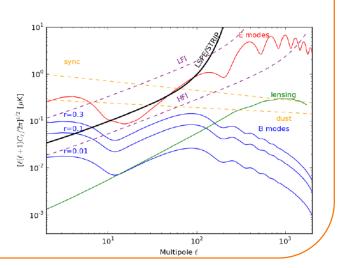
$$C_{\ell}^{EB,\mathrm{o}} = \frac{1}{2} \begin{pmatrix} C_{\ell}^{EE} - C_{\ell}^{BB} \end{pmatrix} \sin(4\alpha)$$
 From theory Brian G. Keating et al. (2013)

(Assumption in the method)

CMB is dominant

- ➤ Only can be used in CMB channel?
- > Need to mask foreground





Cosmic birefringence is negligible

- Create EB correlation
- > Explain afterwards

Review the relation from the beginning

We started from the coefficients of spherical harmonics:

$$E_{\ell,m}^{o} = E_{\ell,m} \cos(2\alpha) - B_{\ell,m} \sin(2\alpha),$$

$$B_{\ell,m}^{o} = E_{\ell,m} \sin(2\alpha) + B_{\ell,m} \cos(2\alpha),$$

From these equations, we find

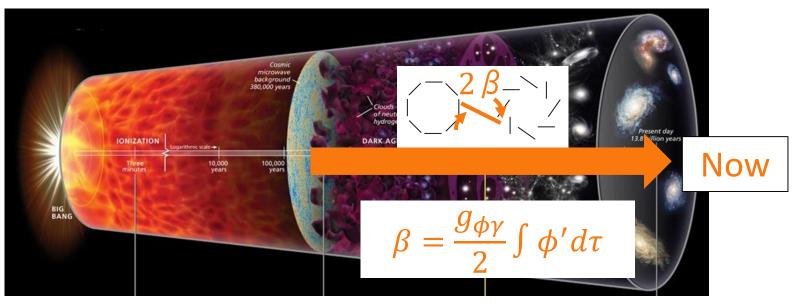
$$C_\ell^{EB,\mathrm{o}} = \frac{1}{2} \left(C_\ell^{EE,\mathrm{o}} - C_\ell^{BB,\mathrm{o}} \right) \tan(4\alpha) + \frac{C_\ell^{EB}}{\cos(4\alpha)} \,.$$
 G. B. Zhao et al. (2015)

Our work

- \triangleright We can estimate α only with observed data
- ➤ If we assume theory CMB power spectra, we can estimate an additional angle!

Cosmological birefringence

During the long travel from recombination era, CMB can be rotated by some physics (e.g. axionic fields)



Credit: Roen Kelly, Discovermagazine

In such case,

- \triangleright Foreground term: rotated only with α
- \triangleright CMB term: rotated with $\alpha + \beta$

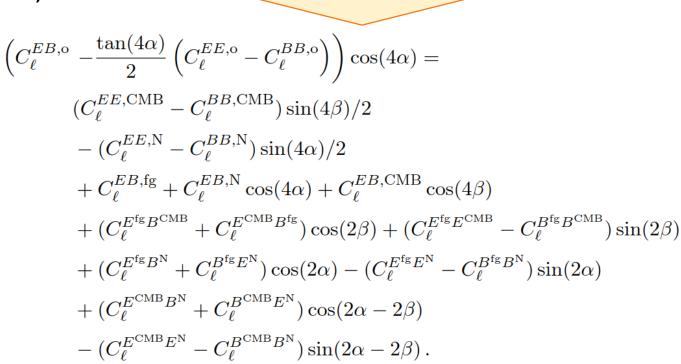
Equations including birefringence rotation:

The coefficients becomes

$$E_{\ell,m}^{\text{o}} = E_{\ell,m}^{\text{fg}} \cos(2\alpha) - B_{\ell,m}^{\text{fg}} \sin(2\alpha) + E_{\ell,m}^{\text{CMB}} \cos(2\alpha + 2\beta) - B_{\ell,m}^{\text{CMB}} \sin(2\alpha + 2\beta)$$

$$B_{\ell,m}^{\text{o}} = E_{\ell,m}^{\text{fg}} \sin(2\alpha) + B_{\ell,m}^{\text{fg}} \cos(2\alpha) + E_{\ell,m}^{\text{CMB}} \sin(2\alpha + 2\beta) + B_{\ell,m}^{\text{CMB}} \cos(2\alpha + 2\beta) + B_{\ell,m}^{\text{CMB}} \cos(2\alpha + 2\beta) + B_{\ell,m}^{\text{CMB}} \cos(2\alpha + 2\beta)$$

From them, we derived



Equations including birefringence rotation:

If we take ensemble average

$$\begin{split} \langle C_\ell^{EB,\mathrm{o}} \rangle = & \frac{\tan(4\alpha)}{2} \left(\langle C_\ell^{EE,\mathrm{o}} \rangle - \langle C_\ell^{BB,\mathrm{o}} \rangle \right) + \frac{\sin(4\beta)}{2\cos(4\alpha)} \left(\langle C_\ell^{EE,\mathrm{CMB}} \rangle - \langle C_\ell^{BB,\mathrm{CMB}} \rangle \right) \\ & + \frac{1}{\cos(4\alpha)} \langle C_\ell^{EB,\mathrm{fg}} \rangle + \frac{\cos(4\beta)}{\cos(4\alpha)} \langle C_\ell^{EB,\mathrm{CMB}} \rangle \,. \end{split} \quad \text{Assume these to be zero Revisit} \end{split}$$

Therefore, we can determine both miscalibration angle and birefringence-rotation angle simultaneously!

Log-likelihood

If we take ensemble average

$$\langle C_{\ell}^{EB,\mathrm{o}} \rangle = \frac{\tan(4\alpha)}{2} \left(\langle C_{\ell}^{EE,\mathrm{o}} \rangle - \langle C_{\ell}^{BB,\mathrm{o}} \rangle \right) + \frac{\sin(4\beta)}{2\cos(4\alpha)} \left(\langle C_{\ell}^{EE,\mathrm{CMB}} \rangle - \langle C_{\ell}^{BB,\mathrm{CMB}} \rangle \right)$$

$$+ \frac{1}{\cos(4\alpha)} \langle C_{\ell}^{EB,\mathrm{fg}} \rangle + \frac{\cos(4\beta)}{\cos(4\alpha)} \langle C_{\ell}^{EB,\mathrm{CMB}} \rangle.$$
 Assume these to be zero Revisit

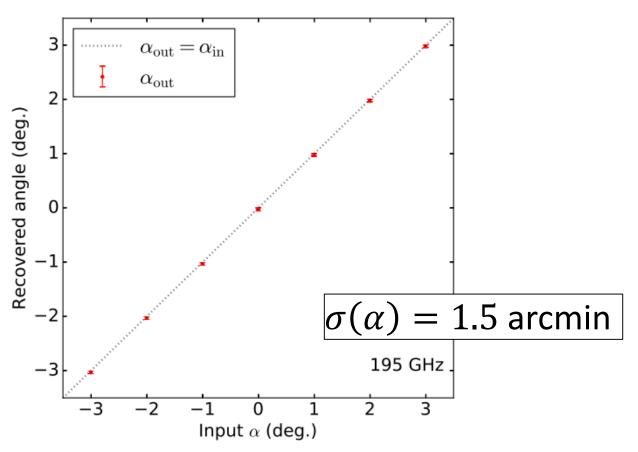


$$-2 \ln \mathcal{L} = \sum_{\ell=2}^{\ell_{\text{max}}} \frac{\left[C_{\ell}^{EB,\text{o}} - \frac{\tan(4\alpha)}{2} \left(C_{\ell}^{EE,\text{o}} - C_{\ell}^{BB,\text{o}} \right) - \frac{\sin(4\beta)}{2\cos(4\alpha)} \left(C_{\ell}^{EE,\text{CMB}} - C_{\ell}^{BB,\text{CMB}} \right) \right]^{2}}{\text{Var} \left(C_{\ell}^{EB,\text{o}} - \frac{\tan(4\alpha)}{2} \left(C_{\ell}^{EE,\text{o}} - C_{\ell}^{BB,\text{o}} \right) \right)}$$

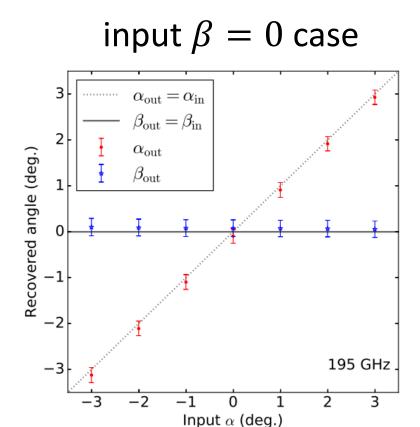
Minimise this log-likelihood to determine α and β

Before the simultaneous determination: α only case

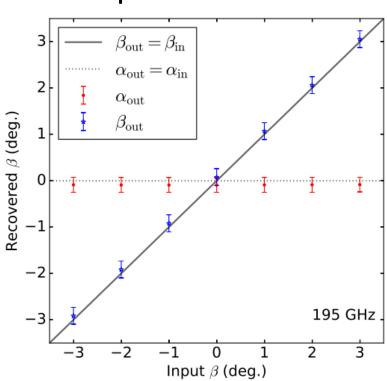
Assuming β =0, we can determine only α as the previous method



Simultaneous determination



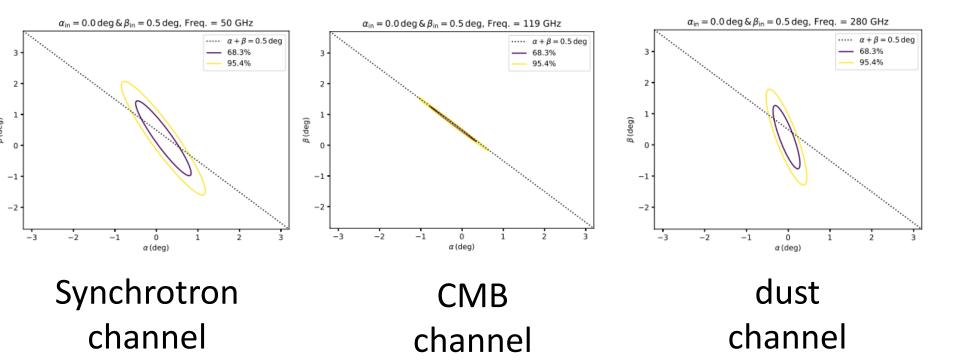
input $\alpha = 0$ case



$$\sigma(\alpha) = 9.6$$
 arcmin and $\sigma(\beta) = 11$ arcmin

Correlation between α and β

$$E_{\ell,m}^{\rm o} = E_{\ell,m}^{\rm fg} \cos(2\alpha) - B_{\ell,m}^{\rm fg} \sin(2\alpha) + E_{\ell,m}^{\rm CMB} \cos(2\alpha + 2\beta) - B_{\ell,m}^{\rm CMB} \sin(2\alpha + 2\beta)$$



- \triangleright CMB has a power to determine $\alpha+\beta$
- \triangleright FG has a power to determine α

Summary

There was a consensus in the CMB community that the measurement of the cosmic birefringence and the polarization angle calibration cannot be done simultaneously.

We have shown that this is not the case.

We can determine the birefringence angle of order 10 arcmin.

This is a great news!

Assumed specification

Table 1: Polarisation sensitivity and beam size of the LiteBIRD telescopes [27]

Frequency (GHz)	Polarisation Sensitivity ($\mu K'$)	Beam Size in FWHM (arcmin)
40	37.5	69
50	24.0	56
60	19.9	48
68	16.2	43
78	13.5	39
89	11.7	35
100	9.2	29
119	7.6	25
140	5.9	23
166	6.5	21
195	5.8	20
235	7.7	19
280	13.2	24
337	19.5	20
402	37.5	17

Foreground EB correlation

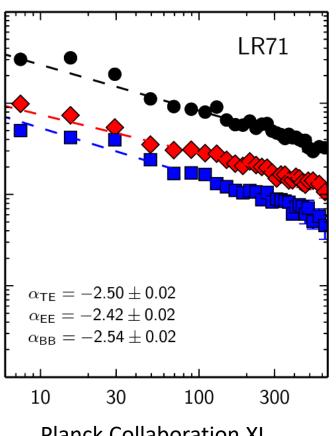
M. H. Abitbol, J. C. Hill, and B. R. Johnson, Mon. Not. Roy. Astron. Soc., 457(2), 1796–1803 (2016), 200 arXiv:1512.06834.

$$C_{\ell}^{dust,XY} = A^{XY} \left(\frac{\ell}{80}\right)^{-2.42} \tag{1}$$

$$C_{\ell,mult}^{dust,XY} = mC_{\ell}^{dust,XY} \tag{2}$$

$$C_{\ell,corr}^{dust,ZB} = f_c \sqrt{C_{\ell,mult}^{dust,ZZ} C_{\ell,mult}^{dust,BB}},$$
(3)

where A^{XY} is the best-fitting amplitude, m is a multiplicative factor, f_c is a correlation fraction, $X,Y \in \{T,E,B\}$ and $Z \in \{T,E\}$.



Planck Collaboration XI (2018), arXiv:1801.04945.

Foreground

$$\langle C_{\ell}^{EB,\mathrm{fg}} \rangle = \frac{f_c \sqrt{\xi}}{1 - \xi} \left(\langle C_{\ell}^{EE,\mathrm{fg}} \rangle - \langle C_{\ell}^{BB,\mathrm{fg}} \rangle \right). \longrightarrow \frac{\sin(4\gamma)}{2} \left(\left\langle C_{\ell}^{EE,fg} \right\rangle - \left\langle C_{\ell}^{BB,fg} \right\rangle \right)$$

$$E_{\ell,m}^{\text{o,fg}} = E_{\ell,m}^{\text{fg}} \cos(2\gamma) - B_{\ell,m}^{\text{fg}} \sin(2\gamma),$$

$$B_{\ell,m}^{\text{o,fg}} = E_{\ell,m}^{\text{fg}} \sin(2\gamma) + B_{\ell,m}^{\text{fg}} \cos(2\gamma),$$

Replace α -> α + γ β ->- γ in birefringence estimation

$$\left(C_{\ell}^{EB,o} - \frac{\tan(4\alpha)}{2} \left(C_{\ell}^{EE,o} - C_{\ell}^{BB,o} \right) \right) \cos(4\alpha + 4\gamma) =$$

$$\left(C_{\ell}^{EE,CMB} - C_{\ell}^{BB,CMB} \right) \sin(-4\gamma)/2$$

$$+ C_{\ell}^{EB,fg} + C_{\ell}^{EB,N} \cos(4\alpha + 4\gamma) + C_{\ell}^{EB,CMB} \cos(-4\gamma)$$

$$- \left(C_{\ell}^{EE,N} - C_{\ell}^{BB,N} \right) \sin(4\alpha + 4\gamma)/2$$

$$+ \left(C_{\ell}^{EfgB^{CMB}} + C_{\ell}^{E^{CMB}B^{fg}} \right) \cos(-2\gamma) + \left(C_{\ell}^{EfgE^{CMB}} - C_{\ell}^{BfgB^{CMB}} \right) \sin(-2\gamma)$$

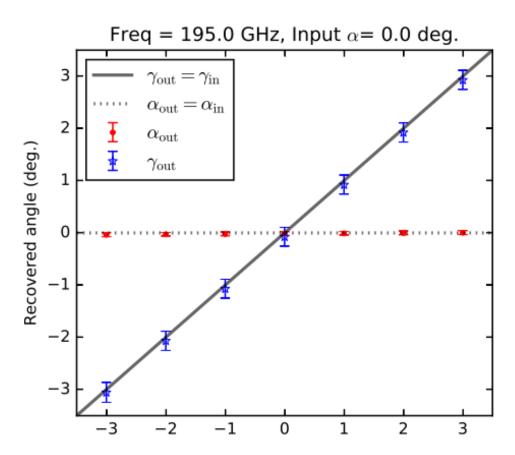
$$+ \left(C_{\ell}^{EfgB^{N}} + C_{\ell}^{BfgE^{N}} \right) \cos(2\alpha + 2\gamma) - \left(C_{\ell}^{EfgE^{N}} - C_{\ell}^{BfgB^{N}} \right) \sin(2\alpha + 2\gamma)$$

$$+ \left(C_{\ell}^{E^{CMB}B^{N}} + C_{\ell}^{B^{CMB}E^{N}} \right) \cos(2\alpha + 4\gamma)$$

$$- \left(C_{\ell}^{E^{CMB}E^{N}} - C_{\ell}^{B^{CMB}B^{N}} \right) \sin(2\alpha + 4\gamma)$$

Results of FG estimation

Replace α -> α + γ β ->- γ in birefringence estimation



Because β is the same in any frequency, we can still determine β even with FG EB correlation.

Variance

$$\begin{aligned} & \operatorname{Var} \left[C_{\ell}^{EB,o} - (C_{\ell}^{EE,o} - C_{\ell}^{BB,o}) \tan(4\alpha) / 2 \right] \\ &= \langle \left[C_{\ell}^{EB,o} - (C_{\ell}^{EE,o} - C_{\ell}^{BB,o}) \tan(4\alpha) / 2 \right]^{2} \rangle - \langle C_{\ell}^{EB,o} - (C_{\ell}^{EE,o} - C_{\ell}^{BB,o}) \tan(4\alpha) / 2 \rangle^{2} \\ &= \frac{1}{2\ell + 1} \langle C_{\ell}^{EE} \rangle \langle C_{\ell}^{BB} \rangle + \frac{\tan^{2}(4\alpha)}{4} \frac{2}{2\ell + 1} \left(\langle C_{\ell}^{EE} \rangle^{2} + \langle C_{\ell}^{BB} \rangle^{2} \right) \\ &- \tan(4\alpha) \frac{2}{2\ell + 1} \langle C_{\ell}^{EB} \rangle \left(\langle C_{\ell}^{EE} \rangle - \langle C_{\ell}^{BB} \rangle \right) + \frac{1}{2\ell + 1} \left(1 - \tan^{2}(4\alpha) \right) \langle C_{\ell}^{EB} \rangle^{2}. \end{aligned} \tag{A1}$$

$$\begin{aligned} &\operatorname{Var} \left[C_{\ell}^{EB,o} - \left(C_{\ell}^{EE,o} - C_{\ell}^{BB,o} \right) \tan(4\alpha) / 2 \right] \\ &\approx \frac{1}{2\ell + 1} C_{\ell}^{EE,o} C_{\ell}^{BB,o} + \frac{\tan^{2}(4\alpha)}{4} \frac{2}{2\ell + 1} \left[\left(C_{\ell}^{EE,o} \right)^{2} + \left(C_{\ell}^{BB,o} \right)^{2} \right] \\ &- \tan(4\alpha) \frac{2}{2\ell + 1} C_{\ell}^{EB,o} \left(C_{\ell}^{EE,o} - C_{\ell}^{BB,o} \right) \right]. \end{aligned}$$