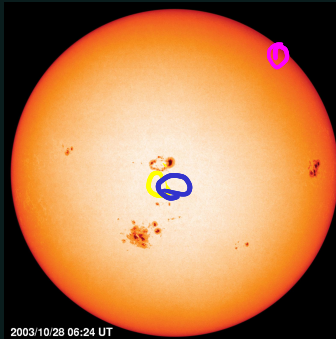


Lecture-4.

Three pictures of the Sun

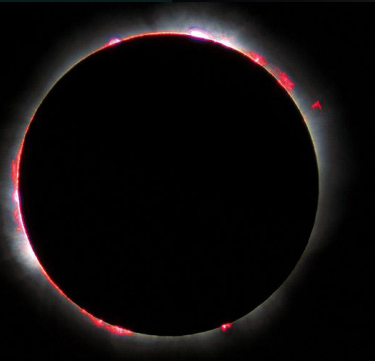
(1) Image of the sun (through a filter).

(2) & (3) are taken during a solar eclipse



(1)

↓
Photosphere



(2)

Chromosphere
(Photosphere is blocked)



(3)

Corona

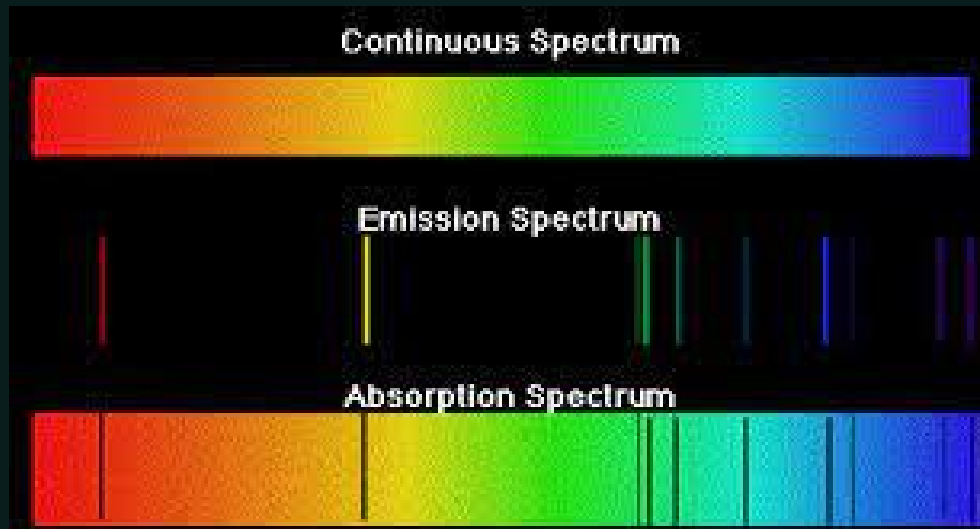
Both the photosphere
& the chromosphere are
eclipsed.

Corona is visible

In the image (1) of photosphere, the central portion is bright while the outskirts are fainter / darker. This phenomenon is called Limb-darkening.

The Spectrum from the Sun

Something about the spectrum



(1)

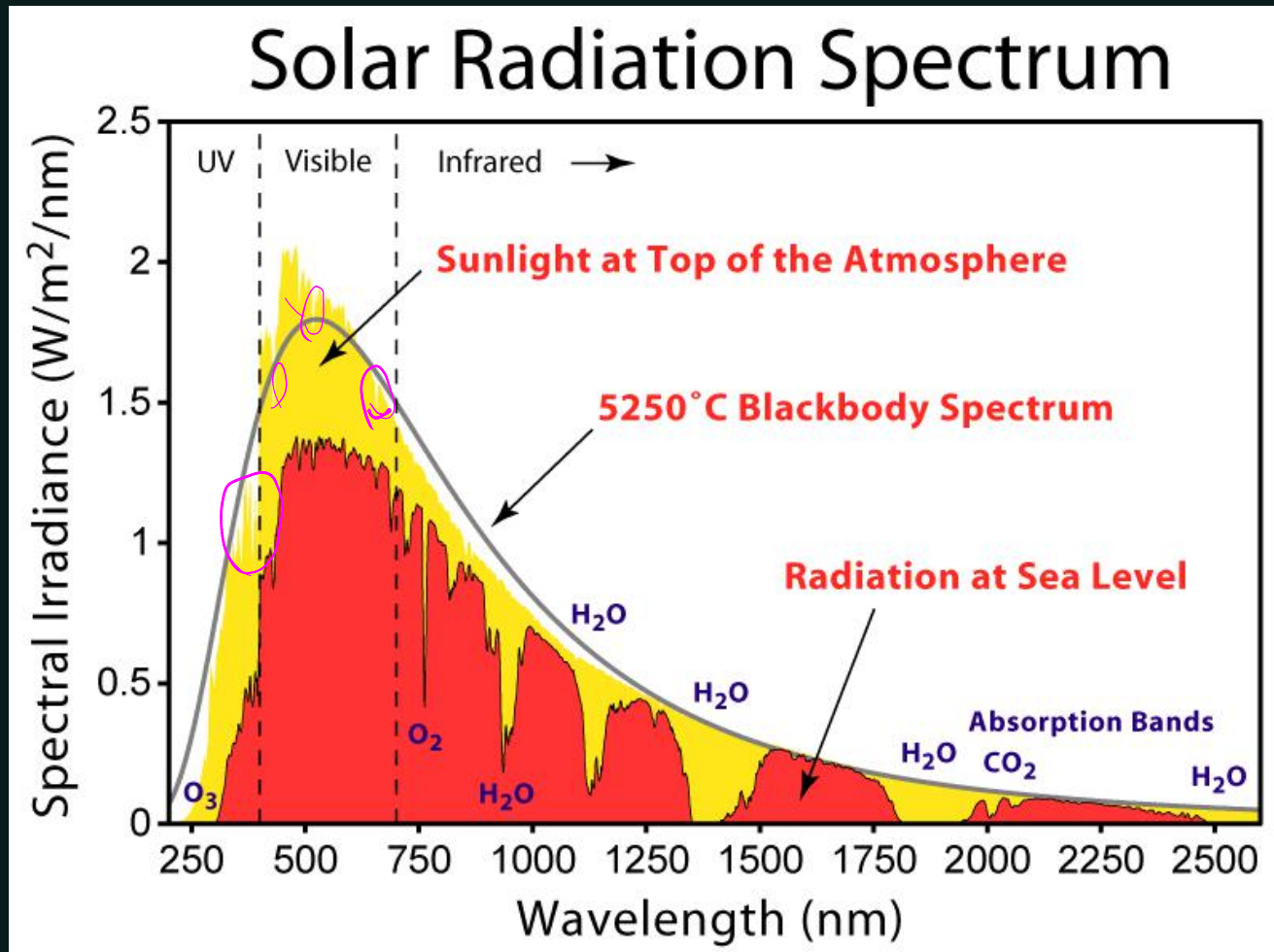
(2) outskirts of the solar disc

(3) → from the centre of the solar disc

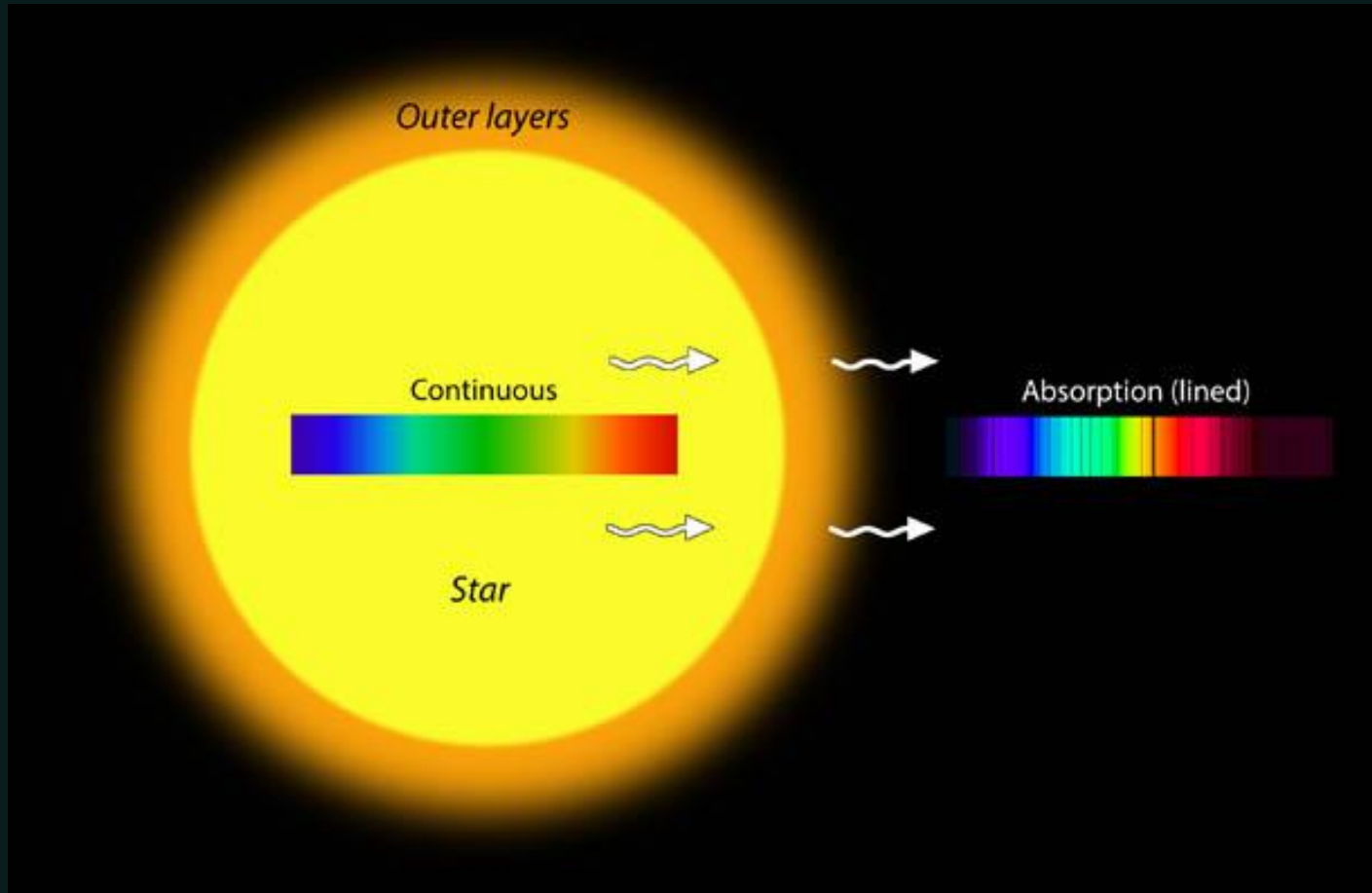
In (2) only a few frequencies/wavelengths are visible
(which are emitted by the source)

In (3) most of the photons are there except a few that are missing / absorbed.

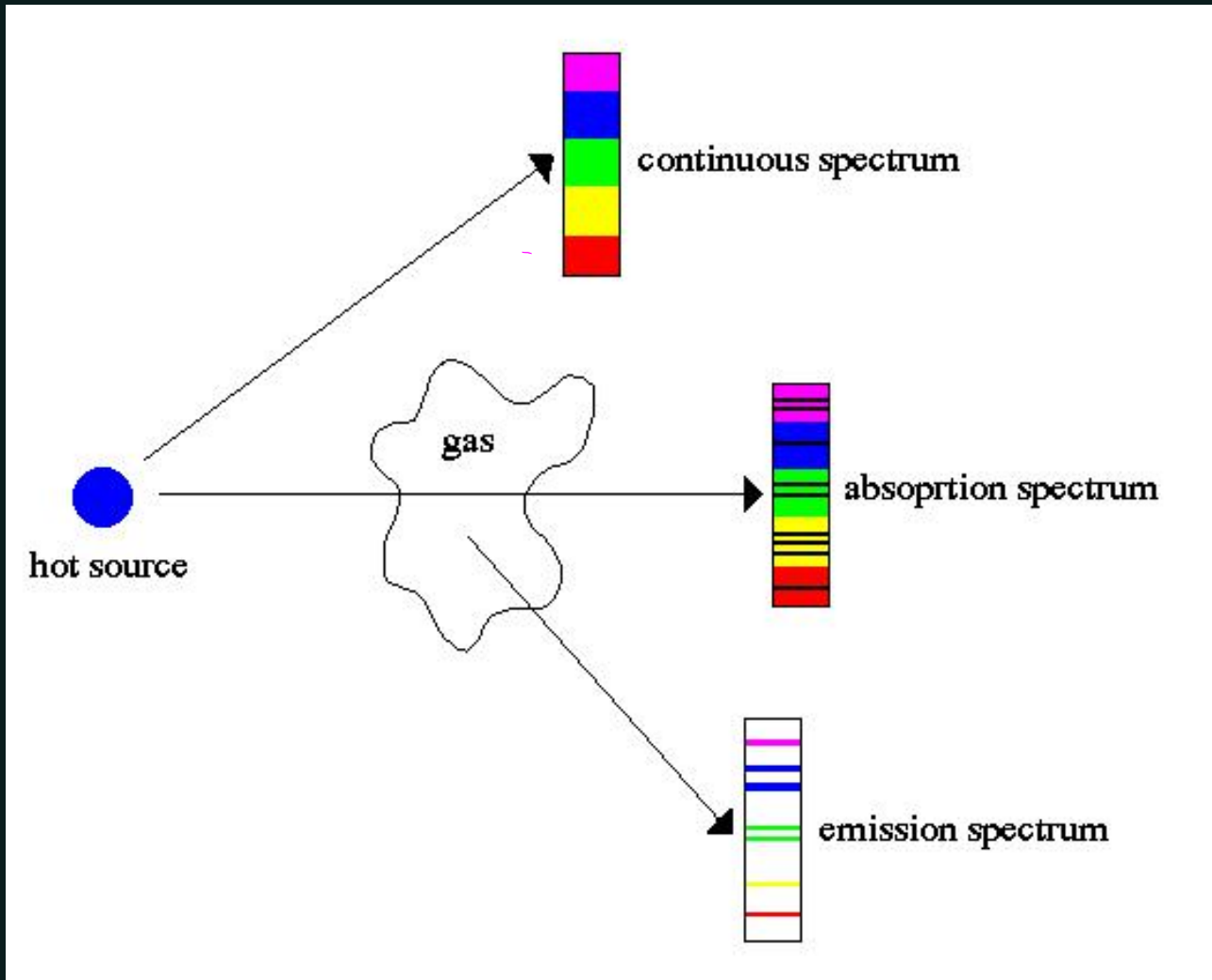
Sun's spectrum is like a Blackbody spectrum



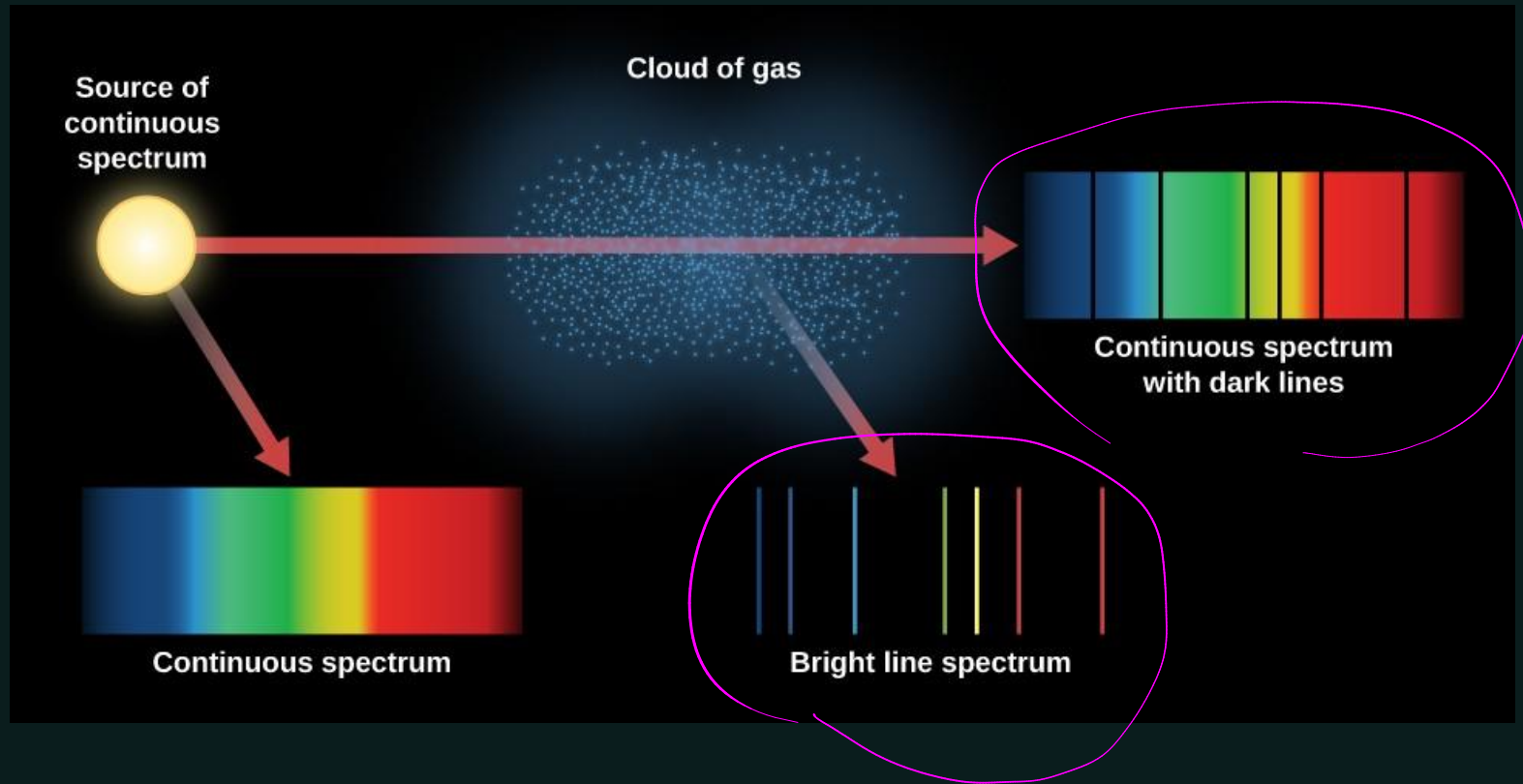
But there are differences



Kirchoff law



Kirchoff Law: spectrum from the Sun

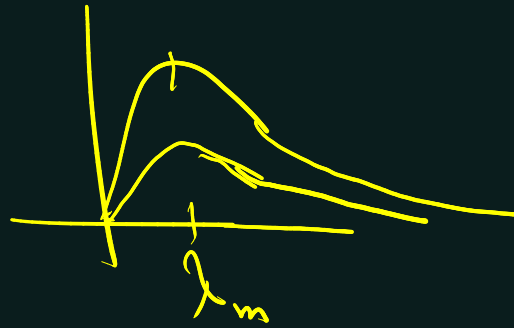


Let's try to describe all this with equations



In the interior of
the sun the photons
of specific wavelengths
are produced but
they get scattered ^{as}
remitted

— Their energies redistribute; they
come in equilibrium with the matter
and all sort of wavelength are present



$$\lambda_m T = \text{constant}$$

$$T = \frac{\text{constant}}{\lambda_m}$$

just by extracting the sample of radiation/photons that are present in a black body (in which the matter & radiation are in equilibrium)

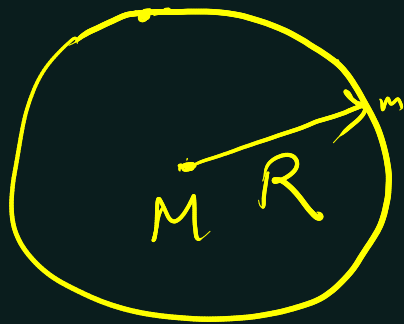
∴ can determine the temperature

Temperature thus measured for the sun is

$$T \approx 5000 \text{ to } 6000^\circ\text{C}$$



Let's do consider the particles in the
Sun.



Sun is a collection of particles under the influence of gravity. Radius = R
 say total mass = M

$$\frac{1}{r} = \frac{GMm^2}{L^2} (1 + \epsilon \cos \theta)$$

Assume that the particle is doing a circular motion $\epsilon = 0$

$$\frac{1}{r} = \frac{GMm^2}{L^2}$$

We want to find out the K.E., P.E.,
and the total energy

Circular orbit

$$\epsilon = \sqrt{1 + \frac{2EL^2}{G^2 M^2 m^3}} = 0$$

$$\frac{2EL^2}{G^2 M^2 m^3} = -1$$

$$\Rightarrow E = - \frac{G^2 M^2 m^3}{2 L^2}$$

$$= - \frac{1}{2} \frac{G M m}{r} \left(\frac{G M m^3}{L^2} \right) r$$

$$E = -\frac{1}{2} \frac{GMm}{r} \left(\frac{GMm^2}{L^2} \right)^{1/2}$$

But $E=0 \Rightarrow \frac{1}{r} = \frac{GMm^2}{L^2} (1 + \cos\theta) = \frac{GMm^2}{L^2}$

$$\Rightarrow \boxed{E = -\frac{1}{2} \frac{GMm}{r}}$$

$$E = K.E + P.E$$

$$P.E = -\frac{GMm}{r}$$

$$E = -\frac{1}{2} \frac{GMm}{r} = \frac{1}{2} \left(-\frac{GMm}{r} \right)$$

$$E = \frac{1}{2} V = \frac{V}{2}$$

$$K.E. = E - V = \frac{V}{2} - V = -\frac{V}{2}$$

$$K.E. = -\frac{1}{2} P.E.$$

$$K.E. = \frac{1}{2} P.E.$$

Virial theorem

If you repeat for elliptical orbit then it is doesn't hold true in the exact energies. But it still holds true if we define averaged over time period

$$\left\{ \begin{array}{l} \langle K.E \rangle_{\text{time period}} = - \frac{1}{2} \langle P.E \rangle \\ E = \frac{1}{2} \langle P.E \rangle \end{array} \right\}$$

Virial theorem

This virial theorem also holds true
for a collection of particles say sun

$$\langle K.E. \rangle = -\frac{1}{2} \langle P.E. \rangle$$

~~~~~  
of all the  
particles in  
the sun

↓  
of all the particles in  
the sun

↓  
But what is this for the gas in the sun.





$$P.E. \quad \sum \frac{G m_i m_j}{r_{ij}}$$

$$P.E. \quad \approx \quad -\frac{3}{5} \frac{G M M}{R}$$

Total P.E.  
of sun

$$\approx \quad -\left(\frac{3}{5}\right) \frac{G M^2}{R}$$

---

Each particle is not moving freely; they are colliding

Net momentum change for a particle colliding  
of the walls of tiny cubical volume in the sun

$$\Delta p = 2mv$$

By  $N$  atoms colliding  $\Delta p_{\text{total}} = 2m v N$

time taken  $\Delta t = \frac{2L}{v}$

Rate of change of momentum

$\Downarrow$   
Force

$$\frac{\Delta p_{\text{total}}}{\Delta t} = \frac{2m v N}{2L/v}$$

$$\text{Force exerted on the wall} = \frac{2m v^2 N}{2L}$$

$$\text{Pressure exerted on the wall} = \frac{F}{A} = \frac{2m v^2 N}{2L \cdot L^2}$$

$$\frac{N}{L^3} = n \quad \text{--- number density}$$

$$P = m v^2 n$$

$$P = m (v_x^2 + v_y^2 + v_z^2) n \approx \frac{1}{2} m v^2 n$$

$$\underline{P = n k_B T = \underbrace{\frac{1}{2} m v^2 n}_{\text{K.E. of all the particles in the sun}}}$$

K.E. of all the  
particles in the sun

$$k_B T \propto \text{K.E.}$$

Kinetic theory  
of  
Gases.

Total K.E. of the sun is roughly  $k_B T N$

Virial Theorem.

$$R \approx 7 \times 10^8 \text{ m}$$

$$K.E. = -\frac{1}{2} P.E.$$

$$Nk_B T = -\frac{1}{2} \left( -\frac{3}{5} \frac{GM^2}{R} \right)$$

$$T = \frac{3}{10} \frac{GM^2}{R N k_B} \rightarrow 1.38 \times 10^{-23}$$

$$N = \frac{2M}{m_p}$$

$$M = 2 \times 10^{30} \text{ kg}$$

$$m_p = 1.67 \times 10^{-27} \text{ kg}$$

$$T = \frac{3}{10} \frac{GM^2}{RNk_B} = \frac{3}{10} \frac{GM^2 m_p}{R \cancel{2M} k_B}$$

$$T = \frac{3}{20} \frac{GM m_p}{k_B R}$$

$$M = 2 \times 10^{30} \text{ kg}$$

$$R = 7 \times 10^8 \text{ m}$$

$$k_B = 1.38 \times 10^{-23}$$

$$m_p = 1.67 \times 10^{-27} \text{ kg}$$

$$T \approx 3 \times 10^6$$

(more than a million K)

So the actual temperature is more than a million kelvin.

{ But we found from the Black-body spectrum calculation, that the temperature is about 5000 K. }

↓↓

This is the surface temperature.

## Problems:

- (1) Prove the virial theorem for the elliptical orbit: (Hint: you will have to calculate the K.E. & P.E. averaged over time period)
- (2) Prove that the total gravitational potential energy of sun is  $-\frac{3GM^2}{5R}$  where  $M$  is the mass of sun and  $R$  the radius



(3) Use the virial theorem argument to calculate the temperature of the sun if

(I) it expands to 10 times its current size

(II) it shrinks to a radius of 10 km

(4) We found that the interior of sun is very hot ( $> \text{million kelvin}$ ) but the surface is not that hot (only around  $5000 \text{ K}$ ). So there must be temperature gradient in successive shells. Using this information; try to explain Limb darkening in words.

