

# Non-Homogeneous PDEs

## Time independent non-homogeneous BC

### Dirichlet BC

- $u_t = \alpha^2 u_{xx}$
- $u(0, t) = u_0, \quad u(L, t) = u_1$
- $u(x, 0) = f(x)$
- $u(x, t) = v(x, t) + \bar{u}(x, t)$

### Steady state problem

- Solution is independent of time
- $u_t = \alpha^2 u_{xx}$
- $u(0, t) = u_0, \quad u(L, t) = u_1$
- $u(x, 0) = f(x)$
- **Step 1** find the steady state solution
- $u_t = u_{xx} = 0 \implies u_\infty(x) = Ax + B$ , the solution in the long run
- $u_\infty(0) = B = u_0 \quad u_\infty(L) = AL + u_0 = u_1$
- $u_\infty(x) = \frac{u_1 - u_0}{L}x + u_0$
- $u(x, t) = v(x, t) + u_\infty(x)$
- $u_t = v_t$
- $u_{xx} = v_{xx} + 0 \implies v_t = v_{xx} \implies \text{PDE}$

### Boundary and initial conditions

- $u(0, t) = v(0, t) + u_\infty(0)$
- $u_0 = v(0, t) + u_0 \implies v(0, t) = 0$
- $u(L, t) = v(L, t) + u_\infty(L)$
- $u_1 = v(L, t) + u_1 \implies v(L, t) = 0$
- $u(x, 0) = v(x, 0) + u_\infty(x)$
- $v(x, 0) = f(x) - u_\infty(x)$

### Solve the PDE

- As done in **separation of variables**
- $v(x, t) = \sum_{n=1}^{\infty} b_n \sin\left(\frac{n\pi x}{L}\right) e^{-a^2\left(\frac{n\pi}{L}\right)^2 t}$

- $b_n = \frac{1}{L} \int_{-L}^L [f(x) - u_\infty(x)] \sin\left(\frac{n\pi x}{L}\right)$
- $u(x, t) = u_\infty(x) + v(x, t)$

## Types of Particular solutions

### Dirichlet

- $u_\infty(x) = Ax + B$

### Mixed type 1/2

- $u_\infty(x) + Ax + B$

### Neumann

- $u_p(x, t) = Ax^2 + Bx + Ct$  particular solution
- $u_t = u_{xx}$
- $u_x(0, t) = q_0, \quad u_x(L, t) = q_1$
- Assume we have a steady state problem
- $0 = u_{xx} \implies u_\infty(x) = Ax + B$
- $u'_\infty(x) = A$
- $u'_\infty(0) = q_0 = A$
- $u'_\infty(L) = q_1 = A$
- This is only possible if  $q_0 = q_1 = A$
- Since Neumann condition does not allow steady state unless  $q_0 = q_1$  let

$$u_p(x, t) = Ax^2 + Bx + Ct$$

- $\frac{\partial u_p}{\partial t} = c$
- $\frac{\partial^2 u_p}{\partial x^2} = 2A$
- So then we must have that  $C = 2A$
- $u_{px}(x, t) = 2Ax + B$
- $u_{px}(0, t) = B = q_0$
- $u_{px}(L, t) = 2AL + B = q_1 \implies A = \frac{q_1 - q_0}{2L}$
- So now we can write  $u_p(x, t) = \frac{q_1 - q_0}{2L} x^2 + q_0 x + \frac{q_1 - q_0}{L} t$

## Eigen Function Expansion

- Given  $u(x, t) = v(x, t) + u_p(x, t)$
- Let  $v(x, t) = \sum_{n=0}^{\infty} V_n(t) b_n \sin\left(\frac{n\pi x}{L}\right)$
- $F(x, t) = \sum_{n=0}^{\infty} S_n(t) \sin\left(\frac{n\pi x}{L}\right)$  this is the source term
- Our aim is to get the coefficients

- $v_t = \sum_{n=0}^{\infty} V'_n(t) \sin\left(\frac{n\pi x}{L}\right)$
- $v_{xx} = \sum_{n=0}^{\infty} -\left(\frac{n\pi}{L}\right)^2 V_n(t) \sin\left(\frac{n\pi x}{L}\right)$
- Remember the solution is  $v_t = v_{xx} + F(x, t) \implies v_t - v_{xx} - F(x, t) = 0$
- So then we have the following
- $\sum_{n=0}^{\infty} \left[ V'_n + \left(\frac{n\pi}{L}\right)^2 V_n - S_n(t) \right] \sin\left(\frac{n\pi x}{L}\right) = 0$
- $S_n(t)$  is known
- $V'_n(t) + \left(\frac{n\pi}{L}\right)^2 V_n(t) = S_n(t)$  which is a first order ode

## Examples

### Example 1

- $u_t = u_{xx}$
- $u_x(0, t) = \ln t \quad u_x(1, t) = t^2$
- $u(x, 0) = f(x)$
- **Find a particular solution**
- $u_p(x, t) = \frac{A(t)}{2}x^2 + B(t)x + C(t)t$
- $\frac{\partial u_p}{\partial x} = A(t)x + B(t)$
- $u_x(0, t) = \ln t = u_{px} = B(t)$
- $u_x(1, t) = t^2 = A(t) + \ln t$
- $A(t) = t^2 - \ln t$
- $u_p(x, t) = \frac{t^2 - \ln t}{2}x^2 + \ln(t)x$  we are ignoring  $C(t)$  since this matches the boundary conditions
- $u(x, t) = v(x, t) = u_p(x, t)$
- $v_t + \frac{(2t - \frac{1}{t})}{2}x^2 + \frac{x}{t} = v_{xx} + t^2 - \ln t$
- $v_t = v_{xx} + t^2 - \ln t - \frac{x}{t} - \frac{2t - \frac{1}{t}}{2}x^2$
- $v_x(0, t) = v_x(1, t) = 0$
- $v(x, 0) = f(x) - u_{\infty}(x, 0)$
- This doesn't work for  $\ln t$  lmao we need to change it to something like  $\ln(t + 1)$

### Example 2

- $u_t = 16u_{xx} + \cos\left(\frac{7\pi x}{2}\right) \quad 0 < x < 2$
- $u_x(0, t) = 1 = u_x(2, t)$
- $u(x, 0) = x^2 - 4$
- $u(x, t) = v(x, t) + u_p(x, t)$
- $u_p(x, t) = \frac{A(t)}{2}x^2 + B(t)x + C(t)$
- $\frac{\partial u_p}{\partial x} = A(t)x + B(t)$
- $u_{px}(0, t) = B = 1$

- $u_{p_x}(2, t) = 1 = 2A + 1 \implies A = 0$
- $u_p(x, t) = x$
- $u(x, t) = v(x, t) + x$ , now plug these into the pde
- $v_t = 16v_{xx} + \cos\left(\frac{7\pi x}{2}\right)$
- $v_x(0, t) = v_x(2, t) = 0$
- $u(x, 0) = v(x, 0) + u_p(x, 0) \implies v(x, 0) = x^2 - 4 - x$
- Let  $v(x, t) = \frac{V_0}{2} + \sum_{n=1}^{\infty} V_n \cos\left(\frac{n\pi x}{2}\right)$
- For a neumann  $\lambda_n \in \left\{0, \left(\frac{n\pi}{L}\right)^2\right\}$
- $\cos\left(\frac{7n\pi}{2}\right) = \frac{S_0}{2} + \sum S_n(t) \cos\left(\frac{n\pi x}{2}\right)$
- $S_n = \frac{2}{L} \int_0^L f(x) \cos\left(\frac{n\pi x}{2}\right)$
- $S_n = \delta_{n7}$
- $v_t = \frac{V_0}{2} + \sum_{n=1}^{\infty} V'_n \cos\left(\frac{7\pi x}{2}\right)$
- $v_{xx} = \sum_{n=1}^{\infty} -\left(\frac{n\pi}{2}\right)^2 V_n \cos\left(\frac{n\pi x}{2}\right)$
- $v_t - 16v_{xx} - \cos\left(\frac{7\pi x}{2}\right) = 0$
- $\left[\frac{V_0}{2} - 16[0] - 0\right] + \sum_{n=1}^{\infty} \left[V'_n + 16\left(\frac{n\pi}{2}\right)^2 V_n = \delta_{n7}\right] \cos\left(\frac{n\pi x}{2}\right)$
- $\frac{\partial V_0}{\partial t} = 0$  (I)
- $V_n + \left(\frac{4n\pi}{2}\right)^2 V_n - \delta_{n7} = 0$  (2)
- From (I) we have that  $V_0 = C_0$
- From (2) we have  $V_{n7} + (2\pi n)^2 V_n = \delta_{n7}$
- Here we must use integrating factor,  $r = e^{\int (2n\pi)^2 dt} = e^{(2n\pi)^2 t}$
- $V_n = e^{-(2n\pi)^2 t} \left[ \int_0^t e^{(2n\pi)^2 t} \delta_{n7} dz + C \right]$
- $V_n = e^{-(2n\pi)^2 t} \left[ \frac{1}{(2n\pi)^2} e^{(2n\pi)^2 t} \delta_{n7} + C \right]$
- $V_n(t) = \frac{1}{(2n\pi)^2} S_{n7} + C e^{-(2n\pi)^2 t}$
- $v(x, t) = \frac{c_0}{2} + \sum_{n=1}^{\infty} \left[ \frac{1}{(2n\pi)^2} \delta_{n7} + C_n e^{-(2n\pi)^2 t} \right] \cos\left(\frac{n\pi x}{2}\right)$
- $v(x, 0) = \frac{c_0}{2} + \sum_{n=1}^{\infty} \left[ \frac{\delta_{n7}}{(2n\pi)^2} + C_n \right] \cos\left(\frac{n\pi x}{2}\right)$
- $a_n = \frac{1}{L} \int_{-L}^L f(x) \cos\left(\frac{n\pi x}{L}\right) dx, \quad L = 2$
- $c_n = a_n - \frac{\delta_{n7}}{(2n\pi)^2}$

### Example 3

- $u_t = u_{xx} + e^{-t} \sin(x), \quad 0 < x < \frac{\pi}{2}$
- $u_0(0, t) = 0, \quad u_x\left(\frac{\pi}{2}, t\right) = e^{-t}$
- $u(x, 0) = x$
- We have mixed boundary conditions
- $\lambda_n - \mu_n^2 \implies \mu_n = 2n + 1$
- $X_n(x) = \sin((2n + 1)x)$
- $u_p(x, t) = A(t)x + B(t)$
- $u_p(0, t) = B(t) = 0$

- $u_{p_x}(x, t) = A(t) \implies A(t) = e^{-t}$
- $v_t = e^{-t}x = v_{xx} + 0 + e^{-t} \sin(x)$
- $v(t) = v_{xx} + e^{-t}[x + \sin x]$
- $u(x, 0) = v(x, 0) + u_p(x, 0)$
- $x = v(x, 0) + x \implies v(x, 0) = 0$
- $v(x, t) = \sum_{n=0}^{\infty} V_n(t) + \sin((2n+1)x)$
- $e^{-t}[x + \sin x] = \sum_{n=0}^{\infty} S_n(t) \sin((2n+1)x)$
- So we know that  $S_n(t) = \frac{2}{\pi} \int_0^{\pi/2} [x + \sin x] \sin((2n+1)x) dx$
- $\int_0^{\pi/2} x \sin[(2n+1)x] dx$
- $dv = \sin[(2n+1)x]$
- $v = -\frac{1}{2n+1} \cos[(2n+1)x]$
- $u = x, \quad du = 1$
- $\int_0^{\pi/2} x \sin[(2n+1)x] dx = -\frac{x}{2n+1} \cos((2n+1)x) \Big|_0^{\pi/2} + \frac{1}{2n+1} + \frac{1}{2n+1} \int_0^{\pi/2} \cos((2n+1)x) dx$
- This gives us  $0 + \frac{(-1)^n}{(2n+1)^2} + \frac{\pi}{4}$
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