

参考:

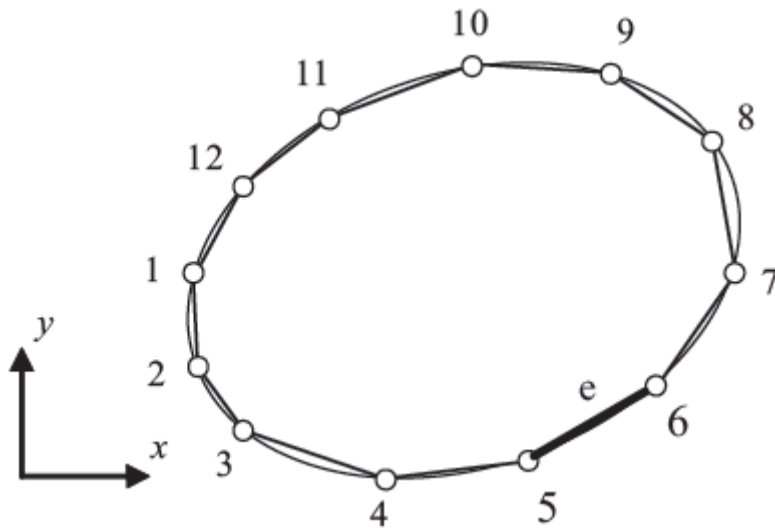
Boundary Element Method

IPCST

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1-D Boundary Elements

- 1-D boundary elements
 - To describe a boundary of a 2-D domain



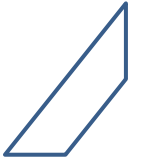
- ✓ Essentially same with 1-D finite elements except their usage
 - Shape functions
 - Gaussian quadrature

$$\mathbf{x} = \sum_i N_i \mathbf{x}_i^e$$

Figure from Beer



2-D or 3-D Elements



- 2-D boundary elements
 - To describe boundaries of 3-D domains or to evaluate 2-D volume integrals
- 3-D cells
 - To evaluate 3-D volume integrals

Iso-parametric Elements

- Same shape functions for coordinates (element shape) and for physical quantity interpolation.
- The physical quantities can be discontinuous at element interfaces.

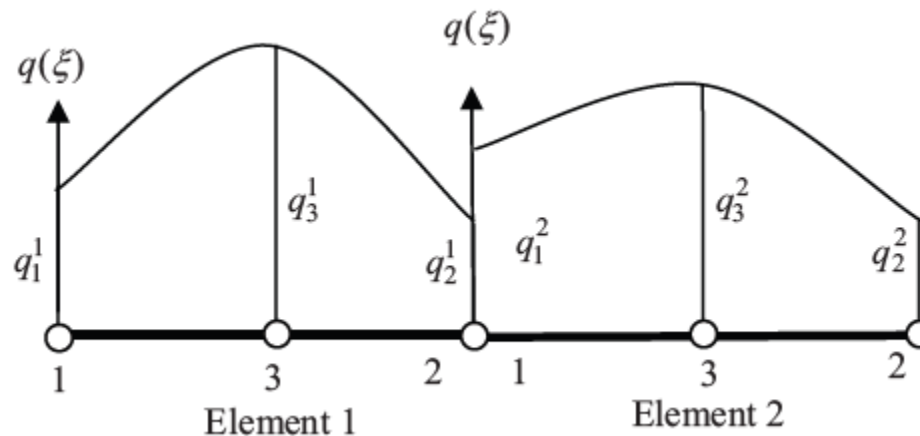


Figure from Beer

Differential Geometry

- Tangential vector

$$\mathbf{v}_\xi = \frac{\partial \mathbf{x}}{\partial \xi} = \sum_i \frac{\partial N_i}{\partial \xi} \mathbf{x}_i^e$$

- Normal vector

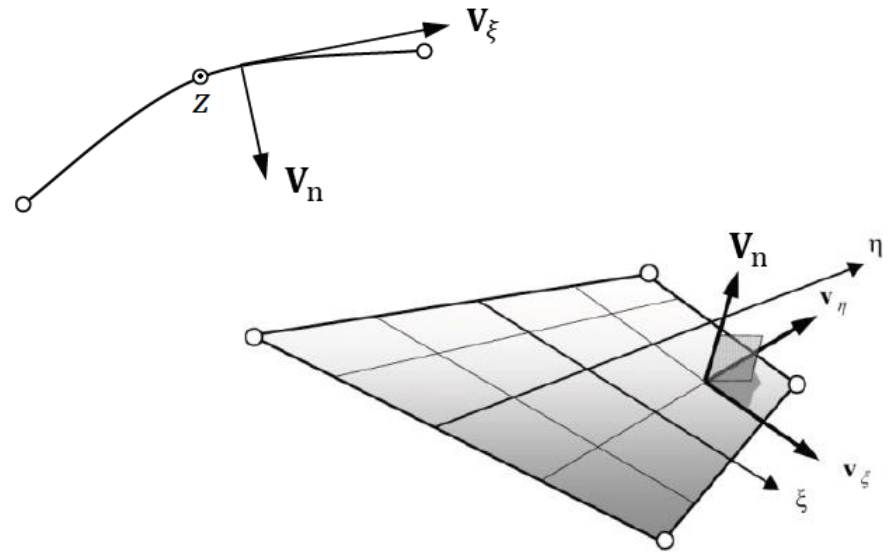
– Line element

$$\mathbf{V}_n = \mathbf{V}_\xi \times \mathbf{n}_z = \left(\frac{dy}{d\xi}, -\frac{dx}{d\xi}, 0 \right)^T$$

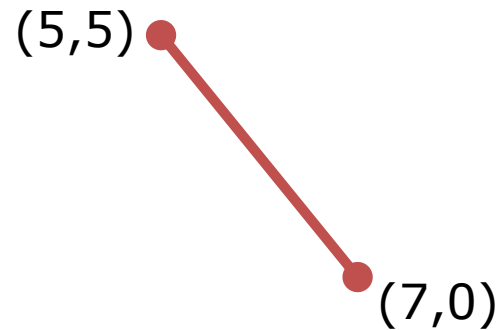
– Surface element

$$\mathbf{V}_n = \mathbf{V}_\xi \times \mathbf{V}_\eta = \left(\frac{\partial y}{\partial \xi} \frac{\partial z}{\partial \eta} - \frac{\partial y}{\partial \eta} \frac{\partial z}{\partial \xi}, \frac{\partial z}{\partial \xi} \frac{\partial x}{\partial \eta} - \frac{\partial z}{\partial \eta} \frac{\partial x}{\partial \xi}, \frac{\partial x}{\partial \xi} \frac{\partial y}{\partial \eta} - \frac{\partial x}{\partial \eta} \frac{\partial y}{\partial \xi} \right)^T$$

$$\diamond V_n = |\mathbf{V}_n| = J \text{ (Jacobian)}$$



Do It Yourself



- Imagine that you should apply BEM and take the above line segment as a (1-D) bar element. Find the relation between the actual coordinates and the natural coordinates of a point on this element.
- Calculate the Jacobian.



Singular Integrals



- Weakly singular integrals
 - Ex.) $\int_0^1 \ln|x - \alpha| dx$ where $0 < \alpha < 1$
$$= \lim_{\varepsilon \rightarrow 0} \int_0^{\alpha - \varepsilon} \ln|x - \alpha| dx + \lim_{\varepsilon \rightarrow 0} \int_{\alpha + \varepsilon}^1 \ln|x - \alpha| dx$$
 - Integrands
 - 2-D: functions of order $\log r$
 - 3-D: functions of order $1/r$
 - These integrals can be handled by the Gaussian quadrature.



Singular Integrals



- Divergent or strongly singular integrals
 - Ex.) $\int_0^1 \frac{1}{x-\alpha} dx$ where $0 < \alpha < 1$
$$= \lim_{\varepsilon \rightarrow 0} \int_0^{\alpha-\varepsilon} \frac{1}{x-\alpha} dx + \lim_{\varepsilon \rightarrow 0} \int_{\alpha+\varepsilon}^1 \frac{1}{x-\alpha} dx$$
 - Integrands
 - 2-D: functions of order $1/r$
 - 3-D: functions of order $1/r^2$
 - These integrals are evaluated as *Cauchy principal values*.
 - Determination of Cauchy principal values is tricky. Sometimes they can be replaced by other equations.



Basic Ideas of BEM



- Use functions satisfying the PDE to approximate the solution
- Then actually, only boundary conditions are needed to be approximated
- Iso-parametric elements + numerical integration



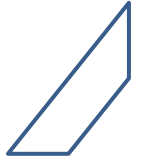
Pros & Cons



- Pros
 - Efficiency \leftarrow relatively small number of nodes
 - Semi-analytical \rightarrow good accuracy
 - Visualization-friendly (ex.: exact contours)
 - Good for extreme systems
 - Ex.) concentrated stress, infinite domain
- Cons
 - Inefficiency \leftarrow handling BEM matrices
 - Fundamental solution is required \rightarrow Rarely applicable to inhomogeneous or non-linear systems



Comparison to FEM



BEM

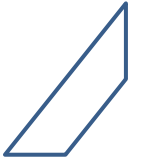
- Discretizing boundary
- Small input data
- Small full matrices
- Better for infinite or semi-infinite domains
- Fundamental solution: required
- Hardly applicable to inhomogeneous or non-linear PDEs

FEM

- Discretizing entire domain
- Rare singularities
- Large sparse matrix
- Better for finite domains
- Fundamental solution: not required
- Applicable to any PDE systems



Application



- Solid Mechanics
 - Fracture, contact
- Acoustics
- Soil-structure, mining & tunnelling
- Fluid mechanics, solid-fluid interaction
- Electromagnetics

Laplace Equation

- Green's second identity

$$\int_D (\phi \Delta \psi - \psi \Delta \phi) dV = \oint_{\partial D} \left(\phi \frac{\partial \psi}{\partial \mathbf{n}} - \psi \frac{\partial \phi}{\partial \mathbf{n}} \right) dS$$

- \mathbf{n} : outward surface normal vector

- Integral equation for Laplace equation

$$\begin{aligned} & \int_D (\varphi(\mathbf{r}) \Delta \psi(\mathbf{p}, \mathbf{r}) - \psi(\mathbf{p}, \mathbf{r}) \Delta \varphi(\mathbf{r})) dV_{\mathbf{r}} \\ &= \oint_{\partial D} \left(\varphi(\mathbf{q}) \frac{\partial \psi(\mathbf{p}, \mathbf{q})}{\partial \mathbf{n}} - \psi(\mathbf{p}, \mathbf{q}) \frac{\partial \varphi(\mathbf{q})}{\partial \mathbf{n}} \right) dS_{\mathbf{q}} \end{aligned}$$

$$\mathbf{p} \in D \cup \partial D; \mathbf{q} \in \partial D; \mathbf{r} \in D$$

$$\Delta \varphi(\mathbf{r}) = 0 \quad \Delta \psi(\mathbf{p}, \mathbf{r}) = -\delta(\mathbf{p} - \mathbf{r})$$

ψ : auxiliary function



Laplace Equation



- Green function
 - Solution of $L\psi(\mathbf{p}, \mathbf{r}) = \delta(\mathbf{p} - \mathbf{r})$
 - L: linear differential operator
 - It is difficult to find ψ if L is non-linear, anisotropic, or inhomogeneous.
 - Green functions for $\Delta\psi(\mathbf{p}, \mathbf{r}) = -\delta(\mathbf{p} - \mathbf{r})$
 - 2-D: $\psi(\mathbf{p}, \mathbf{r}) = -\frac{1}{2\pi} \ln|\mathbf{p} - \mathbf{r}|$
 - 3-D: $\psi(\mathbf{p}, \mathbf{r}) = \frac{1}{4\pi|\mathbf{p} - \mathbf{r}|}$

Laplace Equation

- Boundary integral equation

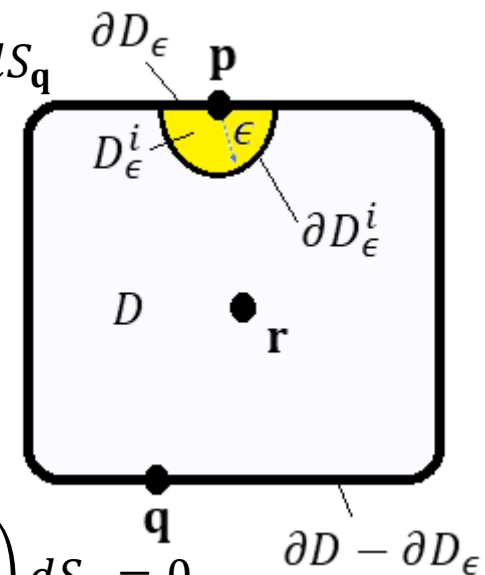
$$\begin{aligned}
 -C(\mathbf{p})\varphi(\mathbf{p}) &= \lim_{\epsilon \rightarrow 0} \int_{\partial D - \partial D_\epsilon} \left(\varphi(\mathbf{q}) \frac{\partial \psi(\mathbf{p}, \mathbf{q})}{\partial \mathbf{n}} - \psi(\mathbf{p}, \mathbf{q}) \frac{\partial \varphi(\mathbf{q})}{\partial \mathbf{n}} \right) dS_{\mathbf{q}} \\
 &+ \lim_{\epsilon \rightarrow 0} \int_{\partial D_\epsilon^i} \left(\varphi(\mathbf{q}) \frac{\partial \psi(\mathbf{p}, \mathbf{q})}{\partial \mathbf{n}} - \psi(\mathbf{p}, \mathbf{q}) \frac{\partial \varphi(\mathbf{q})}{\partial \mathbf{n}} \right) dS_{\mathbf{q}} \\
 &\mathbf{p} \in D \cup \partial D_\epsilon; \mathbf{q} \in \partial D - \partial D_\epsilon \text{ or } \partial D_\epsilon^i
 \end{aligned}$$



$$C(\mathbf{p})\varphi(\mathbf{p}) + \lim_{\epsilon \rightarrow 0} \int_{\partial D - \partial D_\epsilon} \left(\varphi(\mathbf{q}) \frac{\partial \psi(\mathbf{p}, \mathbf{q})}{\partial \mathbf{n}} - \psi(\mathbf{p}, \mathbf{q}) \frac{\partial \varphi(\mathbf{q})}{\partial \mathbf{n}} \right) dS_{\mathbf{q}} = 0$$

$$\mathbf{p} \in D \cup \partial D_\epsilon; \mathbf{q} \in \partial D$$

Geometric coefficient:
 $\frac{1}{2}$ for smooth boundary
 1 for inner domain





Laplace Equation



- Simple discretization

$$\varphi_k = \varphi(\mathbf{p}_k)$$

$$\sum_{k=1}^N H(j, k) \varphi_k = \sum_{k=1}^N G(j, k) \gamma_k$$

$$H(j, k) = \oint_{\partial D_k} \frac{\partial \psi}{\partial \mathbf{n}}(\mathbf{p}_j, \mathbf{q}) dS_{\mathbf{q}} + C(\mathbf{p}_j) \delta_{j,k}$$

$$G(j, k) = \oint_{\partial D_k} \psi(\mathbf{p}_j, \mathbf{q}) dS_{\mathbf{q}}, \quad \gamma_k = \frac{\partial \varphi}{\partial \mathbf{n}}(\mathbf{p}_k)$$

→ $\mathbf{H}\mathbf{h} = \mathbf{G}\mathbf{g}$ form

Laplace Equation

- Discretization with shape functions

$$\varphi(\mathbf{q}) \rightarrow \sum N_k(\mathbf{q})\varphi_k^e, \quad \frac{\partial \varphi(\mathbf{q})}{\partial \mathbf{n}} \rightarrow \sum N_k(\mathbf{q})\gamma_k^e$$

$$\lim_{\epsilon \rightarrow 0} \int_{\partial D - \partial D_\epsilon} \rightarrow \sum_e \int_{\partial D_e}$$

$$\Rightarrow C(\mathbf{p}_j)\varphi(\mathbf{p}_j) + \sum_e \sum_k \varphi_k^e \int_{\partial D_e} N_k(\mathbf{q}) \frac{\partial \psi}{\partial \mathbf{n}}(\mathbf{p}_j, \mathbf{q}) dS_{\mathbf{q}}$$

$$- \sum_e \sum_k \gamma_k^e \int_{\partial D_e} N_k(\mathbf{q}) \psi(\mathbf{p}_j, \mathbf{q}) dS_{\mathbf{q}} = 0$$

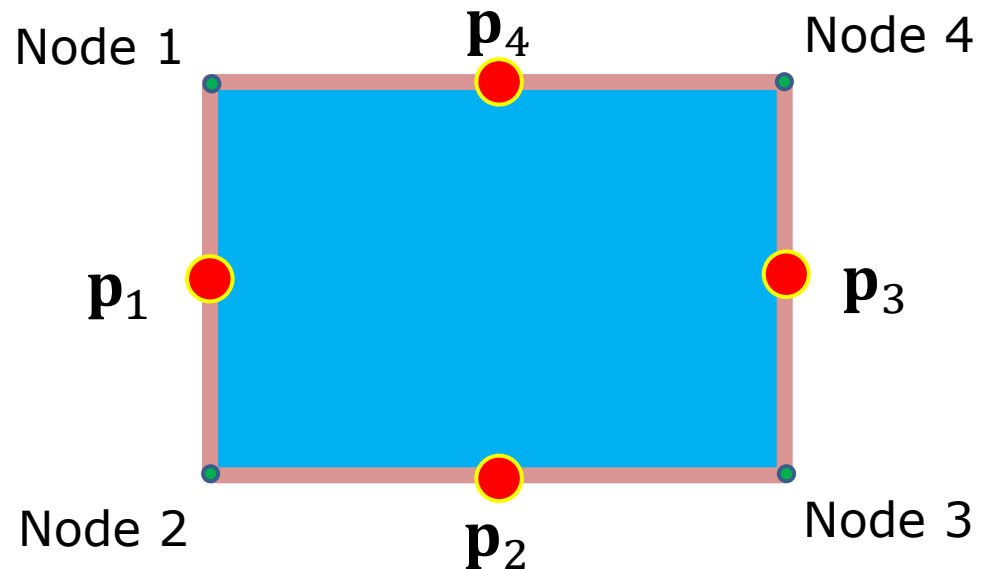
✓ $\varphi(\mathbf{p}_j) = \sum N_k(\mathbf{p}_j)\varphi_k^e$. $\varphi(\mathbf{p}_j) = \varphi_k^e$ if $\mathbf{p}_j = \mathbf{p}_k$ (k th node)

✓ \mathbf{p}_j 's are called collocation points.

Discretization Example

- 2-D rectangle
 - Index j : collocation points ($\mathbf{p}_1 \sim \mathbf{p}_4$)
 - Index k : nodes

✓ # of nodes
= # of shape functions





Numerical Integration



- Gaussian quadrature
 - Usually integration of polynomials.
 - But in BEM, you must integrate $\log r$, r^{-1} , r^{-2} , or r^{-3} terms. \rightarrow often approximated as high order polynomials
 - From the formula of Strout and Secrest (Gaussian Quadrature Formulas, 1966), the upper bound of the integration error for r^{-1} is

$$\varepsilon \leq \frac{4}{(4r/L)^{2N}}$$

- L : element length, N : number of Gauss points.
- Note that $r = |\mathbf{p}_j - \mathbf{q}|$.

Numerical Integration

- Gaussian quadrature
 - The required number of Gauss points depends on R/L where R is the distance between \mathbf{p}_j and its closest element boundary point when \mathbf{p}_j is out of the element.
- ❖ **Minimum R/L for integration order 4 & 5**

N	R/L		
	$O(r^{-1})$	$O(r^{-2})$	$O(r^{-3})$
3	1.4025	2.3187	3.4170
4	0.6735	0.9709	1.2908

- U. Eberwien, C. Duenser, & W. Moser, *Engineering Analysis with Boundary Elements* **29**, 447 (2005).



Numerical Integration



- Gaussian quadrature
 - If the integrand has singularity of $\log r$, you should evaluate the integration analytically or use the Gauss-Laguerre type quadrature.

$$\int_0^1 f(\bar{\xi}) \log \frac{1}{\bar{\xi}} d\bar{\xi} \rightarrow \sum w_i f(\bar{\xi}_i)$$

- Note: $0 < \bar{\xi} < 1$
 - Cf.) $-1 < \xi < 1$ in the Gauss-Legendre type

Numerical Integration

- If \mathbf{p}_j is not located in the integration region, the integrals can be evaluated by using Gaussian quadrature.
 - If \mathbf{p}_j is close to the integration region, subdivision is required to increase R/L . (This is not increasing the number of elements. Don't confuse.)

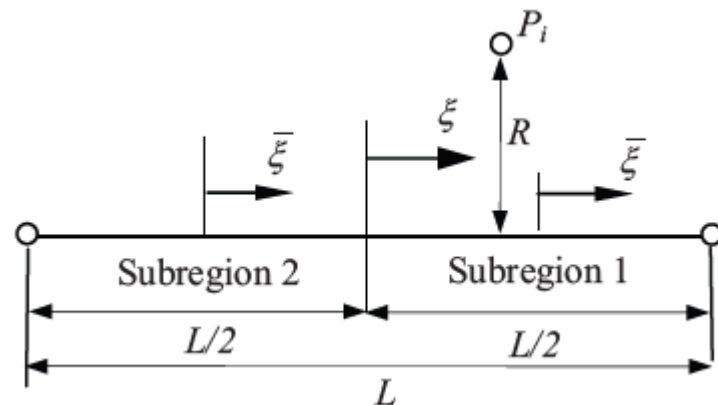
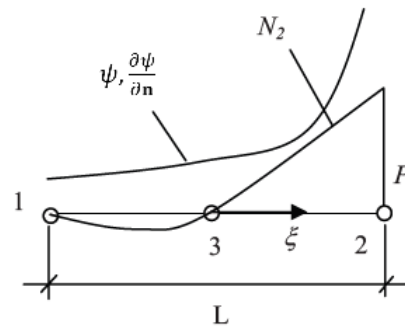
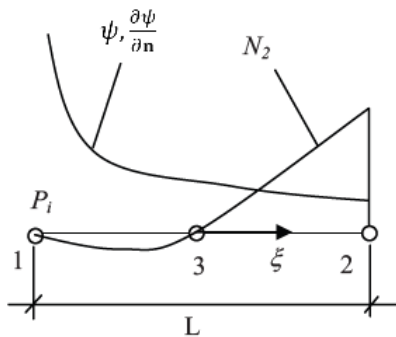


Figure from Beer

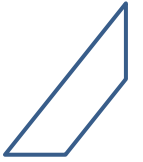
Numerical Integration

- If \mathbf{p}_j is in the integration region, the functions ψ & $\frac{\partial \psi}{\partial \mathbf{n}}$ usually tend to diverge.
 - If the integral is weakly singular or N_k is zero at \mathbf{p}_j , still you can use Gaussian quadrature.
 - Otherwise, you can't choose but use techniques for strongly singular integrals.





BEM Procedure



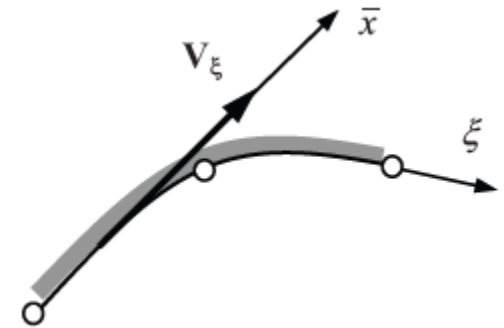
1. Data input & initialization
2. Determination of connectivity of nodes & elements
3. Determination of element boundary conditions
4. Local element equation & integration
5. Assembly & solving
6. Postprocessing

Postprocessing

- Boundary results
 - From the result you've got $\varphi = \sum N_k \varphi_k^e$, you can get its derivatives along the boundary.
 - Remember the tangential vector $\mathbf{v}_\xi = \frac{\partial \mathbf{x}}{\partial \xi} = \sum_i \frac{\partial N_i}{\partial \xi} \mathbf{x}_i^e$

$$\frac{\partial \varphi}{\partial \bar{x}} = \sum_i \frac{\partial N_i}{\partial \xi} \frac{\partial \xi}{\partial \bar{x}} \varphi_k^e$$

- $\frac{\partial \xi}{\partial \bar{x}}$: inverse Jacobian. For 2-D,
$$\frac{\partial \xi}{\partial \bar{x}} = (V_{\xi,x}^2 + V_{\xi,y}^2)^{-1/2}$$





Postprocessing



- Internal domain results

- Remember

$\varphi(\mathbf{p})$

$$= \sum_e \sum_k \gamma_k^e \int_{\partial D_e} N_k(\mathbf{q}) \psi(\mathbf{p}, \mathbf{q}) dS_{\mathbf{q}} - \sum_e \sum_k \varphi_k^e \int_{\partial D_e} N_k(\mathbf{q}) \frac{\partial \psi}{\partial \mathbf{n}}(\mathbf{p}, \mathbf{q}) dS_{\mathbf{q}}$$

- Derivatives at \mathbf{p}

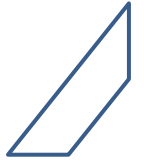
- X-direction, for example,

$\frac{\partial \varphi}{\partial x}(\mathbf{p})$

$$= \sum_e \sum_k \gamma_k^e \int_{\partial D_e} N_k(\mathbf{q}) \frac{\partial \psi}{\partial x}(\mathbf{p}, \mathbf{q}) dS_{\mathbf{q}} - \sum_e \sum_k \varphi_k^e \int_{\partial D_e} N_k(\mathbf{q}) \frac{\partial}{\partial x} \frac{\partial \psi}{\partial \mathbf{n}}(\mathbf{p}, \mathbf{q}) dS_{\mathbf{q}}$$



Corners & Edges



$$C(\mathbf{p})\varphi(\mathbf{p}) = \sum_e \sum_k \gamma_k^e \int_{\partial D_e} N_k(\mathbf{q}) \psi(\mathbf{p}, \mathbf{q}) dS_{\mathbf{q}} - \sum_e \sum_k \varphi_k^e \int_{\partial D_e} N_k(\mathbf{q}) \frac{\partial \psi}{\partial \mathbf{n}}(\mathbf{p}, \mathbf{q}) dS_{\mathbf{q}}$$

- $C(\mathbf{p}) = 1 - \frac{\alpha}{2\pi}$ for sharp corners at \mathbf{p}
 - where α is the interior angle of the corner
- Discontinuous Elements
 - All interpolation nodes are located inside the element.
 - This type of elements are used to avoid indefinite $\frac{\partial \psi}{\partial \mathbf{n}}$ values.

Corners & Edges

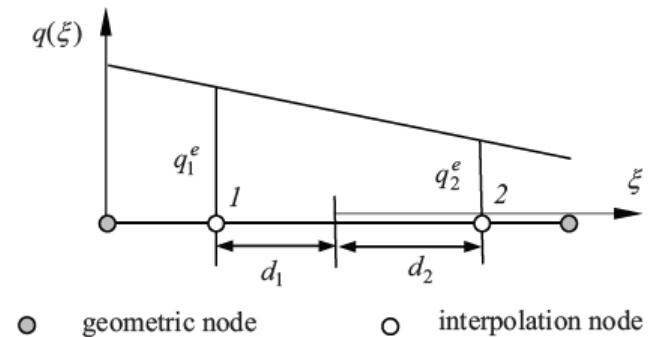
- Discontinuous Elements
 - 1-D linear shape functions

$$N_1(\xi) = \frac{d_1 - \xi}{d_1 + d_2}, N_2(\xi) = \frac{d_2 - \xi}{d_1 + d_2}$$

- $\xi \neq \pm 1$ at interpolation nodes

- Integration

- Same as for continuous elements if \mathbf{p}_j is not located at one of interpolation nodes
- You should divide the integration region at \mathbf{p}_j if \mathbf{p}_j is located at an interpolation node.





References



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- F. París & J. Cañas, Boundary Element Method
- H. Antes, "A Short Course on Boundary Element Methods"
- Wikipedia



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