# Chapter 1

## **Mathematical Preliminaries**

## 1.1 Complexity Theory

Definition of QMA, languages, etcetera.

#### 1.2 Various Mathematical Lemmas

#### 1.2.1 Truncation Lemma

**Lemma 1** (Truncation Lemma). Let H be a Hamiltonian acting on a Hilbertspace  $\mathcal{H}$  and let  $|\Phi\rangle \in \mathcal{H}$  be a normalized state. Let  $\mathcal{K}$  be a subspace of  $\mathcal{H}$ , let P be the projector onto  $\mathcal{K}$ , and let  $\tilde{H} = PHP$  be the Hamiltonian within this subspace. Suppose that, for some T > 0,  $W \in \{H, \tilde{H}\}$ ,  $N_0 \in \mathfrak{h}$ , and  $\delta > 0$ , we have, for all  $0 \le t \le T$ ,

$$e^{-iWt}|\Phi\rangle = |\gamma(t)\rangle + |\epsilon(t)\rangle \text{ with } ||\epsilon(t)\rangle|| \leq \delta$$

and

$$(1-P)H^r|\gamma(t)\rangle = 0 \text{ for all } r \in \{0, 1, \dots, N_0 - 1\}.$$

Then, for all  $0 \le t \le T$ ,

$$\left\| \left( e^{-iHt} - e^{-i\tilde{H}t} \right) |\Phi\rangle \right\| \le \left( \frac{4e\|H\|t}{N_0} + 2 \right) \left( \delta + 2^{-N_0} (1+\delta) \right).$$

**Proposition 1.** Let H be a Hamiltonian acting on a Hilbert space  $\mathcal{H}$ , and let  $|\Phi\rangle \in \mathcal{H}$  be a normalized state. Let  $\mathcal{K}$  be a subspace of  $\mathcal{H}$  such that there exists an  $N_0 \in \natural$  so that for all  $|\alpha\rangle \in \mathcal{K}^{\perp}$  and for all  $n \in \{0, 1, 2, ..., N_0 - 1\}$ ,  $\langle \alpha | H^n | \Phi \rangle = 0$ . Let P be the projector onto  $\mathcal{K}$  and let  $\tilde{H} = PHP$  be the Hamiltonian within this subspace. Then

$$||e^{-it\tilde{H}}|\Phi\rangle - e^{-itH}|\Phi\rangle|| \le 2\left(\frac{e||H||t}{N_0}\right)^{N_0}.$$

*Proof.* Define  $|\Phi(t)\rangle$  and  $|\tilde{\Phi}(t)\rangle$  as

$$|\Phi(t)\rangle = e^{-itH}|\Phi\rangle = \sum_{k=0}^{\infty} \frac{(-it)^k}{k!} H^k |\Phi\rangle \qquad |\tilde{\Phi}(t)\rangle = e^{-it\tilde{H}}|\Phi\rangle = \sum_{k=0}^{\infty} \frac{(-it)^k}{k!} \tilde{H}^k |\Phi\rangle.$$

Note that by assumption,  $\tilde{H}^k|\Phi\rangle = H^k|\Phi\rangle$  for all  $k < N_0$ , and thus the first  $N_0$  terms in the two above sums are equal. Looking at the difference between these two states, we have

$$\begin{aligned} |||\Phi(t)\rangle - |\tilde{\Phi}(t)\rangle|| &= \left\| \sum_{k=0}^{\infty} \frac{(-it)^k}{k!} \left( H^k - \tilde{H}^k \right) |\Phi\rangle \right\| \\ &= \left\| \sum_{k=0}^{N_0 - 1} \frac{(-it)^k}{k!} \left( H^k - \tilde{H}^k \right) |\Phi\rangle - \sum_{k=N_0}^{\infty} \frac{(-it)^k}{k!} \left( H^k - \tilde{H}^k \right) |\Phi\rangle \right\| \\ &\leq \sum_{k=N_0}^{\infty} \frac{t^k}{k!} \left( ||H||^k + ||\tilde{H}||^k \right) \\ &\leq 2 \sum_{k=N_0}^{\infty} \frac{t^k}{k!} ||H||^k \end{aligned}$$

where the last step uses the fact that  $\|\tilde{H}\| \leq \|P\| \|H\| \|P\| = \|H\|$ . Thus for any  $c \geq 1$ , we have

$$\||\Phi(t)\rangle - |\tilde{\Phi}(t)\rangle\| \le \frac{2}{c^{N_0}} \sum_{k=N_0}^{\infty} \frac{(ct)^k}{k!} \|H\|^k$$
  
 $\le \frac{2}{c^{N_0}} \exp(ct\|H\|).$ 

We obtain the best bound by choosing  $c = N_0/||Ht||$ , which gives

$$\||\Phi(t)\rangle - |\tilde{\Phi}(t)\rangle\| \le 2\left(\frac{e\|H\|t}{N_0}\right)^{N_0}$$

as claimed. (If c < 1 then the bound is trivial.)

**Proposition 2.** Let  $U_1, \ldots, U_n$  and  $V_1, \ldots, V_n$  be unitary operators. Then for any  $|\psi\rangle$ ,

$$\left\| \left( \prod_{i=n}^{1} U_{i} - \prod_{i=n}^{1} V_{i} \right) | \psi \rangle \right\| \leq \sum_{j=1}^{n} \left\| (U_{j} - V_{j}) \prod_{i=j-1}^{1} U_{i} | \psi \rangle \right\|. \tag{1.1}$$

*Proof.* The proof is by induction on n. The case n = 1 is obvious. For the induction step, we have

$$\left\| \left( \prod_{i=n}^{1} U_{i} - \prod_{i=n}^{1} V_{i} \right) |\psi\rangle \right\| = \left\| \left( \prod_{i=n}^{1} U_{i} - V_{n} \prod_{i=n-1}^{1} U_{i} + V_{n} \prod_{i=n-1}^{1} U_{i} - \prod_{i=n}^{1} V_{i} \right) |\psi\rangle \right\|$$
(1.2)

$$\leq \left\| (U_n - V_n) \prod_{i=n-1}^1 U_i |\psi\rangle \right\| + \left\| \left( \prod_{i=n-1}^1 U_i - \prod_{i=n-1}^1 V_i \right) |\psi\rangle \right\| \quad (1.3)$$

$$\leq \sum_{j=1}^{n} \left\| (U_j - V_j) \prod_{i=j-1}^{1} U_i |\psi\rangle \right\| \tag{1.4}$$

where the last step uses the induction hypothesis.

Proof of Lemma 1. For  $M \in \natural$  write

$$\begin{split} \|(e^{-iHt} - e^{-i\tilde{H}t})|\Phi\rangle\| &= \left\| \left( \left( e^{-iH\frac{t}{M}} \right)^M - \left( e^{-i\tilde{H}\frac{t}{M}} \right)^M \right) |\Phi\rangle \right\| \\ &\leq \sum_{j=1}^M \left\| \left( e^{-iH\frac{t}{M}} - e^{-i\tilde{H}\frac{t}{M}} \right) e^{-iW(j-1)\frac{t}{M}} |\Phi\rangle \right\| \\ &\leq \sum_{j=1}^M \left\| \left( e^{-iH\frac{t}{M}} - e^{-i\tilde{H}\frac{t}{M}} \right) \left( |\gamma(\frac{(j-1)t}{M})\rangle + |\epsilon(\frac{(j-1)t}{M})\rangle \right) \right\| \\ &\leq 2M\delta + \sum_{j=1}^M \left\| \left( e^{-iH\frac{t}{M}} - e^{-i\tilde{H}\frac{t}{M}} \right) \frac{|\gamma(\frac{(j-1)t}{M})\rangle}{\left\| |\gamma(\frac{(j-1)t}{M})\rangle \right\|} \right\| \left\| |\gamma(\frac{(j-1)t}{M})\rangle \right\| \\ &\leq 2M\delta + 2M \left( \frac{e\|H\|t}{MN_0} \right)^{N_0} (1+\delta) \end{split}$$

where in the second line we have used Proposition ?? and in the last step we have used Proposition ?? and the fact that  $||\gamma(t)\rangle|| \le 1 + \delta$ . Now, for some  $\eta > 1$ , choose

$$M = \left\lceil \frac{\eta e \|H\|t}{N_0} \right\rceil$$

for  $0 < t \le T$  to get

$$\|(e^{-iHt} - e^{-i\tilde{H}t})|\Phi\rangle\| \le 2M \left(\delta + \eta^{-N_0}(1+\delta)\right) \le 2\left(\frac{\eta e\|H\|t}{N_0} + 1\right) \left(\delta + \eta^{-N_0}(1+\delta)\right).$$

The choice  $\eta = 2$  gives the stated conclusion.

Note that it would be slightly better to take a smaller value of  $\eta$ . However, this does not significantly improve the final result; the above bound is simpler and sufficient for our purposes.

### 1.2.2 Nullspace Projection Lemma

**Lemma 2** (Nullspace Projection Lemma). Let  $H_A$  and  $H_B$  be positive semi-definite matrices. Suppose that the nullspace, S, of  $H_A$  is nonempty, and that

$$\gamma(H_B|_S) \ge c > 0$$
 and  $\gamma(H_A) \ge d > 0.$  (1.5)

Then,

$$\gamma(H_A + H_B) \ge \frac{cd}{d + ||H_B||}. (1.6)$$

*Proof.* Let  $|\psi\rangle$  be a normalized state satisfying

$$\langle \psi | H_A + H_B | \psi \rangle = \gamma (H_A + H_B). \tag{1.7}$$

Let  $\Pi_S$  be the projector onto the nullspace of  $H_A$ . First suppose that  $\Pi_S|\psi\rangle = 0$ , in which case

$$\langle \psi | H_A + H_B | \psi \rangle \ge \langle \psi | H_A | \psi \rangle \ge \gamma(H_A)$$
 (1.8)

and the result follows. On the other hand, if  $\Pi_S|\psi\rangle\neq 0$  then we can write

$$|\psi\rangle = \alpha |a\rangle + \beta |a^{\perp}\rangle \tag{1.9}$$

with  $|\alpha|^2 + |\beta|^2 = 1$ ,  $\alpha \neq 0$ , and two normalized states  $|a\rangle$  and  $|a^{\perp}\rangle$  such that  $|a\rangle \in S$  and  $|a^{\perp}\rangle \in S^{\perp}$ . (If  $\beta = 0$  then we may choose  $|a^{\perp}\rangle$  to be an arbitrary state in  $S^{\perp}$  but in the following we fix one specific choice for concreteness.) Note that any state  $|\phi\rangle$  in the nullspace of  $H_A + H_B$  satisfies  $H_A |\phi\rangle = 0$  and hence  $\langle \phi | a^{\perp} \rangle = 0$ . Since  $\langle \phi | \psi \rangle = 0$  and  $\alpha \neq 0$  we also see that  $\langle \phi | a \rangle = 0$ . Hence any state

$$|f(q,r)\rangle = q|a\rangle + r|a^{\perp}\rangle$$
 (1.10)

is orthogonal to the nullspace of  $H_A + H_B$ , and

$$\gamma(H_A + H_B) = \min_{|q|^2 + |r|^2 = 1} \langle f(q, r) | H_A + H_B | f(q, r) \rangle. \tag{1.11}$$

Within the subspace Q spanned by  $|a\rangle$  and  $|a^{\perp}\rangle$ , note that

$$H_A|_Q = \begin{pmatrix} w & v^* \\ v & z \end{pmatrix} \qquad H_B|_Q = \begin{pmatrix} 0 & 0 \\ 0 & y \end{pmatrix} \tag{1.12}$$

where  $w = \langle a|H_B|a\rangle$ ,  $v = \langle a^{\perp}|H_B|a\rangle$ ,  $y = \langle a^{\perp}|H_A|a^{\perp}\rangle$ , and  $z = \langle a^{\perp}|H_B|a^{\perp}\rangle$ , and that we are interested in the smaller eigenvalue of

$$M = H_A|_Q + H_B|_Q = \begin{pmatrix} w & v^* \\ v & y+z \end{pmatrix}.$$
 (1.13)

Letting  $\epsilon_+$  and  $\epsilon_-$  be the two eigenvalues of M with  $\epsilon_+ \geq \epsilon_-$ , note that

$$\epsilon_{+} = ||M|| \le ||H_A|_Q|| + ||H_B|_Q|| \le y + ||H_B|_Q|| \le y + ||H_B||,$$
 (1.14)

where we have used the Cauchy interlacing theorem to note that  $||H_B||_Q|| \leq ||H_B||$ . Additionally, we have that

$$\epsilon_{+}\epsilon_{-} = \det(M) = w(y+z) - |v|^{2} \ge wy \tag{1.15}$$

where we used the fact that  $H_B|_Q$  is positive-semidefinite. Putting this together, we have that

$$\gamma(H_A + H_B) = \min_{|q|^2 + |r|^2 = 1} \langle f(q, r) | H_A + H_B | f(q, r) \rangle = \epsilon_- \ge \frac{wy}{y + ||H_B||}.$$
 (1.16)

As the right hand side increased monotonically with both w and y, and as  $w \ge \gamma(H_B|_S) \ge c$  and  $y \ge \gamma(H_A) \ge d$ , we have

$$\gamma(H_A + H_B) \ge \frac{cd}{d + ||H_B||} \tag{1.17}$$

as required.  $\Box$