## Problema 1

## Parte 1

Le funzioni d'onda  $\psi_S(\vec{r}_1, \vec{r}_2)$  e  $\psi_T(\vec{r}_1, \vec{r}_2)$  sono reali.

$$\psi_S(\vec{r}_1, \vec{r}_2) = C_S(\psi_a(\vec{r}_1)\psi_b(\vec{r}_2) + \psi_b(\vec{r}_1)\psi_a(\vec{r}_2))$$
  
$$\psi_T(\vec{r}_1, \vec{r}_2) = C_T(\psi_a(\vec{r}_1)\psi_b(\vec{r}_2) - \psi_b(\vec{r}_1)\psi_a(\vec{r}_2))$$

$$\langle \psi_S | \psi_S \rangle = 1$$
  $\langle \psi_T | \psi_T \rangle = 1$ 

$$\int d^3r_1 d^3r_2 \psi_S^*(\vec{r}_1, \vec{r}_2) \psi_S(\vec{r}_1, \vec{r}_2) =$$
 
$$\int d^3r_1 d^3r_2 \psi_S(\vec{r}_1, \vec{r}_2) \psi_S(\vec{r}_1, \vec{r}_2) =$$
 
$$C_S^2 \int d^3r_1 d^3r_2 \left( \psi_a(\vec{r}_1) \psi_b(\vec{r}_2) + \psi_b(\vec{r}_1) \psi_a(\vec{r}_2) \right) \left( \psi_a(\vec{r}_1) \psi_b(\vec{r}_2) + \psi_b(\vec{r}_1) \psi_a(\vec{r}_2) \right) =$$
 
$$= C_S^2 \left( \int d^3r_1 d^3r_2 \psi_a^2(\vec{r}_1) \psi_b^2(\vec{r}_2) + \int d^3r_1 d^3r_2 \psi_b^2(\vec{r}_1) \psi_a^2(\vec{r}_2) + 2 \int d^3r_1 d^3r_2 \psi_a(\vec{r}_1) \psi_b(\vec{r}_1) \psi_a(\vec{r}_2) \psi_b(\vec{r}_2) \right)$$

Le funzioni (...) sono già normalizzate, quindi i primi due integrali danno 1.

$$C_S^2 \left( 2 + 2 \int d^3 r_1 \psi_a(\vec{r}_1) \psi_b(\vec{r}_1) \int d^3 r_2 \psi_a(\vec{r}_2) \psi_b(\vec{r}_2) \right) = C_S^2 \left( 2 + 2S^2 \right)$$

Dove abbiamo definito

$$S \equiv \int d^3r \psi_a(\vec{r}) \psi_b(\vec{r})$$

Per  $\psi_T$  i calcoli sono quasi identici, l'unica differenza è che il terzo integrale ha il meno, quindi si trova

$$\int d^3r_1 d^3r_2 \psi_T(\vec{r}_1, \vec{r}_2) \psi_T(\vec{r}_1, \vec{r}_2) = C_T^2 (2 - 2S^2)$$

$$C_S = \frac{1}{\sqrt{2(1+S^2)}}$$

$$C_T = \frac{1}{\sqrt{2(1-S^2)}}$$

Parte 2

$$S = \int d^3r \psi_a(\vec{r}) \psi_b(\vec{r}) = \int d^3r \left( \frac{1}{\sqrt{\pi a^3}} \exp\left( -\frac{|\vec{r} - \vec{r}_a|}{a} \right) \right) \left( \frac{1}{\sqrt{\pi a^3}} \exp\left( -\frac{|\vec{r} - \vec{r}_b|}{a} \right) \right) =$$

$$= \frac{1}{\pi a^3} \int d^3r \exp\left( -\frac{|\vec{r} - \vec{r}_a| + |\vec{r} - \vec{r}_b|}{a} \right)$$

$$\begin{cases} \xi = \frac{|\vec{r} - \vec{r}_a| + |\vec{r} - \vec{r}_b|}{R} \\ \eta = \frac{|\vec{r} - \vec{r}_a| - |\vec{r} - \vec{r}_b|}{R} \\ \phi = \text{prova} \end{cases}$$

dove  $R \equiv |\vec{r}_a - \vec{r}_b|$ 

Lo Jacobiano di questa trasformazione è  $\frac{R^3}{8}(\xi^2-\eta^2)$ 

$$\frac{1}{\pi a^3} \int d^3 r \exp\left(-\frac{|\vec{r} - \vec{r}_a| + |\vec{r} - \vec{r}_b|}{a}\right) = \frac{1}{\pi a^3} \int d^3 r \exp\left(-\frac{R}{a} \frac{|\vec{r} - \vec{r}_a| + |\vec{r} - \vec{r}_b|}{R}\right) =$$

$$= \frac{1}{\pi a^3} \int_0^{2\pi} d\phi \int_{-1}^1 d\eta \int_1^{\infty} d\xi \exp\left(-\frac{R}{a}\xi\right) \frac{R^3}{8} (\xi^2 - \eta^2) =$$

$$= \frac{R^3}{8\pi a^3} \left(\int_0^{2\pi} d\phi \int_{-1}^1 d\eta \int_1^{\infty} d\xi \exp\left(-\frac{R}{a}\xi\right) \xi^2 - \int_0^{2\pi} d\phi \int_{-1}^1 d\eta \eta^2 \int_1^{\infty} d\xi \exp\left(-\frac{R}{a}\xi\right)\right) =$$

$$= \frac{R^3}{8\pi a^3} \left(2\pi \cdot 2 \cdot \int_1^{\infty} d\xi \exp\left(-\frac{R}{a}\xi\right) \xi^2 - 2\pi \cdot \left[\frac{\eta^3}{3}\right]_{-1}^1 \left[\frac{\exp\left(-\frac{R}{a}\xi\right)}{-\frac{R}{a}}\right]_1^{\infty}\right) =$$

$$= \frac{R^3}{8\pi a^3} \left(4\pi \int_1^{\infty} d\xi \exp\left(-\frac{R}{a}\xi\right) \xi^2 - 2\pi \cdot \frac{2}{3} \cdot \frac{\exp\left(-\frac{R}{a}\xi\right)}{\frac{R}{a}}\right) =$$

$$= \frac{R^3}{8\pi a^3} \left(4\pi \cdot \frac{2! \exp\left(-\frac{R}{a}\right)}{\left(\frac{R}{a}\right)^3} \sum_{k=0}^2 \frac{\left(\frac{R}{a}\right)^k}{k!} - \frac{4\pi}{3} \cdot \frac{\exp\left(-\frac{R}{a}\right)}{\frac{R}{a}}\right) =$$

$$= \frac{R^3 \exp\left(-\frac{R}{a}\right)}{2a^3} \left(\frac{2}{\left(\frac{R}{a}\right)^3} \sum_{k=0}^2 \frac{\left(\frac{R}{a}\right)^k}{k!} - \frac{1}{3} \cdot \frac{1}{\frac{R}{a}}\right) =$$

$$= \exp\left(-\frac{R}{a}\right) \left(\sum_{k=0}^2 \frac{\left(\frac{R}{a}\right)^k}{k!} - \frac{1}{6} \cdot \left(\frac{R}{a}\right)^2\right) =$$

$$= \exp\left(-\frac{R}{a}\right) \left(1 + \frac{1}{2} \cdot \left(\frac{R}{a}\right) + \frac{1}{6} \cdot \left(\frac{R}{a}\right)^2\right) =$$

$$= \exp\left(-\frac{R}{a}\right) \left(1 + \frac{1}{2} \cdot \left(\frac{R}{a}\right)\right)$$

Parte 3

$$\begin{split} \frac{\partial}{\partial r_i} \left| \vec{r} - \vec{r}_a \right| &= \frac{r_i - r_{a_i}}{\left| \vec{r} - \vec{r}_a \right|} = \frac{(\vec{r} - \vec{r}_a)_i}{\left| \vec{r} - \vec{r}_a \right|} \\ \nabla \left| \vec{r} - \vec{r}_a \right| &= \frac{\vec{r} - \vec{r}_a}{\left| \vec{r} - \vec{r}_a \right|} \end{split}$$

$$\begin{split} \nabla \exp\left(-\frac{|\vec{r}-\vec{r_a}|}{a}\right) &= \exp\left(-\frac{|\vec{r}-\vec{r_a}|}{a}\right) \nabla \left(-\frac{|\vec{r}-\vec{r_a}|}{a}\right) = \\ &= -\frac{1}{a} \exp\left(-\frac{|\vec{r}-\vec{r_a}|}{a}\right) \frac{\vec{r}-\vec{r_a}}{|\vec{r}-\vec{r_a}|} \end{split}$$

$$\begin{split} \frac{\partial_{r_{1}}}{\partial r_{1}} \left( -\frac{1}{a} \exp\left( -\frac{|\vec{r} - \vec{r}_{a}|}{a} \right) \frac{(\vec{r} - \vec{r}_{a})}{|\vec{r} - \vec{r}_{a}|} \right) = \\ &= -\frac{1}{a} \left[ \frac{(\vec{r} - \vec{r}_{a})_{i}}{|\vec{r} - \vec{r}_{a}|} \frac{\partial}{\partial r_{i}} \exp\left( -\frac{|\vec{r} - \vec{r}_{a}|}{a} \right) + \exp\left( -\frac{|\vec{r} - \vec{r}_{a}|}{a} \right) \frac{\partial_{r_{i}}}{\partial r_{i}} \frac{(\vec{r} - \vec{r}_{a})_{i}}{|\vec{r} - \vec{r}_{a}|} \right] \\ &= -\frac{1}{a} \exp\left( -\frac{|\vec{r} - \vec{r}_{a}|}{a} \right) \left[ \frac{(\vec{r} - \vec{r}_{a})_{i}}{|\vec{r} - \vec{r}_{a}|} \frac{\partial_{r_{i}}}{\partial r_{i}} \left( -\frac{|\vec{r} - \vec{r}_{a}|}{a} \right) + \frac{\partial_{r_{i}}}{(\vec{r} - \vec{r}_{a})_{i}} \frac{(\vec{r} - \vec{r}_{a})_{i}}{|\vec{r} - \vec{r}_{a}|} \right] \\ &= -\frac{1}{a} \exp\left( -\frac{|\vec{r} - \vec{r}_{a}|}{a} \right) \left[ -\frac{1}{a} \frac{(\vec{r} - \vec{r}_{a})_{i}}{|\vec{r} - \vec{r}_{a}|} + \frac{1}{r^{2} - r^{2}a} - (\vec{r} - \vec{r}_{a})_{i} \cdot \frac{(\vec{r} - \vec{r}_{a})_{i}}{|\vec{r} - \vec{r}_{a}|^{2}} \right] \\ &= -\frac{1}{a} \exp\left( -\frac{|\vec{r} - \vec{r}_{a}|}{a} \right) \sum_{i} \left[ -\frac{1}{a} \frac{(\vec{r} - \vec{r}_{a})_{i}}{|\vec{r} - \vec{r}_{a}|^{2}} + \frac{1}{|\vec{r} - \vec{r}_{a}|^{2}} - \frac{(\vec{r} - \vec{r}_{a})_{i}^{2}}{|\vec{r} - \vec{r}_{a}|^{3}} \right] \\ &= -\frac{1}{a} \exp\left( -\frac{|\vec{r} - \vec{r}_{a}|}{a} \right) \sum_{i} \left[ -\frac{1}{a} \frac{(\vec{r} - \vec{r}_{a})_{i}}{|\vec{r} - \vec{r}_{a}|^{2}} + \frac{1}{|\vec{r} - \vec{r}_{a}|^{2}} - \frac{(\vec{r} - \vec{r}_{a})_{i}^{2}}{|\vec{r} - \vec{r}_{a}|^{3}} \right] \\ &= -\frac{1}{a} \exp\left( -\frac{|\vec{r} - \vec{r}_{a}|}{a} \right) \sum_{i} \left[ -\frac{1}{a} \frac{|\vec{r} - \vec{r}_{a}|^{2}}{|\vec{r} - \vec{r}_{a}|^{2}} + \frac{1}{|\vec{r} - \vec{r}_{a}|^{2}} - \frac{(\vec{r} - \vec{r}_{a})_{i}^{2}}{|\vec{r} - \vec{r}_{a}|^{3}} \right] \\ &= -\frac{1}{a} \exp\left( -\frac{|\vec{r} - \vec{r}_{a}|}{a} \right) \left[ -\frac{1}{a} \frac{|\vec{r} - \vec{r}_{a}|^{2}}{|\vec{r} - \vec{r}_{a}|^{2}} + \frac{3}{|\vec{r} - \vec{r}_{a}|^{2}} - \frac{|\vec{r} - \vec{r}_{a}|^{2}}{|\vec{r} - \vec{r}_{a}|^{3}} \right] \\ &= -\frac{1}{a} \exp\left( -\frac{|\vec{r} - \vec{r}_{a}|}{a} \right) \left( -\frac{1}{a} \frac{1}{|\vec{r} - \vec{r}_{a}|} - \frac{2}{|\vec{r} - \vec{r}_{a}|^{2}} - \frac{|\vec{r} - \vec{r}_{a}|^{2}}{|\vec{r} - \vec{r}_{a}|^{3}} \right) \\ &= -\frac{1}{a} \exp\left( -\frac{|\vec{r} - \vec{r}_{a}|}{a} \right) \left( -\frac{1}{a} \frac{1}{|\vec{r} - \vec{r}_{a}|} - \frac{1}{|\vec{r} - \vec{r}_{a}|} - \frac{1}{|\vec{r} - \vec{r}_{a}|} - \frac{1}{|\vec{r} - \vec{r}_{a}|^{2}} \right) \right) \\ &= -\frac{1}{a} \exp\left( -\frac{|\vec{r} - \vec{r}_{a}|}{a} \right) \left( -\frac{1}{a} \frac{1}{|\vec{r} - \vec{r}_{a}|} - \frac{1}{|\vec{r} - \vec{r}_{a}|} - \frac{1}{|\vec{r} - \vec{r}_{a}|} - \frac{1}{|\vec{r} - \vec{r}_{a}|^{2}} \right)$$

Consideriamo solo la parentesi con il laplaciani

$$\begin{split} -\frac{\hbar^2}{2m}(\psi_b(\vec{r}_2)\nabla_1^2\psi_a(\vec{r}_1) + \psi_a(\vec{r}_2)\nabla_1^2\psi_b(\vec{r}_1) + \psi_b(\vec{r}_1)\nabla_2^2\psi_a(\vec{r}_2) + \psi_a(\vec{r}_1)\nabla_2^2\psi_b(\vec{r}_2)) = \\ = -\frac{\hbar^2}{2m}\left(\psi_b(\vec{r}_2)\psi_a(\vec{r}_1)\left(\frac{1}{a^2} - \frac{2}{a\,|\vec{r}_1 - \vec{r}_a|}\right) + \psi_a(\vec{r}_2)\psi_b(\vec{r}_1)\left(\frac{1}{a^2} - \frac{2}{a\,|\vec{r}_1 - \vec{r}_b|}\right) + \psi_b(\vec{r}_1)\psi_a(\vec{r}_2)\left(\frac{1}{a^2} - \frac{2}{a\,|\vec{r}_2 - \vec{r}_a|}\right) + \psi_a(\vec{r}_1)\psi_b(\vec{r}_2)\left(\frac{1}{a^2} - \frac{2}{a\,|\vec{r}_2 - \vec{r}_b|}\right)\right) = \\ = -\frac{\hbar^2}{2m}\left(\psi_a(\vec{r}_1)\psi_b(\vec{r}_2)\left(\frac{2}{a^2} - \frac{2}{a\,|\vec{r}_1 - \vec{r}_a|} - \frac{2}{a\,|\vec{r}_2 - \vec{r}_b|}\right) + \psi_b(\vec{r}_1)\psi_a(\vec{r}_2)\left(\frac{2}{a^2} - \frac{2}{a\,|\vec{r}_1 - \vec{r}_b|} - \frac{2}{a\,|\vec{r}_2 - \vec{r}_a|}\right)\right) = \end{split}$$

Moltiplicando per  $(\psi_a(\vec{r}_1)\psi_b(\vec{r}_2) + \psi_b(\vec{r}_1)\psi_a(\vec{r}_2))$  si trova

$$=-\frac{\hbar^2}{2m}\left(\psi_a(\vec{r}_1)^2\psi_b(\vec{r}_2)^2\left(\frac{2}{a^2}-\frac{2}{a\,|\vec{r}_1-\vec{r}_a|}-\frac{2}{a\,|\vec{r}_2-\vec{r}_b|}\right)+\psi_b(\vec{r}_1)^2\psi_a(\vec{r}_2)^2\left(\frac{2}{a^2}-\frac{2}{a\,|\vec{r}_1-\vec{r}_b|}-\frac{2}{a\,|\vec{r}_2-\vec{r}_a|}\right)\\ +\psi_a(\vec{r}_1)\psi_b(\vec{r}_1)\psi_a(\vec{r}_2)\psi_b(\vec{r}_2)\left(\frac{4}{a^2}-\frac{2}{a\,|\vec{r}_1-\vec{r}_a|}-\frac{2}{a\,|\vec{r}_2-\vec{r}_b|}-\frac{2}{a\,|\vec{r}_1-\vec{r}_b|}-\frac{2}{a\,|\vec{r}_1-\vec{r}_b|}-\frac{2}{a\,|\vec{r}_2-\vec{r}_a|}\right)\right)$$