

## Some Properties of Fourier Transform

### 1 Addition Theorem

If  $g(x) \supset G(s)$  and  $h(x) \supset H(s)$ , and  $a$  and  $b$  are some scalars, then

$$ag(x) + bh(x) \supset aG(s) + bH(s). \quad (1)$$

*Proof.*

$$\begin{aligned} \mathfrak{F}\{ag(x) + bh(x)\} &= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} [ag(x) + bh(x)] e^{-j2\pi sx} \, dx \\ &= a \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} g(x) e^{-j2\pi sx} \, dx + b \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} h(x) e^{-j2\pi sx} \, dx \\ &= aG(s) + bH(s). \end{aligned}$$

□

### 2 Convolution Theorem

If  $g(x) \supset G(s)$  and  $h(x) \supset H(s)$ , then

$$g(x) * h(x) \supset G(s)H(s). \quad (2)$$

*Proof.*

$$\begin{aligned} \mathfrak{F}\{g(x) * h(x)\} &= \int_{-\infty}^{\infty} \left( \int_{-\infty}^{\infty} g(u)h(x-u) \, du \right) e^{-j2\pi sx} \, dx \\ &= \int_{-\infty}^{\infty} g(u) \left( \int_{-\infty}^{\infty} h(x-u) e^{-j2\pi sx} \, dx \right) \, du \\ &= \int_{-\infty}^{\infty} g(u) e^{-j2\pi us} H(s) \, du \\ &= G(s)H(s). \end{aligned}$$

□

### 3 Similarity Theorem

If  $g(x) \supset G(s)$ , then

$$g(ax) \supset \frac{1}{|a|} G\left(\frac{s}{a}\right). \quad (3)$$

*Proof.*

$$\begin{aligned}\int_{-\infty}^{\infty} g(ax) e^{-j2\pi s x} dx &= \frac{1}{|a|} \int_{-\infty}^{\infty} g(ax) e^{-j2\pi(s/a)(ax)} d(ax) \\ &= \frac{1}{|a|} G\left(\frac{s}{a}\right).\end{aligned}$$

□

This means a contraction in one domain gives rise to expansion in another domain, and vice versa.

## 4 Shift Theorem

If  $g(x) \supset G(s)$ , then

$$g(x - a) \supset e^{-j2\pi as} G(s). \quad (4)$$

*Proof.*

$$\begin{aligned}\int_{-\infty}^{\infty} g(x - a) e^{-j2\pi s x} dx &= \int_{-\infty}^{\infty} g(x - a) e^{-j2\pi s(x-a)} e^{-j2\pi sa} d(x - a) \\ &= e^{-j2\pi as} G(s).\end{aligned}$$

□

A shift in position in one domain gives rise to a phase change in another domain. In a similar fashion, we can show that

$$e^{j2\pi ax} g(x) \supset G(s - a). \quad (5)$$

## 5 Modulation Theorem

If  $g(x) \supset G(s)$ , then

$$g(x) \cos(2\pi ax) \supset \frac{1}{2} [G(s - a) + G(s + a)]. \quad (6)$$

*Proof.*

$$\begin{aligned}\mathfrak{F}\{g(x) \cos(2\pi ax)\} &= \frac{1}{2} \int_{-\infty}^{\infty} g(x) e^{j2\pi ax} e^{-j2\pi s x} dx + \frac{1}{2} \int_{-\infty}^{\infty} g(x) e^{-j2\pi ax} e^{-j2\pi s x} dx \\ &= \frac{1}{2} \int_{-\infty}^{\infty} g(x) e^{-j2\pi(s-a)x} dx + \frac{1}{2} \int_{-\infty}^{\infty} g(x) e^{-j2\pi(s+a)x} dx \\ &= \frac{1}{2} [G(s - a) + G(s + a)].\end{aligned}$$

□

Alternatively, if we make use of the fourier transform of a cosine and the convolution theorem in Equation 2,

$$\begin{aligned}\mathfrak{F}\{g(x) \cos(2\pi ax)\} &= \mathfrak{F}\{g(x)\} * \mathfrak{F}\{\cos(2\pi ax)\} \\ &= G(s) * \frac{1}{2} [\delta(s - a) + \delta(s + a)] \\ &= \frac{1}{2} [G(s - a) + G(s + a)].\end{aligned}$$

Multiplication with a cosine has the effect of shifting the spectrum to center on the frequency of the cosine.

## 6 Time Reversal Theorem

If  $g(x) \supset G(s)$ , then

$$g(-x) \supset G(-s). \quad (7)$$

*Proof.*

$$\begin{aligned} \mathfrak{F}\{g(-x)\} &= \int_{-\infty}^{\infty} g(-x) e^{-j2\pi s x} dx \\ &= \int_{\infty}^{-\infty} g(\tilde{x}) e^{-j2\pi(-s)\tilde{x}} (-d\tilde{x}) \quad (\tilde{x} = -x) \\ &= G(-s). \end{aligned}$$

□

## 7 Derivative Theorem

If  $g(x) \supset G(s)$ , then

$$\frac{d}{dx}g(x) \supset j2\pi s G(s). \quad (8)$$

*Proof.*

$$\begin{aligned} g(x) &= \int_{-\infty}^{\infty} G(s) e^{j2\pi s x} ds \\ \frac{d}{dx}g(x) &= \frac{d}{dx} \left( \int_{-\infty}^{\infty} G(s) e^{j2\pi s x} ds \right) \\ &= \int_{-\infty}^{\infty} G(s) \frac{d}{dx} (e^{j2\pi s x}) ds \\ &= \int_{-\infty}^{\infty} [j2\pi s G(s)] e^{j2\pi s x} ds \end{aligned}$$

$$\text{Hence} \quad \mathfrak{F} \left\{ \frac{d}{dx}g(x) \right\} = j2\pi s G(s).$$

□

By extension, we have

$$\frac{d^n}{dx^n}g(x) \supset (j2\pi s)^n G(s). \quad (9)$$

## 8 2D rotation Theorem

If  $g(x, y) \supset G(f_X, f_Y)$ , then

$$g(x \cos \theta + y \sin \theta, -x \sin \theta + y \cos \theta) \supset G(f_X \cos \theta + f_Y \sin \theta, -f_X \sin \theta + f_Y \cos \theta). \quad (10)$$

In words, that means an anti-clockwise rotation of a function by an angle  $\theta$  implies that its Fourier transform is also rotated anti-clockwise by the same angle.

*Proof.* We can define a new coordinate system  $(\tilde{x}, \tilde{y})$ , where

$$\begin{bmatrix} \tilde{x} \\ \tilde{y} \end{bmatrix} = \begin{bmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix}. \quad (11)$$

This is shown in Figure 1. We can also express  $(x, y)$  in terms of the new coordinate system, where

$$\begin{bmatrix} x \\ y \end{bmatrix} = \begin{bmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{bmatrix} \begin{bmatrix} \tilde{x} \\ \tilde{y} \end{bmatrix} = \begin{bmatrix} \tilde{x} \cos \theta - \tilde{y} \sin \theta \\ \tilde{x} \sin \theta + \tilde{y} \cos \theta \end{bmatrix}.$$

Thus,

$$\begin{aligned} \mathfrak{F}\{g(\tilde{x}, \tilde{y})\} &= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} g(\tilde{x}, \tilde{y}) e^{-j2\pi(xf_X + yf_Y)} dx dy \\ &= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} g(\tilde{x}, \tilde{y}) e^{-j2\pi(\tilde{x}f_X \cos \theta - \tilde{y}f_X \sin \theta + \tilde{x}f_Y \sin \theta + \tilde{y}f_Y \cos \theta)} dx dy \\ &= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} g(\tilde{x}, \tilde{y}) e^{-j2\pi[\tilde{x}(f_X \cos \theta + f_Y \sin \theta) + \tilde{y}(-f_X \sin \theta + f_Y \cos \theta)]} dx dy \\ &= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} g(\tilde{x}, \tilde{y}) e^{-j2\pi[\tilde{x}(f_X \cos \theta + f_Y \sin \theta) + \tilde{y}(-f_X \sin \theta + f_Y \cos \theta)]} d\tilde{x} d\tilde{y} \\ &= G(f_X \cos \theta + f_Y \sin \theta, -f_X \sin \theta + f_Y \cos \theta). \end{aligned}$$

In the derivation above, we have made use of the fact that

$$d\tilde{x} d\tilde{y} = \begin{vmatrix} \frac{\partial \tilde{x}}{\partial x} & \frac{\partial \tilde{y}}{\partial x} \\ \frac{\partial \tilde{x}}{\partial y} & \frac{\partial \tilde{y}}{\partial y} \end{vmatrix} dx dy = \begin{vmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{vmatrix} dx dy = dx dy.$$

□

One consequence of the two-dimensional rotation theorem is that if the 2D function is circularly symmetric, its Fourier transform must also be circularly symmetric.

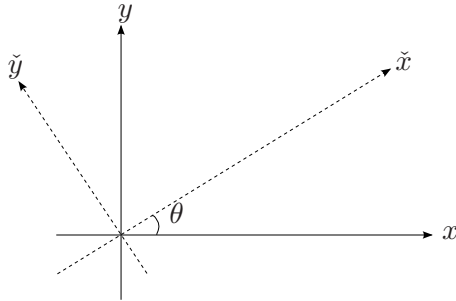


Figure 1: Illustration of a rotation in coordinates.