

1 General Definitions and Tools

1.1 NOTATIONS AND CONVENTIONS

1.1.1 Metric etc.

Minkowski Metric : $\eta^{\mu\nu} = \text{diag}(+, -, -, -)$

Coordinates : $x^\mu = (t, x, y, z)$; Therefore $\partial_\mu = \left(\frac{\partial}{\partial t}, \nabla \right)$

Gamma Matrices : $\{\gamma^\mu, \gamma^\nu\} = 2\eta^{\mu\nu}$; $\gamma_5 = i\gamma^0\gamma^1\gamma^2\gamma^3$

Gamma Combinations : $1, \{\gamma^\mu\}, \{\sigma^{\mu\nu}\}, \{\gamma^\mu\gamma_5\}, \gamma_5$; $\sigma^{\mu\nu} = \frac{i}{2}[\gamma^\mu, \gamma^\nu] = 0 / i\gamma^\mu\gamma^\nu$

1.1.2 Fields

Klein-Gordon Equation : $|\partial_\mu\phi|^2 - m^2|\phi|^2 = 0$

Klein-Gordon Field : $\phi(x) = \int \frac{d^3p}{(2\pi)^3} \frac{1}{\sqrt{2E_p}} \left[a_p e^{-ipx} + b_p^\dagger e^{ipx} \right]$

Dirac Equation : $(i\not{\partial} - m)\psi(x) = 0$

Dirac Field : $\psi(x) = \int \frac{d^3p}{(2\pi)^3} \frac{1}{\sqrt{2E_p}} \sum_{s=1,2} \left[a_p^s u^s(p) e^{-ipx} + b_p^{s\dagger} v^s(p) e^{ipx} \right]$

Gauge Boson : $(\partial^2 + m^2)A^\mu(x) = 0$ (Real Klein-Gordon Equation)

(Before Gauge Fixing) : $A^\mu(x) = \int \frac{d^3p}{(2\pi)^3} \frac{1}{\sqrt{2E_p}} \sum_{r=0,1,2,3} \left[a_p^r \epsilon^r(p) e^{-ipx} + a_p^{r\dagger} \epsilon^{r*}(p) e^{ipx} \right]$

南部-Goldstone Boson : **TODO: .**

Chiral Notation

Gamma Matrices : $\gamma^\mu = \begin{pmatrix} 0 & \sigma^\mu \\ \bar{\sigma}^\mu & 0 \end{pmatrix}$; $\gamma_5 = \begin{pmatrix} -1 & 0 \\ 0 & 1 \end{pmatrix}$

Dirac Field : $\psi = \begin{pmatrix} \psi_L \\ \psi_R \end{pmatrix}$; $\bar{\psi} = \psi^\dagger \gamma^0 = \begin{pmatrix} \psi_R^\dagger & \psi_L^\dagger \end{pmatrix}$

: $u^s(p) = \begin{pmatrix} \sqrt{p \cdot \sigma} \xi^s \\ \sqrt{p \cdot \bar{\sigma}} \xi^s \end{pmatrix}$; $v^s(p) = \begin{pmatrix} \sqrt{p \cdot \sigma} \eta^s \\ -\sqrt{p \cdot \bar{\sigma}} \eta^s \end{pmatrix}$ $[\eta^s = \xi^{-s} := -i\sigma^2(\xi^s)^* = (\xi^2, -\xi^1)]$

Weyl Equations : $i\bar{\sigma} \cdot \partial \psi_L = m\psi_R$; $i\sigma \cdot \partial \psi_L = m\psi_L$

CPT transf. : $P\psi(t, \mathbf{x})P = \eta\gamma^0\psi(t, -\mathbf{x})$ ($|\eta|^2 = 1$)

: $T\psi(t, \mathbf{x})T = \gamma^1\gamma^3\psi(-t, \mathbf{x})$ (ignoring intrinsic phase)

: $C\psi(t, \mathbf{x})C = -i\gamma^2\psi^*(t, \mathbf{x}) = -i(\bar{\psi}\gamma^0\gamma^2)^T$ (")

: $\bar{\psi} \longrightarrow P : \eta^* \bar{\psi} \gamma^0$ $T : -\bar{\psi} \gamma^1 \gamma^3$ $C : i \bar{\psi}^* \gamma^2 = -i(\gamma^0 \gamma^2 \psi)^T$

Electromagnetic Fields : $A^\mu = (\phi, \mathbf{A})$ **【We can invert the signs, but cannot lower the index.】**

: $F_{\mu\nu} = \begin{pmatrix} 0 & \mathbf{E} \\ -\mathbf{E} & B_3 & 0 & -B_1 \\ & -B_2 & B_1 & 0 \end{pmatrix}$

: $F_{\mu\nu}F^{\mu\nu} = -2(\|\mathbf{E}\|^2 - \|\mathbf{B}\|^2)$

1.1.3 Dirac Field Techniques

$$\begin{aligned}
\text{Dirac Equations} & : (\not{p} - m)u^s(p) = 0; \quad (\not{p} + m)v^s(p) = 0 \\
& : \bar{u}^s(p)(\not{p} - m) = 0; \quad \bar{v}^s(p)(\not{p} + m) = 0 \\
\text{Dirac Components} & : \bar{u}^r(p)u^s(p) = 2m\delta^{rs}; \quad u^{r\dagger}(p)u^s(p) = 2E_p\delta^{rs} \\
& : \bar{v}^r(p)v^s(p) = -2m\delta^{rs}; \quad v^{r\dagger}(p)v^s(p) = 2E_p\delta^{rs}; \quad \bar{u}^r(p)v^s(p) = \bar{v}^r(p)u^s(p) = 0 \\
\text{Spin Sums} & : \sum_{\text{spin}} u^s(p)\bar{u}^s(p) = \not{p} + m; \quad \sum_{\text{spin}} v^s(p)\bar{v}^s(p) = \not{p} - m
\end{aligned}$$

1.1.4 CPT Table

	ϕ	A^μ	$\bar{\psi}\psi$	$\bar{\psi}\gamma^\mu\psi$	$\bar{\psi}\sigma^{\mu\nu}\psi$	$\bar{\psi}\gamma^\mu\gamma_5\psi$	$i\bar{\psi}\gamma_5\psi$	∂_μ
P	$\eta\phi$	$\eta-++A^\mu$	+	+---	(+---)(+---)	-+++	-	+---
T	$\zeta\phi$	$\zeta+---A^\mu$	+	+---	-(+---)(+---)	+---	-	-+++
C	$\xi\phi^*$	$\xi+A^{\mu*}$	+	-	-	+	+	+

($\eta\zeta\xi = 1$; especially, photon A^μ is $(\eta, \zeta, \xi) = (-, +, -)$.)

1.1.5 Polarization Sum

Single photon case $M = \epsilon_\mu^*(k)M^\mu$

When Ward identity $k_\mu M^\mu = 0$ is valid,

$$\sum_{\text{pol.}} |M|^2 = \sum_{\text{pol.}} \epsilon_\mu^*(k)\epsilon_\nu(k)M^\mu M^{\nu*} = \eta_{\mu\nu}M^\mu M^{\nu*} \quad (1.1)$$

Double photons case $M = \epsilon_\mu^*(k)\epsilon_\nu'^*(k')M^{\mu\nu}$

When $k_\mu M^{\mu\nu} = k'_\nu M^{\mu\nu} = 0$ is valid,

$$\sum_{\text{pol.}} |M|^2 = \sum_{\text{pol.}} \epsilon_\mu^*(k)\epsilon_\rho(k)\epsilon_\nu'^*(k')\epsilon_\sigma'(k')[M^{\mu\nu}M^{\rho\sigma*} = \eta_{\mu\rho}\eta_{\nu\sigma}M^{\mu\nu}M^{\rho\sigma*} \quad (1.2)$$

【See Sec. A.1 for verbose information.】

1.2 DIRAC'S GAMMA ALGEBRAS

Traces

$$\text{Tr}(\text{any odd \# of } \gamma\text{'s}) = 0 \quad (1.3)$$

$$\text{Tr}(\gamma^\mu\gamma^\nu) = 4\eta^{\mu\nu} \quad (1.4)$$

$$\text{Tr}(\gamma^\mu\gamma^\nu\gamma^\rho\gamma^\sigma) = 4(\eta^{\mu\nu}\eta^{\rho\sigma} - \eta^{\mu\rho}\eta^{\nu\sigma} + \eta^{\mu\sigma}\eta^{\nu\rho}) \quad (1.5)$$

$$\text{Tr}(\gamma_5 \text{ and any odd \# of } \gamma\text{'s}) = 0 \quad (1.6)$$

$$\text{Tr}(\gamma^\mu\gamma^\nu\gamma_5) = 0 \quad (1.7)$$

$$\text{Tr}(\gamma^\mu\gamma^\nu\gamma^\rho\gamma^\sigma\gamma_5) = -4i\epsilon^{\mu\nu\rho\sigma} \quad (1.8)$$

Generally, for some γ -matrices A, B, C, \dots ,

$$\begin{aligned} \text{Tr}(ABCDEF \dots) &= \eta^{AB} \text{Tr}(CDEF \dots) - \eta^{AC} \text{Tr}(BDEF \dots) \\ &+ \eta^{AD} \text{Tr}(BCEF \dots) - \eta^{AE} \text{Tr}(BCDF \dots) + \dots \end{aligned} \quad (1.9)$$

$$\begin{aligned} \text{Tr}(ABCDEF \dots \gamma_5) &= \eta^{AB} \text{Tr}(CDEF \dots \gamma_5) - \eta^{AC} \text{Tr}(BDEF \dots \gamma_5) + \dots \\ &+ \eta^{BC} \text{Tr}(ADEF \dots \gamma_5) - \eta^{BD} \text{Tr}(ACEF \dots \gamma_5) + \dots \\ &+ \eta^{CD} \text{Tr}(ABEF \dots \gamma_5) - \eta^{CE} \text{Tr}(ABDF \dots \gamma_5) + \dots \\ &+ \dots \end{aligned} \quad (1.10)$$

To prove the second equation, we use following technique:

$$\text{Tr}(\gamma^\mu \gamma^\nu \gamma^\rho \gamma^\sigma \dots) = \text{Tr}(\dots \gamma^\sigma \gamma^\rho \gamma^\nu \gamma^\mu); \quad \text{Tr}(\gamma^\mu \gamma^\nu \gamma^\rho \gamma^\sigma \dots \gamma_5) = \text{Tr}(\gamma_5 \dots \gamma^\sigma \gamma^\rho \gamma^\nu \gamma^\mu) \quad (1.11)$$

Contractions

$$\gamma^\mu \gamma_\mu = 4 \quad (1.12)$$

$$\gamma^\mu \gamma^\nu \gamma_\mu = -2\gamma^\nu \quad (1.13)$$

$$\gamma^\mu \gamma^\nu \gamma^\rho \gamma_\mu = 4\eta^{\nu\rho} \quad (1.14)$$

$$\gamma^\mu \gamma^\nu \gamma^\rho \gamma^\sigma \gamma_\mu = -2\gamma^\sigma \gamma^\rho \gamma^\nu \quad (1.15)$$

Generally, for some γ -matrices A, B, C, \dots ,

$$\text{ODD } \# : \quad \gamma^\mu ABC \dots \gamma_\mu = -2(\dots CBA) \quad (1.16)$$

$$\text{EVEN } \# : \quad \gamma^\mu ABC \dots \gamma_\mu = \text{Tr}(ABC \dots) - \text{Tr}(ABC \dots \gamma_5) \cdot \gamma_5 \quad (1.17)$$

Contractions of ϵ 's

$$\epsilon^{\alpha\beta\gamma\delta} \epsilon_{\alpha\beta\gamma\delta} = -24; \quad \epsilon^{\alpha\beta\gamma\mu} \epsilon_{\alpha\beta\gamma\nu} = -6\delta_\nu^\mu; \quad \epsilon^{\alpha\beta\mu\nu} \epsilon_{\alpha\beta\rho\sigma} = -2(\delta_\rho^\mu \delta_\sigma^\nu - \delta_\sigma^\mu \delta_\rho^\nu) \quad (1.18)$$

$$\epsilon^{\mu\alpha\beta\gamma} \epsilon_{\mu\alpha'\beta'\gamma'} = - \left(\delta_{\alpha'}^\alpha \delta_{\beta'}^\beta \delta_{\gamma'}^\gamma + \delta_{\beta'}^\alpha \delta_{\gamma'}^\beta \delta_{\alpha'}^\gamma + \delta_{\gamma'}^\alpha \delta_{\alpha'}^\beta \delta_{\beta'}^\gamma - \delta_{\alpha'}^\alpha \delta_{\gamma'}^\beta \delta_{\beta'}^\gamma - \delta_{\beta'}^\alpha \delta_{\alpha'}^\beta \delta_{\gamma'}^\gamma - \delta_{\gamma'}^\alpha \delta_{\beta'}^\beta \delta_{\alpha'}^\gamma \right) \quad (1.19)$$

$$(p \cdot \sigma)(p \cdot \bar{\sigma}) = p^2$$

- Fierz Transf.
- Noether current
- Gordon Id.
- Majorana Fermions
-
- Feynman Rules(A.1)
- Gamma Algebras(A.3)
- A.4, A.5

2 Standard Model

Gauge Fields We use following notations for the Gauge Group $SU(3) \times SU(2) \times U(1)$:

$$SU(3): \quad G_\mu = G_\mu^a \tau^a; \quad [\tau^a, \tau^b] = i f^{abc} \tau^c, \quad \text{Tr}(\tau^a \tau^b) = \frac{1}{2} \delta^{ab}, \quad (2.1)$$

$$SU(2): \quad W_\mu = W_\mu^a T^a; \quad [T^a, T^b] = i \epsilon^{abc} T^c, \quad \text{Tr}(T^a T^b) = \frac{1}{2} \delta^{ab}, \quad (2.2)$$

$$U(1): \quad B_\mu. \quad (2.3)$$

以下では射影演算子 $P_L^R := \frac{1 \pm \gamma_5}{2}$ をあらわに書くことにする。

2.1 FULL LAGRANGIAN

標準模型の構成粒子から作られる, $SU(3) \times SU(2) \times U(1)$ gauge 対称性を保つ Lagrangian は,

$$\mathcal{L} = -\frac{1}{4} B^{\mu\nu} B_{\mu\nu} - \frac{1}{4} W^{a\mu\nu} W_{\mu\nu}^a - \frac{1}{4} G^{a\mu\nu} G_{\mu\nu}^a \quad (2.4)$$

$$+ \left| \left(\partial_\mu - i g_2 W_\mu - \frac{1}{2} i g_1 B_\mu \right) \phi \right|^2 - V(\phi) \quad (2.5)$$

$$+ \bar{Q} i \gamma^\mu \left(\partial_\mu - i g_3 G_\mu - i g_2 W_\mu - \frac{1}{6} i g_1 B_\mu \right) P_L Q \quad (2.6)$$

$$+ \bar{U} i \gamma^\mu \left(\partial_\mu - i g_3 G_\mu - \frac{2}{3} i g_1 B_\mu \right) P_R U \quad (2.7)$$

$$+ \bar{D} i \gamma^\mu \left(\partial_\mu - i g_3 G_\mu + \frac{1}{3} i g_1 B_\mu \right) P_R D \quad (2.8)$$

$$+ \bar{L} i \gamma^\mu \left(\partial_\mu - i g_2 W_\mu + \frac{1}{2} i g_1 B_\mu \right) P_L L \quad (2.9)$$

$$+ \bar{E} i \gamma^\mu (\partial_\mu + i g_1 B_\mu) P_R E \quad (2.10)$$

$$+ \bar{U} y_u H P_L Q + \bar{D} y_d H^\dagger P_L Q + \bar{E} y_e H^\dagger P_L L + \text{H. c.} \quad (2.11)$$

の形となる。

ただし $F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu - i g [A_\mu, A_\nu] = (\partial_\mu A_\nu^a - \partial_\nu A_\mu^a + g f^{abc} A_\mu^b A_\nu^c) T^a$ である。

これ以外の項が加わらない理由

次元勘定だけからは, これ以外に $\bar{\psi} \gamma_5 \psi$, $\bar{\psi} \gamma_5 \not{D} \psi$, $\epsilon^{\mu\nu\rho\sigma} F_{\mu\nu}^a F_{\rho\sigma}^a$, $\epsilon^{\mu\nu\rho\sigma} D_\mu D_\nu F_{\rho\sigma}^a$ が付け加わりそうに見える。しかし,

$$(\bar{\psi} \gamma_5 \psi)^* = -\bar{\psi} \gamma_5 \psi \quad (2.12)$$

$$(\bar{\psi} \gamma_5 \not{D} \psi)^* = -\bar{\psi} \gamma_5 \not{D} \psi \quad (2.13)$$

であるのでこの2つの項は Lagrangian には寄与せず, また

$$\epsilon^{\mu\nu\rho\sigma} D_\mu D_\nu F_{\rho\sigma}^a = \epsilon^{\mu\nu\rho\sigma} \frac{1}{2} [D_\mu, D_\nu] F_{\rho\sigma}^a = \frac{1}{2} \epsilon^{\mu\nu\rho\sigma} F_{\mu\nu}^a F_{\rho\sigma}^a \quad (2.14)$$

である。ところが $\epsilon F F$ の項は $\psi_R \mapsto e^{i\alpha(x)} \psi_R$ の変換によって打ち消すことができる。これは場の再定義なのか, $U(1)$ gauge 固定なのか。よくわかっていない。

Gauge 場の運動項

$$(\partial A)_{\mu\nu} := \partial_\mu A_\nu - \partial_\nu A_\mu \quad (2.15)$$

と略記することになると，Gauge 場の運動項は

$$\mathcal{L}_{B;\text{kin}} = -\frac{1}{4}(\partial B)(\partial B) \quad (2.16)$$

$$\mathcal{L}_{W;\text{kin}} = -\frac{1}{4}(\partial W^a)(\partial W^a) - g_2 \epsilon^{abc}(\partial_\mu W_\nu^a)W^{\mu b}W^{\nu c} - \frac{g_2^2}{4}(\epsilon^{eab}W_\mu^a W_\nu^b)(\epsilon^{ecd}W^{c\mu}W^{d\nu}) \quad (2.17)$$

$$\mathcal{L}_{G;\text{kin}} = -\frac{1}{4}(\partial G^a)(\partial G^a) - g_3 f^{abc}(\partial_\mu G_\nu^a)G^{\mu b}G^{\nu c} - \frac{g_3^2}{4}(f^{eab}G_\mu^a G_\nu^b)(f^{ecd}G^{c\mu}G^{d\nu}) \quad (2.18)$$

となる。

2.2 HIGGS MECHANISM

Higgs 場が真空期待値 $\langle \phi \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v+h(x) \end{pmatrix}$ ($v, h \in \mathbb{R}$) を持つよう， $\text{SU}(2)$ を gauge 固定する。

$$\mathcal{L}_{\text{Higgs}} = \left| \left(\partial_\mu - i g_2 W_\mu - \frac{1}{2} i g_1 B_\mu \right) \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v+h \end{pmatrix} \right|^2 \quad (2.19)$$

$$= \frac{1}{2}(\partial_\mu h)^2 + \frac{(v+h)^2}{8} \left[g_2^2 W_1^2 + g_2^2 W_2^2 + (g_1 B - g_2 W_3)^2 \right] \quad (2.20)$$

場の norm に注意しながら gauge 場を再定義する。

$$W_\mu^\pm := \frac{1}{\sqrt{2}}(W_\mu^1 \mp i W_\mu^2), \quad Z_\mu^0 := \frac{1}{\sqrt{g_1^2 + g_2^2}}(g_2 W_\mu^3 - g_1 B_\mu), \quad A_\mu := \frac{1}{\sqrt{g_1^2 + g_2^2}}(g_1 W_\mu^3 + g_2 B_\mu) \quad (2.21)$$

その結果

$$\mathcal{L}_{\text{Higgs}} = \frac{1}{2}(\partial_\mu h)^2 + \frac{(v+h)^2}{8} \left[g_2^2 W^+ W^- + g_2^2 W^- W^+ + (g_1^2 + g_2^2)(Z^0)^2 \right], \quad (2.22)$$

$$g_1 B_\mu = |e| A - \frac{|e|s}{c} Z, \quad (2.23)$$

$$g_2 W_\mu = \frac{g_2}{\sqrt{2}} (W_\mu^+ T^+ + W_\mu^- T^-) + \left(\frac{|e|c}{s} Z_\mu^0 + |e| A_\mu \right) T^3 \quad (2.24)$$

となる。ここで $T^\pm = T^1 \pm i T^2$ であり，また Weinberg 角 θ_w と素電荷 e を導入した：

$$|e| := \frac{g_1 g_2}{\sqrt{g_1^2 + g_2^2}}; \quad s := \sin \theta_w, \quad c := \cos \theta_w; \quad g_1 = \frac{|e|}{c}, \quad g_2 = \frac{|e|}{s}. \quad (2.25)$$

Higgs potential

Higgs 項は， $\text{SU}(2)$ 対称性より

$$V(\phi) = -\mu^2(\phi^\dagger \phi) + \lambda|\phi^\dagger \phi|^2 \quad (2.26)$$

の形に限られる。故に先述のように $\text{SU}(2)$ を固定すると

$$V(h) = \frac{1}{4}\lambda h^4 + v\lambda h^3 + \mu^2 h^2 \quad (2.27)$$

$$= \frac{1}{4}\lambda h^4 + \sqrt{\frac{\lambda}{2}} m_h h^3 + \frac{1}{2}\mu^2 h^2 \quad (2.28)$$

となり，Higgs の質量は $m_h = \frac{\mu}{\sqrt{2}}$ で与えられる。

Gauge 項

Gauge 項は

$$\mathcal{L}_{\text{gauge}} = -\frac{1}{4} [(\partial W^3)(\partial W^3) + (\partial B)(\partial B) + 2(\partial W^+)(\partial W^-) + G^{\alpha\mu\nu} G_{\mu\nu}^a] \quad (2.29)$$

$$+ \frac{ig_2}{2} [(W^+ W^-)(\partial W^3) + (W^3 W^+)(\partial W^-) + (W^- W^3)(\partial W^+)] \quad (2.30)$$

$$+ \frac{g_2^2}{4} [(W^+ W^-)(W^+ W^-) + 2(W^3 W^+)(W^- W^3)] \quad (2.31)$$

$$= -\frac{1}{4} [(\partial Z^0)(\partial Z^0) + (\partial A)(\partial A) + 2(\partial W^+)(\partial W^-) + G^{\alpha\mu\nu} G_{\mu\nu}^a] \quad (2.32)$$

$$+ \frac{i|e|c}{2s} [W^+ W^- (\partial Z^0) + Z^0 W^+ (\partial W^-) + W^- Z^0 (\partial W^+)] \quad (2.33)$$

$$+ \frac{i|e|}{2} [W^+ W^- (\partial A) + A W^+ (\partial W^-) + W^- A (\partial W^+)] \quad (2.34)$$

$$+ \frac{|e|^2}{4s^2} W^+ W^+ W^- W^- + \frac{|e|^2 c^2}{2s^2} W^+ W^- Z^0 Z^0 \quad (2.35)$$

$$+ \frac{|e|^2 c}{s} W^+ W^- Z^0 A + \frac{|e|^2}{2} W^+ W^- A A \quad (2.36)$$

となる。

湯川項

湯川項は，SU(2) の脚を露わに書くと

$$\mathcal{L}_{\text{湯川}} = \bar{U} y_u H P_L Q + \bar{D} y_d H^\dagger P_L Q + \bar{E} y_e H^\dagger P_L L + \text{H. c.} \quad (2.37)$$

$$= \bar{U} y_u \epsilon^{\alpha\beta} H^\alpha P_L Q^\beta + \bar{D} y_d H^{\dagger\alpha} P_L Q^\alpha + \bar{E} y_e H^{\dagger\alpha} P_L L^\alpha + \text{H. c.} \quad (2.38)$$

$$= (v + h) (\bar{U} y_u P_L Q^1 + \bar{D} y_d P_L Q^2 + \bar{E} y_e P_L L^2) + \text{H. c.} \quad (2.39)$$

となる（ただし $\epsilon^{\alpha\beta}$ の符号は y_u に吸収させた）。

2.3 FULL LAGRANGIAN AFTER HIGGS MECHANISM

ここでは簡単のため P_L などを省略する。全 Lagrangian は

$$\mathcal{L} = \mathcal{L}_{\text{gauge}} \quad (2.40)$$

$$+ m_W^2 W^+ W^- + \frac{m_Z^2}{2} Z^2 \quad (2.41)$$

$$\text{【Higgs Field】} + (\partial_\mu h)^2 - \frac{1}{2} \mu^2 h^2 - \sqrt{\frac{\lambda}{2}} m_h h^3 - \frac{1}{4} \lambda h^4 \quad (2.42)$$

$$+ \frac{v g_2^2}{4} W^+ W^- h + \frac{v(g_1^2 + g_2^2)}{8} Z^2 h \quad (2.43)$$

$$+ \frac{g_2^2}{4} W^+ W^- h^2 + \frac{g_1^2 + g_2^2}{8} Z^2 h^2 \quad (2.44)$$

$$+ h \bar{U} y_u Q^1 + h \bar{D} y_d Q^2 + h \bar{E} y_e L^2 + \text{H. c.} \quad (2.45)$$

$$\text{【SU(3) および微分項】} + \bar{Q} (i \not{\partial} + g_3 \not{G}) Q + \bar{U} (i \not{\partial} + g_3 \not{G}) U + \bar{D} (i \not{\partial} + g_3 \not{G}) D \quad (2.46)$$

$$+ \bar{L} (i \not{\partial}) L + \bar{E} (i \not{\partial}) E \quad (2.47)$$

$$\text{【W boson】} + \bar{Q} \frac{g_2}{\sqrt{2}} (W^+ T^+ + W^- T^-) Q + \bar{L} \frac{g_2}{\sqrt{2}} (W^+ T^+ + W^- T^-) L \quad (2.48)$$

$$\text{【A&Z}^0 \text{ boson】} + \bar{Q} \left[\left(T^3 + \frac{1}{6} \right) |e| A + \left(\frac{|e|c}{s} T^3 - \frac{|e|s}{6c} \right) Z^0 \right] Q \quad (2.49)$$

$$+ \bar{U} \left(\frac{2}{3} |e| A - \frac{2|e|s}{3c} Z \right) U \quad (2.50)$$

$$+ \bar{D} \left(-\frac{1}{3} |e| A + \frac{|e|s}{3c} Z \right) D \quad (2.51)$$

$$+ \bar{L} \left[\left(T^3 - \frac{1}{2} \right) |e| A + \left(\frac{|e|c}{s} T^3 + \frac{|e|s}{2c} \right) Z^0 \right] L \quad (2.52)$$

$$+ \bar{E} \left(-|e| A + \frac{|e|s}{c} Z \right) E \quad (2.53)$$

$$\text{【湯川項】} + v \bar{U} y_u Q^1 + v \bar{D} y_d Q^2 + v \bar{E} y_e L^2 + \text{H. c.} \quad (2.54)$$

である。ただしここで

$$m_W := \frac{g_2 v}{2}, \quad m_Z := \frac{\sqrt{g_1^2 + g_2^2}}{2} v \quad (2.55)$$

を導入した。

2.4 MASS EIGENSTATES

湯川行列を対角化し、質量の固有状態を得ることを考える。

湯川行列 $Y := vy$ に対して特異値分解^{*1}を行う。即ち、2 つの unitary 行列 Ψ, Φ および $m_i \geq 0$ によって

$$Y = \Phi^\dagger \begin{pmatrix} m_1 & 0 & 0 \\ 0 & m_2 & 0 \\ 0 & 0 & m_3 \end{pmatrix} \Psi =: \Phi^\dagger M \Psi \quad (2.56)$$

^{*1} $A^\dagger A$ および AA^\dagger は Hermite 行列であるため、unitary 行列により対角化可能であり、固有値は全て非負である。

と展開する。これを用いて

$$Q^1 \mapsto \Psi_u^\dagger Q^1, \quad Q^2 \mapsto \Psi_d^\dagger Q^2, \quad L \mapsto \Psi_e^\dagger L, \quad U \mapsto \Phi_u^\dagger U, \quad D \mapsto \Phi_d^\dagger D, \quad E \mapsto \Phi_e^\dagger E \quad (2.57)$$

と置き換えてやることによって、湯川項が対角化できる：

$$\mathcal{L}_{\text{湯川}} = \bar{U} M_u Q^1 + D M_d Q^2 + E M_e L^2 + \text{H. c.} \quad (2.58)$$

しかし、 Q^1 と Q^2 を別の方法で変換したため、 W boson との結合の項に歪みが生じる：

$$\mathcal{L}_{QQW} = \frac{g_2}{\sqrt{2}} (\bar{Q}^1 \Psi_u \quad \bar{Q}^2 \Psi_d) \begin{pmatrix} 0 & W^+ \\ W^- & 0 \end{pmatrix} \begin{pmatrix} \Psi_u^\dagger Q^1 \\ \Psi_d^\dagger Q^2 \end{pmatrix} \quad (2.59)$$

$$= \frac{g_2}{\sqrt{2}} [\bar{Q}^2 W^- X Q^1 + \bar{Q}^1 W^+ X^\dagger Q^2] \quad (2.60)$$

ここで $X := \Psi_d \Psi_u^\dagger$ である。

この項は明らかに flavor violating であり、また CP violating でもある。

CP 変換により、spinor は

$$\psi \mapsto -i\eta^* (\bar{\psi} \gamma^2)^T, \quad \bar{\psi} \mapsto i\eta (\gamma^2 \psi)^T \quad (2.61)$$

のように変換される。

\mathcal{L}_{QQW} 以外の項は、例えば

$$\bar{Q}(i\not{\partial}) P_L Q \mapsto (\gamma^2 Q)^T (i\not{\partial}^P) P_L (\bar{Q} \gamma^2)^T \quad (2.62)$$

$$= i(\gamma^2 Q)^T (\partial_\mu^P \bar{Q} \gamma^2 P_L \gamma^{\mu T})^T \quad (2.63)$$

$$= -i(\partial_\mu^P \bar{Q} \gamma^2 P_L \gamma^{\mu T} \gamma^2 Q) \quad (2.64)$$

$$= i(\bar{Q} \gamma^2 P_L (\gamma^2 \gamma^0 \gamma^\mu \gamma^0 \gamma^2) \gamma^2 \partial_\mu^P Q) \quad (2.65)$$

$$= i(\bar{Q} \gamma^\mu \partial_\mu P_L Q) \quad (2.66)$$

のように CP 不変であるが、 \mathcal{L}_{QQW} の項は

$$\bar{Q}^2 W^- X P_L Q^1 \mapsto (\gamma^2 Q^2)^T (-W^{+P}) X P_L (\bar{Q}^1 \gamma^2)^T \quad (2.67)$$

$$= -W_\mu^{+P} (\gamma^2 Q^2)^T (\bar{Q}^1 X^\dagger \gamma^2 P_L \gamma^{\mu T})^T \quad (2.68)$$

$$= (\bar{Q}^1 X^\dagger W^+ P_L Q^2) \quad (2.69)$$

のように変換する。 CP 変換により結合定数の符号が変わるべきであることも踏まえると CP の保存は $X^T = X^\dagger$ と同値である。即ち CP の保存は、 X が実行列であることと同値である。

以上より，標準模型の Lagrangian は

$$\mathcal{L} = \mathcal{L}_{\text{gauge}} \quad (2.70)$$

$$\text{【質量項】} + m_W^2 W^+ W^- + \frac{m_Z^2}{2} Z^2 \quad (2.71)$$

$$+ \bar{U} M_u P_L Q^1 + \bar{D} M_d P_L Q^2 + \bar{E} M_e P_L L^2 + \text{H. c.} \quad (2.72)$$

$$\text{【Higgs Field】} + (\partial_\mu h)^2 - \frac{1}{2} \mu^2 h^2 - \sqrt{\frac{\lambda}{2}} m_h h^3 - \frac{1}{4} \lambda h^4 \quad (2.73)$$

$$\text{【Higgs との結合】} + \frac{v g_2^2}{4} W^+ W^- h + \frac{v (g_1^2 + g_2^2)}{8} Z^2 h \quad (2.74)$$

$$+ \frac{g_2^2}{4} W^+ W^- h^2 + \frac{g_1^2 + g_2^2}{8} Z^2 h^2 \quad (2.75)$$

$$+ \frac{1}{v} \bar{U} M_u P_L Q^1 h + \frac{1}{v} \bar{D} M_d P_L Q^2 h + \frac{1}{v} \bar{E} M_e P_L L^2 h + \text{H. c.} \quad (2.76)$$

$$\text{【SU(3) および微分項】} + \bar{Q} (i \not{\partial} + g_3 \not{G}) P_L Q + \bar{U} (i \not{\partial} + g_3 \not{G}) P_R U + \bar{D} (i \not{\partial} + g_3 \not{G}) P_R D \quad (2.77)$$

$$+ \bar{L} (i \not{\partial}) P_L L + \bar{E} (i \not{\partial}) P_R E \quad (2.78)$$

$$\text{【W boson】} + \frac{g_2}{\sqrt{2}} [\bar{Q}^2 W^- X P_L Q^1 + \bar{Q}^1 W^+ X^\dagger P_L Q^2] \quad \text{【CP and flavor violating!】} \quad (2.79)$$

$$+ \bar{L} \frac{g_2}{\sqrt{2}} (W^+ T^+ + W^- T^-) P_L L \quad (2.80)$$

$$\text{【A\&Z}^0 \text{ boson】} + \bar{Q} \left[\left(T^3 + \frac{1}{6} \right) |e| A + \left(\frac{|e|c}{s} T^3 - \frac{|e|s}{6c} \right) Z^0 \right] P_L Q \quad (2.81)$$

$$+ \bar{U} \left(\frac{2}{3} |e| A - \frac{2|e|s}{3c} Z \right) P_R U \quad (2.82)$$

$$+ \bar{D} \left(-\frac{1}{3} |e| A + \frac{|e|s}{3c} Z \right) P_R D \quad (2.83)$$

$$+ \bar{L} \left[\left(T^3 - \frac{1}{2} \right) |e| A + \left(\frac{|e|c}{s} T^3 + \frac{|e|s}{2c} \right) Z^0 \right] P_L L \quad (2.84)$$

$$+ \bar{E} \left(-|e| A + \frac{|e|s}{c} Z \right) P_R E \quad (2.85)$$

となる。

2.5 CHIRAL NOTATION

以上の Lagrangian を chiral 表示で表すと、まず最初は

$$\mathcal{L} = (\text{Higgs terms}) + (\text{Gauge fields strength}) \quad (2.86)$$

$$+ Q_L^\dagger i \bar{\sigma}^\mu \left(\partial_\mu - i g_3 G_\mu - i g_2 W_\mu - \frac{1}{6} i g_1 B_\mu \right) Q_L \quad (2.87)$$

$$+ U_R^\dagger i \sigma^\mu \left(\partial_\mu - i g_3 G_\mu - \frac{2}{3} i g_1 B_\mu \right) U_R \quad (2.88)$$

$$+ D_R^\dagger i \sigma^\mu \left(\partial_\mu - i g_3 G_\mu + \frac{1}{3} i g_1 B_\mu \right) D_R \quad (2.89)$$

$$+ L_L^\dagger i \bar{\sigma}^\mu \left(\partial_\mu - i g_2 W_\mu + \frac{1}{2} i g_1 B_\mu \right) L_L \quad (2.90)$$

$$+ E_R^\dagger i \sigma^\mu (\partial_\mu + i g_1 B_\mu) E_R \quad (2.91)$$

$$+ U_R^\dagger y_u H Q_L + D_R^\dagger y_d H^\dagger Q_L + E_R^\dagger y_e H^\dagger L_L + \text{H. c.} \quad (2.92)$$

$$= (\text{Higgs terms}) + (\text{Gauge fields strength}) \quad (2.93)$$

$$+ i Q_L^\dagger \bar{\sigma}^\mu \partial_\mu Q_L + i U_R \bar{\sigma}^\mu \partial_\mu U_R^\dagger + i D_R \bar{\sigma}^\mu \partial_\mu D_R^\dagger + i L_L^\dagger \bar{\sigma}^\mu \partial_\mu L_L + i E_R \bar{\sigma}^\mu \partial_\mu E_R^\dagger \quad (2.94)$$

$$+ g_3 \left(Q_L^\dagger \bar{\sigma}^\mu G_\mu Q_L + U_R^\dagger \bar{\sigma}^\mu G_\mu U_R + D_R^\dagger \bar{\sigma}^\mu G_\mu D_R \right) \quad (2.95)$$

$$+ g_2 \left(Q_L^\dagger \bar{\sigma}^\mu W_\mu Q_L + L_L^\dagger \bar{\sigma}^\mu W_\mu L_L \right) \quad (2.96)$$

$$+ g_1 \left(\frac{1}{6} Q_L^\dagger \bar{\sigma}^\mu B_\mu Q_L + \frac{2}{3} U_R^\dagger \bar{\sigma}^\mu B_\mu U_R - \frac{1}{3} D_R^\dagger \bar{\sigma}^\mu B_\mu D_R - \frac{1}{2} L_L^\dagger \bar{\sigma}^\mu B_\mu L_L - E_R^\dagger \bar{\sigma}^\mu B_\mu E_R \right) \quad (2.97)$$

$$+ U_R^\dagger y_u H Q_L + D_R^\dagger y_d H^\dagger Q_L + E_R^\dagger y_e H^\dagger L_L + \text{H. c.} \quad (2.98)$$

であり、そして最終的には

$$\mathcal{L} = (\text{Gauge bosons and Higgs}) \quad (2.99)$$

$$+ i Q_L^\dagger \bar{\sigma}^\mu \partial_\mu Q_L + i U_R \bar{\sigma}^\mu \partial_\mu U_R^\dagger + i D_R \bar{\sigma}^\mu \partial_\mu D_R^\dagger + i L_L^\dagger \bar{\sigma}^\mu \partial_\mu L_L + i E_R \bar{\sigma}^\mu \partial_\mu E_R^\dagger \quad (2.100)$$

$$+ g_3 \left(Q_L^\dagger \bar{\sigma}^\mu G_\mu Q_L + U_R^\dagger \bar{\sigma}^\mu G_\mu U_R + D_R^\dagger \bar{\sigma}^\mu G_\mu D_R \right) \quad (2.101)$$

$$+ m_u (u_R^\dagger u_L + u_L^\dagger u_R) + (\text{quarks}) + m_e (e_R^\dagger e_L + e_L^\dagger e_R) + (\text{leptons}) \quad (2.102)$$

$$+ \frac{m_u}{v} (u_R^\dagger u_L + u_L^\dagger u_R) h + (\text{quarks}) + \frac{m_e}{v} (e_R^\dagger e_L + e_L^\dagger e_R) h + (\text{leptons}) \quad (2.103)$$

$$+ \frac{g_2}{\sqrt{2}} \left[(d_L^\dagger \ s_L^\dagger \ b_L^\dagger) \bar{\sigma}^\mu W_\mu^- X \begin{pmatrix} u_L \\ c_L \\ t_L \end{pmatrix} + (u_L^\dagger \ c_L^\dagger \ t_L^\dagger) \bar{\sigma}^\mu W_\mu^+ X^\dagger \begin{pmatrix} d_L \\ s_L \\ b_L \end{pmatrix} \right] \quad (2.104)$$

$$+ \frac{g_2}{\sqrt{2}} \left[\nu_e^\dagger \bar{\sigma}^\mu W_\mu^+ e_L + e_L^\dagger \bar{\sigma}^\mu W_\mu^- \nu_e \right] \quad (2.105)$$

$$+ |e| \left[\frac{2}{3} u_L^\dagger \bar{\sigma}^\mu A_\mu u_L - \frac{1}{3} d_L^\dagger \bar{\sigma}^\mu A_\mu d_L + \frac{2}{3} u_R^\dagger \sigma^\mu A_\mu u_R - \frac{1}{3} d_R^\dagger \sigma^\mu A_\mu d_R + (\text{quarks}) \right. \quad (2.106)$$

$$\left. - e_L^\dagger \bar{\sigma}^\mu A_\mu e_L - e_R^\dagger \sigma^\mu A_\mu e_R + (\text{leptons}) \right] \quad (2.107)$$

$$+ \frac{|e|s}{c} \left[\left(\frac{c^2}{2s^2} - \frac{1}{6} \right) u_L^\dagger \bar{\sigma}^\mu Z_\mu u_L - \left(\frac{c^2}{2s^2} + \frac{1}{6} \right) d_L^\dagger \bar{\sigma}^\mu Z_\mu d_L - \frac{2}{3} u_R^\dagger \sigma^\mu Z_\mu u_R + \frac{1}{3} d_R^\dagger \sigma^\mu Z_\mu d_R \right. \quad (2.108)$$

$$\left. + \left(\frac{c^2}{2s^2} + \frac{1}{2} \right) \nu_e^\dagger \bar{\sigma}^\mu Z_\mu \nu_e - \left(\frac{c^2}{2s^2} - \frac{1}{2} \right) e_L^\dagger \bar{\sigma}^\mu Z_\mu e_L + e_R^\dagger \sigma^\mu Z_\mu e_R + (\text{others}) \right] \quad (2.109)$$

となる。

実験値

実験値としては, $v = (\sqrt{2}G_F)^{-1/2} \sim 246 \text{ GeV} \sim 174\sqrt{2} \text{ GeV}$ および $M_w \sim 37.3 \text{ GeV} / \sin \theta_w$ が知られている。**TODO: 拡充**

3 楊-Mills Theory

3.1 U(1) THEORY

3.1.1 General SU(N)

$$\mathcal{L} = \quad (3.1)$$

$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu + \frac{g}{i} [A_\mu, A_\nu] \quad (3.2)$$

$$(3.3)$$

Gauge Transformation

For any Lie group G ,

$$V : \mathbb{R}^{1,3} \rightarrow G \quad (3.4)$$

$$A_\mu \mapsto V \left(A_\mu + \frac{i}{g} \partial_\mu \right) V^{-1} \quad (3.5)$$

$$F_{\mu\nu} \mapsto V F_{\mu\nu} V^{-1} \quad (3.6)$$

If the gauge group G is **compact**, it has a finite-dimensional unitary representation.

$$\text{For } t^a : \text{hermitian representation,} \quad (3.7)$$

$$[t^a, t^b] = i f^{ab}_c t^c \quad \text{and} \quad f \in \mathbb{R} \quad (3.8)$$

$$0 = f^D_{ab} f^E_{Dc} + f^D_{ca} f^E_{Db} + f^D_{bc} f^E_{Da} \quad (3.9)$$

$$V = \exp[i \alpha^a T^a] \quad \text{for } \alpha^a \in \mathbb{R} \quad (3.10)$$

$$(3.11)$$

with generators $\{T^a\}$ written in an hermitian representation,

$$A_\mu \mapsto V \left(A_\mu + \frac{1}{g} (\partial_\mu \alpha^a) T^a \right) V^{-1} \quad (3.12)$$

$$\simeq A_\mu + \frac{1}{g} (\partial_\mu \alpha^a) T^a + i \alpha^a A_\mu^b [T^a, T^b] \quad (3.13)$$

$$D_\mu = \partial_\mu - i g A_\mu^a T^a \quad (\text{for appropriate representation}) \quad (3.14)$$

4 Spinor

$\eta^{\mu\nu} = (-, +, +, +)$ case

Grassmann Number : $(ab)^\dagger = b^\dagger a^\dagger$ for $a, b \in \mathbb{G}$

: \implies for $a, b \in \mathbb{G}^{\mathbb{R}}, ab \in \mathbf{i} \mathbb{G}^{\mathbb{R}}$

γ matrix : $\{\gamma^\mu, \gamma^\nu\} = 2\eta^{\mu\nu} \cdot \mathbf{1}$

: $\gamma^{\mu\nu} = \frac{1}{2}(\gamma^\mu \gamma^\nu - \gamma^\nu \gamma^\mu)$ etc...

: $(\mathbf{i} \gamma^0)^\dagger := \mathbf{i} \gamma^0, \quad \gamma^{i\dagger} := \gamma^i$

Dirac Conjugate : $\bar{\psi} = \mathbf{i} \psi^\dagger \gamma^0$

5 Mathematics

5.1 GROUP THEORY

5.1.1 Lie Group and Lie Algebra

- G が group である ... 積が定義されており, 積閉・単位元・逆元の 3 条件を満たす。
- 群 G が Lie group である ... G が同時に C^∞ 多様体であり, 積演算と逆元写像が共に C^∞ 級である。
- Lie 群 G が COMPLEX Lie group である ... 積演算と逆元写像が共に正則写像である。
- Vector 空間 g が Lie algebra である ... 括弧積が定義されており, 線型性・反対称性・Jacobi 恒等式を満たす。
- Lie 群 G の単位元における接空間を, G の Lie algebra \mathfrak{g} という。
 - \mathfrak{g} は G の左不変な vector 場全体である。
 - \mathfrak{g} は vector 場の括弧積の下で Lie algebra となる。
- G として有限次元 Lie 群を考えると,
 - その Lie 代数の基底 B_i に対して structure constant c が $[B_i, B_j] = c_{ij}^k B_k$ として定義できる。

* * *

- Compact Lie 群は線型 Lie 群である。
- G として Linear group $GL(n; \mathbb{R})$ を考えると,
 - その Lie 代数は n 次実正方形行列全体となる。
 - Vector 場の括弧積は commutation relation $[X, Y] = XY - YX$ となる。
- Lie 群は, $GL(n; \mathbb{C})$ の部分 Lie 群と局所同型になるような位相群でかつ連結成分が高々可算個であるものである。

以下では, Lie 群として $GL(n; \mathbb{R})$ の部分群を考えることにし, Lie 代数の元を行列により表現する。

5.1.2 Matrix Representation

- Lie 群 G の Lie 代数の基底の組を, G の generators と言う。
- $GL(n; \mathbb{R})$ の元は n 次元行列で表せる。
- Lie 群 G の生成子 $\{T_i\}$ に対し, 以下の 2 つは共に G の単位元近傍の局所座標系を与える。

$$(x_1, \dots, x_m) \mapsto e^{x_1 T_1 + \dots + x_m T_m} \qquad (x_1, \dots, x_m) \mapsto e^{x_1 T_1} \dots e^{x_m T_m} \quad (5.1)$$

- Lie 群 G が compact である ...
 1. 多様体 G が compact である。 **TODO: これは何故同値なのか?**
 2. G の生成子 $\{T_i\}$ を, $\text{Tr}(T_i T_j) = k \delta_{ij}$ かつ $k > 0$ となるように取り替えることができる。

【この基底の下では構造定数が完全反対称になる。】

- Compact 群 G は, unitary representation を持つ。
故に, 単位元の近傍では有限個の Hermitian matrix T^i と parameters $x^i \in \mathbb{R}$ により, G の元を

$$e^{i x^i T^i} \quad (5.2)$$

と表すことが出来る。

5.1.3 結論

Compact Lie 群の元のうち, 単位元近傍にあるもの V は,
Hermitian Representation

$$V = \exp(\mathrm{i} x^i T^i) \quad \text{where} \quad T^i : \text{Hermitian Matrix}, \quad x^i \in \mathbb{R}, \\ [T^i, T^j] = \mathrm{i} f^{ijk} T^k, \quad \mathrm{Tr}(T^i T^j) = \lambda \delta^{ij} > 0; \quad f \in \mathbb{R}$$

Real Representation

$$V = \exp(x^i R^i) \quad \text{where} \quad R^i : \text{Real Matrix}, \quad x^i \in \mathbb{R}, \\ [R^i, R^j] = -f^{ijk} R^k, \quad \mathrm{Tr}(R^i R^j) = -\lambda \delta^{ij} < 0; \quad f \in \mathbb{R}$$

と表すことが出来る。

6 Supersymmetry

6.1 GENERAL RELATIONS

Coordinates

$$y := x + i\theta\sigma\bar{\theta}, \quad y^+ := x - i\theta\sigma\bar{\theta} \quad (6.1)$$

6.2 CHIRAL SUPERFIELDS

Definition

$$\bar{D}_{\dot{\alpha}}\Phi = 0 \quad (6.2)$$

Explicit Expression

$$\Phi = \phi(y) + \sqrt{2}\theta\psi(y) + \theta\theta F(y) \quad (6.3)$$

$$= \phi(x) + i\theta\sigma^\mu\bar{\theta}\partial_\mu\phi(x) + \frac{1}{4}\theta\theta\bar{\theta}\bar{\theta}\partial^2\phi(x) + \sqrt{2}\theta\psi(x) - \frac{i}{\sqrt{2}}\theta\theta\partial_\mu\psi(x)\sigma^\mu\bar{\theta} + \theta\theta F(x) \quad (6.4)$$

$$\Phi^\dagger = \phi^*(x) - i\theta\sigma^\mu\bar{\theta}\partial_\mu\phi^*(x) + \frac{1}{4}\theta\theta\bar{\theta}\bar{\theta}\partial^2\phi^*(x) + \sqrt{2}\bar{\theta}\bar{\psi}(x) + \frac{i}{\sqrt{2}}\bar{\theta}\bar{\theta}\sigma^\mu\partial_\mu\bar{\psi}(x) + \bar{\theta}\bar{\theta}F^*(x) \quad (6.5)$$

Changing Bases

$$\phi(y) = \phi(x) + i\theta\sigma^\mu\bar{\theta}\partial_\mu\phi(x) + \frac{1}{4}\theta\theta\bar{\theta}\bar{\theta}\partial^2\phi(x) \quad (6.6)$$

$$= \phi(y^+) + 2i\theta\sigma^\mu\bar{\theta}\partial_\mu\phi(y^+) + \theta\theta\bar{\theta}\bar{\theta}\partial^2\phi(y^+) \quad (6.7)$$

$$\phi(y^+) = \phi(x) - i\theta\sigma^\mu\bar{\theta}\partial_\mu\phi(x) + \frac{1}{4}\theta\theta\bar{\theta}\bar{\theta}\partial^2\phi(x) \quad (6.8)$$

$$= \phi(y) - 2i\theta\sigma^\mu\bar{\theta}\partial_\mu\phi(y) + \theta\theta\bar{\theta}\bar{\theta}\partial^2\phi(y) \quad (6.9)$$

$$\phi(x) = \phi(y) - i\theta\sigma^\mu\bar{\theta}\partial_\mu\phi(y) + \frac{1}{4}\theta\theta\bar{\theta}\bar{\theta}\partial^2\phi(y) \quad (6.10)$$

$$= \phi(y^+) + i\theta\sigma^\mu\bar{\theta}\partial_\mu\phi(y^+) + \frac{1}{4}\theta\theta\bar{\theta}\bar{\theta}\partial^2\phi(y^+) \quad (6.11)$$

Product of Chiral Superfields

$$\begin{aligned} \Phi_i^\dagger\Phi_j \text{ (in } x\text{-basis)} &= \phi_i^*\phi_j + \sqrt{2}\phi_i^*\theta\psi_j + \sqrt{2}\bar{\theta}\bar{\psi}_i\phi_j + \theta\theta\phi_i^*F_j + \bar{\theta}\bar{\theta}F_i^*\phi_j \\ &\quad + i\theta\sigma^\mu\bar{\theta}(\phi_i^*\partial_\mu\phi_j - \partial_\mu\phi_i^*\phi_j) + 2\bar{\theta}\bar{\psi}_i\theta\psi_j \\ &\quad - \frac{i}{\sqrt{2}}\theta\theta(\phi_i^*\partial_\mu\psi_j - \partial_\mu\phi_i^*\psi_j)\sigma^\mu\bar{\theta} + \sqrt{2}\theta\theta\bar{\theta}\bar{\psi}_iF_j \\ &\quad + \frac{i}{\sqrt{2}}\bar{\theta}\bar{\theta}\theta\sigma^\mu(\partial_\mu\bar{\psi}_i\phi_j - \bar{\psi}_i\partial_\mu\phi_j) + \sqrt{2}\bar{\theta}\bar{\theta}F_i^*\theta\psi_j \end{aligned} \quad (6.12)$$

$$\begin{aligned} &+ \theta\theta\bar{\theta}\bar{\theta}\left[F_i^*F_j + \frac{1}{4}\phi_i^*\partial^2\phi_j + \frac{1}{4}\partial^2\phi_i^*\phi_j - \frac{1}{2}\partial_\mu\phi_i^*\partial_\mu\phi_j + \frac{i}{2}\partial_\mu\bar{\psi}_i\bar{\sigma}^\mu\psi_j - \frac{i}{2}\bar{\psi}_i\bar{\sigma}^\mu\partial_\mu\psi_j\right] \\ &\rightsquigarrow \phi_i^*\phi_j + \sqrt{2}\phi_i^*\theta\psi_j + \sqrt{2}\bar{\theta}\bar{\psi}_i\phi_j + \theta\theta\phi_i^*F_j + \bar{\theta}\bar{\theta}F_i^*\phi_j \\ &\quad + i\theta\sigma^\mu\bar{\theta}(\phi_i^*\partial_\mu\phi_j - \partial_\mu\phi_i^*\phi_j) + 2\bar{\theta}\bar{\psi}_i\theta\psi_j \\ &\quad + \sqrt{2}\theta\theta\bar{\theta}(\bar{\psi}_iF_j - i\bar{\sigma}^\mu\psi_j\partial_\mu\phi_i^*) + \sqrt{2}\bar{\theta}\bar{\theta}\theta(\psi_jF_i^* - i\sigma^\mu\bar{\psi}_i\partial_\mu\phi_j) \\ &\quad + \theta\theta\bar{\theta}\bar{\theta}[F_i^*F_j - \partial_\mu\phi_i^*\partial_\mu\phi_j - i\bar{\psi}_i\bar{\sigma}^\mu\partial_\mu\psi_j] \end{aligned} \quad (6.13)$$

$$\Phi_i \Phi_j (\text{in } y\text{-basis}) = \phi_i \phi_j + \sqrt{2}\theta [\psi_i \phi_j + \phi_i \psi_j] + \theta\theta [\phi_i F_j + F_i \phi_j - \psi_i \psi_j] \quad (6.14)$$

$$\begin{aligned} \Phi_i \Phi_j \Phi_k (\text{in } y\text{-basis}) &= \phi_i \phi_j \phi_k + \sqrt{2}\theta [\psi_i \phi_j \phi_k + \phi_i \psi_j \phi_k + \phi_i \phi_j \psi_k] \\ &\quad + \theta\theta [F_i \phi_j \phi_k + \phi_i F_j \phi_k + \phi_i \phi_j F_k - \psi_i \psi_j \phi_k - \psi_i \phi_j \psi_k - \phi_i \psi_j \psi_k] \end{aligned} \quad (6.15)$$

Note that products of chiral superfields $\Phi_1 \Phi_2 \dots$ are again chiral superfields.

Superpotential

$$W = \int d^2\theta \left[\lambda_i \Phi_i + \frac{1}{2} m_{ij} \Phi_i \Phi_j + \frac{1}{3} y_{ijk} \Phi_i \Phi_j \Phi_k \right] \quad (6.16)$$

$$\begin{aligned} &= \lambda_i F_i + \frac{1}{2} m_{ij} (\phi_i F_j + F_i \phi_j - \psi_i \psi_j) \\ &\quad + \frac{1}{3} y_{ijk} (F_i \phi_j \phi_k + \phi_i F_j \phi_k + \phi_i \phi_j F_k - \psi_i \psi_j \phi_k - \psi_i \phi_j \psi_k - \phi_i \psi_j \psi_k) \\ &\quad (\text{In } x\text{-basis, since we have omitted all } \bar{\theta}\text{s.}) \end{aligned} \quad (6.17)$$

6.3 VECTOR SUPERFIELDS

6.3.1 Abelian Case

General Definitions

Vector Superfields : $V = V^\dagger$

Gauge Transf. : $V \rightarrow V + \Phi + \Phi^\dagger$

Field Strength : $W_\alpha = -\frac{1}{4} \bar{D} \bar{D} D_\alpha V$; $\bar{W}_{\dot{\alpha}} = -\frac{1}{4} D D \bar{D}_{\dot{\alpha}} V$

Lagrangian : $\mathcal{L} = \frac{1}{4} \left(W^\alpha W_\alpha \Big|_{\theta\theta} + \bar{W}_{\dot{\alpha}} \bar{W}^{\dot{\alpha}} \Big|_{\bar{\theta}\bar{\theta}} \right)$

Explicit Expression

$$\begin{aligned} V &= C(x) + i\theta\chi(x) - i\bar{\theta}\bar{\chi}(x) \\ &\quad + \frac{i}{2}\theta\theta [M(x) + iN(x)] - \frac{i}{2}\bar{\theta}\bar{\theta} [M(x) - iN(x)] - \theta\sigma^\mu\bar{\theta}A_\mu(x) \\ &\quad + i\theta\theta\bar{\theta} \left[\bar{\lambda}(x) + \frac{i}{2}\bar{\sigma}^\mu\partial_\mu\chi(x) \right] - i\bar{\theta}\bar{\theta}\theta \left[\lambda(x) + \frac{i}{2}\sigma^\mu\partial_\mu\bar{\chi}(x) \right] \\ &\quad + \frac{1}{2}\theta\theta\bar{\theta}\bar{\theta} \left[D(x) + \frac{1}{2}\partial^2 C(x) \right] \end{aligned} \quad (6.18)$$

In Wess-Zumino gauge,

$$V \rightarrow -\theta\sigma^\mu\bar{\theta}A_\mu(x) + i\theta\theta\bar{\theta}\bar{\lambda}(x) - i\bar{\theta}\bar{\theta}\theta\lambda(x) + \frac{1}{2}\theta\theta\bar{\theta}\bar{\theta}D(x) \quad (6.19)$$

$$= -\theta\sigma^\mu\bar{\theta}A_\mu(y) + i\theta\theta\bar{\theta}\bar{\lambda}(y) - i\bar{\theta}\bar{\theta}\theta\lambda(y) + \frac{1}{2}\theta\theta\bar{\theta}\bar{\theta} [D(y) - i\partial_\mu A^\mu(y)] \quad (6.20)$$

$$= -\theta\sigma^\mu\bar{\theta}A_\mu(y^+) + i\theta\theta\bar{\theta}\bar{\lambda}(y^+) - i\bar{\theta}\bar{\theta}\theta\lambda(y^+) + \frac{1}{2}\theta\theta\bar{\theta}\bar{\theta} [D(y^+) + i\partial_\mu A^\mu(y^+)] \quad (6.21)$$

Field Strength

$$W_\alpha = -\frac{1}{4} \bar{D} \bar{D} e^{2gV} D_\alpha e^{-2gV} \quad \text{where } V = V^a T^a \quad (6.22)$$

Field Strength

Defining all component fields as including generators and coupling constants,

$$\begin{aligned} e^{\pm 2gV^a T^a} &\rightarrow e^{\pm 2V} \\ &= 1 \mp 2\theta\sigma^\mu\bar{\theta}A_\mu(x) \mp 2i [\bar{\theta}\bar{\theta}\theta\lambda(x) - \theta\theta\bar{\theta}\bar{\lambda}(x)] + \theta\theta\bar{\theta}\bar{\theta} [-A^\mu(x)A_\mu(x) - D(x)] \end{aligned} \quad (6.23)$$

Therefore, in y^+ -basis,

$$\begin{aligned} D_\alpha e^{-2V} &= \frac{\partial}{\partial\theta^\alpha} \left\{ 1 + 2\theta\sigma^\mu\bar{\theta}A_\mu + 2i [\bar{\theta}\bar{\theta}\theta\lambda - \theta\theta\bar{\theta}\bar{\lambda}] + \theta\theta\bar{\theta}\bar{\theta} [A^\mu A_\mu - D - i\partial_\mu A^\mu] \right\} \\ &= 2(\sigma^\mu\bar{\theta})_\alpha A_\mu + 2i\bar{\theta}\bar{\theta}\lambda_\alpha - 4i\theta_\alpha\bar{\theta}\bar{\lambda} + 2\theta_\alpha\bar{\theta}\bar{\theta} [A^\mu A_\mu - D - i\partial_\mu A^\mu] \end{aligned} \quad (6.24)$$

$$\begin{aligned} e^{2V} D_\alpha e^{-2V} &= \left\{ 1 - 2\theta\sigma^\mu\bar{\theta}A_\mu - 2i [\bar{\theta}\bar{\theta}\theta\lambda - \theta\theta\bar{\theta}\bar{\lambda}] \right\} D_\alpha e^{-2V} \\ &= 2(\sigma^\mu\bar{\theta})_\alpha A_\mu + 2i\bar{\theta}\bar{\theta}\lambda_\alpha - 4i\theta_\alpha\bar{\theta}\bar{\lambda} + 2\theta_\alpha\bar{\theta}\bar{\theta} [A^\mu A_\mu - D - i\partial_\mu A^\mu] \\ &\quad - 2\theta\sigma^\mu\bar{\theta}A_\mu [2(\sigma^\mu\bar{\theta})_\alpha A_\mu - 4i\theta_\alpha\bar{\theta}\bar{\lambda}] \\ &\quad + 2i\theta\theta\bar{\theta}\bar{\lambda} \cdot 2(\sigma^\mu\bar{\theta})_\alpha A_\mu \\ &= 2(\sigma^\mu\bar{\theta})_\alpha A_\mu - 4i\theta_\alpha\bar{\theta}\bar{\lambda} + 2\theta_\alpha\bar{\theta}\bar{\theta} [A^\mu A_\mu - D - i\partial_\mu A^\mu] + 2i\bar{\theta}\bar{\theta}\lambda_\alpha \\ &\quad - 2A_\mu A_\nu \bar{\theta}\bar{\theta}\epsilon_{\alpha\gamma}(\theta\sigma^\mu\bar{\sigma}^\nu)^\gamma - 4iA_\mu\theta\theta\bar{\theta}\bar{\theta}(\sigma^\mu\bar{\lambda})_\alpha \quad (\text{in } y^+\text{-basis}) \end{aligned} \quad (6.25)$$

$$\begin{aligned} &= 2(\sigma^\mu\bar{\theta})_\alpha A_\mu - 2i\bar{\theta}\bar{\theta}\epsilon_{\alpha\gamma}(\theta\sigma^\nu\bar{\sigma}^\mu)^\gamma\partial_\nu A_\mu - 4i\theta_\alpha\bar{\theta}\bar{\lambda} - 2\theta\theta\bar{\theta}\bar{\theta}(\sigma^\mu\partial_\mu\bar{\lambda})_\alpha \\ &\quad + 2\theta_\alpha\bar{\theta}\bar{\theta} [A^\mu A_\mu - D - i\partial_\mu A^\mu] + 2i\bar{\theta}\bar{\theta}\lambda_\alpha \\ &\quad - 2A_\mu A_\nu \bar{\theta}\bar{\theta}\epsilon_{\alpha\gamma}(\theta\sigma^\mu\bar{\sigma}^\nu)^\gamma - 4iA_\mu\theta\theta\bar{\theta}\bar{\theta}(\sigma^\mu\bar{\lambda})_\alpha \quad (\text{in } y\text{-basis}) \end{aligned} \quad (6.26)$$

In y -basis, $\bar{D}\bar{D} = 4 \cdot \frac{\partial}{\partial(\bar{\theta}\bar{\theta})}$. Therefore,

$$W_\alpha = -\frac{1}{4}\bar{D}\bar{D}e^{2V}D_\alpha e^{-2V} \quad (6.27)$$

$$\begin{aligned} &= 2(A_\mu A_\nu + i\partial_\nu A_\mu)(\sigma^\nu\bar{\sigma}^\mu\theta)_\alpha + 2\theta\theta(\sigma^\mu\partial_\mu\bar{\lambda})_\alpha \\ &\quad - 2\theta_\alpha [A^\mu A_\mu - D - i\partial_\mu A^\mu] - 2i\lambda_\alpha + 4iA_\mu\theta\theta(\sigma^\mu\bar{\lambda})_\alpha \quad (\text{in } y\text{-basis}) \end{aligned} \quad (6.28)$$

$$\begin{aligned} &= 2(A_\mu A_\nu + i\partial_\nu A_\mu)(\sigma^\nu\bar{\sigma}^\mu\theta)_\alpha + 2\theta\theta(\sigma^\mu\partial_\mu\bar{\lambda})_\alpha \\ &\quad - 2\theta_\alpha [A^\mu A_\mu - D - i\partial_\mu A^\mu] - 2i\lambda_\alpha + 4iA_\mu\theta\theta(\sigma^\mu\bar{\lambda})_\alpha \\ &\quad - i\theta\theta(\sigma^\nu\bar{\sigma}^\mu\sigma^\rho\bar{\theta})_\alpha(\partial_\rho A_\mu A_\nu + A_\mu\partial_\rho A_\nu + i\partial_\nu\partial_\rho A_\mu) \\ &\quad + i\theta\theta(\sigma^\rho\bar{\theta})_\alpha(\partial_\rho A^\mu A_\mu + A^\mu\partial_\rho A_\mu - \partial_\rho D - i\partial_\mu\partial_\rho A^\mu) \quad (\text{in } x\text{-basis}) \end{aligned} \quad (6.29)$$

$$\begin{aligned} W^\alpha &= 2(A_\mu A_\nu + i\partial_\nu A_\mu)(\theta\sigma^\mu\bar{\sigma}^\nu)^\alpha - 2\theta\theta(\partial_\mu\bar{\lambda}\bar{\sigma}^\mu)^\alpha \\ &\quad - 2\theta^\alpha [A^\mu A_\mu - D - i\partial_\mu A^\mu] - 2i\lambda^\alpha - 4iA_\mu\theta\theta(\bar{\lambda}\bar{\sigma}^\mu)^\alpha \quad (\text{in } y\text{-basis}) \end{aligned} \quad (6.30)$$

$$\begin{aligned}
W^\alpha W_\alpha \Big|_{\theta\theta} &= 4(A_\mu A_\nu + i\partial_\nu A_\mu)(A_\rho A_\sigma + i\partial_\sigma A_\rho)(\theta\sigma^\mu\bar{\sigma}^\nu\sigma^\sigma\bar{\sigma}^\rho\theta) \\
&\quad - 4(A_\mu A_\nu + i\partial_\nu A_\mu)[A^\mu A_\mu - D - i\partial_\mu A^\mu](\theta\sigma^\mu\bar{\sigma}^\nu\theta) \\
&\quad + 4i\theta\theta(\partial_\mu\bar{\lambda}\bar{\sigma}^\mu\lambda) \\
&\quad - 4[A^\mu A_\mu - D - i\partial_\mu A^\mu](A_\mu A_\nu + i\partial_\nu A_\mu)(\theta\sigma^\nu\bar{\sigma}^\mu\theta) \\
&\quad + 4\theta\theta[A^\mu A_\mu - D - i\partial_\mu A^\mu][A^\mu A_\mu - D - i\partial_\mu A^\mu] \\
&\quad + 4i[A^\mu A_\mu - D - i\partial_\mu A^\mu]\theta\lambda \\
&\quad - 4i(\lambda\sigma^\nu\bar{\sigma}^\mu\theta)(A_\mu A_\nu + i\partial_\nu A_\mu) \\
&\quad - 4i\theta\theta(\lambda\sigma^\mu\partial_\mu\bar{\lambda}) \\
&\quad + 8\theta\theta\lambda^\alpha A_\mu(\sigma^\mu\bar{\lambda})_\alpha \\
&\quad - 8\theta\theta A_\mu(\bar{\lambda}\bar{\sigma}^\mu\lambda) \\
&= 4(\theta\sigma^\mu\bar{\sigma}^\nu\sigma^\sigma\bar{\sigma}^\rho\theta)[A_\mu A_\nu + i\partial_\nu A_\mu][A_\rho A_\sigma + i\partial_\sigma A_\rho] \\
&\quad - 4(\theta\sigma^\mu\bar{\sigma}^\nu\theta)\left\{A^\mu A_\mu - D - i\partial_\mu A^\mu, A_\mu A_\nu + i\partial_\nu A_\mu\right\} \\
&\quad + 4i\theta\theta(\partial_\mu\bar{\lambda}\bar{\sigma}^\mu\lambda - \lambda\sigma^\mu\partial_\mu\bar{\lambda}) \\
&\quad + 4\theta\theta[A^\mu A_\mu - D - i\partial_\mu A^\mu][A^\mu A_\mu - D - i\partial_\mu A^\mu] \\
&\quad + 8\theta\theta\lambda^\alpha A_\mu(\sigma^\mu\bar{\lambda})_\alpha \\
&\quad - 8\theta\theta A_\mu(\bar{\lambda}\bar{\sigma}^\mu\lambda) \quad (\text{in } y\text{-basis}) \\
&\rightsquigarrow 4(\eta^{\mu\nu}\eta^{\sigma\rho} - \eta^{\mu\sigma}\eta^{\nu\rho} + \eta^{\mu\rho}\eta^{\nu\sigma})[A_\mu A_\nu + i\partial_\nu A_\mu][A_\rho A_\sigma + i\partial_\sigma A_\rho] \\
&\quad - 4\left\{A^\mu A_\mu - D - i\partial_\mu A^\mu, A^\nu A_\nu + i\partial_\nu A^\nu\right\} \\
&\quad + 4[A^\mu A_\mu - D - i\partial_\mu A^\mu][A^\mu A_\mu - D - i\partial_\mu A^\mu] \\
&\quad + 4i(\partial_\mu\bar{\lambda}\bar{\sigma}^\mu\lambda - \lambda\sigma^\mu\partial_\mu\bar{\lambda}) + 8\lambda^\alpha A_\mu(\sigma^\mu\bar{\lambda})_\alpha - 8A_\mu(\bar{\lambda}\bar{\sigma}^\mu\lambda) \\
&= -2F^{\mu\nu}F_{\mu\nu} + 4(D + 2i\partial_\mu A^\mu)(D + 2i\partial_\mu A^\mu) \\
&\quad + 4i(\partial_\mu\bar{\lambda}\bar{\sigma}^\mu\lambda - \lambda\sigma^\mu\partial_\mu\bar{\lambda}) + 8\lambda^\alpha A_\mu(\sigma^\mu\bar{\lambda})_\alpha - 8A_\mu(\bar{\lambda}\bar{\sigma}^\mu\lambda), \tag{6.31}
\end{aligned}$$

$$\begin{aligned}
\bar{W}_{\dot{\alpha}}\bar{W}^{\dot{\alpha}} \Big|_{\bar{\theta}\bar{\theta}} &= -2F^{\mu\nu}F_{\mu\nu} + 4(D - 2i\partial_\mu A^\mu)(D - 2i\partial_\mu A^\mu) \\
&\quad + 4i(\partial_\mu\lambda\sigma^\mu\bar{\lambda} - \bar{\lambda}\bar{\sigma}^\mu\partial_\mu\lambda) + 8\bar{\lambda}^\alpha A_\mu(\bar{\sigma}^\mu\lambda)_\alpha - 8A_\mu(\lambda\sigma^\mu\bar{\lambda}), \tag{6.32}
\end{aligned}$$

where

$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu + i[A_\mu, A_\nu] \tag{6.33}$$

Defining the generators as ^{*2}

$$T^{a\dagger} = T^a, \quad [T^a, T^b] = if^{abc}T^c, \quad \text{Tr } T^a T^b = k\delta^{ab} \quad (k > 0), \quad \text{Tr}(\{T^a, T^b\}T^c) = 0, \tag{6.34}$$

$$\begin{aligned}
\frac{1}{16kg^2} \text{Tr} \left(W^\alpha W_\alpha \Big|_{\theta\theta} + \bar{W}_{\dot{\alpha}}\bar{W}^{\dot{\alpha}} \Big|_{\bar{\theta}\bar{\theta}} \right) &= -\frac{1}{4}F^{a\mu\nu}F_{\mu\nu}^a + \frac{1}{2}D^a D^a - 2\partial_\mu A^{a\mu}\partial_\nu A^{a\nu} \\
&\quad - \frac{1}{2}i(\lambda\sigma^\mu D_\mu\bar{\lambda} + \bar{\lambda}\bar{\sigma}^\mu D_\mu\lambda) \tag{6.35}
\end{aligned}$$

Here

$$F_{\mu\nu}^a = \partial_\mu A_\nu^a - \partial_\nu A_\mu^a - f^{abc}A_\mu^b A_\nu^c \tag{6.36}$$

$$D_\mu\lambda^a = \partial_\mu\lambda^a - gf^{abc}A_\mu^b\lambda^c \tag{6.37}$$

$$D_\mu\bar{\lambda}^a = \partial_\mu\bar{\lambda}^a - gf^{abc}A_\mu^b\bar{\lambda}^c \tag{6.38}$$

^{*2} The last one is anomaly-free condition, yielding $\text{Tr } T^a T^b T^c = ikgf^{abc}/2$.

Interaction Terms

$$e^{-2gV^aT^a} = 1 + 2g\theta\sigma^\mu\bar{\theta}A_\mu^aT^a + 2ig\bar{\theta}\bar{\theta}\theta\lambda^aT^a - 2ig\theta\theta\bar{\theta}\bar{\lambda}^aT^a - \theta\theta\bar{\theta}\bar{\theta} [g^2A^{a\mu}A_\mu^bT^aT^b + gD^aT^a] \quad (6.39)$$

$$\begin{aligned} \Phi^\dagger e^{-2gV^aT^a} \Phi \Big|_{\theta^4} &= \phi^* \theta\theta\bar{\theta}\bar{\theta} [-g^2A^{a\mu}A_\mu^bT^aT^b - gD^aT^a] \phi \\ &\quad + \sqrt{2}\phi^* [2ig\bar{\theta}\bar{\theta}\theta\lambda^aT^a] \theta\psi + \sqrt{2}\bar{\theta}\bar{\psi} [-2ig\theta\theta\bar{\theta}\bar{\lambda}^aT^a] \phi \\ &\quad + i\theta\sigma^\mu\bar{\theta} [\phi^*(2g\theta\sigma^\mu\bar{\theta}A_\mu^aT^a)\partial_\mu\phi - \partial_\mu\phi^*(2g\theta\sigma^\mu\bar{\theta}A_\mu^aT^a)\phi] \\ &\quad + 2\bar{\theta}\bar{\psi} [2g\theta\sigma^\mu\bar{\theta}A_\mu^aT^a] \theta\psi \\ &\quad + \theta\theta\bar{\theta}\bar{\theta} \left[F^*F + \frac{1}{4}\phi^*\partial^2\phi + \frac{1}{4}\partial^2\phi^*\phi - \frac{1}{2}\partial_\mu\phi^*\partial_\mu\phi + \frac{i}{2}\partial_\mu\bar{\psi}\bar{\sigma}^\mu\psi - \frac{i}{2}\bar{\psi}\bar{\sigma}^\mu\partial_\mu\psi \right] \\ &\rightsquigarrow F^*F - D_\mu\phi^*D^\mu\phi - i\bar{\psi}\bar{\sigma}^\mu D_\mu\psi \\ &\quad - gD^a(\phi^*T^a\phi) - \sqrt{2}ig\lambda^a(\phi^*T^a\psi) + \sqrt{2}ig\bar{\lambda}^a(\bar{\psi}T^a\phi) \end{aligned} \quad (6.40)$$

$$D_\mu\phi = (\partial_\mu - igA^{a\mu}T^a)\phi \quad (6.41)$$

$$D_\mu\phi^* = \partial_\mu\phi^* + igA_\mu^a(\phi^*T^a) \quad (6.42)$$

$$D_\mu\psi = (\partial_\mu - igA_\mu^aT^a)\psi \quad (6.43)$$

7 Supergravity

7.1 MINIMAL SUGRA LAGRANGIAN

Minimal SUGRA Lagrangian is constructed from supergravity multiplet $(e_a{}^\mu, \psi_\mu^\alpha, B_\mu, F_\phi)$.

$$\mathcal{L} = -\frac{M^2}{2}eR + e\epsilon^{\mu\nu\rho\sigma}\bar{\psi}_\mu\bar{\sigma}_\nu D_\rho\psi_\sigma \quad (7.1)$$

where

$$D_\mu\psi_\nu := \partial_\mu\psi_\nu + \frac{1}{2}\omega_\mu{}^{ab}\sigma_{ab}\psi_\nu \quad [\omega_\mu{}^{ab} : \text{“spin 接続”}] \quad (7.2)$$

$$e := \det e_a{}^\mu \quad (7.3)$$

$$M := 1/\sqrt{8\pi G} \quad (\text{Reduced Planck mass}) \quad (7.4)$$

$$R := e_a{}^\mu e_b{}^\nu R_{\mu\nu}{}^{ab} \quad (7.5)$$

$$R_{\mu\nu}{}^{ab} := \partial_\mu\omega_\nu{}^{ab} - \partial_\nu\omega_\mu{}^{ab} - \omega_\mu{}^{ac}\omega_{\nu c}{}^b + \omega_\nu{}^{ac}\omega_{\mu c}{}^b. \quad (7.6)$$

7.2 GENERAL SUGRA LAGRANGIAN

The components of general SUGRA Lagrangian is

$$\Phi_i = (\psi_i, \chi_i^\alpha, F_i), \quad V^{(a)} = (A_\mu^{(a)}, \lambda^{\alpha(a)}, D^{(a)}), \quad G = (e_\mu{}^a, \psi_\mu^\alpha, B_\mu, F_\phi), \quad (7.7)$$

and described with following functions:

- Kähler potential $K(\Phi, \Phi^*)$
 - Real function of chiral multiplets.
 - In global SUSY, $\int d^4\theta K$ yields kinetic terms of the chiral multiplet.
 - “Minimal Kähler” is (if no gauge interaction) $K = \Phi\Phi^\dagger$, which is

$$\int d^4\theta \Phi\Phi^* = \partial_\mu\phi^*\partial_\mu\phi + i\bar{\chi}\bar{\sigma}^\mu\partial_\mu\chi + F^*F. \quad (7.8)$$

- Super Potential $W(\Phi)$
- Gauge kinetic term $f_{(a)(b)}(\Phi)$
 - Some function which satisfies $f_{(a)(b)} = f_{(b)(a)}$.
 - $(a), (b), \dots$ are indices for adjoint representation of gauge group.
 - Minimal one is $f_{(a)(b)} \propto \delta_{(a)(b)}$.

$$\begin{aligned}
\mathcal{L} = & -\frac{1}{2}eR + eg_{ij*}D_\mu\phi^iD^\mu\phi^{*j} - \frac{1}{2}eg^2D_{(a)}D^{(a)} \\
& + i eg_{ij*}\bar{\chi}^j\bar{\sigma}^\mu D_\mu\chi^i + e\epsilon^{\mu\nu\rho\sigma}\bar{\psi}_\mu\bar{\sigma}_\nu D_\rho\psi_\sigma \\
& - \frac{1}{4}ef^R_{(ab)}F^{(a)}_{\mu\nu}F^{\mu\nu(b)} + \frac{1}{8}e\epsilon^{\mu\nu\rho\sigma}f^I_{(ab)}f^{(a)}_{\mu\nu}f^{(b)}_{\rho\sigma} \\
& + \frac{i}{2}e\left[\lambda_{(a)}\sigma^\mu D_\mu\bar{\lambda}^{(a)} + \bar{\lambda}_{(a)}\bar{\sigma}^\mu D_\mu\lambda^{(a)}\right] - \frac{1}{2}f^I_{(ab)}D_\mu\left[e\lambda^{(a)}\sigma^\mu\bar{\lambda}^{(b)}\right] \\
& + \sqrt{2}egg_{ij*}X_{(a)}^{*j}\chi^i\lambda^{(a)} + \sqrt{2}egg_{ij*}X_{(a)}^i\bar{\chi}^j\bar{\lambda}^{(a)} \\
& - \frac{i}{4}\sqrt{2}eg\partial_i f_{(ab)}D^{(a)}\chi^i\lambda^{(b)} + \frac{i}{4}\sqrt{2}eg\partial_{i*}f_{(ab)}^*D^{(a)}\bar{\chi}^i\bar{\lambda}^{(b)} \\
& - \frac{1}{4}\sqrt{2}e\partial_i f_{(ab)}\chi^i\sigma^{\mu\nu}\lambda^{(a)}F_{\mu\nu}^{(b)} - \frac{1}{4}\sqrt{2}e\partial_{i*}f_{(ab)}^*\bar{\chi}^i\bar{\sigma}^{\mu\nu}\bar{\lambda}^{(a)}F_{\mu\nu}^{(b)} \\
& + \frac{1}{2}egD_{(a)}\psi_\mu\sigma^\mu\bar{\lambda}^{(a)} - \frac{1}{2}egD_{(a)}\bar{\psi}_\mu\bar{\sigma}^\mu\lambda^{(a)} \\
& - \frac{1}{2}\sqrt{2}eg_{ij*}D_\nu\phi^{*j}\chi^i\sigma^\mu\bar{\sigma}^\nu\psi_\mu - \frac{1}{2}\sqrt{2}eg_{ij*}D_\nu\phi^i\bar{\chi}^j\bar{\sigma}^\mu\sigma^\nu\bar{\psi}_\mu \\
& - \frac{i}{4}e\left[\psi_\mu\sigma^{\nu\rho}\sigma^\mu\bar{\lambda}_{(a)} + \bar{\psi}_\mu\bar{\sigma}^{\nu\rho}\bar{\sigma}^\mu\lambda_{(a)}\right]\left[F_{\nu\rho}^{(a)} + \hat{F}_{\nu\rho}^{(a)}\right] \\
& + \frac{1}{4}eg_{ij*}\left[i\epsilon^{\mu\nu\rho\sigma}\psi_\mu\sigma_\nu\bar{\psi}_\rho + \psi_\mu\sigma^\sigma\bar{\psi}^\mu\right]\chi^i\sigma_\sigma\bar{\chi}^i \\
& - \frac{1}{8}e\left[g_{ij*}g_{kl*} - 2R_{ij*kl*}\right]\chi^i\chi^k\bar{\chi}^j\bar{\chi}^l \\
& + \frac{1}{16}e\left[2g_{ij*}f^R_{(ab)} + f^{R(cd)-1}\partial_i f_{(bc)}\partial_{j*}f_{(ad)}^*\right]\bar{\chi}^j\bar{\sigma}^\mu\chi^i\bar{\lambda}^{(a)}\bar{\sigma}_\mu\lambda^{(b)} \\
& + \frac{1}{8}e\nabla_i\partial_j f_{(ab)}\chi^i\chi^j\lambda^{(a)}\lambda^{(b)} + \frac{1}{8}e\nabla_{i*}\partial_{j*}f_{(ab)}^*\bar{\chi}^i\bar{\chi}^j\bar{\lambda}^{(a)}\bar{\lambda}^{(b)} \\
& + \frac{1}{16}ef^{R(cd)-1}\partial_i f_{(ac)}\partial_j f_{(bd)}\chi^i\lambda^{(a)}\chi^j\lambda^{(b)} \\
& + \frac{1}{16}ef^{R(cd)-1}\partial_{i*}f_{(ac)}^*\partial_{j*}f_{(bd)}^*\bar{\chi}^i\bar{\lambda}^{(a)}\bar{\chi}^j\bar{\lambda}^{(b)} \\
& - \frac{1}{16}eg^{ij*}\partial_i f_{(ab)}\partial_{j*}f_{(cd)}^*\lambda^{(a)}\lambda^{(b)}\bar{\lambda}^c\bar{\lambda}^{(d)} \\
& + \frac{3}{16}e\lambda_{(a)}\sigma^\mu\bar{\lambda}^{(a)}\lambda_{(b)}\sigma_\mu\bar{\lambda}^{(b)} \\
& + \frac{i}{4}\sqrt{2}e\partial_i f_{(ab)}\left[\chi^i\sigma^{\mu\nu}\lambda^{(a)}\psi_\mu\sigma_\nu\bar{\lambda}^{(b)} - \frac{1}{4}\bar{\psi}_\mu\bar{\sigma}^\mu\chi^i\lambda^{(a)}\lambda^{(b)}\right] \\
& + \frac{i}{4}\sqrt{2}e\partial_{i*}f_{(ab)}^*\left[\bar{\chi}^i\bar{\sigma}^{\mu\nu}\bar{\lambda}^{(a)}\bar{\psi}_\mu\bar{\sigma}_\nu\lambda^{(b)} - \frac{1}{4}\psi_\mu\sigma^\mu\bar{\chi}^i\bar{\lambda}^{(a)}\bar{\lambda}^{(b)}\right] \\
& - ee^{K/2}\left[W^*\psi_\mu\sigma^{\mu\nu}\psi_\nu + W\bar{\psi}_\mu\bar{\sigma}^{\mu\nu}\bar{\psi}_\nu\right] \\
& + \frac{i}{2}\sqrt{2}ee^{K/2}\left[D_iW\chi^i\sigma^\mu\bar{\psi}_\mu + D_{i*}W^*\bar{\chi}^i\bar{\sigma}^\mu\psi_\mu\right] \\
& - \frac{1}{2}ee^{K/2}\left[D_iD_jW\chi^i\chi^j + D_{i*}D_{j*}W^*\bar{\chi}^i\bar{\chi}^j\right] \\
& + \frac{1}{4}ee^{K/2}g^{ij*}\left[D_{j*}W^*\partial_i f_{(ab)}\lambda^{(a)}\lambda^{(b)} + D_iW\partial_{j*}f_{(ab)}^*\bar{\lambda}^{(a)}\bar{\lambda}^{(b)}\right] \\
& - ee^K\left[g^{ij*}(D_iW)(D_{j*}W^*) - 3W^*W\right]
\end{aligned} \tag{7.9}$$

付録 A Verbose Notes

A.1 POLARIZATION SUM

Firstly we focus on the single photon case $M = \epsilon_\mu^*(k)\epsilon_\nu'^*(k')M^{\mu\nu}$. Here we set $k = (E, 0, 0, E)$, and $\epsilon = (0, 1, 0, 0) \oplus (0, 0, 1, 0)$. Then

$$\sum_{\text{pol.}} |M|^2 = \sum_{\text{pol.}} \epsilon_\mu^*(k)\epsilon_\nu(k)M^\mu M^{\nu*} = |M^1|^2 + |M^2|^2, \quad (\text{A.1})$$

while

$$\eta_{\mu\nu}M^\mu M^{\nu*} = |M^1|^2 + |M^2|^2 \quad (\text{A.2})$$

for Ward identity $k_\mu M^\mu = 0$. Therefore the replacement

$$\sum_{\text{pol.}} \epsilon_\mu \epsilon_\nu' \rightarrow \eta_{\mu\nu} \quad (\text{A.3})$$

is valid.

Secondly we think about the double photons case ^{*3} $M = \epsilon_\mu^*(k)\epsilon_\nu'^*(k')M^{\mu\nu}$. Here we set

$$k = (E, 0, 0, E) \quad \epsilon = (0, 1, 0, 0) \oplus (0, 0, 1, 0) \quad (\text{A.4})$$

$$k' = (E, 0, 0, -E) \quad \epsilon' = (0, \cos \theta, \sin \theta, 0) \oplus (0, -\sin \theta, \cos \theta, 0). \quad (\text{A.5})$$

Then doing some simple calculations, we can get

$$\sum_{\text{pol.}} |M|^2 = |M^{11}|^2 + |M^{12}|^2 + |M^{21}|^2 + |M^{22}|^2. \quad (\text{A.6})$$

Nevertheless, naïve replacement does not work, because our Ward identities

$$k_\mu \epsilon_\nu'^*(k')M^{\mu\nu} = \epsilon_\mu^*(k)k_\nu' M^{\mu\nu} = 0 \quad (\text{A.7})$$

obviously does not help us. If we can omit ϵ s from these identities, that is if

$$k_\mu M^{\mu\nu} = k_\nu' M^{\mu\nu} = 0, \quad (\text{A.8})$$

we can recover validity of the replacement:

$$\eta_{\mu\rho}\eta_{\nu\sigma}M^{\mu\nu}M^{\rho\sigma*} = -\eta_{\nu\sigma}(M^{1\nu}M^{1\sigma*} + M^{2\nu}M^{2\sigma*}) \quad (\text{A.9})$$

$$= |M^{11}|^2 + |M^{12}|^2 + |M^{21}|^2 + |M^{22}|^2. \quad (\text{A.10})$$

Then what's happening? Why this replacement is not valid? Actually our new conditions (A.8) seem to guarantee that we are summing not only “physical” but also “unphysical” polarizations. Meanwhile if we use some physical condition such as $\epsilon \cdot k = 0$, (A.8) break down while Ward identities (A.7) are still valid.

Now let's check what is happening from another viewpoint. First we suppose M satisfies our new conditions (A.8), and define $\widetilde{M}^{\mu\nu}$ and \widetilde{M} as

$$\widetilde{M}^{\mu\nu} := M^{\mu\nu} + k^\mu p^\nu + p'^\mu k'^\nu, \quad (\text{A.11})$$

$$\widetilde{M} := \epsilon_\mu^*(k)\epsilon_\nu'^*(k')\widetilde{M}^{\mu\nu}. \quad (\text{A.12})$$

^{*3} This part is derived from 濱口幸一's notebook.

This alternative amplitude satisfies Ward identities (since photon is massless and $\epsilon \cdot k = 0$), and furthermore $\widetilde{M} = M$. Therefore \widetilde{M} is physically identical to M . However technically these are very different, just because we cannot perform our “naïve replacement” for this \widetilde{M} :

$$\eta_{\mu\rho}\eta_{\nu\sigma}\widetilde{M}^{\mu\nu}\widetilde{M}^{\rho\sigma*} = \eta_{\mu\rho}\eta_{\nu\sigma} (M^{\mu\nu} + k^\mu p^\nu + p'^\mu k'^\nu) (M^{\rho\sigma*} + k^\rho p^{\sigma*} + p'^{\rho*} k'^\sigma) \quad (\text{A.13})$$

$$= \sum_{\text{pol.}} |M|^2 + [(k \cdot p'^*)(k' \cdot p) + \text{H. c.}]. \quad (\text{A.14})$$

After all, we have obtained following expression:

$$\begin{aligned} \sum_{\text{pol.}} |\widetilde{M}|^2 &= \sum_{\text{pol.}} |M|^2 \quad (\text{Furthermore } \widetilde{M} = M) \\ &= \sum_{\text{pol.}} |\epsilon_\mu^*(k) \epsilon_\nu'^*(k') M^{\mu\nu}|^2 = \sum_{\text{pol.}} |\epsilon_\mu^*(k) \epsilon_\nu'^*(k') \widetilde{M}^{\mu\nu}|^2 \\ &= \eta_{\mu\rho}\eta_{\nu\sigma} M^{\mu\nu} M^{\rho\sigma*} \\ &\neq \eta_{\mu\rho}\eta_{\nu\sigma} \widetilde{M}^{\mu\nu} \widetilde{M}^{\rho\sigma*} = \sum_{\text{pol.}} |\widetilde{M}|^2 + [(k \cdot p'^*)(k' \cdot p) + \text{H. c.}]. \end{aligned} \quad (\text{A.15})$$