

### Question 1 Particle on a spring

A particle of mass  $m$  is constrained to move along a line in the  $x$  direction. It's attached to a spring with constant  $k$  that is connected to a point at a distance  $L$  perpendicular to the line. The natural length of the spring is  $L_0 = L/2$ .

(a)

Determine the system's potential energy  $U$  in terms of the particle's position  $x = x(t)$

The potential energy is

$$U = \frac{1}{2}k(\Delta L)^2 = \frac{1}{2}k\left(\sqrt{L^2 + x^2} - \frac{L}{2}\right)^2$$

(b)

Find the position  $x = x_{\min}$  where the potential energy has a minimum

From the intuition, the minimum potential energy happens when  $x = 0$  (the spring is least stretched). The calculation also agrees on that:

$$U' = \frac{1}{2}k \cdot 2 \left( \sqrt{L^2 + x^2} - \frac{L}{2} \right) \cdot \frac{1}{2}(L^2 + x^2)^{-\frac{1}{2}} \cdot 2x = kx \cdot \left( 1 - \frac{L}{2\sqrt{L^2 + x^2}} \right)$$

is equal to zero only when  $x = 0$  and that

$$U'' = k\left(1 - \frac{L}{2\sqrt{L^2 + x^2}}\right) + kx\left(-\frac{L}{2} \cdot -\frac{1}{2}(L^2 + x^2)^{-3/2}\right)$$

with  $U''(0) = \frac{1}{2}k > 0$  show that the point is indeed the minimum.

(c)

Determine the effective spring constant  $k_{\text{eff}} = U''(x = x_{\min})$

From the previous calculation, we know that  $k_{\text{eff}} = U''(0) = \frac{1}{2}k$

(d)

Determine the system's oscillation frequency

The oscillation frequency  $\omega = \sqrt{\frac{k_{\text{eff}}}{m}} = \sqrt{\frac{k}{2m}}$

### Question 2 Simple Pendulum

(A pendulum, with a massless rod of length  $L$  and a point bob of mass  $m$ , swings in a gravitational field  $\vec{g}$  as shown. Assume the pendulum to be near the Earth's surface such that  $g = \text{const}$ . The pendulum has energy  $E$ , that is not necessarily small. We choose the reference frame such that the potential energy vanishes at  $\theta = 0$ , i.e.,  $U(\theta = 0) = 0$

(a)

Determine the pendulum maximum angular speed  $\dot{\theta}$ , in terms of  $E$ ,  $L$ ,  $m$  and  $g$

The total energy  $E$  is equal to

$$E = U + T = mgL(1 - \cos \theta) + \frac{1}{2}mL^2\dot{\theta}^2$$

Thus

$$\dot{\theta} = \sqrt{\frac{2E}{mL^2} - \frac{2g}{L}(1 - \cos \theta)}$$

and maximum is achieved when setting  $\cos \theta = 1$ .

$$\boxed{\dot{\theta} = \sqrt{\frac{2E}{mL^2}}} \quad (1)$$

(b)

Determine the maximum amplitude,  $\theta_{\max}$ , of the oscillation in terms of  $E$ ,  $L$ ,  $m$ ,  $g$

From the question (a), we could rearrange the equation to

$$\theta = \arccos\left(1 - \frac{E}{mgL} - \frac{L}{2g}\dot{\theta}^2\right)$$

and the maximum is achieved when setting  $\dot{\theta} = 0$ .

$$\boxed{\theta = \arccos\left(1 - \frac{E}{mgL}\right)} \quad (2)$$

(c)

How large must the energy  $E$  be such that the pendulum whirls around rather than just oscillates?

For pendulum to whirls around, the maximum  $\theta$  must be at least  $\pi$  (see the graph). Thus

$$\begin{aligned} \pi &= \arccos\left(1 - \frac{E}{mgL}\right) \\ -1 &= 1 - \frac{E}{mgL} \\ \boxed{E} &= 2mgL \end{aligned} \quad (3)$$

### Question 3 Compound pendulum

A compound pendulum is formed of a rigid rod of length  $2a$ , with one mass  $m_1 = m$  at its tip and another mass  $m_2 = 2m$  at its midpoint as shown.

Determine the pendulum's period for small oscillations, i.e.,  $\theta(t)$  small.

Hint: one approach is start by constructing the formulas for its kinetic energy as a function of  $\dot{\theta}$  and for its potential energy as a function of  $\theta$

The potential energy is

$$\begin{aligned}
U &= U_{\text{small}} + U_{\text{big}} \\
&= m_1 g \cdot 2a(1 - \cos \theta) + m_2 g \cdot a(1 - \cos \theta) \\
&= 4mga(1 - \cos \theta)
\end{aligned}$$

and the kinetic energy is

$$\begin{aligned}
T &= T_{\text{small}} + T_{\text{big}} \\
&= \frac{1}{2} m_1 (2a\dot{\theta})^2 + \frac{1}{2} m_2 (a\dot{\theta})^2 \\
&= \frac{1}{2} m \cdot 4a^2 \dot{\theta}^2 + \frac{1}{2} \cdot 2m \cdot a^2 \dot{\theta}^2 \\
&= 3ma^2 \dot{\theta}^2
\end{aligned}$$

so, the total energy is

$$E = U + T = 4mga(1 - \cos \theta) + 3ma^2 \dot{\theta}^2$$

and since the total energy is conserved

$$\frac{dE}{dt} = 4mga \sin \theta \cdot \dot{\theta} + 3ma^2 \cdot 2\dot{\theta} \cdot \ddot{\theta} = 0$$

and thus

$$4mga \sin \theta + 6ma^2 \ddot{\theta} = 0$$

since  $\theta$  is small, we use the approximation that  $\sin \theta \approx \theta$

$$4mga\theta + 6ma^2 \ddot{\theta} = 0$$

which we recognize that

$$\begin{aligned}
M_{\text{eff}} &= 6ma^2 \\
K_{\text{eff}} &= 4mga
\end{aligned}$$

and thus the

$$\omega = \sqrt{\frac{K_{\text{eff}}}{M_{\text{eff}}}} = \sqrt{\frac{4mga}{6ma^2}} = \sqrt{\frac{2g}{3a}}$$

and thus the period is

$$\boxed{P = \frac{2\pi}{\omega} = 2\pi \sqrt{\frac{3a}{2g}}} \quad (4)$$

#### Question 4

A pendulum, with a massless rod of length  $L$  and a point bob of mass  $m$  at its tip, oscillates in an oil bath. The drag force acting on the bob is  $\vec{F} = -c\vec{v}$  and is directed opposite to the bob's velocity vector  $\vec{v}$ . You may ignore buoyancy and assume motion near the Earth's surface such that  $g = \text{const}$

(a)

Derive the differential equation describing the bob's motion  $\theta(t)$  for an arbitrary amplitude  $\theta_{\max}$  (but such that the bob remains submerged in the oil).

Doing some forces analysis:

$$I\ddot{\theta} = -mgL \sin \theta - cL(L\dot{\theta})$$

that is

$$mL^2\ddot{\theta} + cL^2\dot{\theta} + mgL \sin \theta = 0$$

(b)

Linearize the differential equation that you obtained in (a) and check it for obvious errors: is the damping positive? Is the effective stiffness (or "effective spring constant") positive? What is the damped frequency of the vibration?

Assume that the system is underdamped.

Assume that  $\theta$  is small, so  $\sin \theta \approx \theta$ . So, the equation is

$$mL^2\ddot{\theta} + cL^2\dot{\theta} + mgL\theta = 0$$

which is equivalent to solving

$$\ddot{\theta} + \frac{c}{m}\dot{\theta} + \frac{g}{L}\theta = 0 \quad (5)$$

The  $M_{\text{eff}} = 1$ ,  $C_{\text{eff}} = \frac{c}{m}$ ,  $K_{\text{eff}} = \frac{g}{L}$ . and thus

$$\omega_n = \sqrt{\frac{K_{\text{eff}}}{M_{\text{eff}}}} = \sqrt{\frac{g}{L}}$$

and thus the damping zeta  $\zeta$  is

$$\zeta = \frac{C_{\text{eff}}}{2M_{\text{eff}}\omega_n} = \frac{c}{2m} \sqrt{\frac{L}{g}}$$

by assumption,  $c > 0$ . and obviously  $L > 0$  and  $m > 0$ . Thus, the  $\zeta > 0$ .

The effective spring constant is just  $K_{\text{eff}} = g/L > 0$ , (or  $mgL$  for original equation) is obviously greater than zero.

By assumption of that system is underdamped, the damped frequency  $\omega_d$  is

$$\omega_d = \omega_n \sqrt{1 - \zeta^2} = \sqrt{\frac{g}{L} \left(1 - \frac{c^2}{4m^2} \cdot \frac{L}{g}\right)} = \sqrt{\frac{g}{L} - \frac{c^2}{4m^2}}$$

(c)

Setting the initial conditions  $\theta(t=0) = \theta_0$  and  $\dot{\theta}(t=0) = \Omega$ , find the trajectory  $\theta(t)$

By the Lecture note 13, Equation (43), the solution has the form of

$$\theta(t) = \exp(-\zeta\omega_n t)(A \cos(\omega_d t) + B \sin(\omega_d t))$$

It's derivative is

$$\dot{\theta}(t) = -\zeta\omega_n \exp(-\zeta\omega_n t)(A \cos(\omega_d t) + B \sin(\omega_d t)) + \exp(-\zeta\omega_n t)(-A\omega_d \sin(\omega_d t) + B\omega_d \cos(\omega_d t))$$

By the initial conditions

$$\theta(0) = 1(A \cdot 1 + B \cdot 0) = \theta_0$$

$$\dot{\theta}(0) = -A\zeta\omega_n + B\omega_d = \Omega$$

and thus

$$A = \theta_0 \quad B = \frac{\Omega + \theta_0 \zeta \omega_n}{\omega_d}$$

and the solution is thus

$$\theta(t) = \exp(-\zeta\omega_n t) \left( \theta_0 \cos(\omega_d t) - \frac{\Omega}{\zeta\omega_n \omega_d} \sin(\omega_d t) \right) \quad (6)$$

where the  $\omega_n$ ,  $\omega_d$ , and  $\zeta$  is calculated in (b)

### Question 5

Computer Algebra programs like mathematica are useful tools to do lengthy computations. Here you will practice working with mathematica using the problem in Question 4, In particular,

(a)

Verify that your solution in (4c) does indeed solve the differential equation of (4b) by inserting your solution

We define some constants first

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In[1]:= Meff = 1; Ceff = c / m; Keff = g / L
Out[1]=  $\frac{g}{L}$ 

In[2]:= Wn = Sqrt[Keff / Meff]
Out[2]=  $\sqrt{\frac{g}{L}}$ 

In[3]:= Ze = Ceff / (2 * Meff * Wn)
Out[3]=  $\frac{c}{2 \sqrt{\frac{g}{L}} m}$ 

In[4]:= Wd = Wn * Sqrt[1 - Ze ^ 2]
Out[4]=  $\sqrt{\frac{g}{L}} \sqrt{1 - \frac{c^2 L}{4 g m^2}}$ 

In[5]:= A = ThetaZero
Out[5]= ThetaZero

In[6]:= B = (Omega + ThetaZero * Ze * Wn) / Wd
Out[6]=  $\frac{\Omega + \frac{c \text{ThetaZero}}{2 m}}{\sqrt{\frac{g}{L}} \sqrt{1 - \frac{c^2 L}{4 g m^2}}}$ 

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and write the  $\theta(t)$  function, put it back into (5)

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In[7]:= Theta[t_] = Exp[-Ze*Wn*t] * (A * Cos[Wd*t] + B * Sin[Wd*t])

Out[7]= e^{\frac{-c t}{2 m}} \left( \text{ThetaZero} \cos \left[ \sqrt{\frac{g}{L}} \sqrt{1 - \frac{c^2 L}{4 g m^2}} t \right] + \frac{\left( \text{Omega} + \frac{c \text{ThetaZero}}{2 m} \right) \sin \left[ \sqrt{\frac{g}{L}} \sqrt{1 - \frac{c^2 L}{4 g m^2}} t \right]}{\sqrt{\frac{g}{L}} \sqrt{1 - \frac{c^2 L}{4 g m^2}}} \right)

In[8]:= Simplify[Theta''[t] + c/m*Theta'[t] + g/L*Theta[t]]

Out[8]= 0

```

we get 0 as expected, which shows that our solution is correct.

(b)

Plot the motion  $\theta(t)$  in the interval  $t = 0\text{s}, \dots, 10\text{s}$ . Set the parameters  $\theta_0 = 0.1$ ,  $L = 1\text{m}$ ,  $m = 0.5\text{kg}$ ,  $c = 1\text{kg/s}$ ,  $g = 9.81\text{m/s}^2$ ,  $\Omega = 20\text{Hz}$

Hint: it is useful to introduce the dimensionless time parameter  $\bar{t} = t/s$ , and plot  $\theta(\bar{t})$

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In[16]:= ThetaZero = 0.1; L = 1; m = 0.5; c = 1; g = 9.81; Omega = 20

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Out[16]= 20

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In[17]:= Theta[t]

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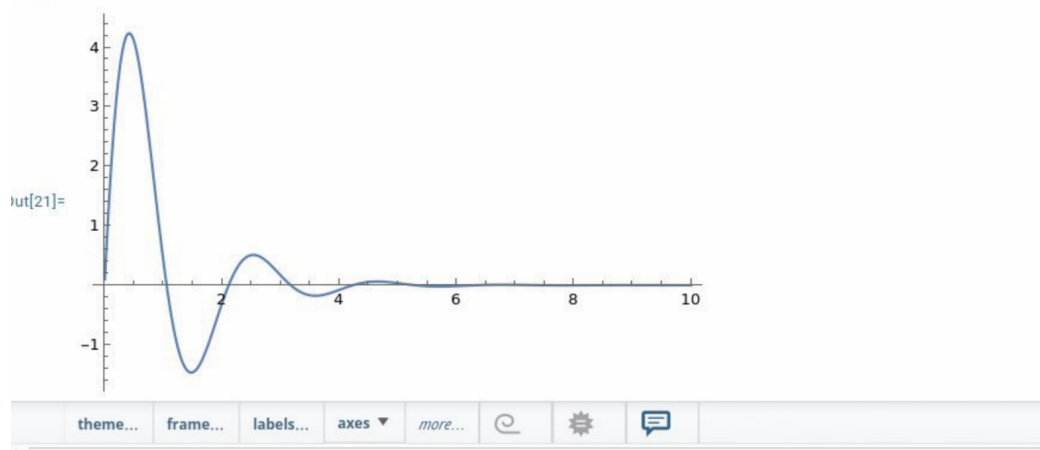
Out[17]= e^{-1. t} (0.1 Cos[2.96816 t] + 6.77186 Sin[2.96816 t])

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In[21]:= Plot[Theta[t], {t, 0, 10}, PlotRange -> All]

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Giving value to the constants, and plot the  $\theta(\bar{t})$ . As we could see, this is indeed a damped oscillation.