



# Crusher

 Date	@March 7, 2025
 tags	1D

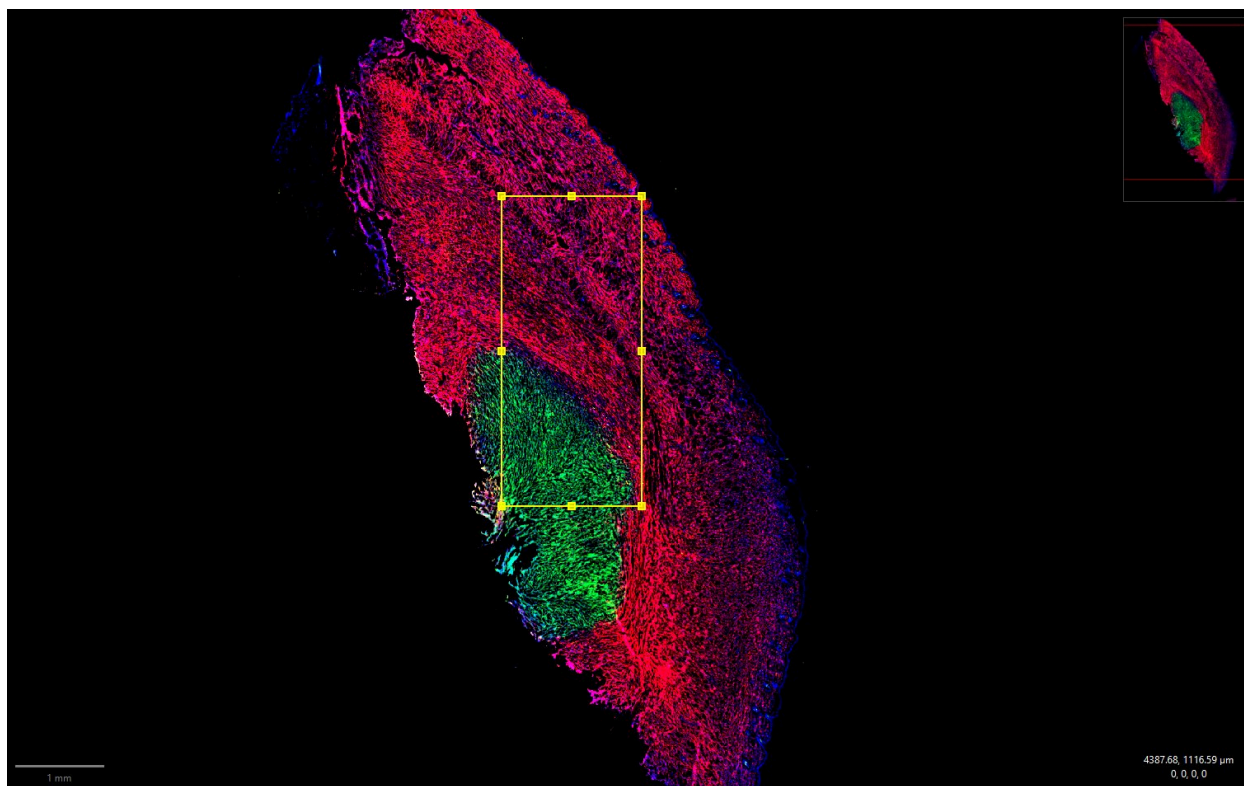
- Decided to go back to 1D for now.
- Motivation was because of what appears to be a clear separation of the T-cells in the 10mm tumor.
- Added the 2D data to the 1D simulation repository.
- Created a crusher file which project all points onto some line I defined.
- Plots all the points as well as the line to get an idea of what the projection will look like.
- Also plots a pdf of all the projected points.



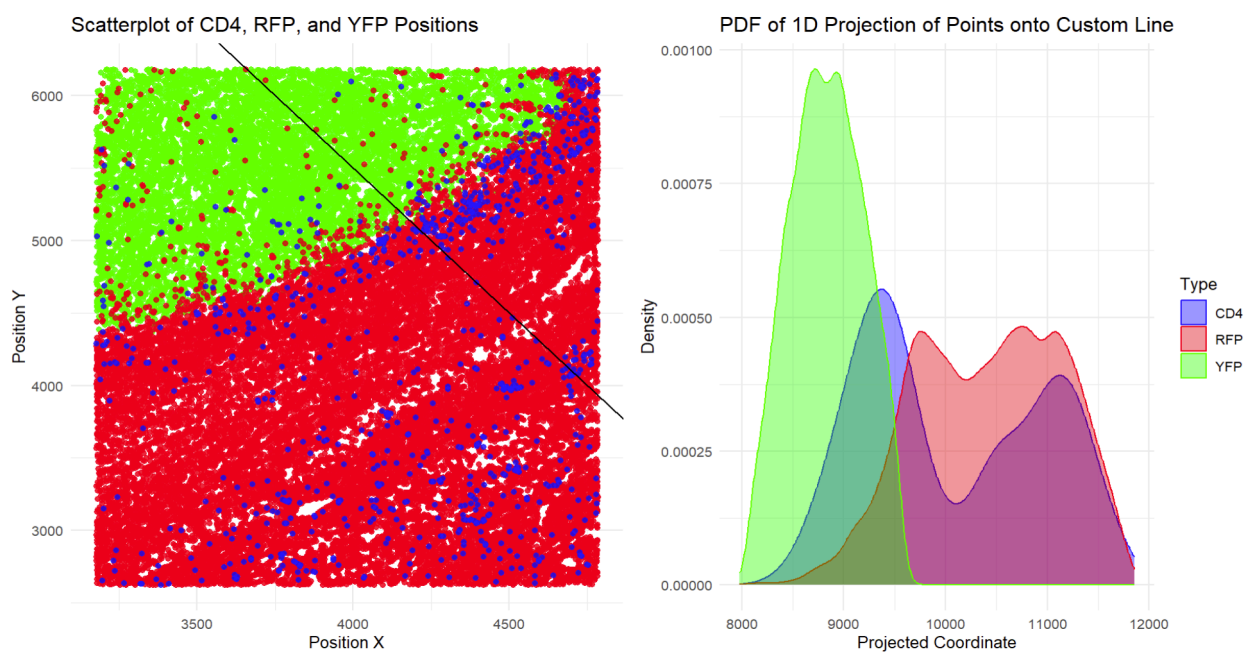
## Info

- There's a weird warning, saying that a reached elapsed time limit.
- Takes a weird amount of time to load.
- Jiggling the plot window seems to make it load better, idk why.

## 10mm\_tumor:






- Chose  $y = -2m + 13500$



- Seems pretty legit for now.

# 1D simulation pt2

 Date	@March 7, 2025
 Files & media	<a href="#"><u>chemoking_gradient_calc.pdf</u></a>
 tags	1D

- Going back because modeling the T-cell movement in 2D seems to be a little weird right now, hopefully will go back eventually.
- Will be making a change where T-cells will have a velocity that will be influenced by the gradient.

## The Setup:

- $b_{min}, b_{max}$  : smallest and largest positions observed across the entire dataset.
- Let  $X_i(0)$  denote the position of t-cell  $i$  at time step 0:

$$X_i(0) \sim Uniform(b_{min}, b_{max})$$

- Let  $V_i(0)$  denote the velocity of t-cell  $i$  at time step 0:

$$V_i(0) \sim Uniform(-1, 1)$$

- Let  $X_i(t), V_i(t)$  denote the position and velocity of T-cell  $i$  at time step  $t$

$$X_i(t) = X_i(t-1) + V_i(t-1)$$

## The Behaviors:



#### Note

- These are subject to change

## Chemokine Gradient

- Let  $c(x)$  denote the concentration of chemokine at position  $x$ :
  - For  $n$  cancer cells,  $x_i$  denotes the position and  $M_i$  denotes the heat of cancer cell  $i$
  - $D$  denotes the diffusion constant of the chemokines.
  - $\delta$  denotes the decay constant of the chemokines

$$c(x) = \sum_{i=1}^n \frac{M_i}{2\sqrt{D\delta}} e^{-\sqrt{\delta/D}(|x-x_i|)}$$

- Can see the derivation for it in the pdf ``chemokine_gradient_calc.pdf``.
- Can pull out some constants for faster computation

$$c(x) = \frac{1}{2\sqrt{D\delta}} \sum_{i=1}^n M_i e^{-\sqrt{\delta/D}(|x-x_i|)}$$

- Let  $\nabla c(x)$  denote the concentration gradient at position  $x$ 
  - Will point in the direction of greater concentration.

$$\nabla c(x) = -\frac{1}{2D} \sum_{i=1}^n M_i e^{-\sqrt{\delta/D}(|x-x_i|)} \cdot \text{sgn}(x - x_i)$$

- Additionally, let  $|\nabla c(x)|$  denote the magnitude of said gradient.
- Not sure what to call this, will refer to it as the **logistic function**
- Will be used to determine how the gradient will influence the velocity.

$$\sigma_{\alpha}(y) = \frac{e^{\alpha y}}{1 + e^{\alpha y}}$$

- The larger that  $\alpha$  is, the stronger that the gradient will influence the velocity.
- Will also have a speed term  $v$

#### Chemokine Visualization 10mm Tumor

### **Behavior 0: Logistic Always Update**

- The Cancer Cells will always listen to the gradient.
- So
- If  $\nabla c(X_i(t)) > 0$

$$V_i(t) = \sigma_{\alpha}(|\nabla c(X_i(t))|) \cdot v$$

- If  $\nabla c(X_i(t)) < 0$

$$V_i(t) = -\sigma_{\alpha}(|\nabla c(X_i(t))|) \cdot v$$

### **Behavior 1: Probability becomes the velocity**

- The logistic function will be used as a probability that the T-cells respond to the gradient.
- let  $p \sim \text{uniform}(0, 1)$
- The gradient will affect the velocity if

$$p < \sigma_\alpha(|\nabla c(X_i(t))|)$$

- The velocity update:

- If  $\nabla c(X_i(t)) > 0$

$$V_i(t) = \sigma_\alpha(|\nabla c(X_i(t))|) \cdot v$$

- If  $\nabla c(X_i(t)) < 0$

$$V_i(t) = -\sigma_\alpha(|\nabla c(X_i(t))|) \cdot v$$

- When the gradient isn't affected:

$$V_i(t) = V_i(t - 1)$$

## results

# Chemokine Visualization 10mm Tumor

- Green Represents YFP Cancer Cells Distribution
- Red Represents RFP Cancer Cells Distribution
- Black Represents the Concentration/Gradient scale is on the right



## Takeaways

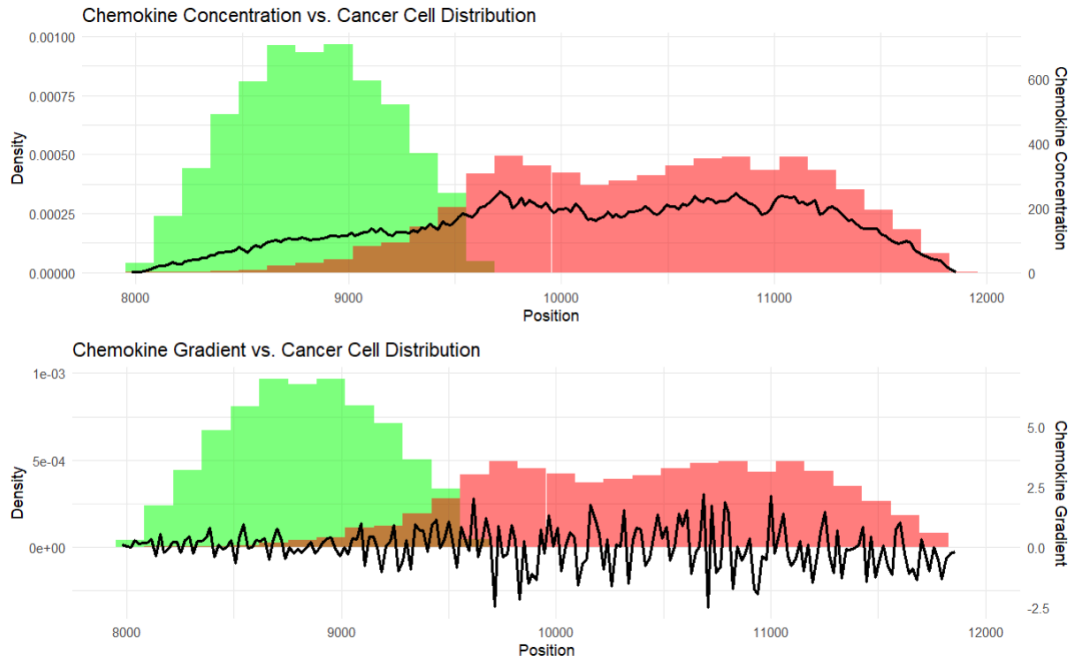
- For the RFP, YFP heat ratio, the smaller the ratio the longer it will take for a T-cell to exit it's region.
  - I think anything above a 1-2 ratio and below like 10 is reasonable.
- For Diffusion/Decay:
  - Higher diffusion/Smaller Decay should be used if we want the T-cells to reach the two noticed clusters in Crusher for the 10mm Tumor
  - Smaller Diffusion/Greater Decay should be used to have the T-cells be more spread out.
  - I think Diffusion 10/ Decay 0.1, is good if we want T-cells to be more spread out
  - I think Diffusion 10/ Decay 0.01, is good if we want T-cells to focus more at those peaks.

## Modifying YFP RFP Heat Ratios

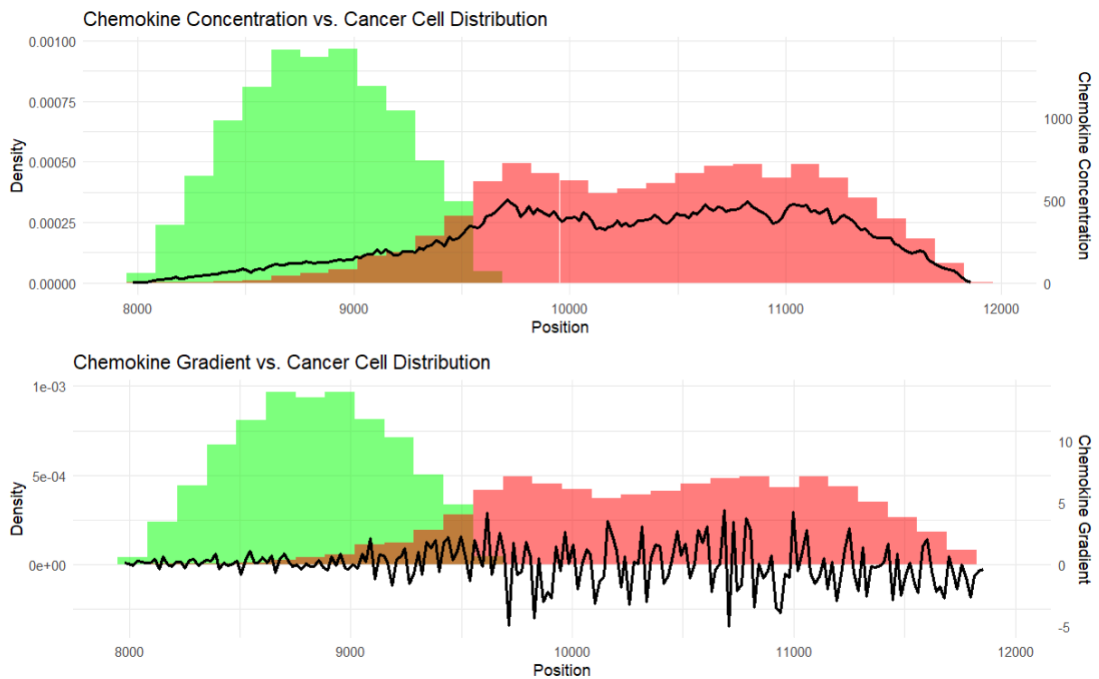
- Diffusion = 10
- Delta = 0.1

# Modifying YFP RFP Heat Ratios

- Diffusion = 10
- Delta = 0.1
- YFP Heat = 1
- RFP Heat = 2

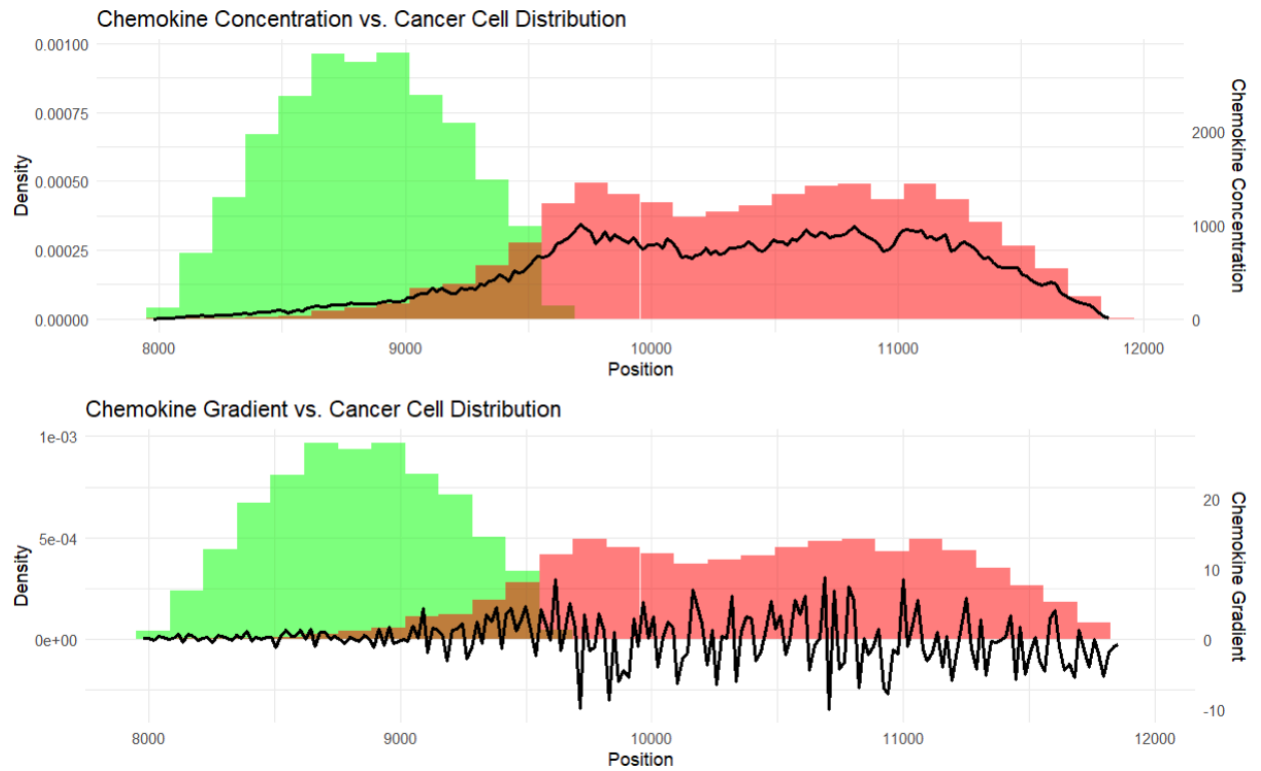


- Diffusion = 10
- Delta = 0.1
- YFP Heat = 1
- RFP Heat = 4



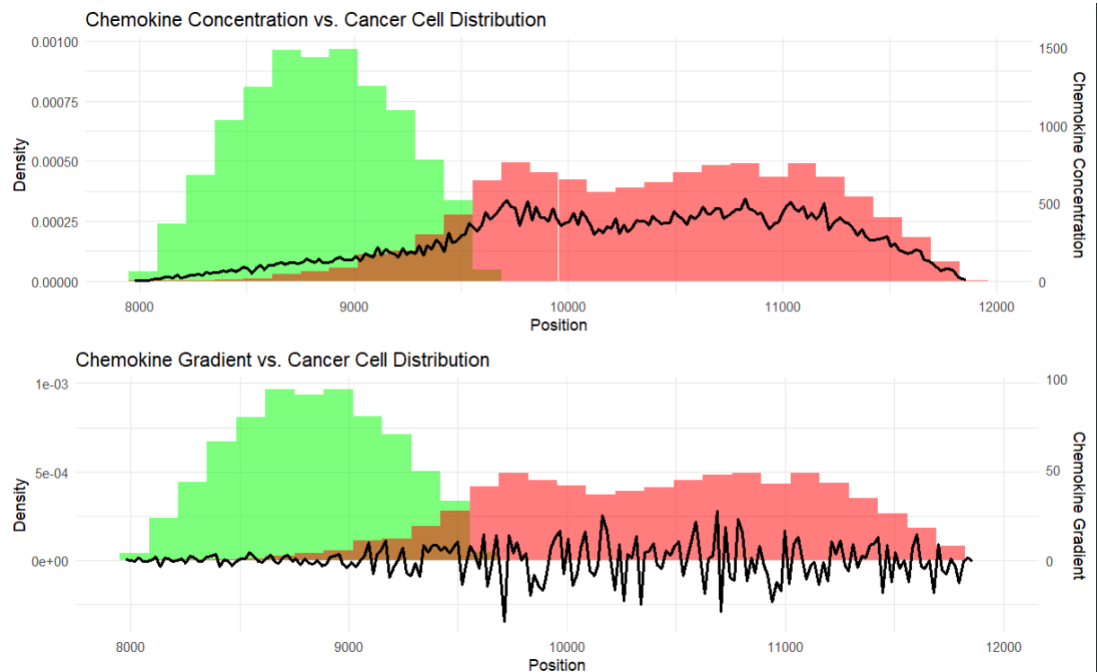


- Diffusion = 10
- Delta = 0.1
- YFP Heat = 1
- RFP Heat = 8

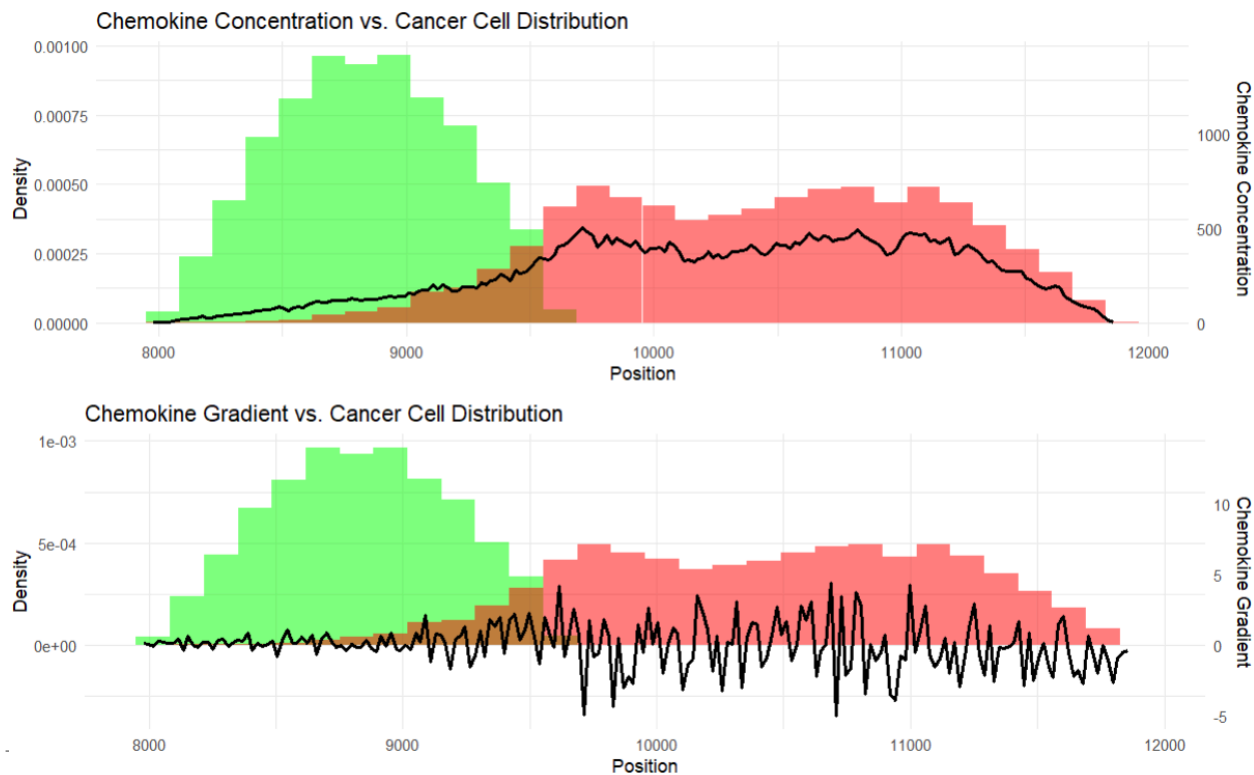


# Modifying Diffusion

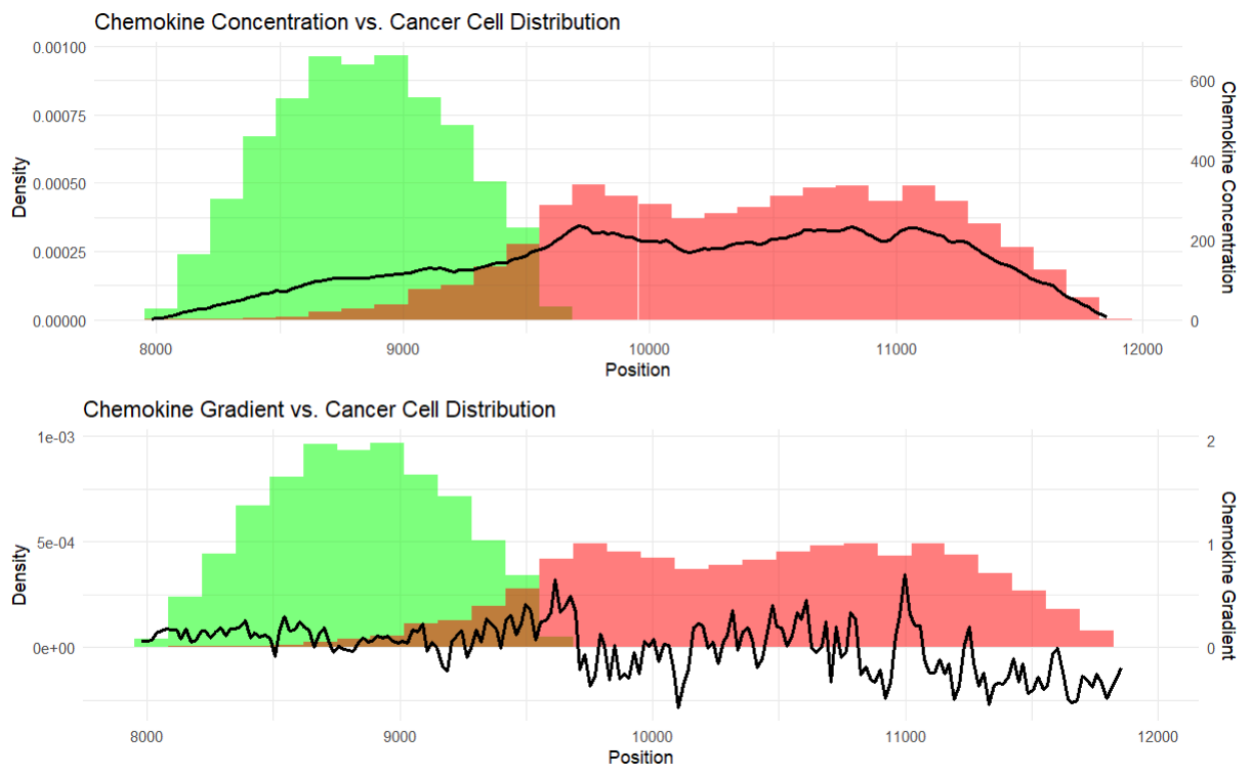
- Diffusion = 1
- Delta = 0.1
- YFP Heat = 1
- RFP Heat = 4



- Diffusion = 10
- Delta = 0.1
- YFP Heat = 1
- RFP Heat = 4

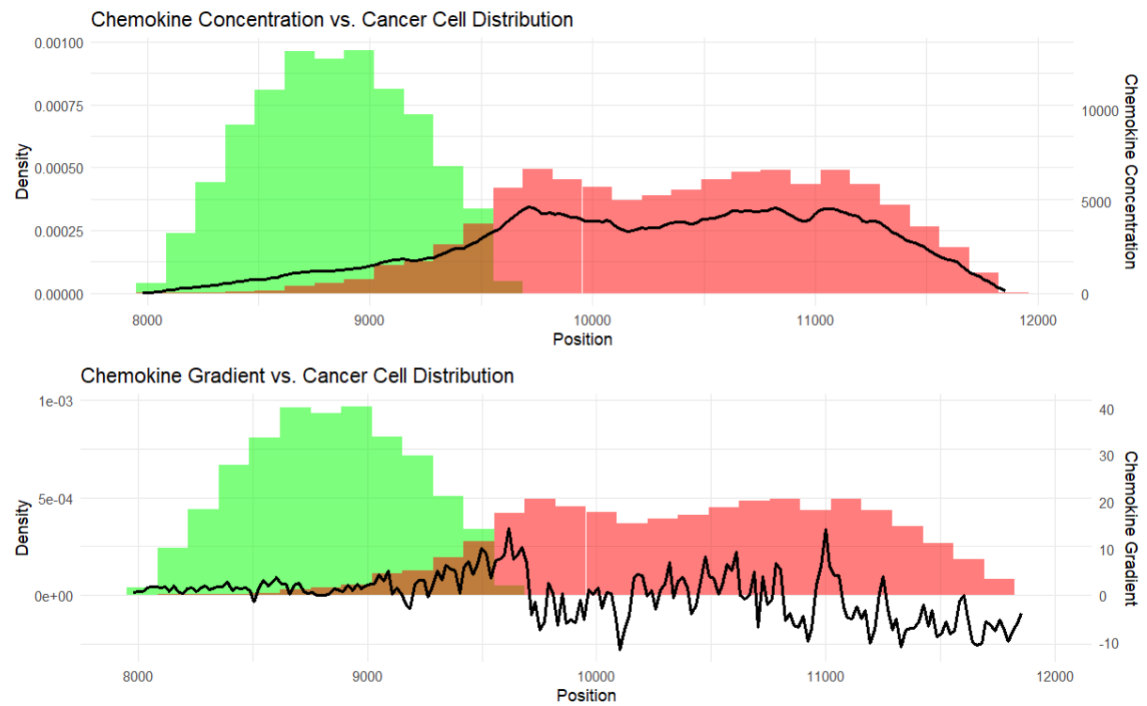


- Diffusion = 100
- Delta = 0.1
- YFP Heat = 1
- RFP Heat = 4

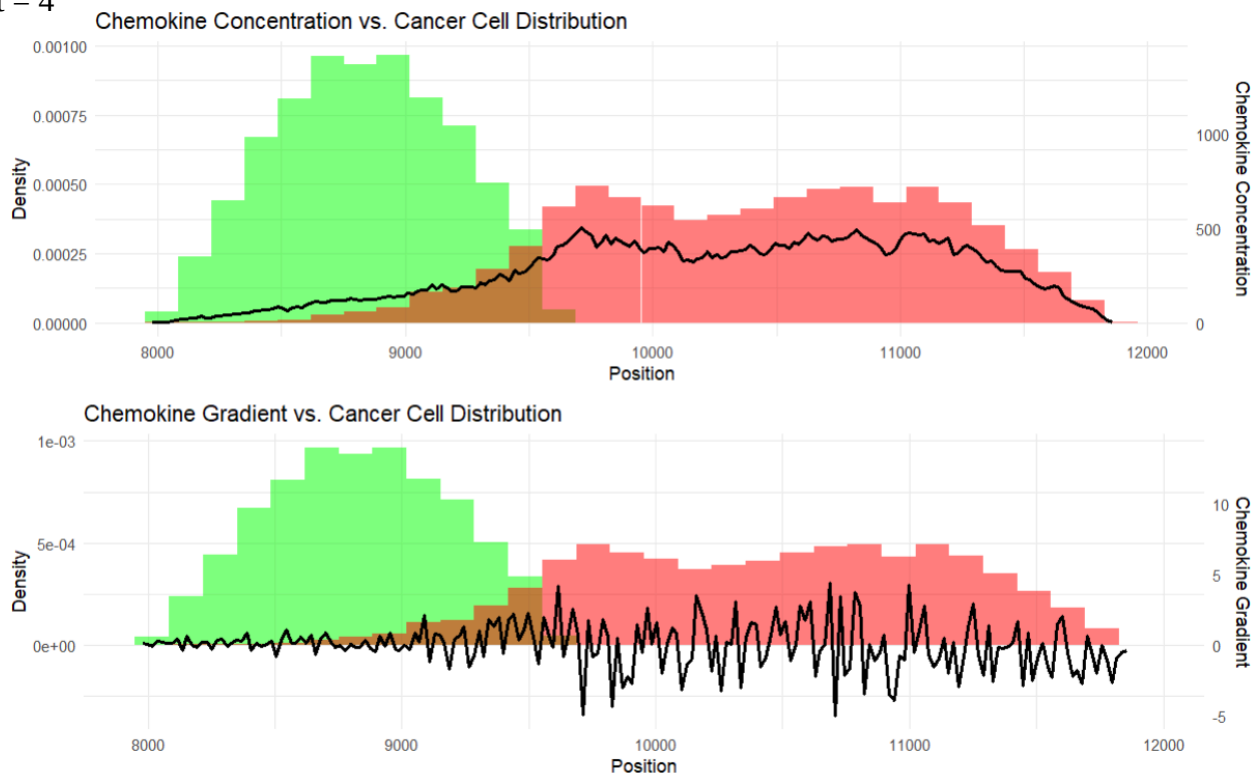


# Modifying Decay

- Diffusion = 10
- Delta = 0.01
- YFP Heat = 1
- RFP Heat = 4

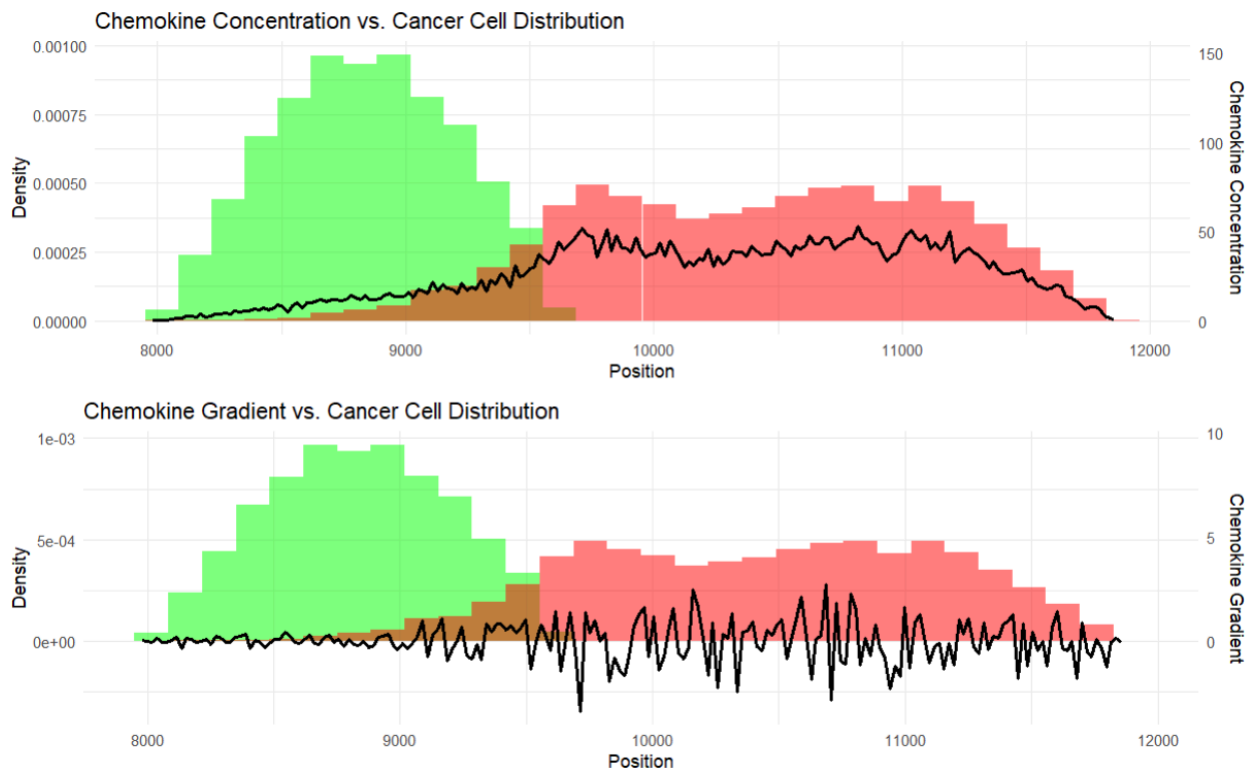


- Diffusion = 10
- Diffusion = 10
  - Delta = 1
- Delta = 0.1
  - YFP Heat = 1
- YFP Heat = 1
  - RFP Heat = 4
- RFP Heat = 4





- Diffusion = 10=
- Delta = 1=
- YFP Heat = 1=
- RFP Heat = 4

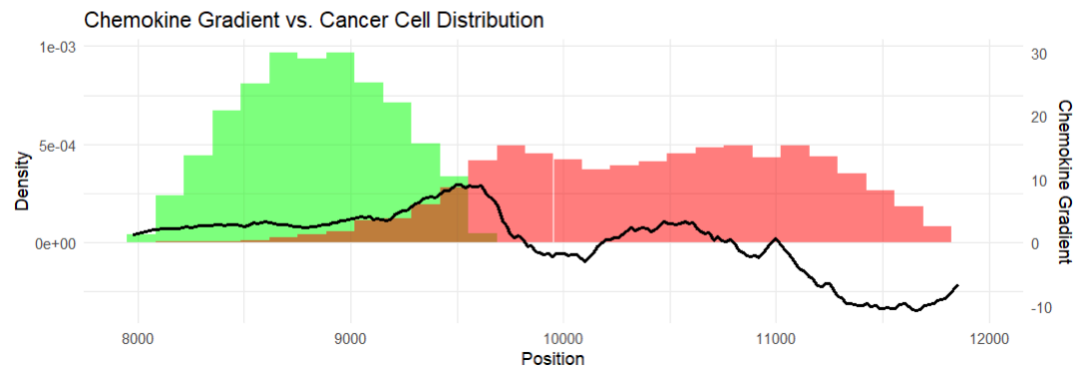
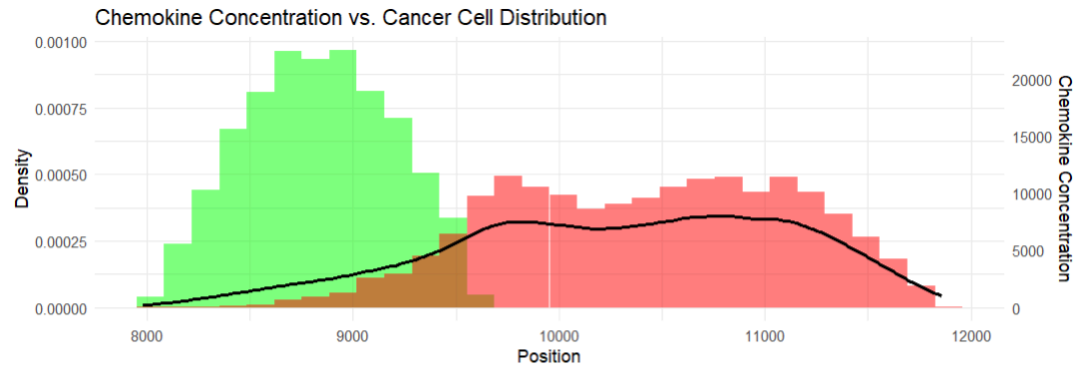


# Paper Constants

## Simulation Constant Searching

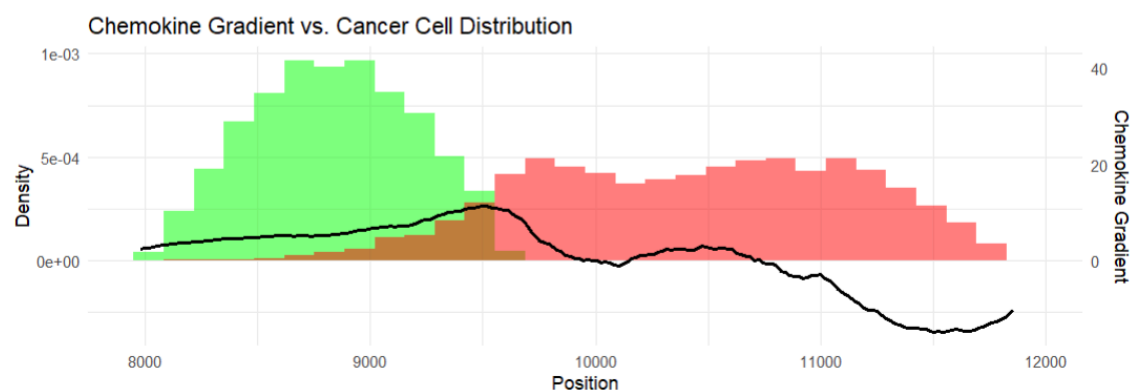
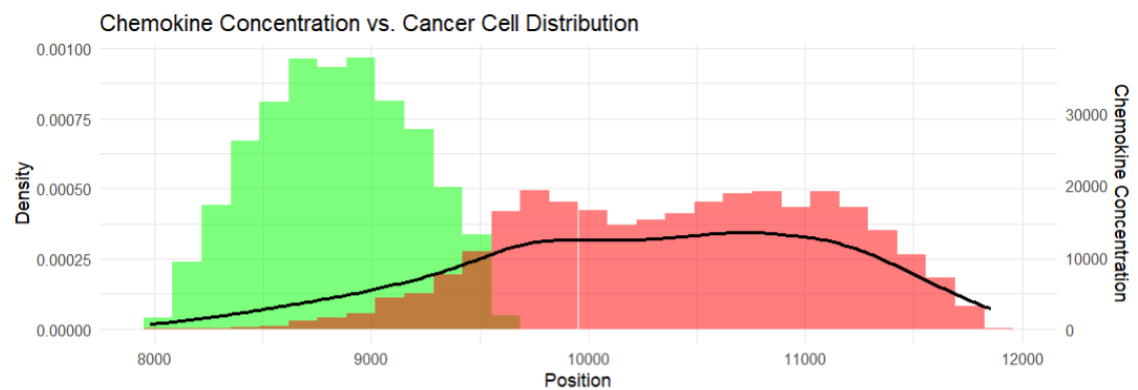
### CCL21

- $d_{1/2} = 128.0$
- $K = 0.00542$
- YFP Heat = 1
- RFP Heat = 4



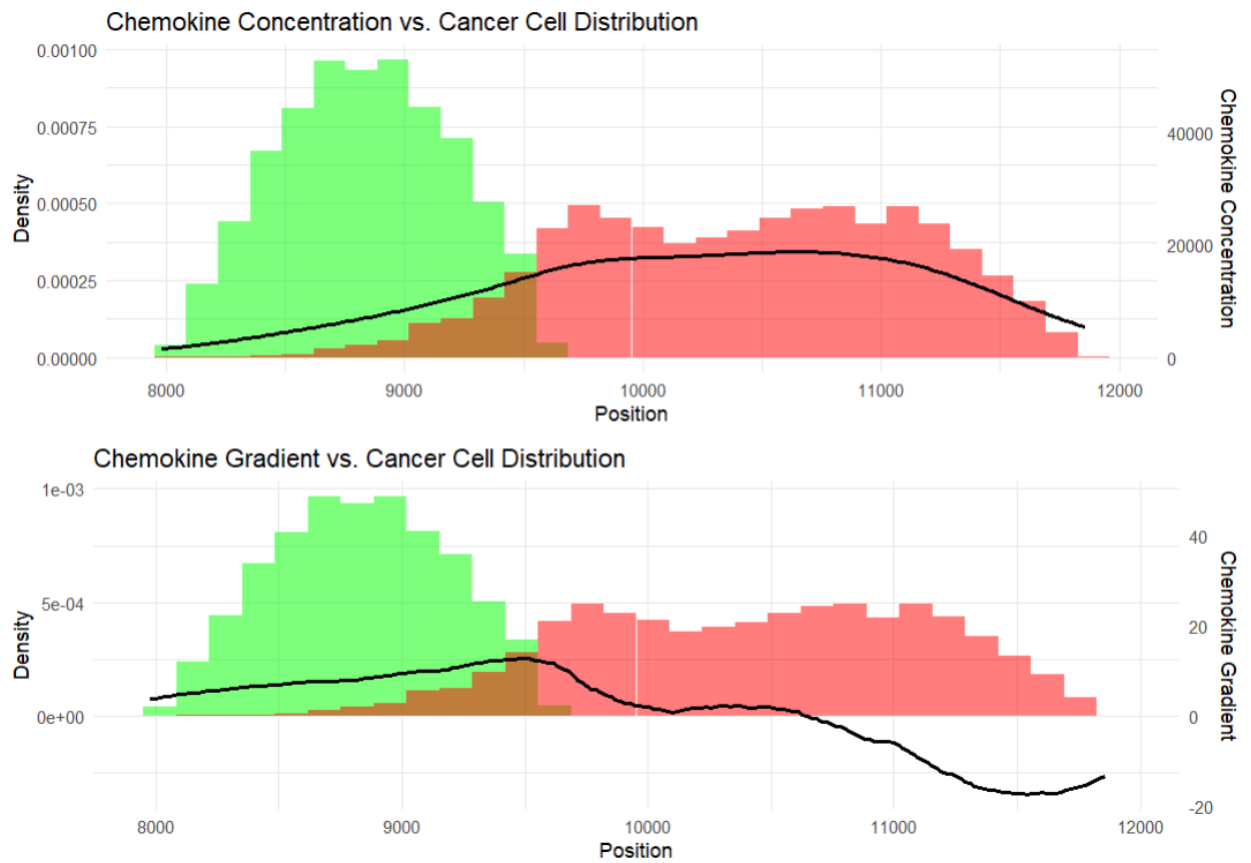
### CCL19

- $d_{1/2} = 226.4$
- $K = 0.00306$
- YFP Heat = 1
- RFP Heat = 4



## CCL21trunc

- $d_{1/2} = 338.6$
- $K = 0.00205$
- YFP Heat = 1
- RFP Heat = 4





# results

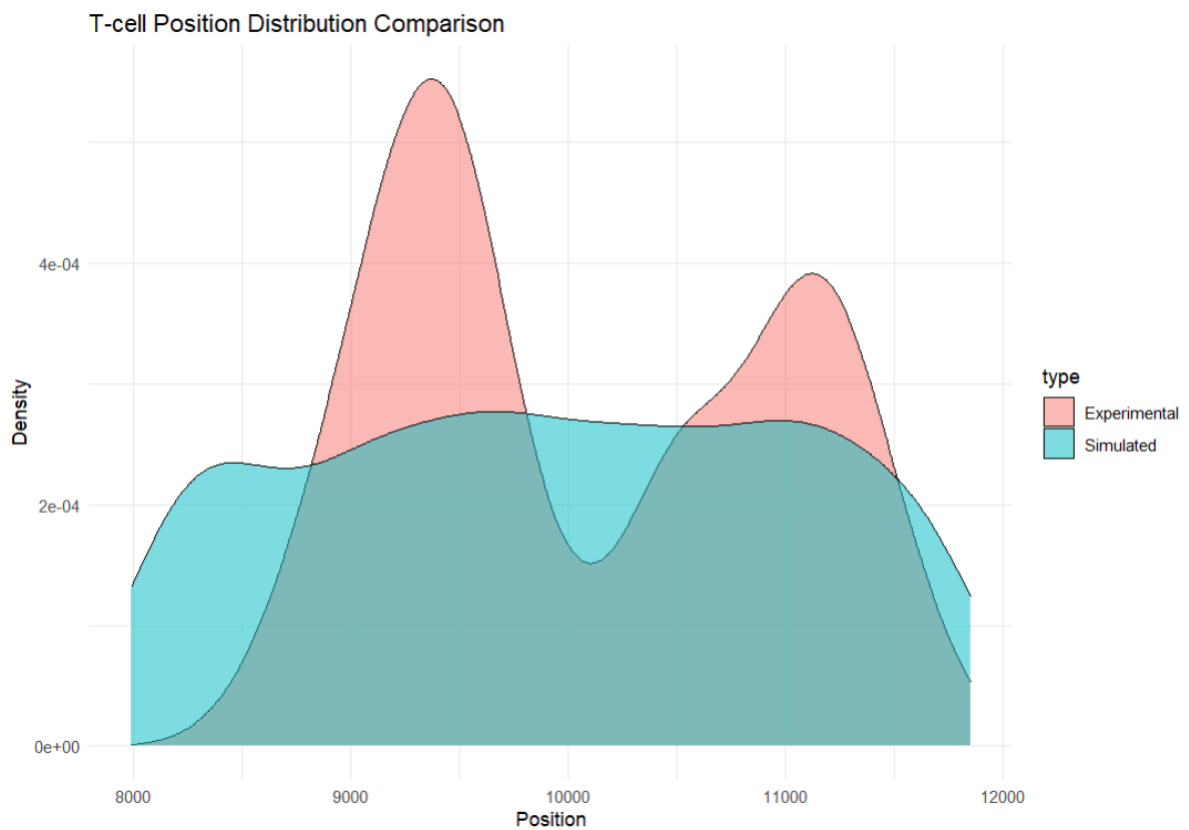
## Info

- Anything greater than  $\alpha = 0$  does not appear to have change.

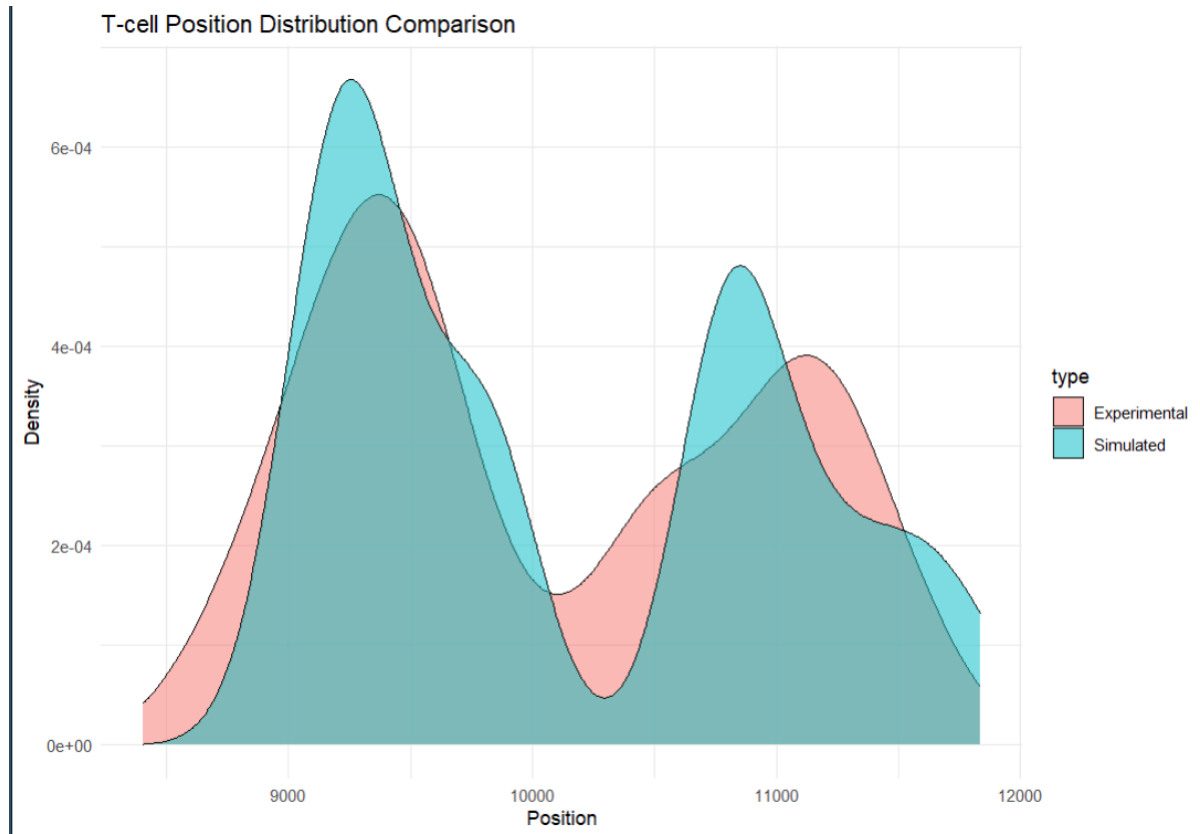
## CCL21

### Behavior 0

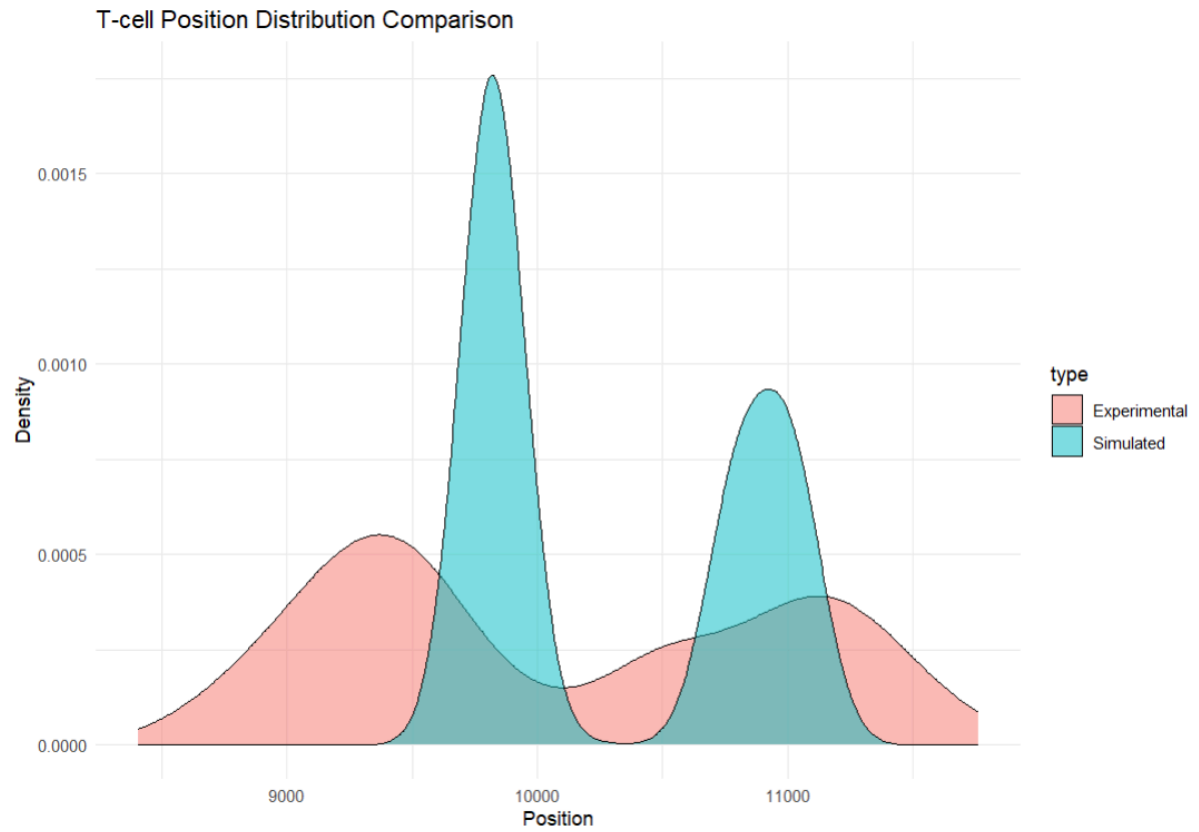
- $\alpha = -10$
- speed = 2
- num-steps = 10000



- $\alpha = -1$
- speed = 2
- num-steps = 10000

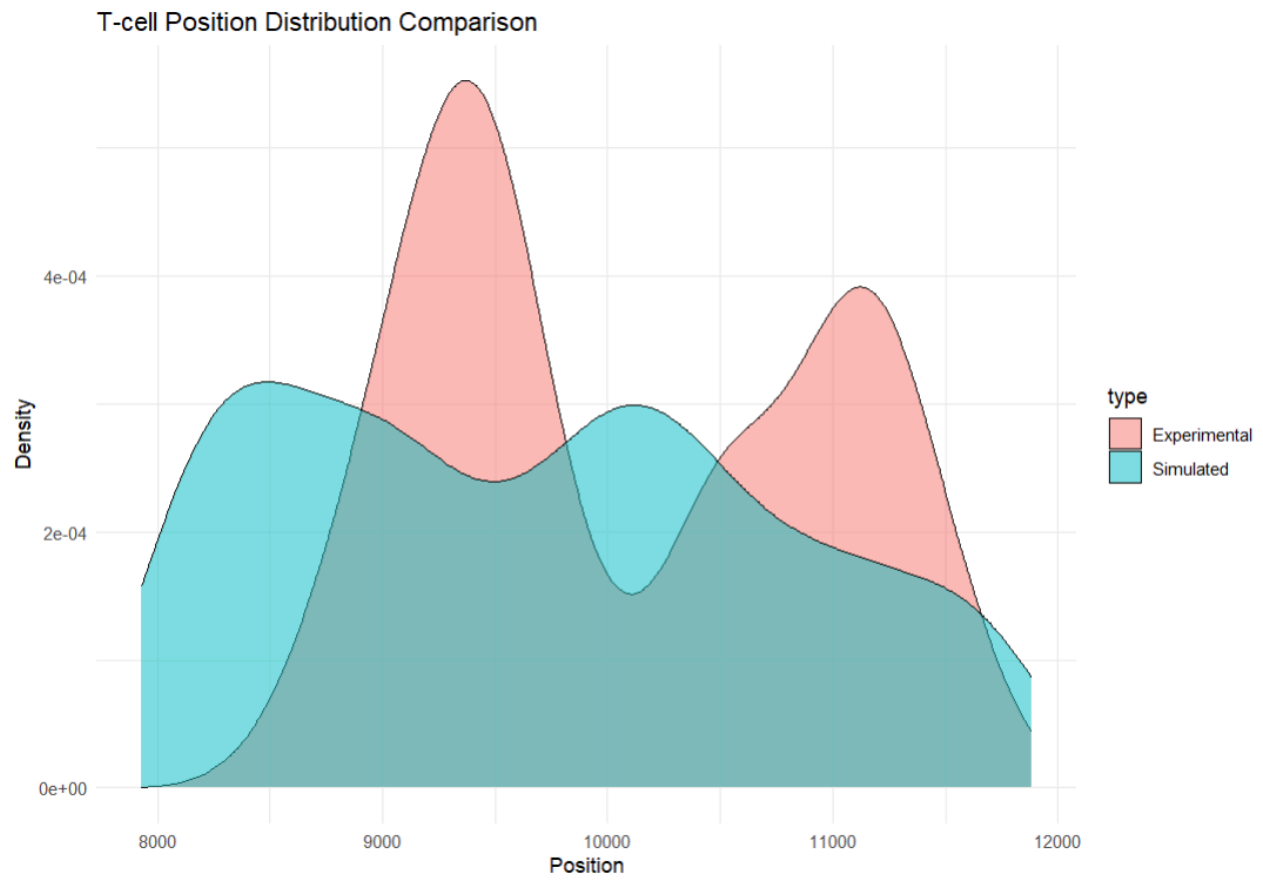


- $\alpha = 1e-6$
- speed = 2
- num-steps = 10000

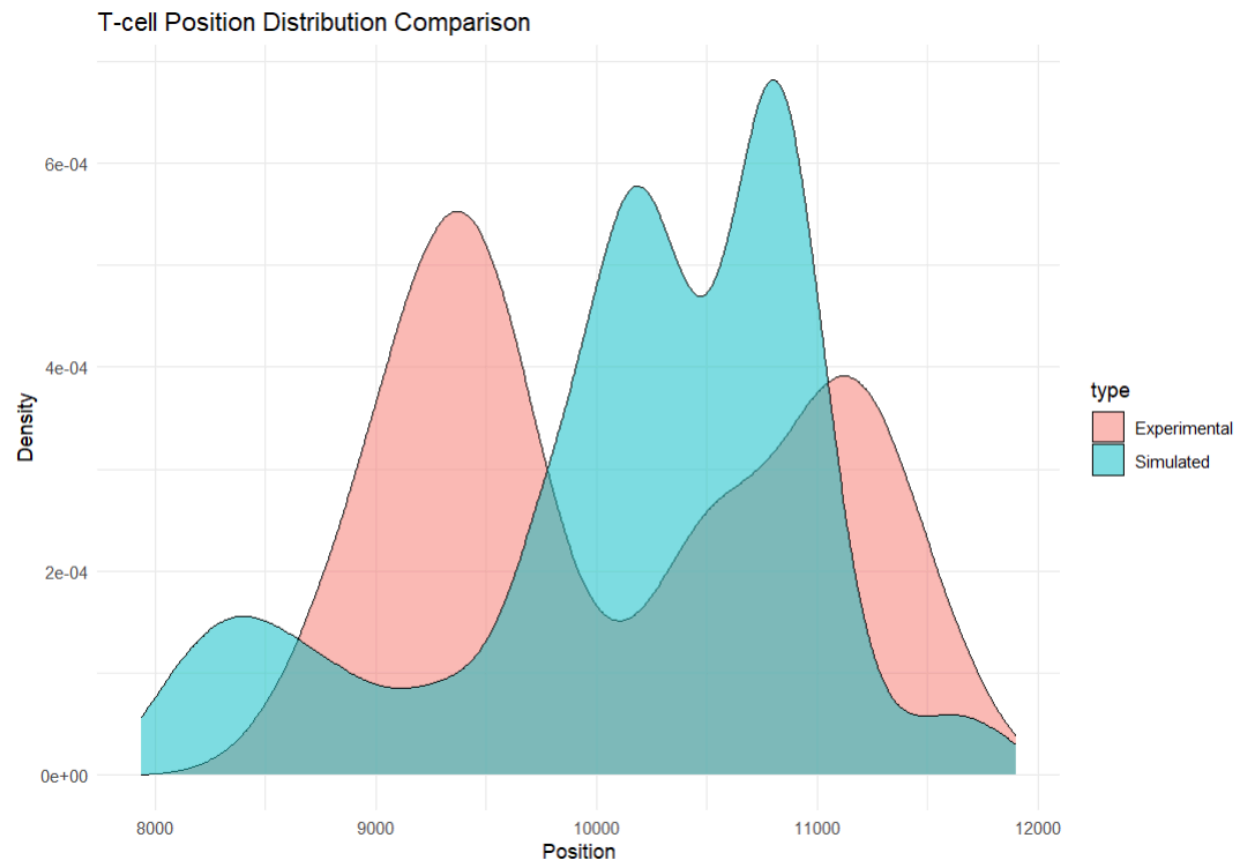


## Behavior 1

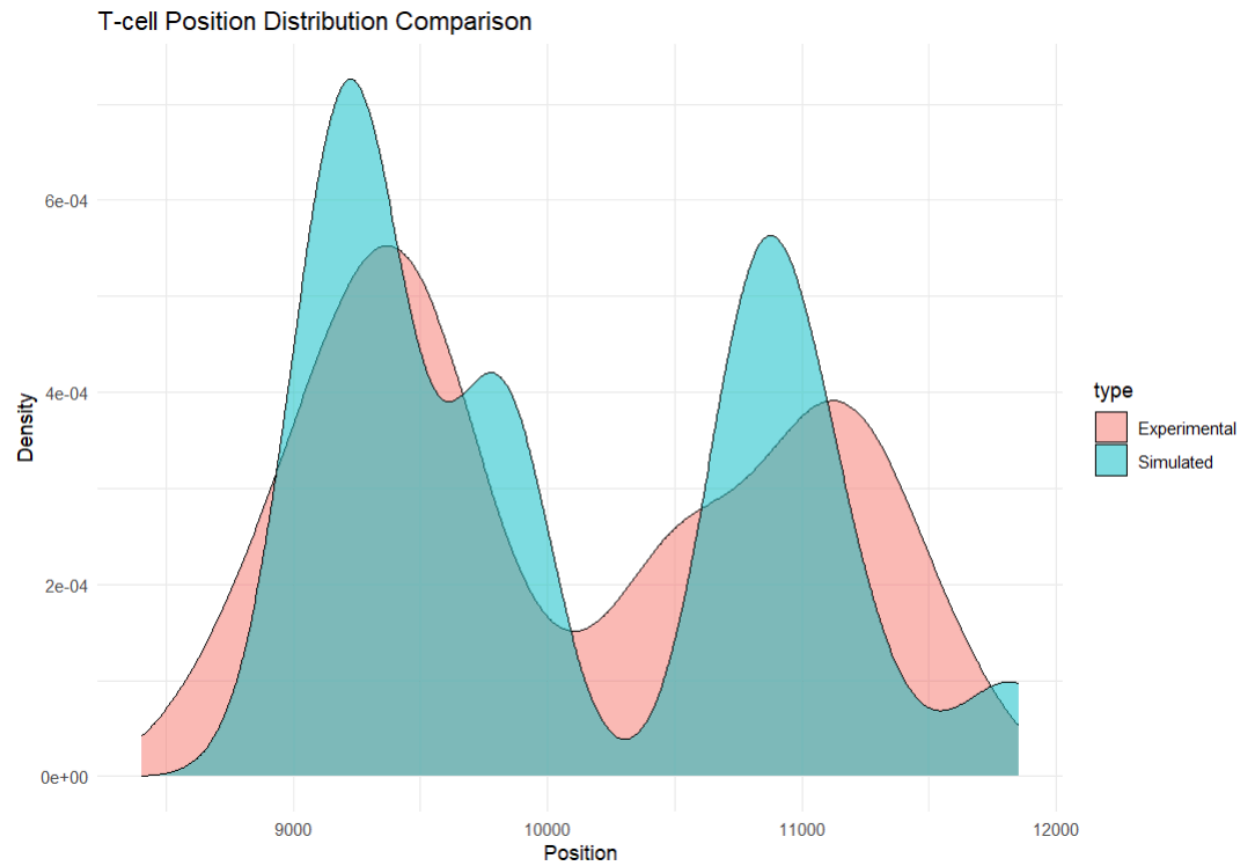
- $\alpha = -100$
- $\text{speed} = 2$
- $\text{num-steps} = 10000$



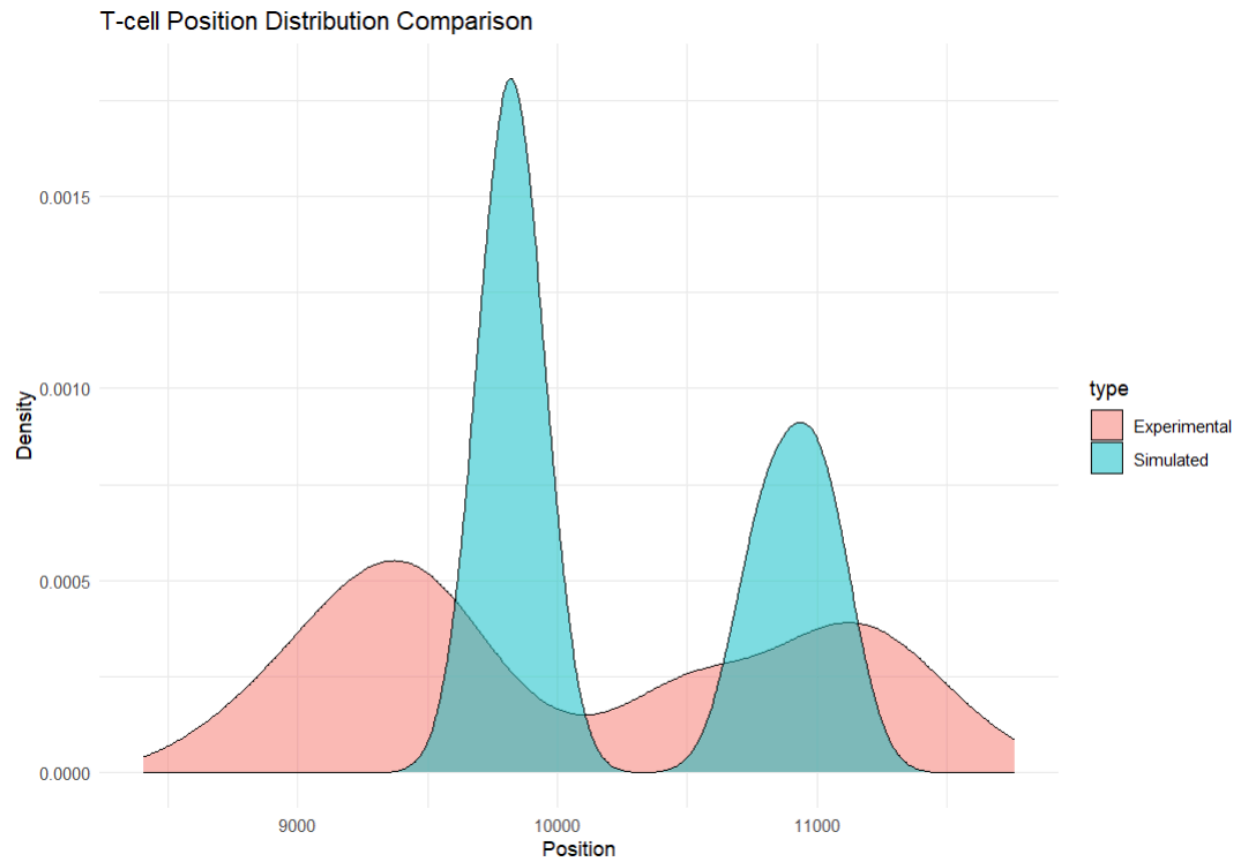
- $\alpha = -10$
- $\text{speed} = 2$
- $\text{num-steps} = 10000$



- $\alpha = -1$
- $\text{speed} = 2$
- $\text{num-steps} = 10000$



- $\alpha = 1e-6$



# Steady-State Diffusion with Degradation in 2D

## 1 Introduction

We aim to derive and solve the steady-state diffusion equation for chemokine concentration in a two-dimensional space with degradation. We assume a system where chemokines diffuse outward from cancer cells while simultaneously degrading at a constant rate.

## 2 Deriving the 2D Diffusion Equation

### 2.1 Fick's First Law

Fick's First Law states that the diffusion flux  $\mathbf{J}$  is proportional to the negative gradient of the concentration:

$$\mathbf{J} = -D\nabla c, \quad (1)$$

where:

- $\mathbf{J} = (J_x, J_y)$  is the diffusion flux (amount of substance per unit area per unit time),
- $D$  is the diffusion coefficient,
- $c(x, y, t)$  is the chemokine concentration,
- $\nabla c$  is the concentration gradient.

In Cartesian coordinates, this expands to

$$J_x = -D \frac{\partial c}{\partial x}, \quad J_y = -D \frac{\partial c}{\partial y}. \quad (2)$$

### 2.2 Fick's Second Law

Applying the principle of mass conservation, we obtain Fick's Second Law:

$$\frac{\partial c}{\partial t} = -\nabla \cdot \mathbf{J}. \quad (3)$$

Expanding the divergence in two dimensions:

$$\frac{\partial c}{\partial t} = D \left( \frac{\partial^2 c}{\partial x^2} + \frac{\partial^2 c}{\partial y^2} \right). \quad (4)$$

If chemokines degrade at a constant rate  $\delta$ , we introduce a decay term:

$$\frac{\partial c}{\partial t} = D \left( \frac{\partial^2 c}{\partial x^2} + \frac{\partial^2 c}{\partial y^2} \right) - \delta c. \quad (5)$$



### 2.3 Steady-State Approximation

For a steady-state solution, we set  $\partial c / \partial t = 0$ :

$$D \left( \frac{\partial^2 c}{\partial x^2} + \frac{\partial^2 c}{\partial y^2} \right) = \delta c. \quad (6)$$

Dividing by  $D$ :

$$\frac{\partial^2 c}{\partial x^2} + \frac{\partial^2 c}{\partial y^2} = \frac{\delta}{D} c. \quad (7)$$

## 3 Conversion to Polar Coordinates

Since the problem is radially symmetric, we use polar coordinates  $(r, \theta)$ , where

$$x = r \cos \theta, \quad y = r \sin \theta.$$

The Laplacian in polar coordinates (assuming no  $\theta$  dependence) simplifies to

$$\frac{\partial^2 c}{\partial r^2} + \frac{1}{r} \frac{\partial c}{\partial r} = \frac{\delta}{D} c.$$

Define

$$\lambda^2 = \frac{\delta}{D} \implies \lambda = \sqrt{\frac{\delta}{D}}.$$

The steady-state radial equation then becomes:

$$\frac{d^2 c}{dr^2} + \frac{1}{r} \frac{dc}{dr} - \lambda^2 c = 0.$$

## 4 Substitution $\rho = \lambda r$

To transform the above ODE into the standard modified Bessel equation of order 0, we set

$$\rho = \lambda r \implies r = \frac{\rho}{\lambda}.$$

Using the chain rule:

1. First derivative:

$$\frac{d}{dr} = \frac{d\rho}{dr} \frac{d}{d\rho} = \lambda \frac{d}{d\rho}.$$

2. Second derivative:

$$\frac{d^2}{dr^2} = \frac{d}{dr} \left( \lambda \frac{d}{d\rho} \right) = \lambda^2 \frac{d^2}{d\rho^2}.$$

Substitute these into

$$\frac{d^2 c}{dr^2} + \frac{1}{r} \frac{dc}{dr} - \lambda^2 c = 0.$$

We get:

$$\lambda^2 \frac{d^2 c}{d\rho^2} + \frac{1}{r} \left( \lambda \frac{dc}{d\rho} \right) - \lambda^2 c = 0.$$

Since  $r = \rho/\lambda$ , we have  $\frac{1}{r} = \frac{\lambda}{\rho}$ , so

$$\lambda \frac{1}{r} \frac{dc}{d\rho} = \lambda^2 \frac{1}{\rho} \frac{dc}{d\rho}.$$

Hence,

$$\lambda^2 \left[ \frac{d^2c}{d\rho^2} + \frac{1}{\rho} \frac{dc}{d\rho} - c \right] = 0.$$

Dividing by  $\lambda^2$  (which is nonzero), we obtain the **modified Bessel equation** of order 0:

$$\frac{d^2c}{d\rho^2} + \frac{1}{\rho} \frac{dc}{d\rho} - c = 0.$$

#### 4.1 Solution via Modified Bessel's Equation

This is recognized as a modified Bessel's equation of order zero. By letting  $\rho = \lambda r$ , we rewrite the ODE in the standard form of the modified Bessel equation:

$$\frac{d^2c}{d\rho^2} + \frac{1}{\rho} \frac{dc}{d\rho} - c = 0. \quad (8)$$

The general solution is a linear combination of  $I_0(\rho)$  and  $K_0(\rho)$ , the modified Bessel functions of the first and second kind of order zero :

$$c(\rho) = A I_0(\rho) + B K_0(\rho). \quad (9)$$

Reverting to  $r$ , we have

$$c(r) = A I_0(\lambda r) + B K_0(\lambda r). \quad (10)$$

#### 4.2 Applying Boundary Conditions

**Decay at Infinity.** As  $r \rightarrow \infty$ ,  $I_0(\lambda r)$  grows exponentially, while  $K_0(\lambda r)$  decays exponentially . To ensure  $c(r) \rightarrow 0$  at large  $r$ , we must set  $A = 0$ . Hence, the solution simplifies to

$$c(r) = B K_0(\lambda r). \quad (11)$$