



Green Functions of Nonperturbtive QCD

(Part I)

Si-xue Qin

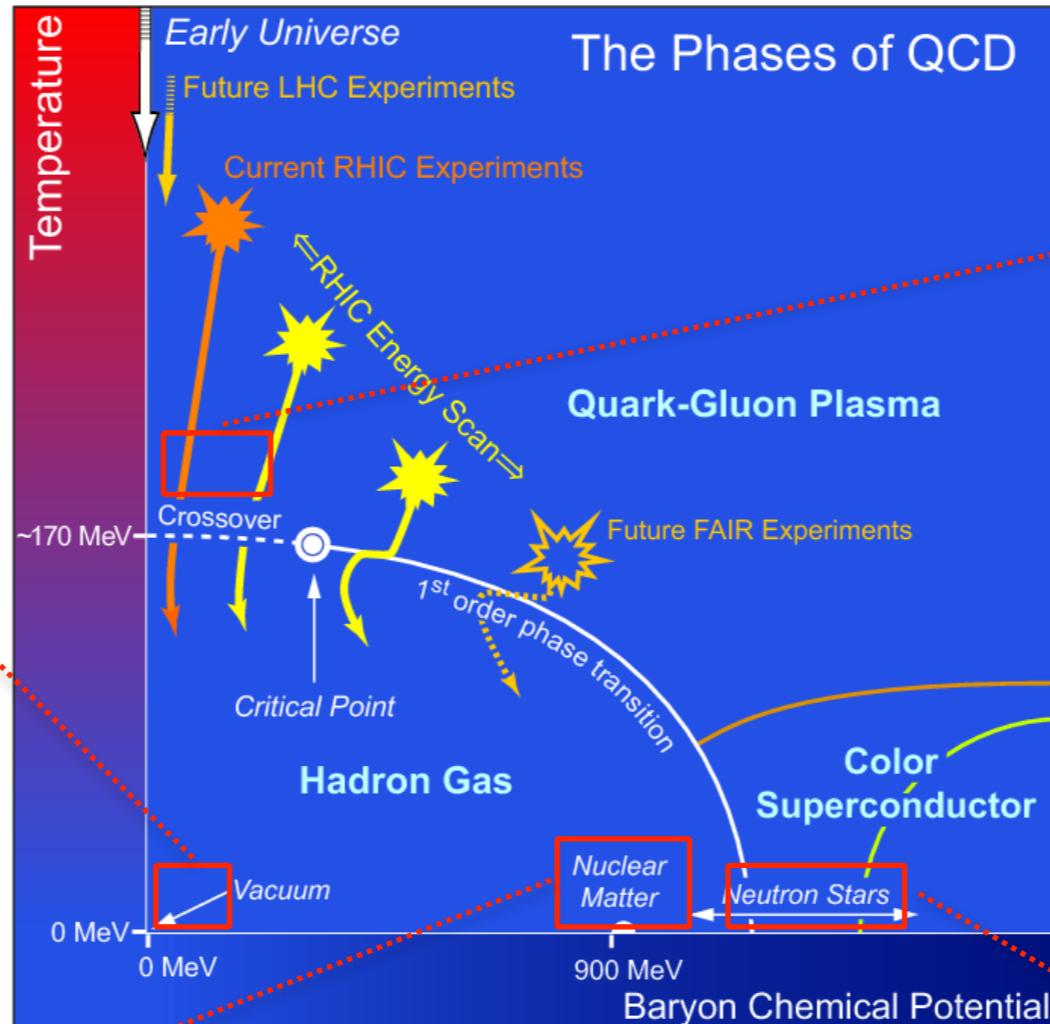
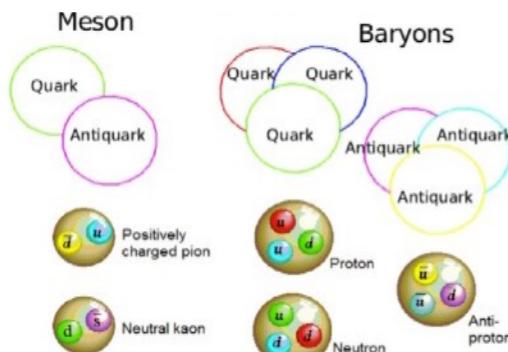
(秦思学)

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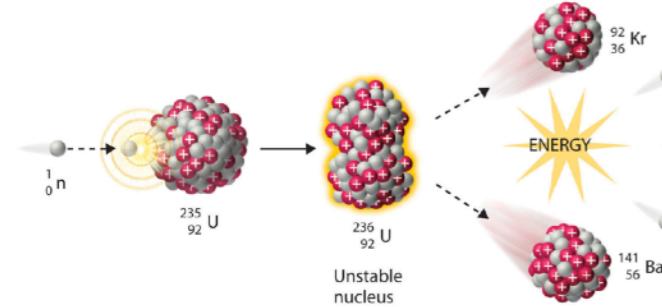
(重庆大学 物理学院)

Frontier

Hadron

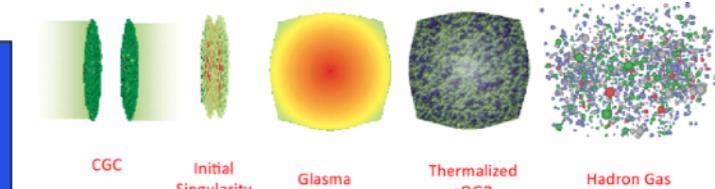


Nucleus

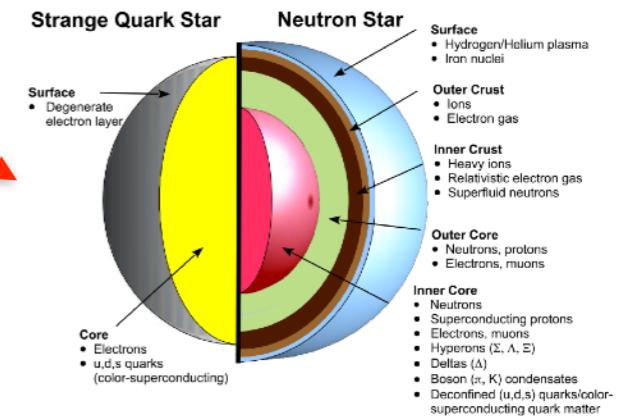


Few-body

QGP

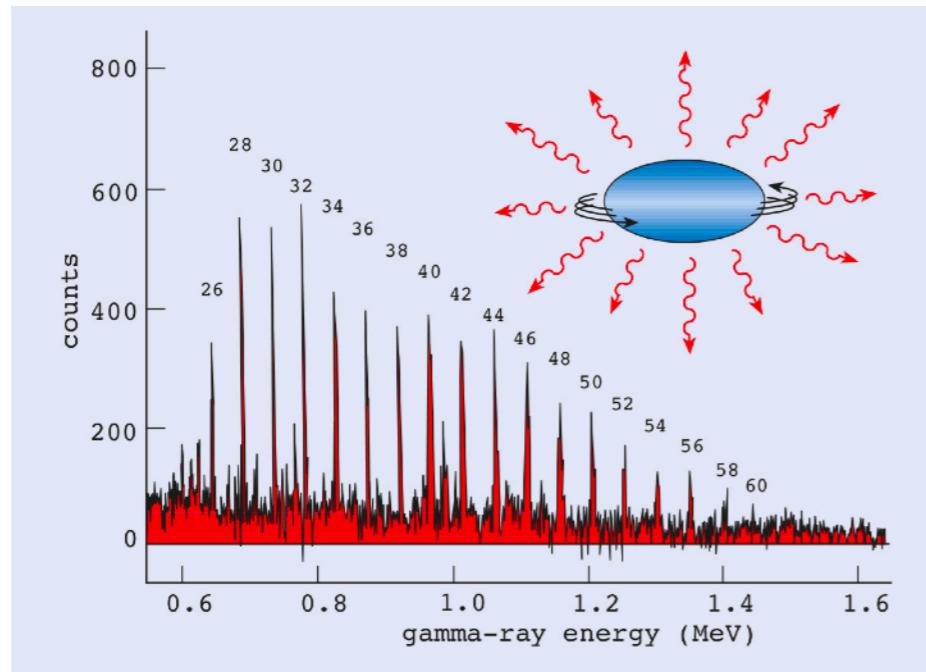


Compact Star



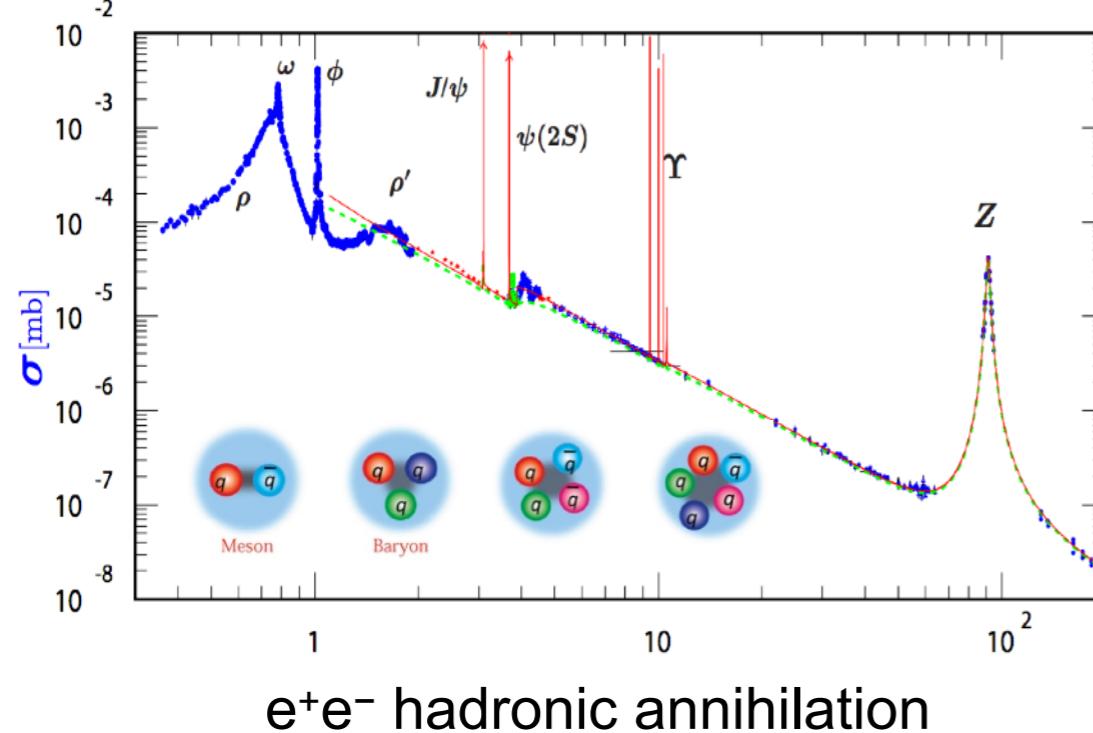
Many-body

Frontier: Experiments

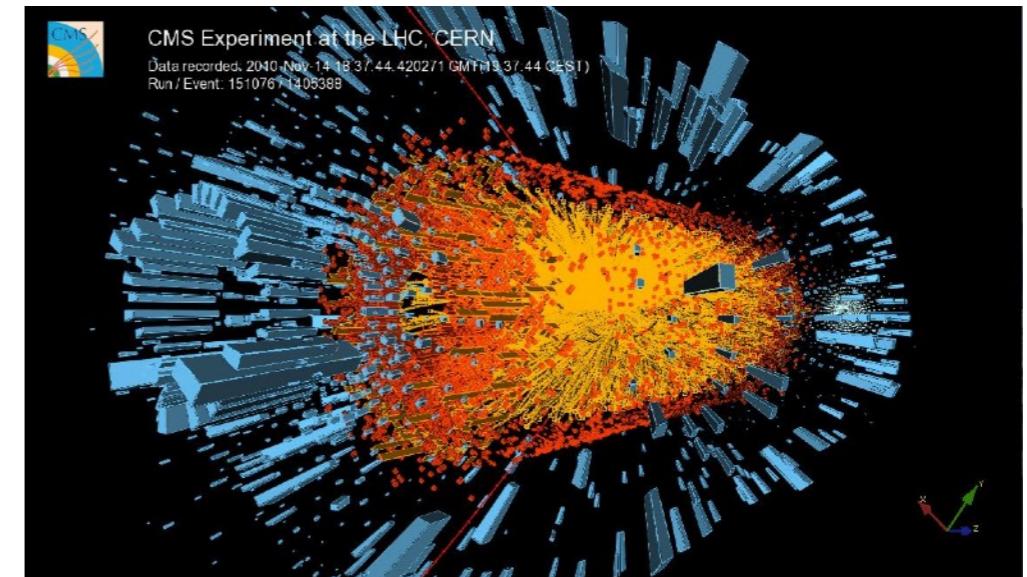


Novel states of nuclei

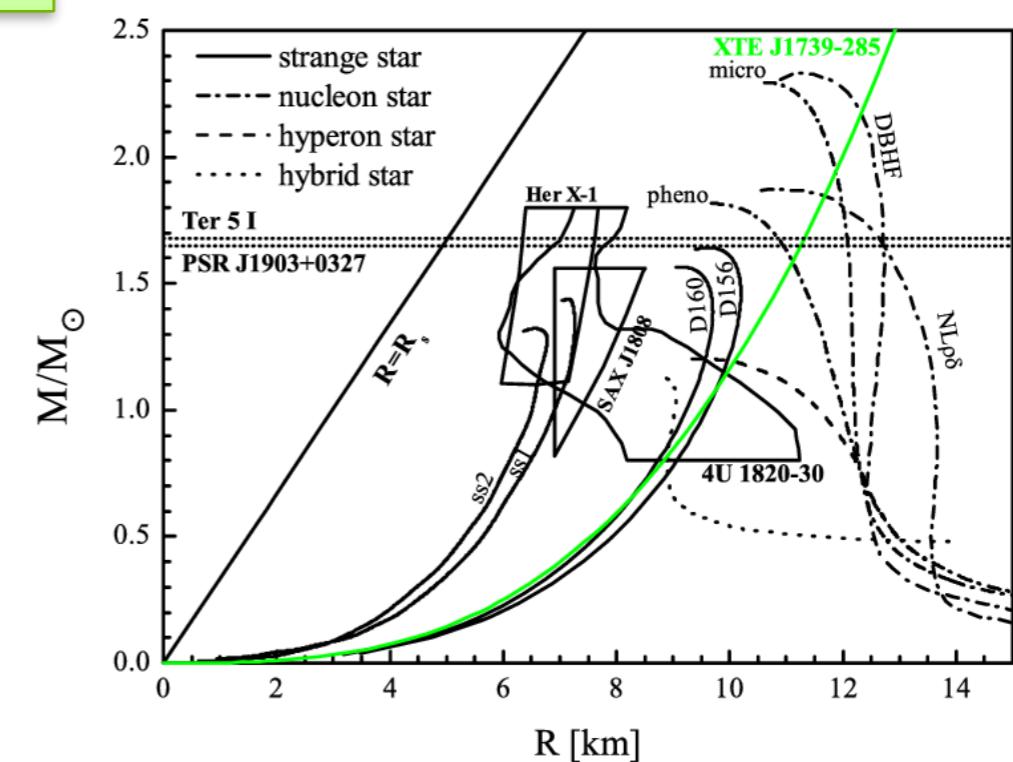
Solve QCD



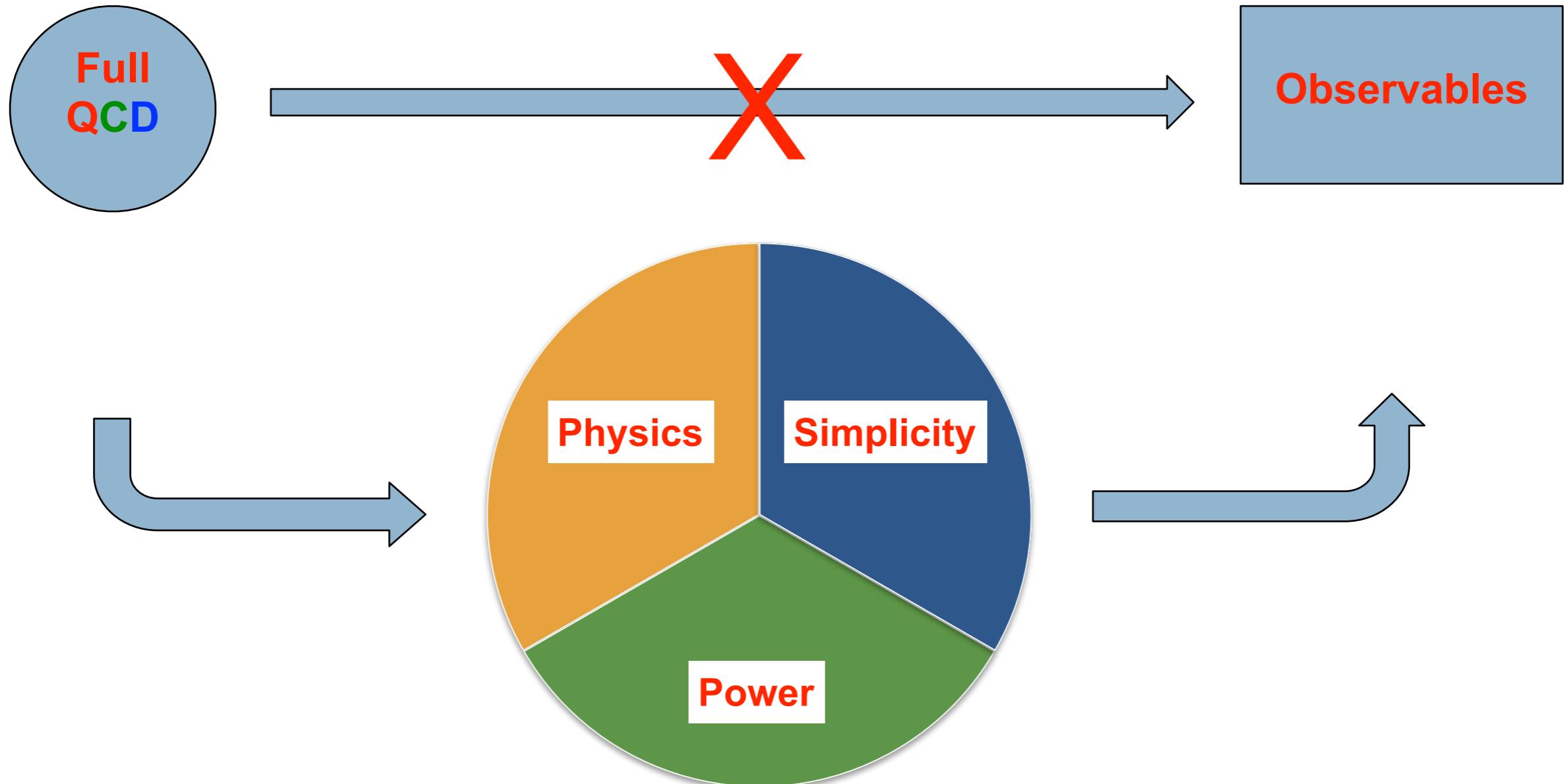
e^+e^- hadronic annihilation



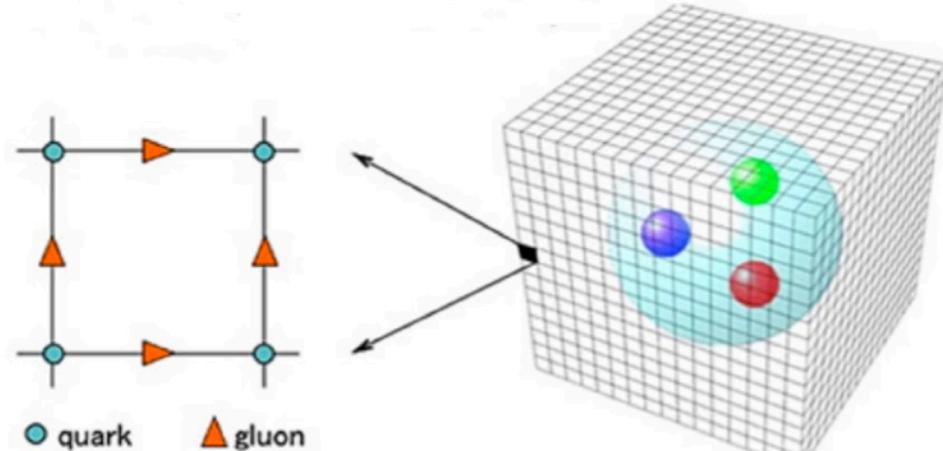
relativistic heavy-ion collision



mass-radius relation of compact stars



Lattice QCD, NJL, Sum rules, AdS/QCD, LFQCD, Effective theories, ..., Functional methods



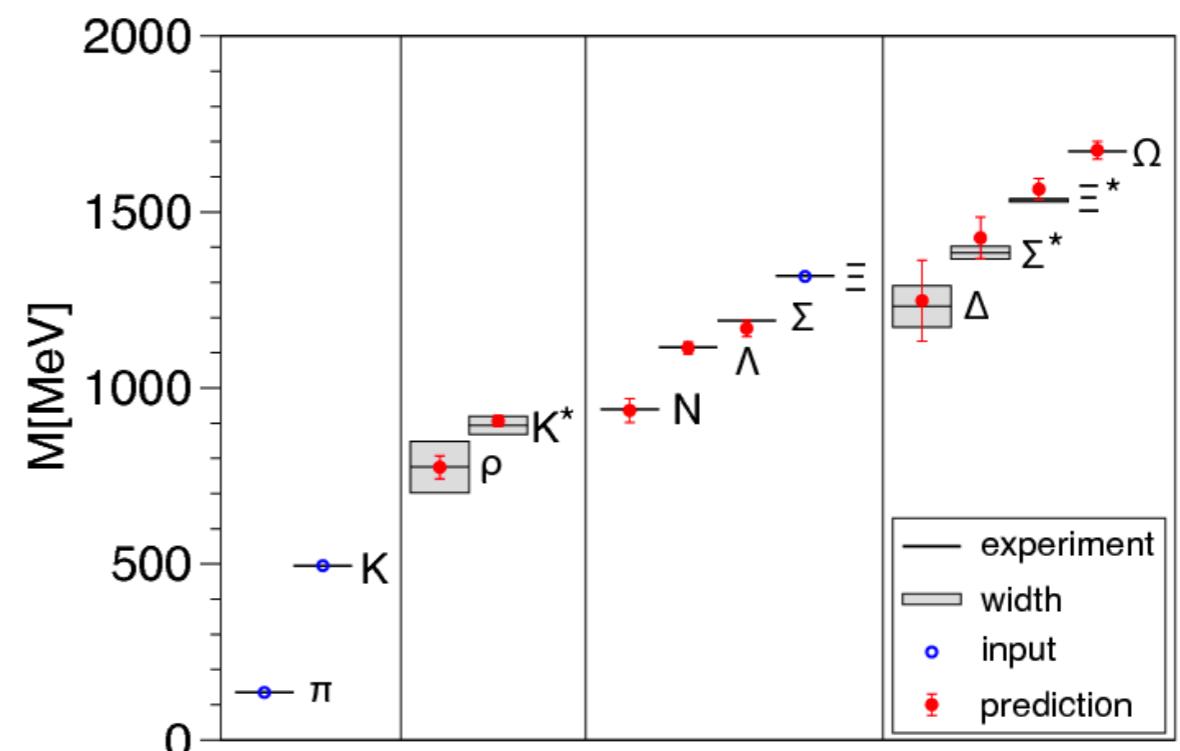
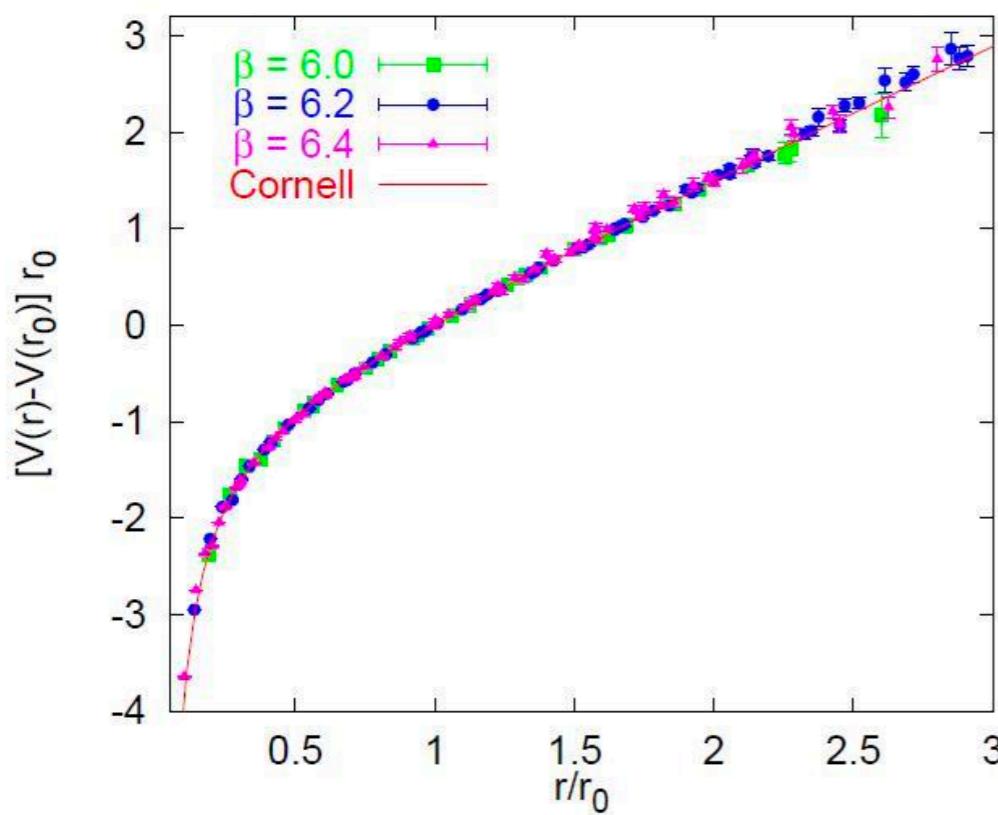
$$S = \int d^4x \left(\frac{1}{4} F^2 - \bar{q} M q \right)$$

$$Z = \int D\Lambda \det M e^{-S}$$

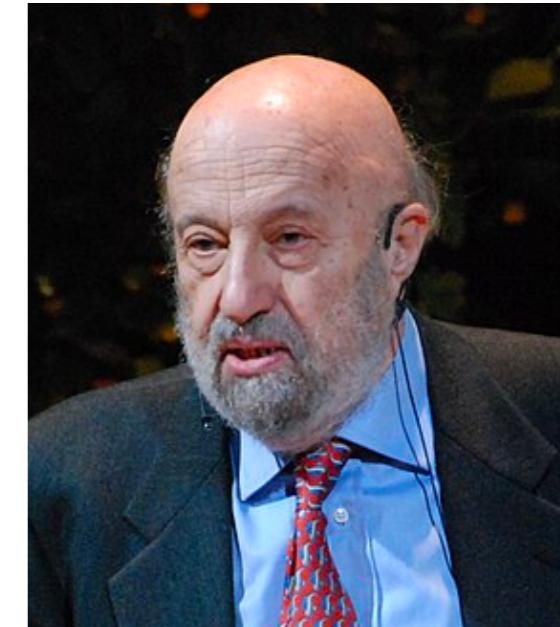
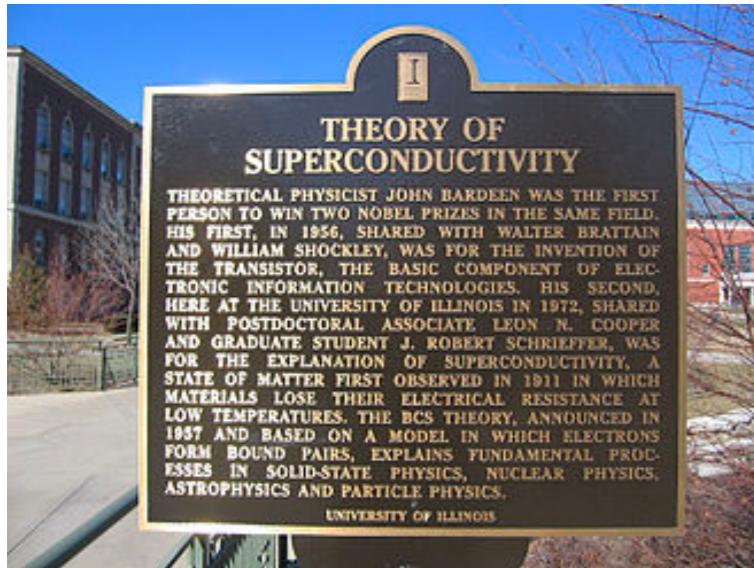
sign problem



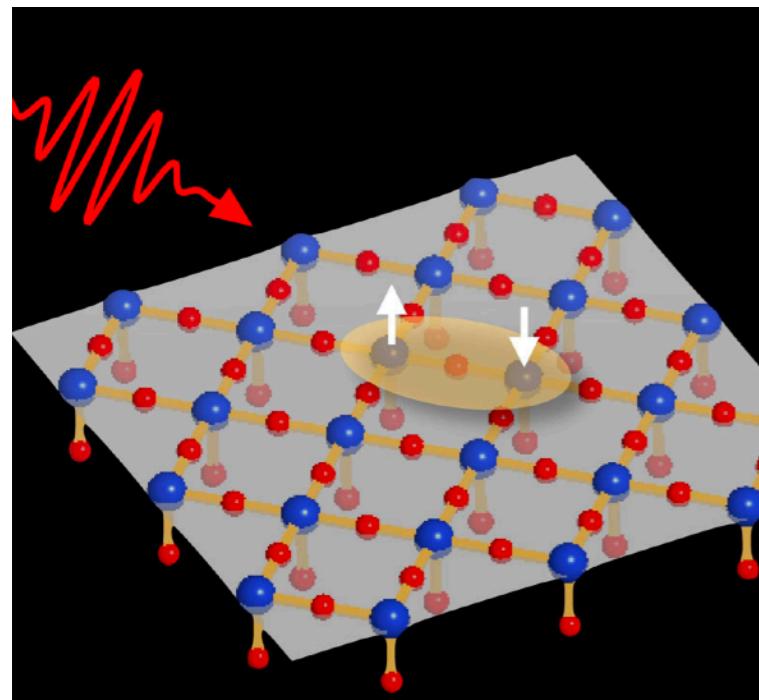
Kenneth Wilson



Frontier: Nambu—Jona-Lasinio model



Yoichiro Nambu Giovanni Jona-Lasinio



$$\mathcal{L} = i\bar{q}_L \gamma \cdot \partial q_L + i\bar{q}_R \gamma \cdot \partial q_R + \lambda(\bar{q}_L q_R)(\bar{q}_R q_L)$$

NJL: χ -symmetry

$$\Phi(\vec{x}) = \frac{1}{N_c} \text{Tr} \left\{ T \exp \left[i \int_0^\beta d\tau A_4(\vec{x}, \tau) \right] \right\}$$

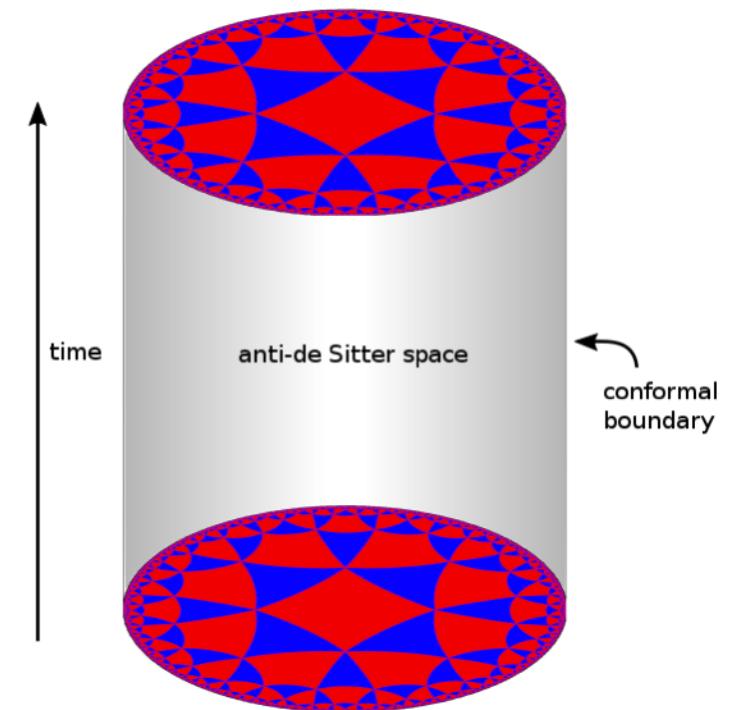
$$|\Phi(\vec{x})| = \exp(-\beta F)$$

Gauge invariant,
Not center invariant

PNJL: χ -symmetry & confinement

Mapping: String theory vs QFT

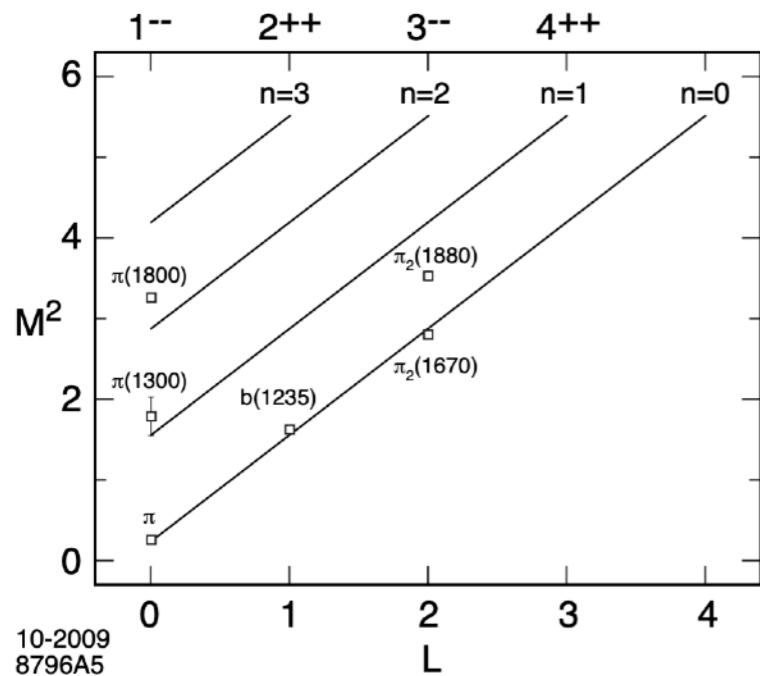
String theory on $AdS_5 \times S^5 \simeq \mathcal{N} = 4, SU(N)$ gauge theory in 4D



AdS/CFT: strong/weak coupling duality

$$\lambda = g_{YM}^2 N \quad \lambda \sim \left(\frac{\ell_{AdS}}{\ell_{string}} \right)^4 \quad \frac{\ell_{AdS}^{d-1}}{G_N} \sim \left(\frac{\ell_{AdS}}{\ell_P} \right)^{d-1} \sim N^2$$

AdS/QCD: holographic theory



$$\frac{\eta}{s} = \frac{1}{4\pi}$$

No proof



Juan Maldacena

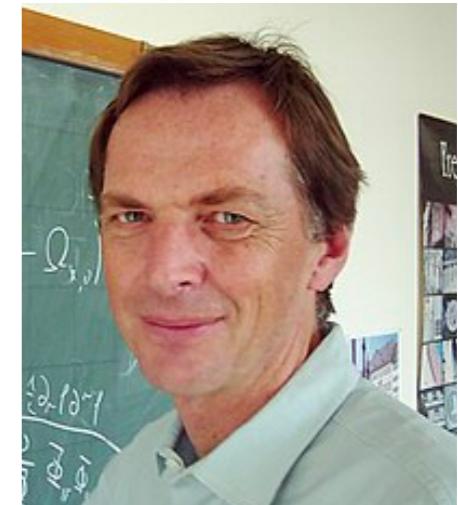
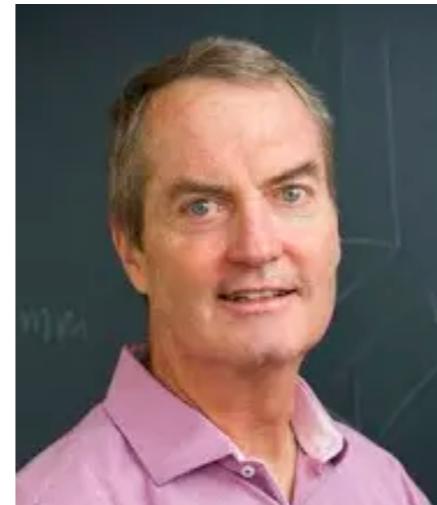
Frontier: Functional methods

Startpoint: QFT partition function **solvable** and **traceable** approximation to QCD

$$Z_{QFT}[\phi, j] = \int D\phi e^{-S[\phi, j]}$$

Differential version:

Functional renormalization equation

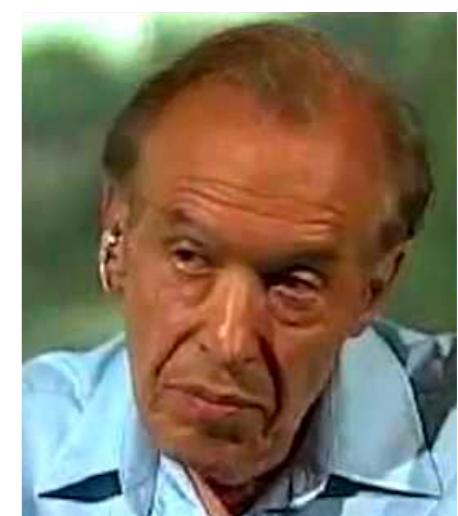


Joseph Polchinski Christof Wetterich

Integral version:

Dyson-Schwinger equation

NOT limited in HEP

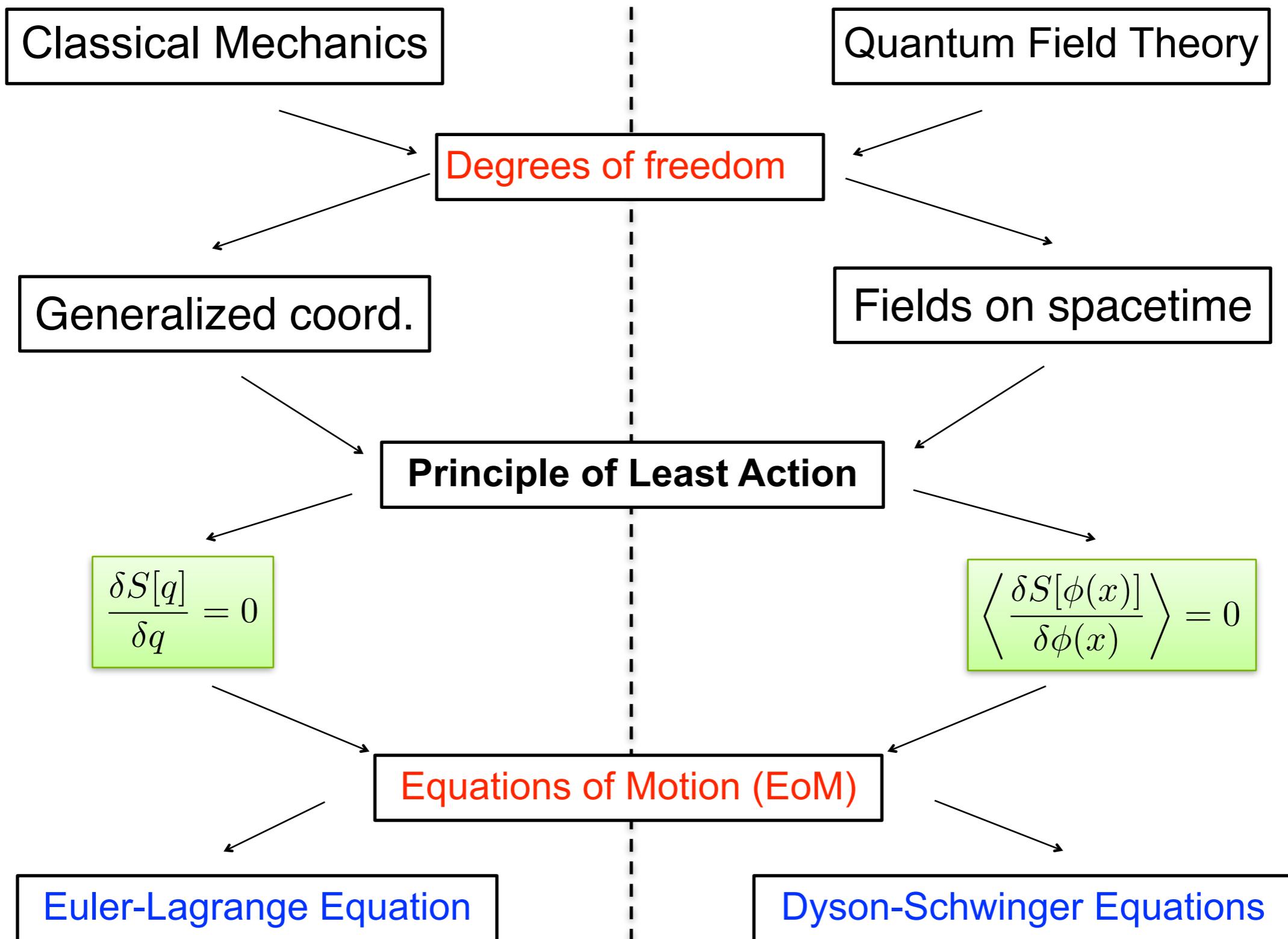


Freeman Dyson

Julian Schwinger

Framework

Framework: Equation of motion



Framework: Green function

Properties of Green function:

See, e.g., arXiv:hep-ph/0208074, arXiv:0909.0703

- express particle scattering
- encode all observables
- use Feynman diagrams

$$G^{(n)}(x_1, \dots, x_n) = \langle \Omega | T[\phi(x_1) \dots \phi(x_n)] | \Omega \rangle = \frac{\int [D\phi] e^{-S[\phi]} \phi(x_1) \dots \phi(x_n)}{\int [D\phi] e^{-S[\phi]}}$$

Generating function:

- Partition function of QFT

$$Z[J] = \int [D\phi] e^{-S[\phi] + \int d^4x J(x)\phi(x)}$$

$$G^{(n)}(x_1, \dots, x_n) = \frac{1}{Z[0]} \left. \frac{\delta^n Z[J]}{\delta J(x_1) \dots \delta J(x_n)} \right|_{J=0}$$

- Identity of generating function:

$$0 = \int [D\phi] \frac{\delta}{\delta \phi(x)} \exp \left\{ -S[\phi] + \int d^4x J(x)\phi(x) \right\}$$

$$0 = \int [D\phi] \left\{ -\frac{\delta S[\phi]}{\delta \phi(x)} + J(x) \right\} \exp \left\{ -S[\phi] + \int d^4x J(x)\phi(x) \right\}$$

- Generating equation of DSE:

$$\left\langle \frac{\delta S[\phi]}{\delta \phi(x)} \right\rangle_J = \frac{1}{Z[J]} \frac{\delta S[\delta/\delta J]}{\delta \phi(x)} Z[J] = J(x)$$

Framework: Symmetry breaking

- Classical theory:

$$S[\phi] \xrightarrow{\text{EoM}} O$$

- Quantum theory:

$$Z[\phi, j] = \int [D\phi] e^{-S[\phi] + j\phi} \xrightarrow{\text{EoM}} \{O, O', \dots\}$$

Symmetric	$S[\phi]$	$[D\phi]$	$\{O\}$
Anomalous	$S[\phi]$	$\cancel{[D\phi]}$	$\{O\}$
Spontaneous	$S[\phi]$	$[D\phi]$	$\{O, O', \dots\}$

Framework: Symmetry breaking

Symmetry	Anomaly	Fate
$SU(3)_C$	no	unbroken
$SU(N)_V \times U(1)_V$	no	unbroken
$U(1)_A$	yes	not a symmetry
$SU(N)_A$	no	spontaneously broken
Scale Invariance	yes	not a symmetry
CP	yes?	not a symmetry?

QCD Lagrangian:

$$\mathcal{L}_{\text{QCD}} = \bar{\psi}_i (i(\gamma^\mu D_\mu)_{ij} - m\delta_{ij})\psi_j - \frac{1}{4}G_{\mu\nu}^a G_a^{\mu\nu}$$

Quark gap equation:

- ▶ Functional differential of action:

$$\frac{\delta S}{\delta \bar{\psi}} = (i\cancel{D} - m_0)\psi + g A_\mu^a \gamma^\mu t^a \psi$$

- ▶ Generating equation of DSE:

$$\frac{1}{Z} \left\{ (i\cancel{D} - m_0) \frac{\delta}{\delta \bar{\eta}} + g \frac{\delta}{\delta J_a^\mu} \gamma^\mu t^a \frac{\delta}{\delta \bar{\eta}} \right\} \cdot Z = \eta(x)$$

- ▶ Extraction of two-point function:

$$(i\cancel{D} - m) \frac{\delta^2 W}{\delta \bar{\eta}_x \delta \eta_y} \Big|_{j=0} - g \gamma^\mu t^a \frac{\delta^3 W}{\delta J_{ax}^\mu \delta \bar{\eta}_x \delta \eta_y} \Big|_{j=0} = \delta(x - y)$$

$$Z[J] = \sum_{N=0}^{\infty} \frac{(-1)^N}{N!} W^N[J] = e^{-W[J]}$$

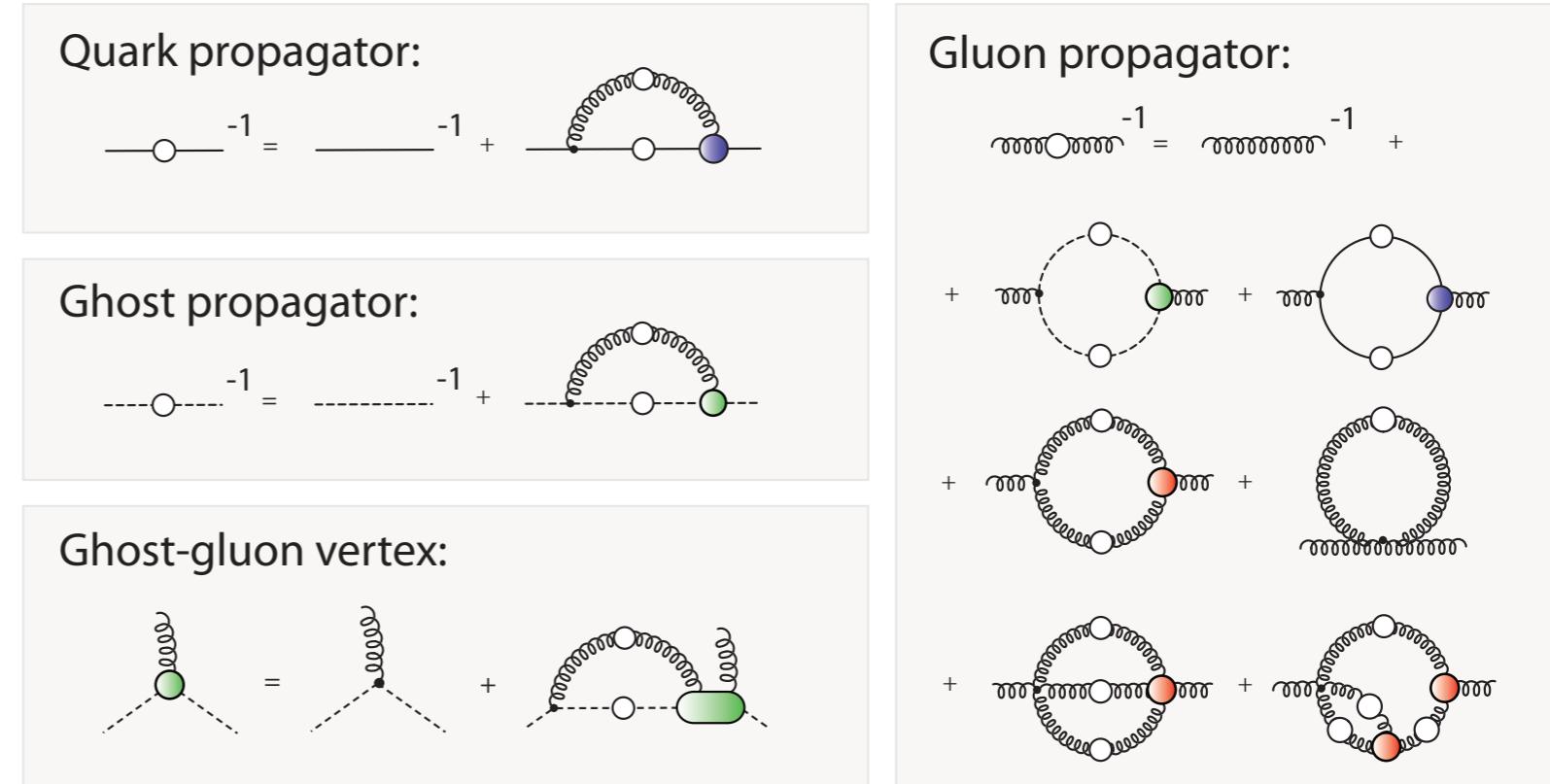
Quark propagator:



Framework: Disadvantage vs Advantage

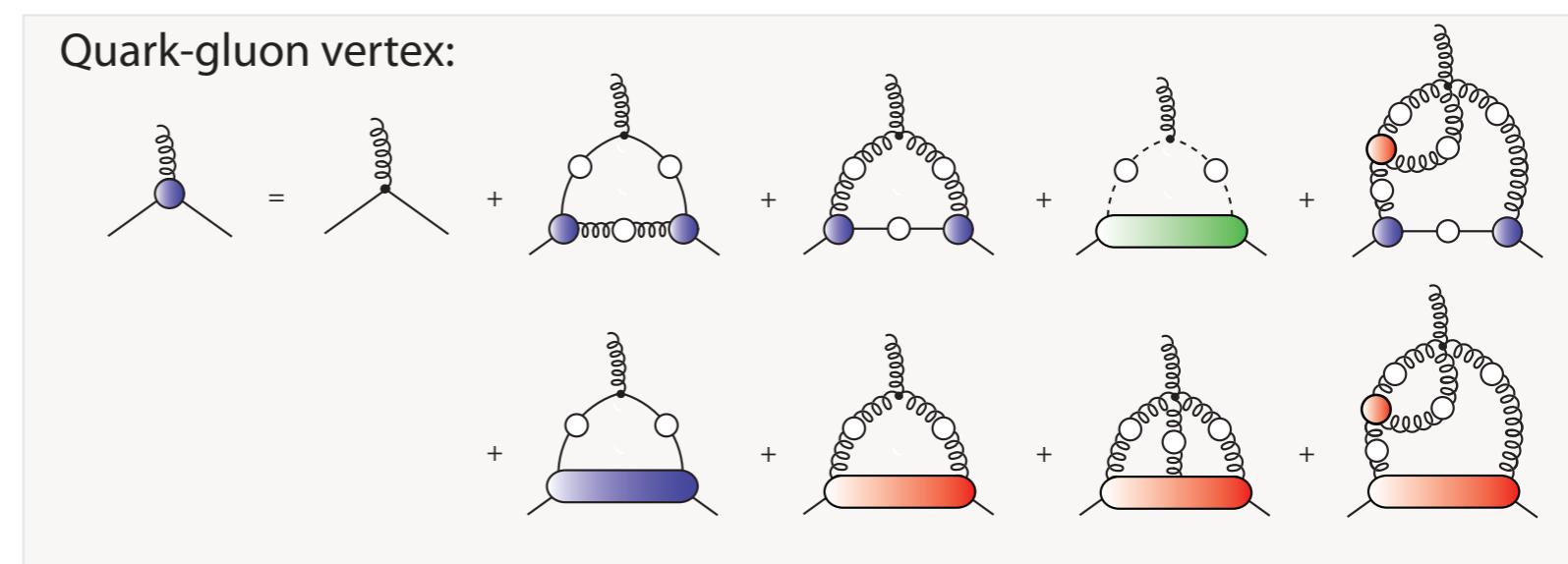
Disadvantage:

- ◆ Most equations are very complicated.
- ◆ Multi-order Green functions couple together.



Advantage:

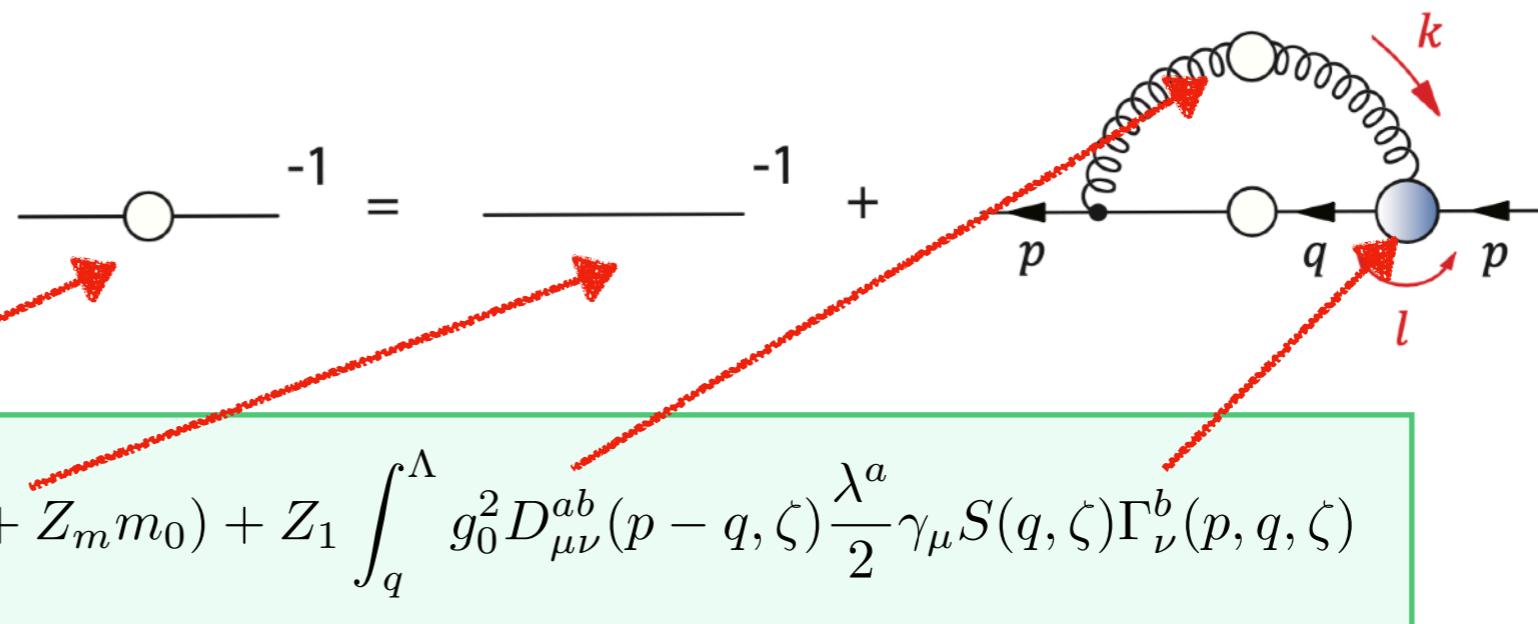
- ◆ Solid connection to QCD
 - ◆ Controllable complexities
- Modeling** **Truncation**



Formalism

Quark gap equation:

- Quark self-energy



Solution structure:

- Decomposition of Lorentz structure
- Equations of dressing functions

$$S(p, \zeta) = \frac{1}{i\cancel{p}A(p^2, \zeta^2) + B(p^2, \zeta^2)} = \frac{Z(p^2, \zeta^2)}{i\cancel{p} + M(p^2)}$$

$$A(p^2) = 1 + \frac{1}{4p^2} \text{tr}[-i\cancel{p}\Sigma(p^2)], \quad B(p^2) = \frac{1}{4} \text{tr}[\Sigma(p^2)]$$

$$A_x = 1 + \frac{1}{6\pi^3} \int_0^\infty y dy \frac{\Theta_A(x, y; A_x, B_x, A_y, B_y)}{y A_y^2 + B_y^2}$$

$$B_x = m + \frac{1}{6\pi^3} \int_0^\infty y dy \frac{\Theta_B(x, y; A_x, B_x, A_y, B_y)}{y A_y^2 + B_y^2}$$

Formalism: One-body equation

Non-linear equation:

- Discrete momenta

$$\{p_i^2\} = \{10^{t_i}\}, \quad t_i \in [-\Lambda, \Lambda]$$

$$\vec{G}(\vec{R}) = \vec{R} - \vec{F}(\vec{R}), \quad H_{ij} = \delta_{ij} - \frac{\partial F_i(\vec{R})}{\partial R_j}$$

Trivial iteration, Newton's iteration, Broyden's iteration, ...

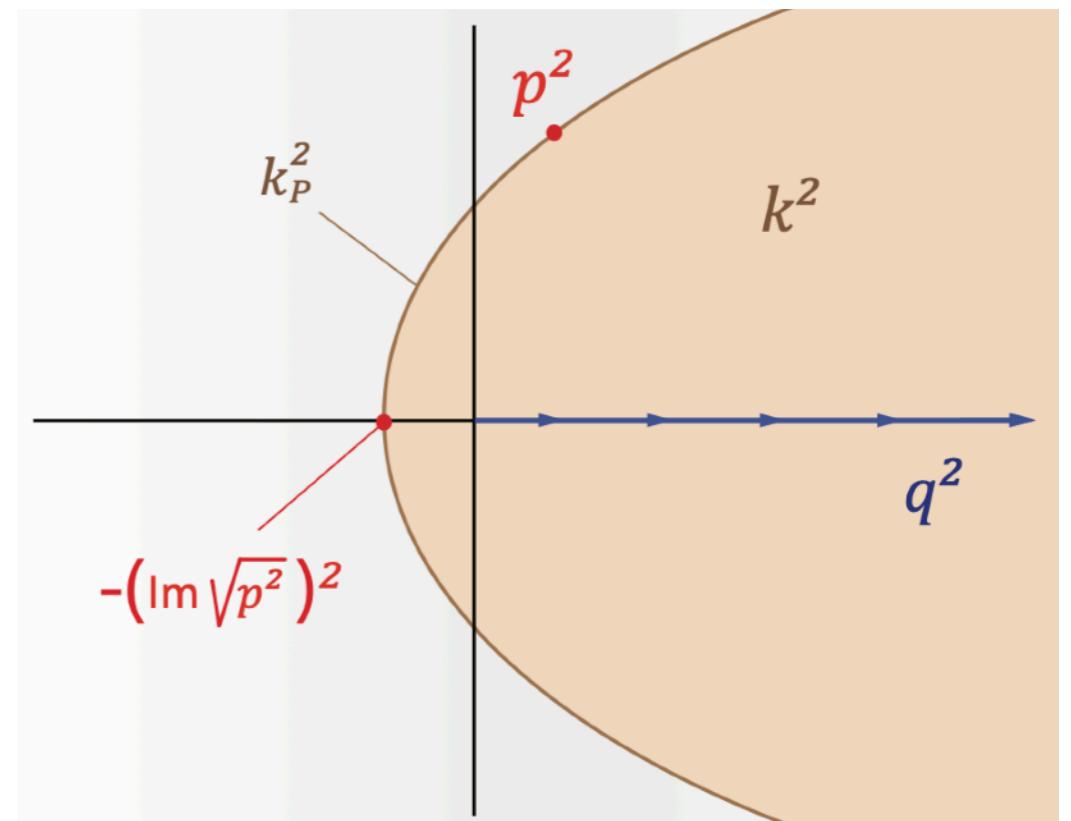
- Complex momenta

$$P^2 = -M^2, \quad P_\mu = (\vec{P}, P_4), \quad P_4 = i\sqrt{\vec{P}^2 + M^2}$$

$$q = (\vec{q}, q_4 + iM), \quad \Re(q^2) + M^2 \geq \frac{\Im^2(q^2)}{M^2}$$

- Complex integral

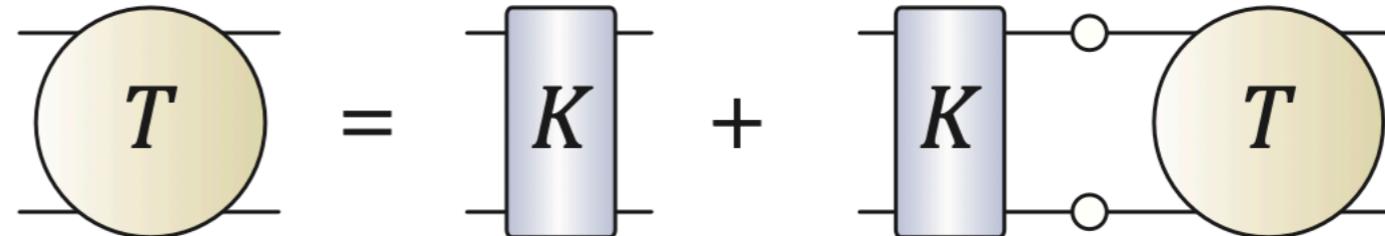
$$F(q'^2) = \frac{\oint_C \frac{F(q^2)}{q^2 - q'^2} dq^2}{\oint_C \frac{1}{q^2 - q'^2} dq^2}$$



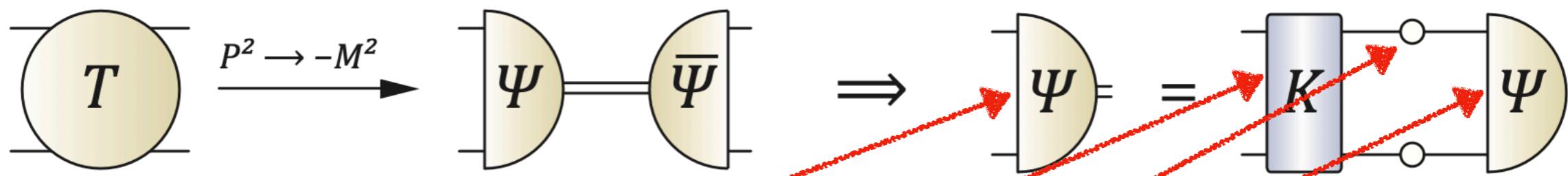
Formalism: Two-body equation

Bethe-Salpeter equation:

- DSE of T-matrix



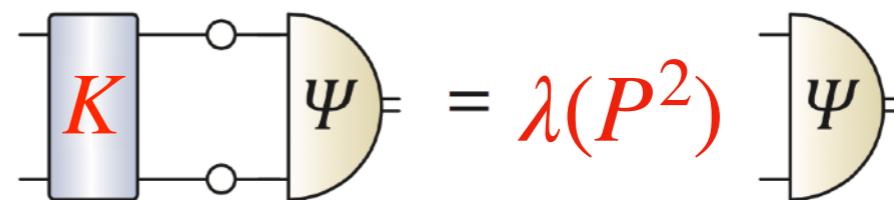
- Bound-state pole



$$\Gamma(k, P)_{\alpha\beta} = \int_q^\Lambda K_{\alpha\beta}^{\rho\sigma}(q, k, P) [S(q_+) \Gamma(q, P) S(q_-)]_{\rho\sigma}$$

- Schrödinger equation

$$H |\psi\rangle = E |\psi\rangle$$



Formalism: Two-body equation

Linear equation:

$$\text{[Diagram: A gray rectangle with two horizontal lines entering and two exiting, followed by a bracket containing } \left(\begin{array}{c} \bullet \\ \bullet \end{array} \right)^{-1} - \text{ a white rectangle labeled } K \text{, followed by another gray rectangle with two horizontal lines}] = \text{ [Diagram: A gray rectangle with two horizontal lines]} = \text{ [Diagram: A gray rectangle with two horizontal lines]}$$

► BSA normalization

$$\lim_{\text{on-shell}} \frac{i}{p^2 - m^2 + i\epsilon} = \text{[Diagram: A vertex with two incoming lines and one outgoing line, followed by a bracket containing } \left(\begin{array}{c} \bullet \\ \bullet \end{array} \right)^{-1} - \text{ a white rectangle labeled } K \text{, followed by a vertex with one incoming line and two outgoing lines]} = 1$$

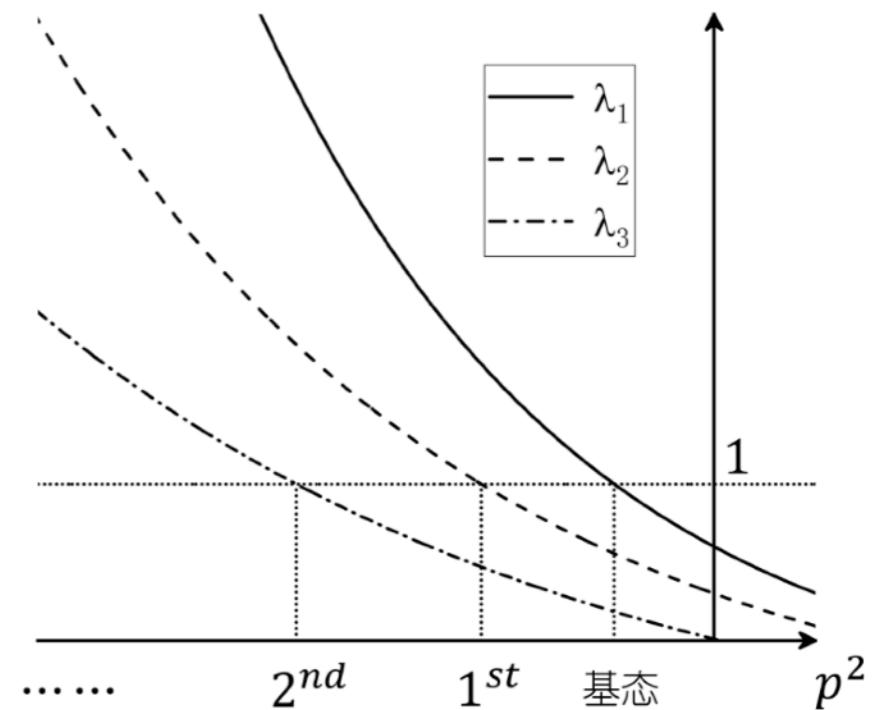
► Lorentz structure

$$\Gamma(P, q) \rightarrow \{\mathbb{1}, \hat{P}, \not{q}, \hat{P}\not{q}\} \times f_i(P^2, q^2, P \cdot q)$$

$$\lambda(P^2) \vec{F}_X(P^2) = \mathbf{K}_X(P^2) \vec{F}_X(P^2)$$

$$\Gamma_\mu(P, q) \rightarrow \{\mathbb{1}, \hat{P}, \not{q}, \hat{P}\not{q}\} \times \{\gamma_\mu, P_\mu, q_\mu\} \times f_i(P^2, q^2, P \cdot q)$$

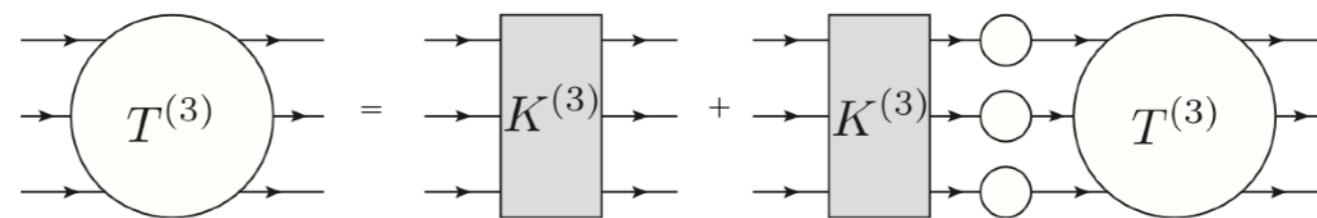
$J = 0$	$J = 1$
$\tau_1 = \mathbb{1}$	$\tau_1^\mu = \gamma^\mu$
$\tau_2 = \hat{P}$	$\tau_2^\mu = \gamma^\mu \hat{P}$
$\tau_3 = z \not{q}_T$	$\tau_3^\mu = i \hat{q}^\mu$
$\tau_4 = i [\not{q}, \hat{P}]$	$\tau_4^\mu = z \hat{q}^\mu \hat{P}$
	$\tau_5^\mu = z (\gamma^\mu \not{q}_T - \hat{q}^\mu)$
	$\tau_6^\mu = i \hat{q}^\mu \hat{P} - \frac{i}{2} \gamma^\mu [\not{q}, \hat{P}]$
	$\tau_7^\mu = \hat{q}^\mu \not{q}_T - \frac{1}{3} \hat{q}_T^2 \gamma^\mu$
	$\tau_8^\mu = \frac{1}{3} \hat{q}_T^2 \gamma^\mu \hat{P} - \frac{1}{2} \hat{q}^\mu [\not{q}, \hat{P}]$



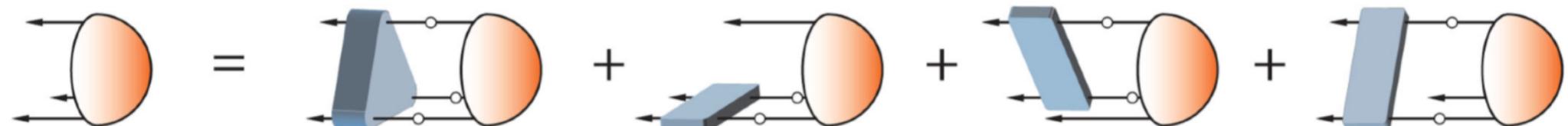
Formalism: Three-body equation

Faddeev equation:

- DSE of T-matrix



- Scatter kernel



Amplitude structure:

$$\Psi = \Psi_{\text{Color}} \otimes \Psi_{\text{Flavor}} \otimes \Psi_{\text{Dirac}}$$

Dirac (J): $4 \times 4 \times 4 \times 4$ Flavor (I): $2 \times 2 \times 2 \times 2$ Color (0): $3 \times 3 \times 3$

- Flavor structure

$2 \otimes 2 \otimes 2$	T	T_z	$\frac{3}{2}$	$\frac{1}{2}$	$-\frac{1}{2}$	$-\frac{3}{2}$
4	$\frac{3}{2}$	uuu	$\frac{1}{\sqrt{3}} (uud + udu + duu)$	$\frac{1}{\sqrt{3}} (udd + dud + ddu)$	ddd	
2	$\frac{1}{2} (\mathsf{F}_S)$		$-\frac{1}{\sqrt{6}} (udu + duu - 2uud)$	$\frac{1}{\sqrt{6}} (udd + dud - 2ddu)$		
2	$\frac{1}{2} (\mathsf{F}_A)$		$\frac{1}{\sqrt{2}} (udu - duu)$	$\frac{1}{\sqrt{2}} (udd - dud)$		

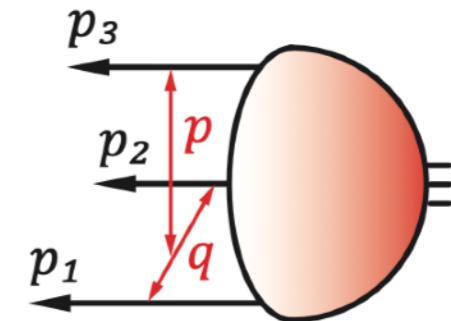
Formalism: Three-body equation

► Dirac structure

$$W_\mu = \frac{1}{2} \epsilon_{\mu\nu\alpha\beta} P^\nu J^{\alpha\beta}$$

Pauli-Lubanski operator

$$J = S + L, \quad L = L_p + L_q$$



s	l	#
3/2	2	4
3/2	2	4
3/2	2	4
3/2	2	4

s	l	#
3/2	2	4
3/2	1	4
3/2	1	4
3/2	1	4

s	l	#
1/2	0	8
1/2	1	8
1/2	1	8
1/2	1	8

J=1/2: 64

J=3/2: 128

► Momentum structure

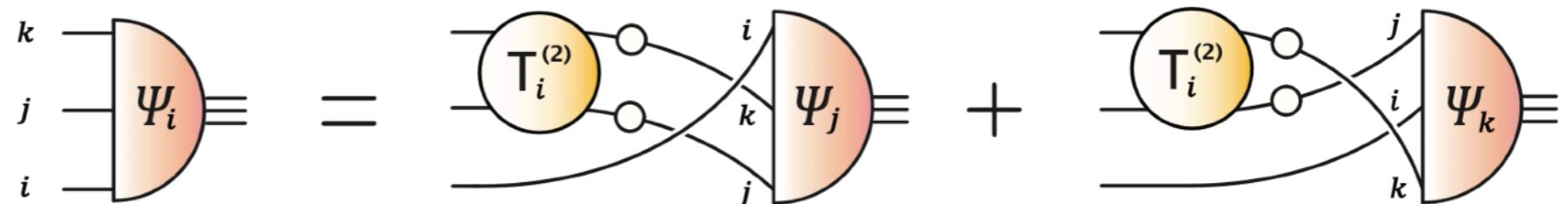
$$(P, q, p) \rightarrow (q^2, p^2, P \cdot q, P \cdot p, p \cdot q)$$

rest frame: 5-dim, 100+

Formalism: Three-body equation

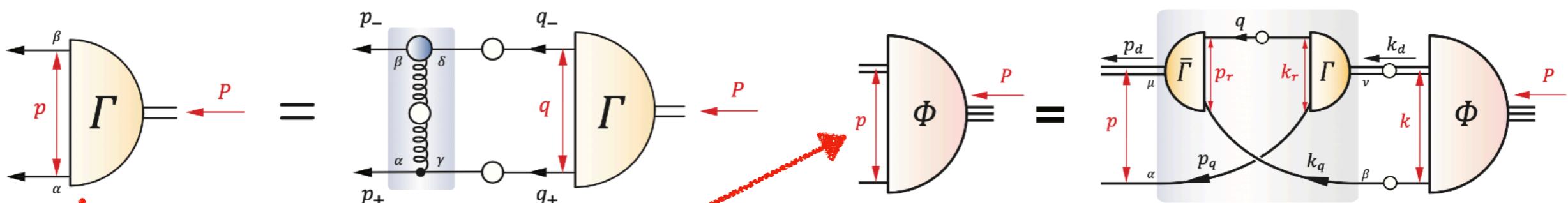
Diquark picture:

- Two-body matrix



- Diquark correlation

2 x two-body equation



$$\Gamma(k, P)_{\alpha\beta} = \int_q^\Lambda K_{\alpha\beta}^{\rho\sigma}(q, k, P) [S(q_+) \Gamma(q, P) S^T(q_-)]_{\rho\sigma}$$

SC - PS AV - VC

$$\Phi_{\alpha\beta}^a(p, P) = \int_k^\Lambda \{ \Gamma^b(k_r, k_d) S^T(q) \bar{\Gamma}^a(p_r, -p_d) S(k_q) \Phi^c(k, P) \}_{\alpha\beta} D^{bc}(k_d)$$

$$\sim \frac{1}{M^2}$$

Formalism: Summary of equations

- **One-body gap equation**



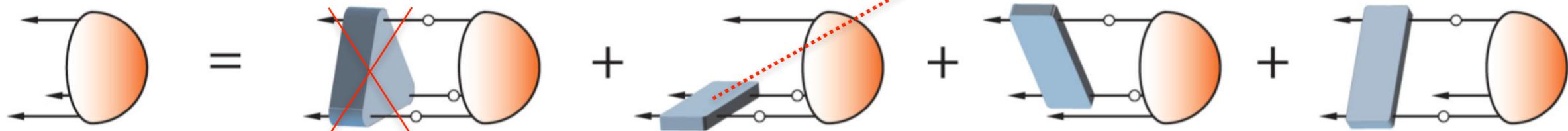
Gluon propagator

- **Two-body Bethe-Salpeter equation**



Scattering Kernel

- **Three-body Faddeev equation**



Trial

Trial: Gluon model

Tensor structure

$$D_{\mu\nu}^{ab}(k) = \frac{\delta^{ab}}{k^2} \left[\left(\delta_{\mu\nu} - \frac{k_\mu k_\nu}{k^2} \right) \mathcal{D}(k^2) + \xi \frac{k_\mu k_\nu}{k^2} \right]$$

$$g^2 D_{\mu\nu}(k) = \left(\delta_{\mu\nu} - \frac{k_\mu k_\nu}{k^2} \right) \mathcal{G}(k^2)$$

Simple model

$$\mathcal{G}(k^2) = \frac{1}{M_G^2} \theta(\Lambda^2 - k^2)$$

contact

$$\mathcal{G}(k^2) = 4\pi^2 \eta^2 \delta^4(k)$$

constant

Realistic model

$$\mathcal{G}(k^2) = \mathcal{G}_{\text{IR}}(k^2) + \frac{4\pi}{k^2} \alpha_{\text{pQCD}}(k^2)$$

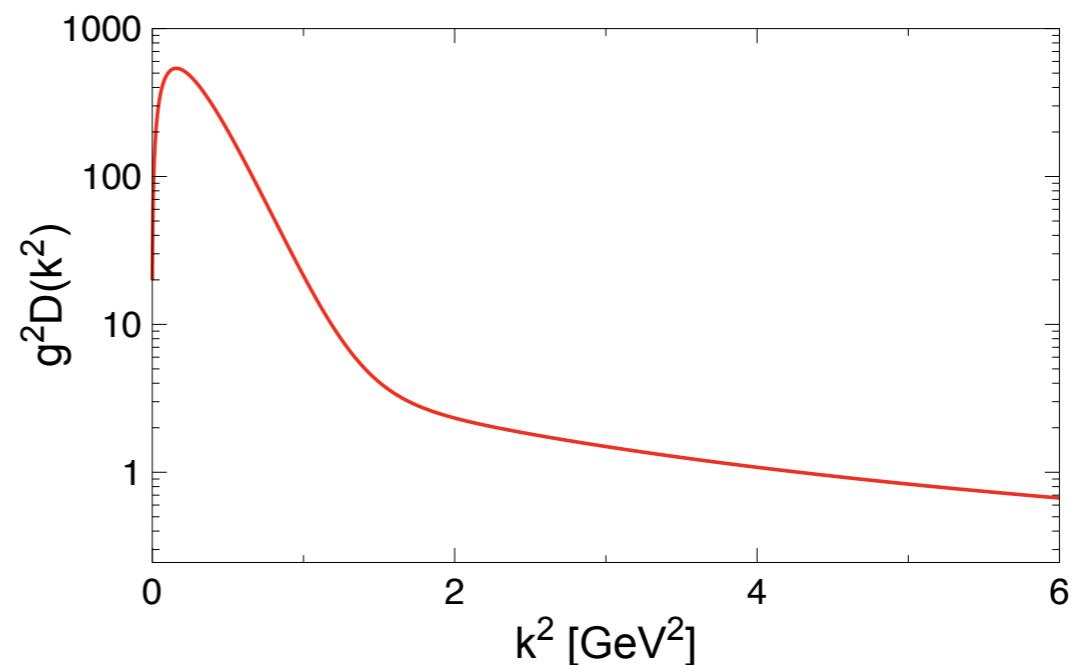
$$\alpha_{\text{pQCD}}(k^2) = \frac{2\pi\gamma_m(1 - e^{-k^2/4m_t^2})}{\ln[\tau + (1 + k^2/\Lambda_{\text{QCD}}^2)^2]}$$

Linear

$$\mathcal{G}_{\text{IR}}(k^2) = 4\pi^2 \frac{\chi^2}{k^4 + \Delta}$$

Exponential

$$\mathcal{G}_{\text{IR}}(k^2) = \frac{4\pi^2}{\omega^6} D k^2 e^{-k^2/\omega^2}$$



Trial: Vertex model

► Dirac structure

$$\Gamma_\mu(k, Q) = \{\gamma_\mu, k_\mu, Q_\mu\} \times \{1, \not{k}, \not{Q}, \not{k}\not{Q}\}$$

► Simple model

$$\Gamma_\mu^{\text{Bare}}(k, Q) = \gamma_\mu$$

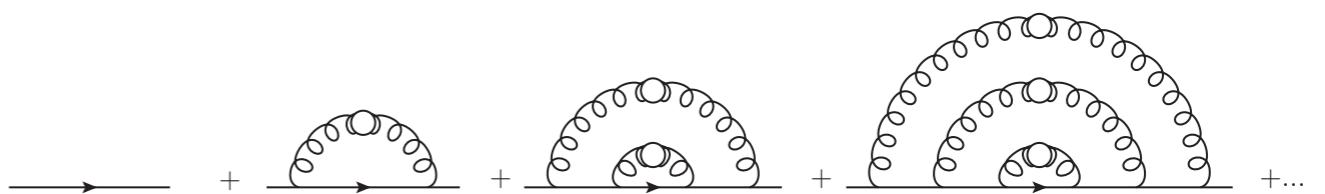
► Fredholm equation

$$A(p^2, \zeta^2) = Z_2 + \frac{4}{3} \int_q^\Lambda \mathcal{G}(k^2) \frac{A(q^2, \zeta^2)}{q^2 A^2(q^2, \zeta^2) + B^2(q^2, \zeta^2)} \left[\frac{p \cdot q}{p^2} + \frac{2(k \cdot p)(k \cdot q)}{p^2 k^2} \right]$$

$$B(p^2, \zeta^2) = Z_4 m_0 + \frac{4}{3} \int_q^\Lambda \mathcal{G}(k^2) \frac{3B(q^2, \zeta^2)}{q^2 A^2(q^2, \zeta^2) + B^2(q^2, \zeta^2)}$$

► Rainbow approximation

$$g^2 D_{\mu\nu}(k) \Gamma_\nu(k, Q) = \left(\delta_{\mu\nu} - \frac{k_\mu k_\nu}{k^2} \right) \not{\gamma}_\nu \mathcal{G}(k^2)$$



Trial: Quark solution

- Using MOM renormalization scheme:

$$S^{-1}(\zeta) = i\cancel{\not{q}} + m_f$$

$$\Sigma(p) = i\cancel{p}[A'(p^2) - 1] + B'(p^2)$$

- Subtraction scheme (I):

$$S^{-1}(p) = Z_2(i\cancel{p} + Z_m m_0) + \Sigma(p)$$

$$S^{-1}(\zeta) = i\cancel{\not{q}}[Z_2 + A'(\zeta^2) - 1] + Z_2 Z_m m_0 + B'(\zeta^2)$$

$$Z_2 = 2 - A'(\zeta^2), \quad Z_2 Z_m m_0 = m_f - B'(\zeta^2)$$

- Subtraction scheme (II):

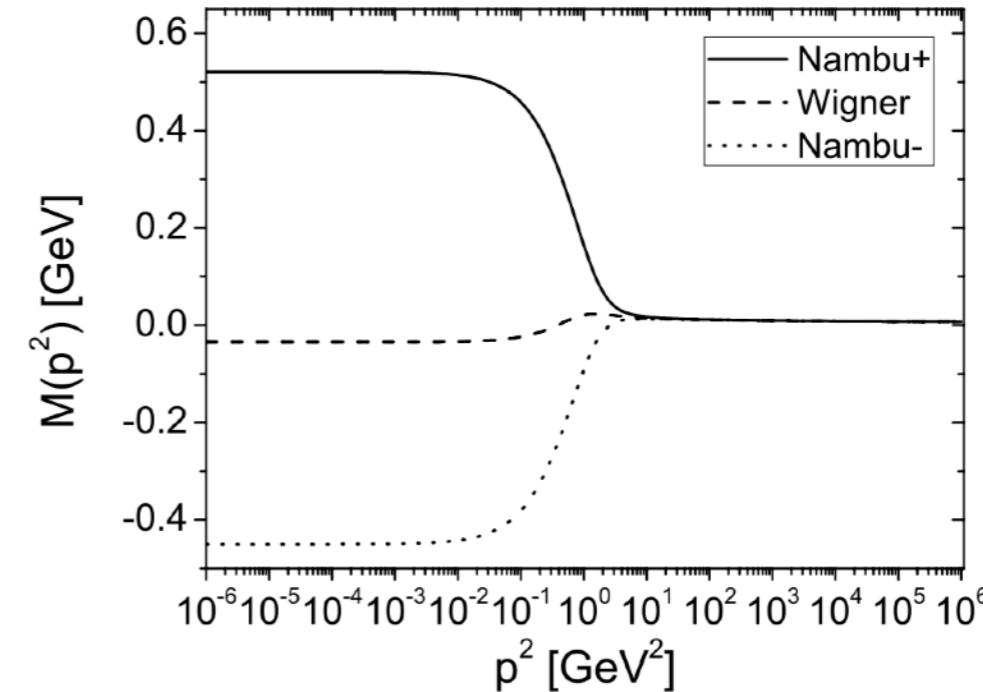
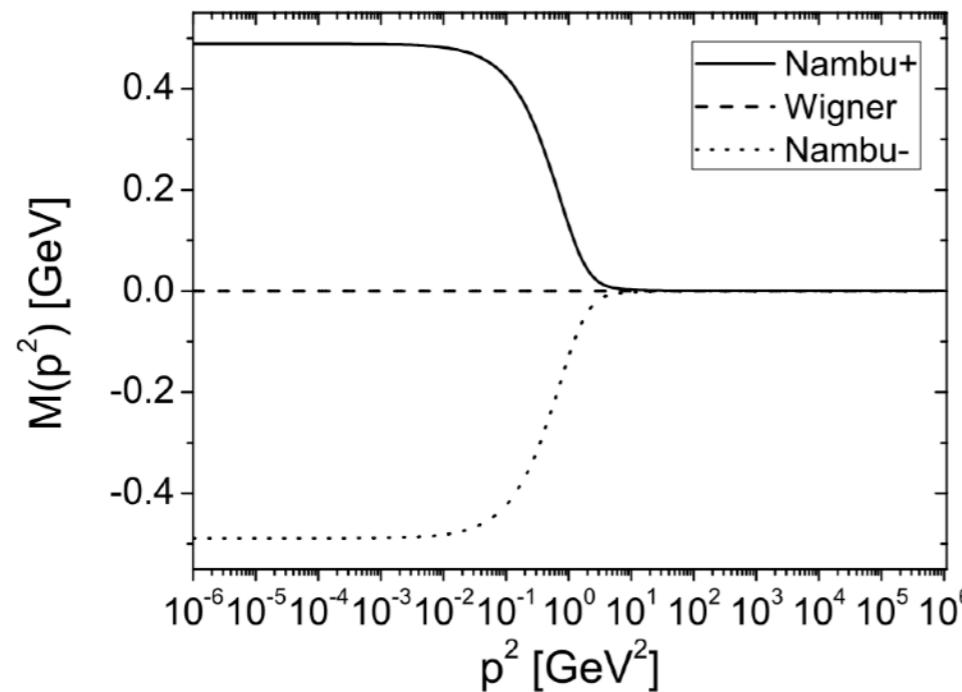
$$S^{-1}(p) = Z_2(i\cancel{p} + Z_m m_0) + Z_2 \Sigma(p)$$

$$S^{-1}(\zeta) = i\cancel{\not{q}}[Z_2 A'(\zeta^2)] + Z_2 Z_m m_0 + Z_2 B'(\zeta^2)$$

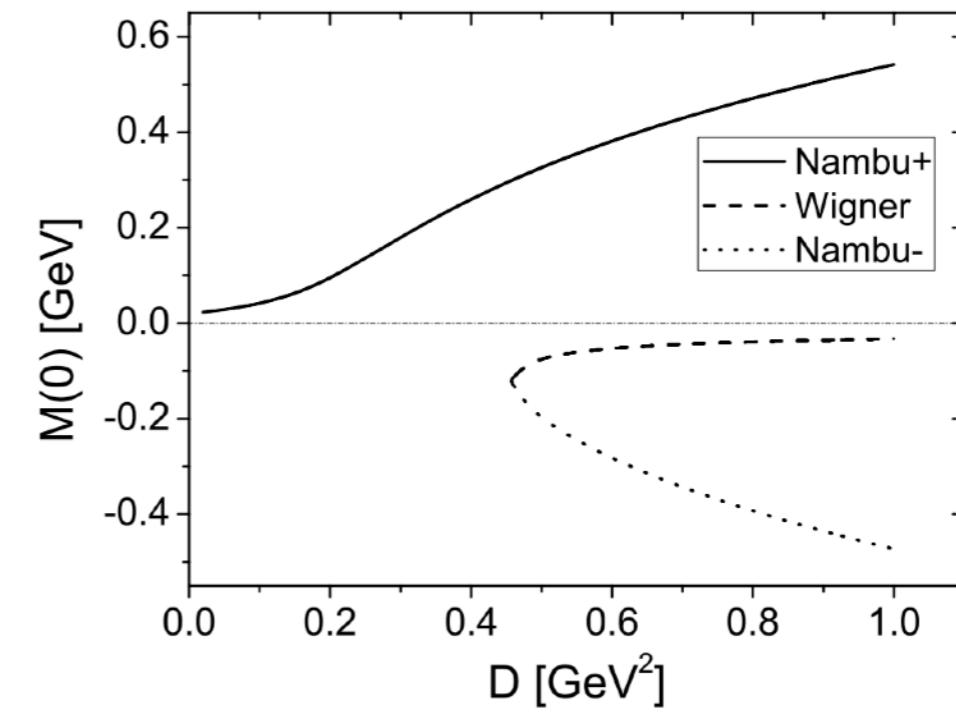
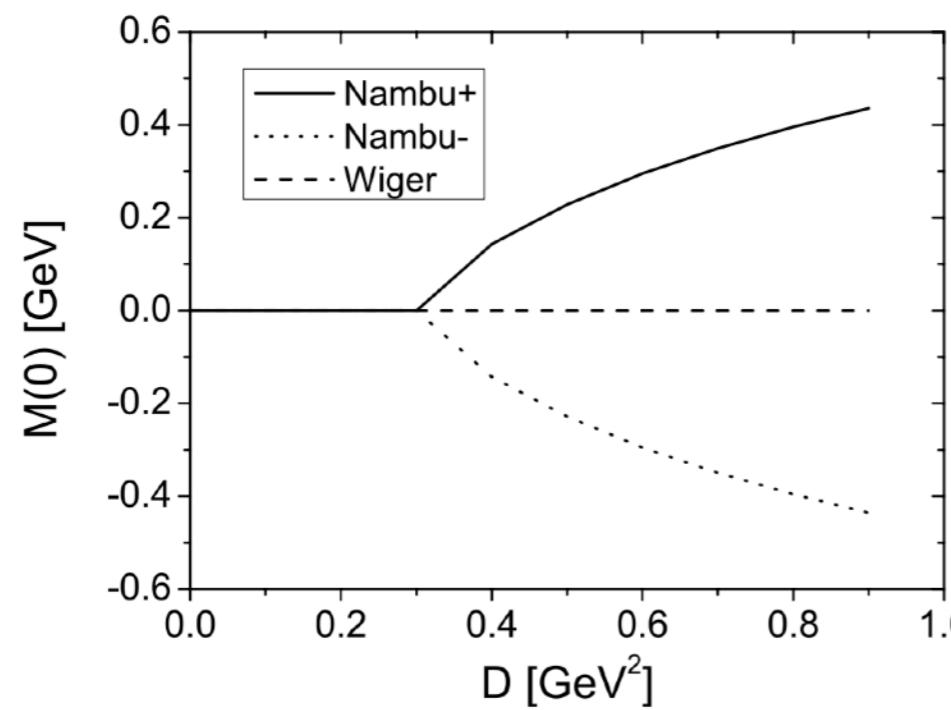
$$Z_2 = \frac{1}{A'(\zeta^2)}, \quad Z_2 Z_m m_0 = m_f - \frac{B'(\zeta^2)}{A'(\zeta^2)}$$

Trial: Quark solution

► Mass function



► Critical coupling



Trial: Quark solution

► Chiral condensate

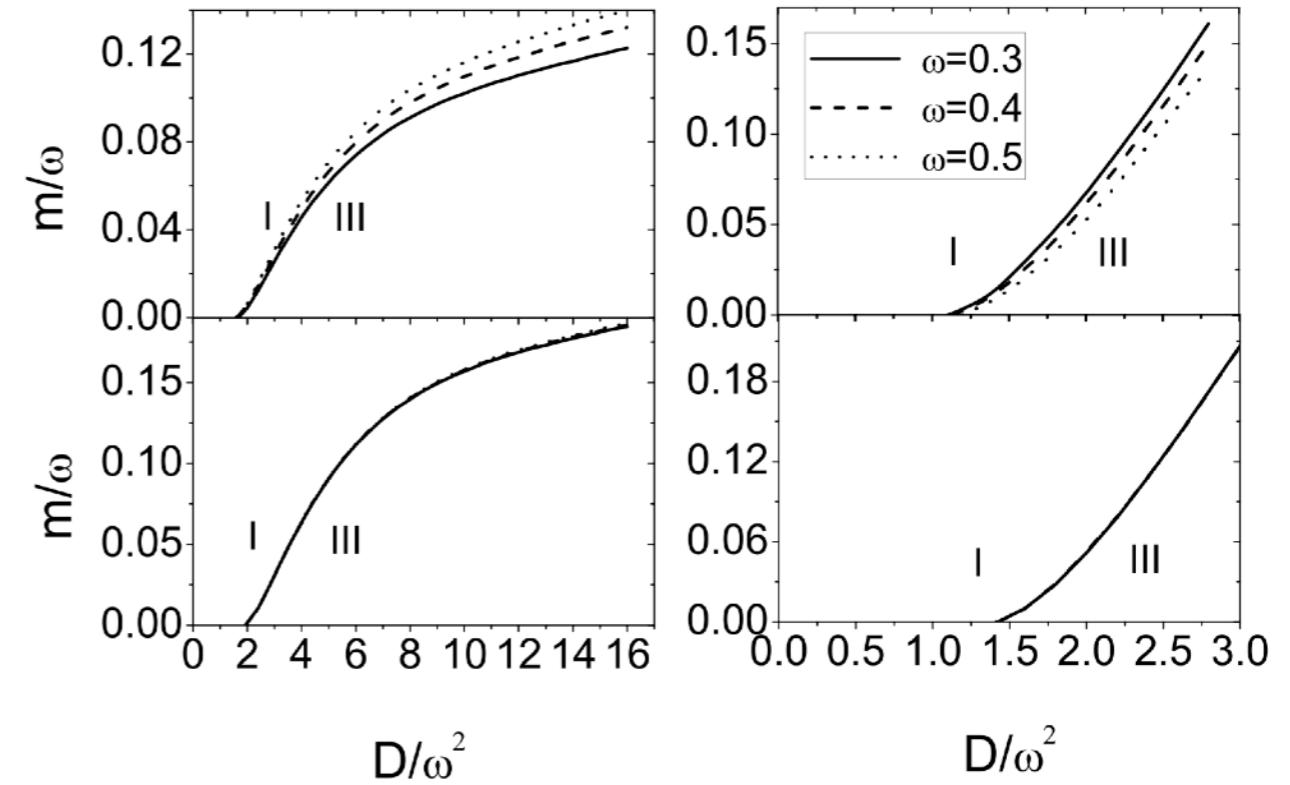
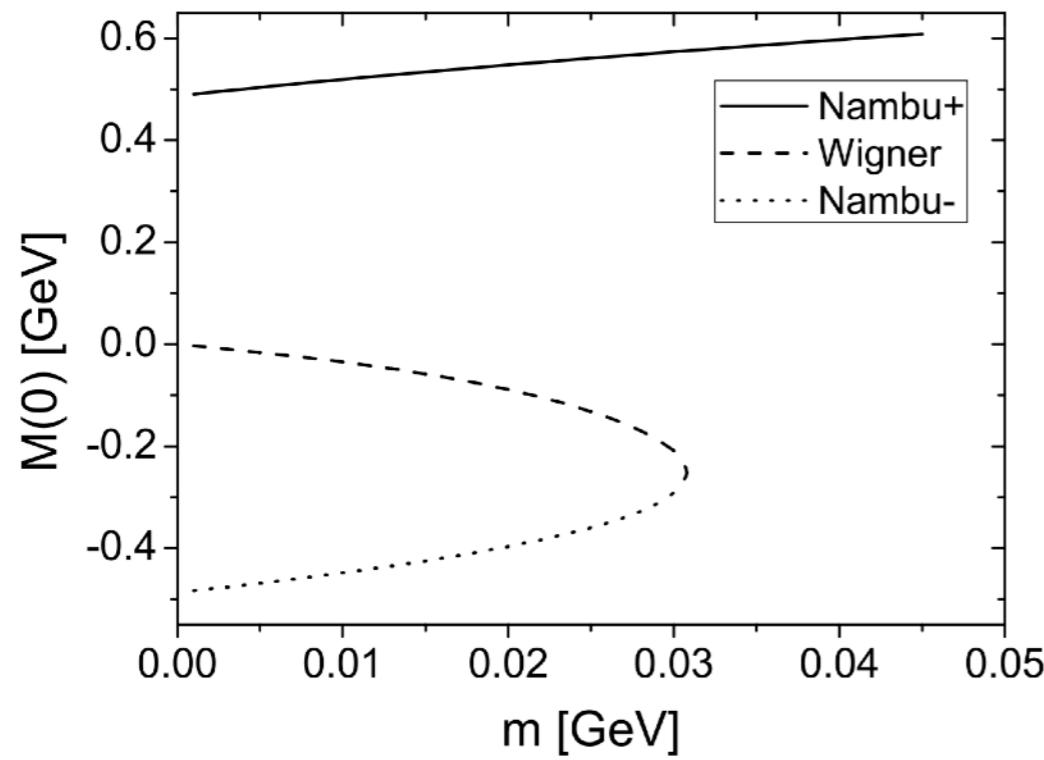
$$-\langle \bar{q}q \rangle^0 = \text{tr}\langle \bar{q}q \rangle = \lim_{\Lambda \rightarrow \infty} \int_q \text{tr}S(q) = N_c \lim_{\Lambda \rightarrow \infty} \int_q \frac{Z(q^2)M(q^2)}{q^2 + M^2(q^2)}$$

$$\langle \Omega | \bar{q}q | \Omega \rangle = \langle \Omega | \bar{q}_L q_R | \Omega \rangle + \langle \Omega | \bar{q}_R q_L | \Omega \rangle$$

$$q(x) \rightarrow q'(x) = e^{i\gamma_5 \theta} q(x), \quad \bar{q}(x) \rightarrow \bar{q}'(x) = \bar{q}(x) e^{i\gamma_5 \theta}$$

$$S'(p) = \int d^4x e^{ip \cdot (x-y)} \langle \Omega' | q(x) \bar{q}(y) | \Omega' \rangle = -i \not{p} \sigma_v(p^2) + e^{2i\gamma_5 \theta} \sigma_s(p^2)$$

► Critical current mass



Trial: Scattering kernel

► Dirac structure

$$K_{\alpha\beta}^{\rho\sigma}(q, k, P) = K_{\rho\sigma}^L(q, k, P) \otimes K_{\alpha\beta}^R(q, k, P)$$

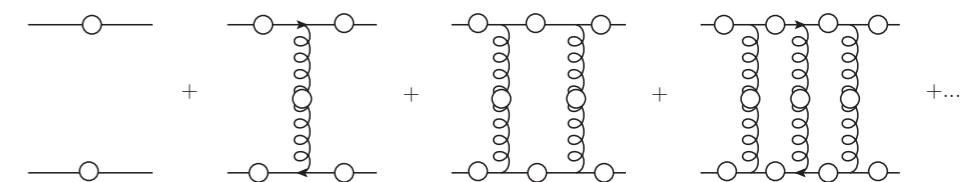
256 → 64

► Simple model

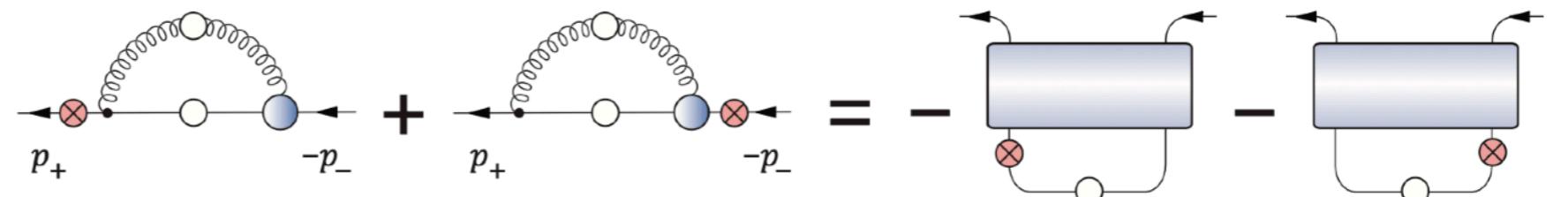
$$K_{\alpha\beta, BD}^{\rho\sigma, AC}(q, k, P) = -g^2 D_{\mu\nu}(k - q) \left[\frac{\lambda^a}{2} \right]_{AC} \left[\frac{\lambda^a}{2} \right]_{BD} [\gamma_\mu]_{\alpha\rho} [\gamma_\nu]_{\beta\sigma}$$

► Ladder approximation

$$\Gamma_X^{fg}(k, P) = - \int_q^\Lambda g^2 D_{\mu\nu}(k - q) \frac{\lambda^a}{2} \gamma_\mu S_f(q_+) \Gamma_X^{fg}(q, P) S_g(q_-) \frac{\lambda^a}{2} \gamma_\nu$$



Ward identities

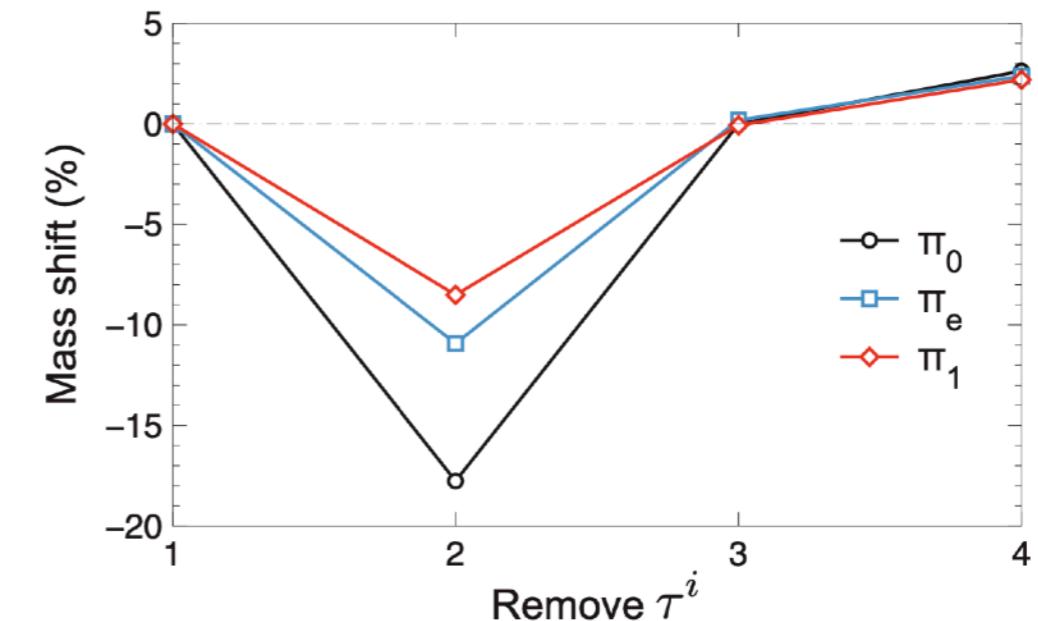
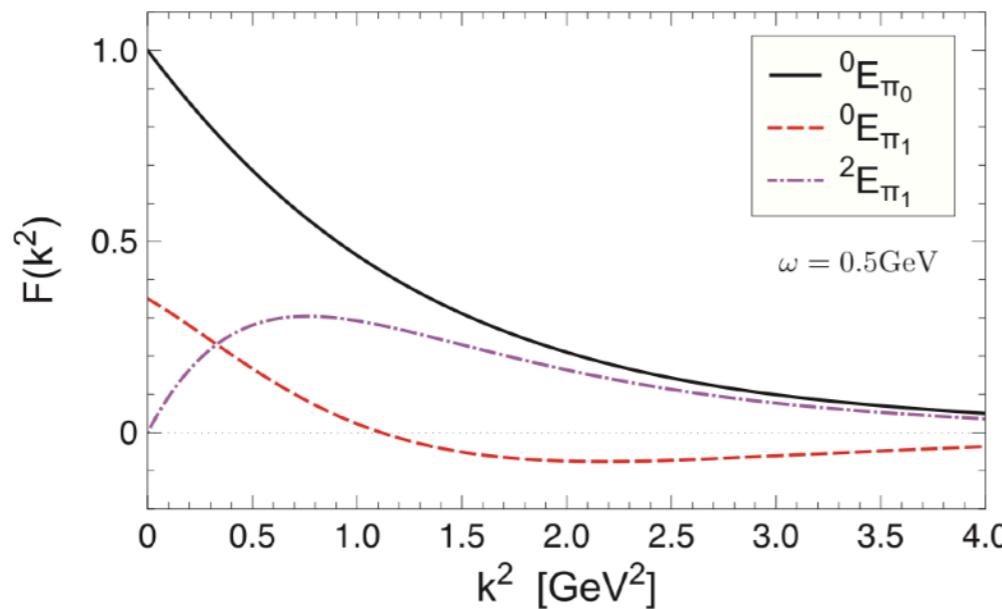


Trial: Meson solution

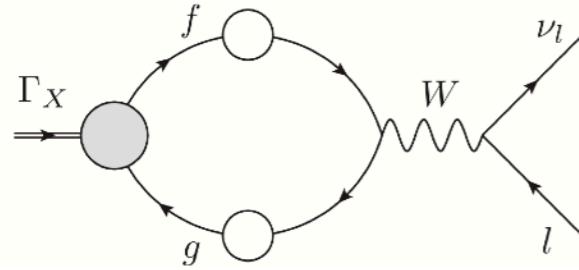
BSE amplitude

$$\gamma_5 \mathcal{E}(k, P) + \gamma_5 \not{k} \mathcal{F}(k, P) + \gamma_5 \not{k} G(k, P) + \gamma_5 [\not{k}, \not{P}] H(k, P)$$

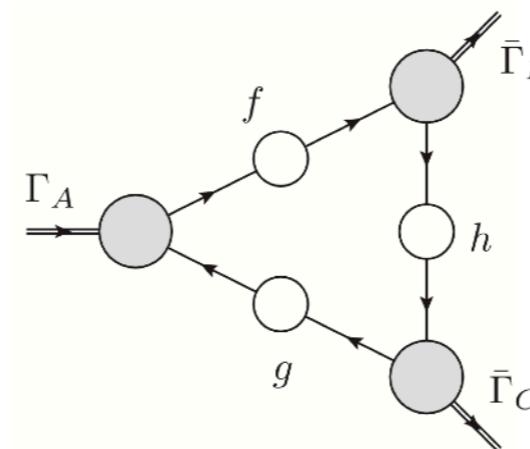
$$P \neq (-1)^{L+1}$$



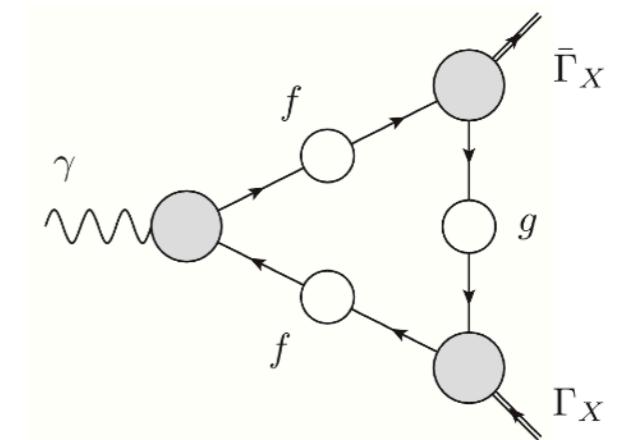
Meson observable



leptonic decay



strong decay



EM structure

Trial: Meson solution

► Comparison with experiments

Summary of light meson results

$m_{u=d} = 5.5 \text{ MeV}$, $m_s = 125 \text{ MeV}$ at $\mu = 1 \text{ GeV}$

Pseudoscalar (PM, Roberts, PRC56, 3369)

	expt.	calc.
$-\langle \bar{q}q \rangle_\mu^0$	$(0.236 \text{ GeV})^3$	$(0.241^\dagger)^3$
m_π	0.1385 GeV	0.138^\dagger
f_π	0.0924 GeV	0.093^\dagger
m_K	0.496 GeV	0.497^\dagger
f_K	0.113 GeV	0.109

Charge radii (PM, Tandy, PRC62, 055204)

r_π^2	0.44 fm^2	0.45
$r_{K^+}^2$	0.34 fm^2	0.38
$r_{K^0}^2$	-0.054 fm^2	-0.086

$\gamma\pi\gamma$ transition (PM, Tandy, PRC65, 045211)

$g_{\pi\gamma\gamma}$	0.50	0.50
$r_{\pi\gamma\gamma}^2$	0.42 fm^2	0.41

Weak K_{l3} decay (PM, Ji, PRD64, 014032)

$\lambda_+(e3)$	0.028	0.027
$\Gamma(K_{e3})$	$7.6 \cdot 10^6 \text{ s}^{-1}$	7.38
$\Gamma(K_{\mu 3})$	$5.2 \cdot 10^6 \text{ s}^{-1}$	4.90

Vector mesons

(PM, Tandy, PRC60, 055214)

$m_{\rho/\omega}$	0.770 GeV	0.742
$f_{\rho/\omega}$	0.216 GeV	0.207
m_{K^*}	0.892 GeV	0.936
f_{K^*}	0.225 GeV	0.241
m_ϕ	1.020 GeV	1.072
f_ϕ	0.236 GeV	0.259

Strong decay (Jarecke, PM, Tandy, PRC67, 035202)

$g_{\rho\pi\pi}$	6.02	5.4
$g_{\phi KK}$	4.64	4.3
$g_{K^* K\pi}$	4.60	4.1

Radiative decay

(PM, nucl-th/0112022)

$g_{\rho\pi\gamma}/m_\rho$	0.74	0.69
$g_{\omega\pi\gamma}/m_\omega$	2.31	2.07
$(g_{K^* K\gamma}/m_K)^+$	0.83	0.99
$(g_{K^* K\gamma}/m_K)^0$	1.28	1.19

Scattering length

(PM, Cotanch, PRD66, 116010)

a_0^0	0.220	0.170
a_0^2	0.044	0.045
a_1^1	0.038	0.036

Tandy @ Beijing Lectures 2010

◆ **Frontier**: theoretical and experimental topics, non-perturbative approaches

◆ **Framework**: QFT equation of motion, QCD's DSE, advantages and disadvantages

◆ **Formalism**: one-body, two-body, and three-body equations, basic ingredients

◆ **Trial**: simple approximations for basic ingredients, quark solutions, meson solutions