第五章 傅里叶积分变换

5.1 傅里叶积分变换定义

5.1.1 知识回顾

傅里叶级数展开(周期函数) 傅里叶级数展开 $^{\text{\tiny MR}\,l \to \infty}$ (非周期函数) 傅里叶积分表达

往复的周期 把1往两边拉伸

以
$$2l$$
 为周期的函数 $f(x)$ 在区间 $[-l,l]$ 上可表达: $f(x) = \frac{a_0}{2} + \sum_{n=1}^{\infty} \left(a_n \cos \frac{n\pi x}{l} + b_n \sin \frac{n\pi x}{l} \right)$ $\left(a_n = \frac{1}{l} \int_{-l}^{l} f(t) \cos \frac{n\pi t}{l} dt, n = 0,1,2,\cdots \right)$

$$\begin{cases} a_n = \frac{1}{l} \int_{-l}^{l} f(t) \cos \frac{n\pi t}{l} dt, n = 0, 1, 2, \cdots \\ b_n = \frac{1}{l} \int_{-l}^{l} f(t) \sin \frac{n\pi t}{l} dt, n = 1, 2, \cdots \end{cases}$$
 考虑 a_n, b_n 的表达式,得到:

$$f(x) = \frac{1}{2l} \int_{-l}^{l} f(t) dt + \sum_{n=1}^{\infty} \frac{1}{l} \int_{-l}^{l} f(t) \left(\cos \frac{n\pi t}{l} \cos \frac{n\pi x}{l} + \sin \frac{n\pi t}{l} \sin \frac{n\pi x}{l} \right) dt$$
$$= \frac{1}{2l} \int_{-l}^{l} f(t) dt + \sum_{n=1}^{\infty} \frac{1}{l} \int_{-l}^{l} f(t) \cos \left[\frac{n\pi}{l} (x - t) \right] dt \qquad \stackrel{\text{height}}{=} 2\pi \alpha_n = \frac{n\pi}{l}, \Delta \alpha_n = \frac{\pi}{l}$$

$$f(x) = \frac{1}{2l} \int_{-l}^{l} f(t) dt + \frac{1}{\pi} \sum_{n=1}^{\infty} \left\{ \int_{-l}^{l} f(t) \cos \left[\alpha_n(x-t) \right] dt \right\} \Delta \alpha_n$$

假定
$$f(x)$$
 在 $(-\infty, +\infty)$ 上**绝对可积**, $\int_{-\infty}^{+\infty} |f(t)| dt < +\infty$, 则 $\frac{1}{2l} \int_{-l}^{l} f(t) dt = 0$

$$f(x) = \frac{1}{\pi} \int_0^{+\infty} \{ \int_{-\infty}^{+\infty} f(t) \cos[\alpha(x-t)] dt \} d\alpha \quad \text{为} f(x)$$
的傅里叶积分公式。

对积分表达进一步处理,利用欧拉公式: $e^{i\beta} = \cos \beta + i\sin \beta$

$$f(x) = \frac{1}{2\pi} \int_0^{+\infty} \int_{-\infty}^{+\infty} f(t) e^{i\alpha(x-t)} dt d\alpha + \frac{1}{2\pi} \int_0^{+\infty} \int_{-\infty}^{+\infty} f(t) e^{-i\alpha(x-t)} dt d\alpha$$

$$f(x) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} f(t) e^{i\alpha(x-t)} dt d\alpha$$
 $f(x)$ 的傅里叶积分公式

$$f(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} e^{i\alpha x} \left\{ \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} f(t) e^{-i\alpha t} dt \right\} d\alpha \qquad$$
 记为 $F(\alpha)$,称为 $f(x)$ 的傅里叶变换 $\mathcal{F}[f(x)]$

5.1.2 傅里叶积分变换的定义

基本条件 设 f(x) 在 $(-\infty, +\infty)$ 上的任一有限区间上满足 Dirichlet 条件,在 $(-\infty, +\infty)$ 上绝对可积

傅里叶逆变换 $f(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} F(\alpha) e^{i\alpha x} d\alpha$ 为 $F(\alpha)$ 的傅里叶**逆变换**,记为 $f(x) = \mathcal{F}^{-1}[F(\alpha)]$

关系 "信号"
$$f(x)$$
 $\xrightarrow{\mathscr{F}}$ $F(\alpha)$ 频谱分布函数 "反演" \mathscr{F}^{-1}

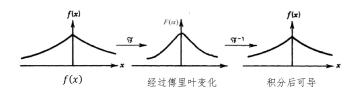
注意 傅里叶积分变换来自于非周期傅里叶函数的积分表达,**两者成对出现**

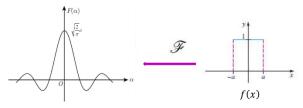
示例: 求函数的傅里叶变换,其中 a 为常数: $(1)f(x) = e^{-a|x|};$ $(2)f(x) = \begin{cases} 1, |x| \leq a, \\ 0, |x| > a. \end{cases}$

①
$$F(\alpha) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} e^{-a|x|} e^{-i\alpha x} dx = \frac{1}{\sqrt{2\pi}} \left\{ \int_{-\infty}^{0} e^{x(a-i\alpha)} dx + \int_{0}^{+\infty} e^{x(-a-i\alpha)} dx \right\} = \sqrt{\frac{2}{\pi}} \frac{a}{a^2 + \alpha^2}$$
 上式中: $\frac{1}{\sqrt{2\pi}} \frac{1}{a-i\alpha} \int_{-\infty}^{0} de^{x(a-i\alpha)}$ 可以继续积分

② 当
$$\alpha \neq 0$$
 时 $F(\alpha) = \frac{1}{\sqrt{2\pi}} \int_{-a}^{a} e^{-i\alpha x} dx$ 可以写为正余弦形式 = $\sqrt{\frac{2}{\pi}} \frac{\sin{(a\alpha)}}{\alpha}$

当
$$\alpha=0$$
 时 $\lim_{\alpha\to 0}F(\alpha)=\frac{2}{\sqrt{\pi}}\lim_{\alpha\to 0}\frac{\sin(a\alpha)}{\alpha}=\sqrt{\frac{2}{\pi}}\alpha=F(0)=\frac{1}{\sqrt{2\pi}}\int_{-a}^{a}dx$ 验证极限值等于函数值,连续





5.2 傅里叶积分变换的性质

5.2.1 线性性质

内容 傅里叶变换和逆变换都是线性变换,即 (积分满足线性性质)

$$\begin{split} \mathcal{F}[c_1f_1(x) + c_2f_2(x)] &= c_1\mathcal{F}[f_1(x)] + c_2\mathcal{F}[f_2(x)] = c_1F_1(\alpha) + c_2F_2(\alpha) \\ \mathcal{F}^{-1}[c_1F_1(\alpha) + c_2F_2(\alpha)] &= c_1\mathcal{F}^{-1}[F_1(\alpha)] + c_2\mathcal{F}^{-1}[F_2(\alpha)] = c_1f_1(x) + c_2f_2(x) \end{split}$$

5.2.2 位移性质

内容 设 x_0 为任意实常数,则 $\mathbf{\mathcal{F}}[f(x \pm x_0)] = e^{\pm i\alpha x_0}\mathbf{\mathcal{F}}[f(x)]$

设:
$$f(x) \xrightarrow{\mathcal{F}} F(\alpha)$$
, 则 $f(x-c) \xrightarrow{\mathcal{F}} e^{-i\alpha c} F(\alpha)$

证明
$$\mathcal{F}[f(x-c)] = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} f(x-c) e^{-i\alpha x} dx \xrightarrow{\Leftrightarrow \xi = x-c} \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} f(\xi) e^{-i\alpha \xi} e^{-i\alpha c} d\xi$$

5.2.3 相似性质 (伸缩性质)

内容 设 c 为任意非零实常数,则 $\mathcal{F}[f(cx)] = \frac{1}{|c|} F(\frac{\alpha}{c})$

证明
$$\mathcal{F}[f(cx)] = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} f(cx) e^{-i\alpha x} dx \xrightarrow{\frac{2\xi = cx}{c} (c>0)} \frac{1}{c} \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} f(\xi) e^{-i\frac{\alpha}{c}\xi} d\xi = \frac{1}{c} F\left(\frac{\alpha}{c}\right)$$

$$\xrightarrow{\frac{2\xi = cx}{c} (c<0)} \frac{1}{c} \frac{1}{\sqrt{2\pi}} \int_{+\infty}^{-\infty} f(\xi) e^{-i\frac{\alpha}{c}\xi} d\xi = \frac{1}{c} F\left(\frac{\alpha}{c}\right) = -\frac{1}{c} \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} f(\xi) e^{-i\frac{\alpha}{c}\xi} d\xi$$
 合并后为绝对值

5.2.4 微分性质

内容 若当 $|x| \to \infty$ 时, $f(x) \to 0$ (函数在无穷远处为零) $f^{(k)}(x) \to 0$ $(k = 1, 2, \dots, n - 1)$,则 $\mathcal{F}[f'(x)] = i\alpha \mathcal{F}[f(x)], \quad \mathcal{F}[f''(x)] = (i\alpha)^2 \mathcal{F}[f(x)], \dots, \mathcal{F}[f^{(n)}(x)] = (i\alpha)^n \mathcal{F}[f(x)]$

变化前为求导,变化后为幂的形式

证明
$$\mathcal{F}[f'(x)] = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} f'(x) \, e^{-i\alpha x} dx \xrightarrow{\text{分部积分}} \frac{1}{\sqrt{2\pi}} f(x) e^{-i\alpha x} \Big|_{-\infty \text{此项为零}}^{\infty} - (-i\alpha) \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} f(x) e^{-i\alpha x} dx = i\alpha \mathcal{F}[f(x)]$$
在像原函数空间一阶求导,等价于在该处乘上一个ia

5.2.5 积分性质

$$\mathcal{F}\left[\int_{x_0}^x f(\xi) d\xi\right] = \frac{1}{i\alpha} \mathcal{F}[f(x)]$$

证明
$$\left(\int_{x_0}^x f(\xi) d\xi \right)' = f(x) \xrightarrow{\underline{\text{mżd}}\underline{\#}\underline{\mathbb{P}}\underline{H}} \mathcal{F} \left[\left(\int_{x_0}^x f(\xi) d\xi \right)' \right] = F(\alpha) \qquad \mathcal{F} \left[\left(\int_{x_0}^x f(\xi) d\xi \right)' \right] = (i\alpha) \mathcal{F} \left[\int_{x_0}^x f(\xi) d\xi \right]'$$

5.2.6 乘多项式性质

$$\mathcal{F}[xf(x)] = i\frac{\mathrm{d}F(\alpha)}{\mathrm{d}\alpha}, \dots, \mathcal{F}[x^n f(x)] = i^n \frac{\mathrm{d}^n F(\alpha)}{\mathrm{d}\alpha^n}$$

证明
$$\mathcal{F}[xf(x)] = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} x f(x) e^{-i\alpha x} dx = -\frac{1}{i} \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} \frac{d}{d\alpha} (f(x)e^{-i\alpha x}) dx = i \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} \frac{d}{d\alpha} (f(x)e^{-i\alpha x}) dx$$
$$= i \frac{d}{d\alpha} \left(\frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} (f(x)e^{-i\alpha x}) dx \right) = i \frac{dF(\alpha)}{d\alpha}$$

5.2.7 卷积定理

卷积运算 定义 f(x) 和 g(x) 的卷积运算为 $f(x)*g(x)=\frac{1}{\sqrt{2\pi}}\int_{-\infty}^{+\infty}f(x-t)g(t)dt=\frac{1}{\sqrt{2\pi}}\int_{-\infty}^{+\infty}f(t)g(x-t)dt$

性质 $\mathcal{F}[(f*g)(x)] = F(\alpha)G(\alpha) \qquad \qquad \mathcal{F}^{-1}[F(\alpha)G(\alpha)] = (f*g)(x) \quad \text{后者常用 (频谱域→物理空间)}$

示例 1. 计算下列二个函数的卷积: $\begin{cases} f(x) = x \\ g(x) = e^{-x^2} \end{cases}$

$$(f*g)(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} (x-\xi)e^{-\xi^2}d\xi = \frac{x}{\sqrt{2}}$$
 (已知结论 $\int_{-\infty}^{\infty} e^{-\xi^2}d\xi = \sqrt{\pi}$)

2. 计算卷积。定义 $g_n(x)$: $g_n(x) = \begin{cases} n\sqrt{\frac{\pi}{2}}, & |x| \leq \frac{1}{n} \\ 0, &$ 其它.

一个新的函数 $F(x) = \int_{-\infty}^{x} f(t)dt$.试计算 $f(x) * g_n(x)$

3. 计算函数 $f(x) = e^{-\frac{a}{2}x^2}$ 的傅里叶变换

根据定义有 $F(\alpha) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{-\frac{ax^2}{2}} e^{-i\alpha x} dx$ (直接计算复杂) 进而得到 $F(0) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{-\frac{ax^2}{2}} dx = \frac{1}{\sqrt{a}}$ 又有: $f(x) = e^{-\frac{a}{2}x^2} \rightarrow f'(x) + axf(x) = 0$ 两边傅里叶变换 $i\alpha F(\alpha) + ai\frac{dF(\alpha)}{d\alpha} = 0$ 微分方程 $F(\alpha) = Ce^{-\frac{a^2}{2a}}$

可得: $F(\alpha) = \frac{1}{\sqrt{a}}e^{-\frac{\alpha^2}{2a}}$ 即有: $e^{-\frac{a}{2}\chi^2} \stackrel{\mathcal{F}}{\to} \frac{1}{\sqrt{a}}e^{-\frac{\alpha^2}{2a}}$ 可用公式 $(a \to 2a)$: $e^{-a\chi^2} \stackrel{\mathcal{F}}{\Longrightarrow} \frac{1}{\sqrt{2a}}e^{-\frac{\alpha^2}{4a}}$

例如,计算 $e^{-\tilde{a}^2\alpha^2t} \overset{\mathcal{F}^{-1}}{\Longrightarrow}$? 有: $\sqrt{2a}e^{-ax^2} \overset{\mathcal{F}}{\Longrightarrow} e^{-\frac{\alpha^2}{4a}}$ 则: $\begin{cases} \frac{1}{4a} = \tilde{a}^2t \\ a = \frac{1}{4\tilde{a}^2t} \end{cases}$ 则逆变换 $\frac{1}{\tilde{a}}\frac{1}{\sqrt{2t}}e^{-\frac{x^2}{4\tilde{a}^2t}}$

5.3 傅里叶积分变换在求解偏微分方程初值问题中的应用

5.3.1 偏导数的傅里叶变换

条件 对于多变量函数 u(x,t)的偏导数 $u_x(x,t)$, $u_{xx}(x,t)$, $u_t(x,t)$, $u_{tt}(x,t)$ 针对x做变换,固定t

变换
$$\mathcal{F}[u_x] = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} u_x(x,t) e^{-i\alpha x} dx = i\alpha \mathcal{F}[u] \qquad \qquad \mathcal{F}[u_{xx}] = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} u_{xx}(x,t) e^{-i\alpha x} dx = -\alpha^2 \mathcal{F}[u]$$

$$\mathcal{F}[u_t] = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} u_t(x,t) e^{-i\alpha x} dx = \frac{\partial}{\partial t} \mathcal{F}[u] \qquad \qquad \mathcal{F}[u_{tt}] = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} u_{tt}(x,t) e^{-i\alpha x} dx = \frac{\partial^2}{\partial t^2} \mathcal{F}[u]$$

5.3.2 求解热传导方程的初值问题

问题
$$\begin{cases} u_t = a^2 u_{xx} + f(x,t) & -\infty < x < +\infty, t > 0 \\ u(x,0) = \phi(x) & -\infty < x < +\infty \end{cases}$$

$$\mathcal{F}[u(x,t)] = \mathbf{U}(\alpha,t)$$

问题转化 対 $u(x,t),f(x,t)$ 和 $\phi(x)$ 关于 x 进行傅里叶变换 ⇒ $\mathcal{F}[f(x,t)] = \mathbf{F}(\alpha,t)$
 $\mathcal{F}[\phi(x)] = \mathbf{\Phi}(\alpha)$

原问题转为:
$$\begin{cases} U_t(\alpha,t) = a^2(i\alpha)^2 U(\alpha,t) + F(\alpha,t) \\ U(\alpha,0) = \Phi(\alpha) \end{cases} \quad U_t + a^2 \alpha^2 U = F(\alpha,t) \overset{\text{等价}}{\Longrightarrow} y'(x) + ay(x) = f(x)$$

常数变易法:
$$\xrightarrow{\text{两边乘}e^{ax}} (y(x)e^{ax})' = f(x)e^{ax} \xrightarrow{\text{积} f(x)} y(x) = y(0)e^{-ax} + \int_0^x f(\xi)e^{-a(x-\xi)}d\xi$$

代入原方程相关内容:
$$U(\alpha,t) = \Phi(\alpha)e^{-\alpha^2a^2t} + \int_0^t F(\alpha,\tau)e^{-\alpha^2a^2(t-\tau)}d\tau$$

$$= \phi(x) * \mathcal{F}^{-1} \left[e^{-\alpha^2 a^2 t} \right] + \int_0^t f(x,\tau) * \mathcal{F}^{-1} \left[e^{-\alpha^2 a^2 (t-\tau)} \right] d\tau \xrightarrow{\mathcal{F}^{-1} \left[e^{-a^2 \alpha^2 t} \right] = \frac{1}{a\sqrt{2t}} e^{\frac{x^2}{4a^2 t}}} \text{ $\frac{x}{2}$}$$

$$= \frac{1}{2a\sqrt{\pi t}} \int_{-\infty}^{+\infty} e^{\frac{(x-\xi)^2}{4a^2t}} \phi(\xi) d\xi + \int_{0}^{t} \int_{-\infty}^{+\infty} \frac{f(\xi,\tau)}{2a\sqrt{\pi(t-\tau)}} e^{\frac{(x-\xi)^2}{4a^2(t-\tau)}} d\xi d\tau \xrightarrow{\text{\bf a} \text{ in } \text{\bf a} \text{ in }$$

$$u(x,t) = \int_{-\infty}^{+\infty} G(x-\xi,t)\phi(\xi)d\xi + \int_{0}^{t} \int_{-\infty}^{+\infty} G(x-\xi,t-\tau)f(\xi,\tau)d\xi d\tau$$

$$u(x,t) = \phi(x) * G(x,t) + \int_0^t f(x,\tau) * G(x,t-\tau) d\tau$$

注意 ① G(x,t)热核函数或高斯核函数: t=0时刻,x=0处的单位强度的热源在其它位置时刻的响应函数

② 对于热传导初值问题 $\begin{cases} u_t = a^2 u_{xx}, \\ u(x,0) = \varphi(x) \end{cases}$ 其解的表达式为:

$$u(x,t) = \int_{-\infty}^{+\infty} G(x-\xi,t)\phi(\xi)d\xi = \phi(x) * G(x,t)$$

③ 对于热传导初值问题的解,在下列极限意义下满足初始条件: $u(x,0) = \varphi(x)$

