

¹³⁴⁸ 9.21 UNDERSTANDING THE ¹³⁴⁹ “RENORMALIZATION GROUP” ¹³⁵⁰

One might wonder if the "Renormalization Group" (RG) is, in fact, a mathematical group. The name is a famous misnomer, and understanding why it's both wrong and right reveals the entire conceptual leap from historical QFT to the modern, Wilsonian framework.

¹³⁵⁴ THE CENTRAL QUESTION: A GROUP OR NOT?

In the modern, arguably more general sense, the Renormalization Group is **not a mathematical group**, but a **semigroup**.

The confusion arises because the name is a historical artifact from an older, more formal perspective. To see the difference, we must distinguish between two different ways of thinking about the RG:

- 1360 1. The modern (Wilsonian) picture of a *physical flow*.
- 1361 2. The historical (QFT) picture of a *re-parameterization*.

¹³⁶² THE MODERN ANSWER: A SEMIGROUP OF PHYSICAL ¹³⁶³ FLOWS

When we think of the RG in the modern sense (thanks to Ken Wilson), we are describing a physical process of "coarse-graining" or "integrating out" high-energy (short-distance) degrees of freedom. We are literally averaging over and discarding information.

A mathematical **group** must satisfy four axioms:

- 1369 1. **Closure:** (Transform A) + (Transform B) = (New Transform C). (*This holds.*)
- 1370 2. **Associativity:** (A+B)+C = A+(B+C). (*This holds.*)
- 1371 3. **Identity Element:** A transformation that does nothing. (*This holds.*)
- 1372 4. **Inverse Element:** For every transformation, there must be an inverse that perfectly undoes it. (***This fails.***)

The Wilsonian RG fails on **Invertibility**.

The act of integrating out high-energy modes is an **irreversible** process. You are averaging over details. Once you've coarse-grained a system, you have *lost information* about the precise, short-distance microscopic arrangement.

You cannot run the RG "backwards" from a low-energy effective theory to uniquely discover the one, true high-energy UV theory. In fact, the entire concept of **universality** rests on this fact: many different UV theories can "flow" to the *same* low-energy IR theory.

1384 This is a **many-to-one mapping**. Because the transformations are not
1385 invertible, the RG is not a group. It is a semigroup (or more specifically, a
1386 monoid, as it has an identity element). This irreversibility is what gives the
1387 RG "flow" its famous "arrow of scale."

1388 A DIFFERENT ASPECT: A LIE GROUP OF 1389 RE-PARAMETERIZATION

1390 So, why the name? The name is a historical artifact from the 1950s,
1391 coined by Ernst Stueckelberg and André Petermann (the "Renorm-gruppe")
1392 and further developed by Gell-Mann and Low.
1393 Their "group" was not about integrating out physical modes. It was about
1394 the **ambiguity** in the perturbative renormalization procedure. (This is how
1395 we introduced renormalization in the last part of Chapter II of our classes.)

- 1396 • In perturbative QFT, to get rid of infinities, we introduce an
1397 arbitrary, unphysical energy scale, μ , often called the
1398 "renormalization scale" or "subtraction point."
- 1399 • All our renormalized couplings (g), masses (m), and field
1400 normalizations (Z) become functions of this arbitrary scale: $g(\mu)$,
1401 $m(\mu)$, $Z(\mu)$.
- 1402 • Physical observables (like a scattering cross-section) *cannot* depend
1403 on this arbitrary choice of μ .
- 1404 • Therefore, a change in our choice of scale (from μ_1 to μ_2) must be
1405 perfectly compensated by a corresponding change in the couplings
1406 and fields to keep the physics invariant.

1407 This set of transformations *is* a group.

- 1408 • **Closure:** A transformation $\mu_1 \rightarrow \mu_2$ followed by $\mu_2 \rightarrow \mu_3$ is just the
1409 transformation $\mu_1 \rightarrow \mu_3$.
- 1410 • **Associativity:** This holds.
- 1411 • **Identity:** The transformation $\mu_1 \rightarrow \mu_1$.
- 1412 • **Inverse:** The transformation from $\mu_2 \rightarrow \mu_1$ is the *exact* inverse of
1413 $\mu_1 \rightarrow \mu_2$.

1414 This forms a continuous, invertible **Lie group**. The famous
1415 Callan-Symanzik and Gell-Mann-Low equations are the differential
1416 expressions (i.e., the Lie algebra) of this group. They describe how the
1417 couplings must "run" with scale μ to preserve the invariance of the
1418 physics.

1419 RECONCILIATION: IT IS UP TO THE PHYSICS CONTEXT

1420 One can rightly ask: "Even if I am given a single UV theory, like the
1421 Standard Model, its flow is described by a deterministic differential

equation (the Callan-Symanzik equation). Can't I just run this equation backwards and find the UV theory? Isn't that an inverse?"

That's correct. If you are given **one single, well-defined theory**, the flow is just a **re-parameterization** of that one theory. You can slide your "viewing window" μ up and down, and the transformations are invertible. In this limited sense, for a *single theory*, the flow is a group.

The Wilsonian revolution was a philosophical shift. We stopped thinking about a *single theory* (in particular, for condensed matter theorists) and started thinking about the **space of all possible theories** (so-called "theory space").

The physical act of coarse-graining is a *flow* in this giant space. Each point in the space is a different Lagrangian (a different set of all possible coupling constants).

In this picture, we see universality. For example:

- An Ising model on a square lattice.
- An Ising model on a triangular lattice.
- A real-world liquid-gas system near its critical point.

These are three *wildly different* microscopic (UV) theories. But as we apply the Wilsonian RG (zoom out, coarse-grain, and rescale), they all flow toward the *exact same* point in theory space: the "Ising fixed point."

This is the **many-to-one** mapping that is the hallmark of a semigroup. The reason we prefer this picture is that it *explains the physics*. It explains *why* our low-energy world looks simple.

The irreversible, information-losing, semigroup nature of the RG flow is what washes away all the complex, unknown details of Planck-scale physics, leaving us with the simple, renormalizable, effective field theories (like the Standard Model) that we use to describe our world.