



中国科学技术大学  
University of Science and Technology of China

# Spin Related Nonadiabatic Molecular Dynamics by Hefei-NAMD

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# OUTLINE

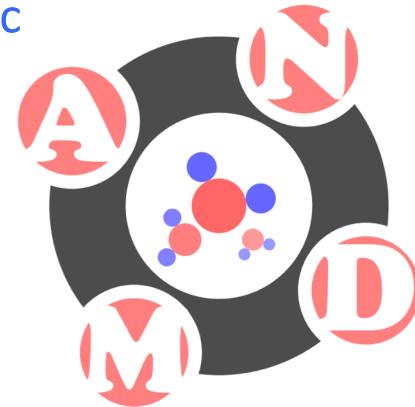
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## METHODOLOGY

- ◆ Time-dependent Schrödinger equation with SOC
- ◆ Two representations: spin-adiabatic & spin-diabatic
- ◆ Work flow of NAMD with SOC

## PRACTICE

- ◆ Different calculations of two representations
- ◆ VASP patches for outputting SOC matrix elements
- ◆ Running NAMD and post data processing



# Electronic Time-dependent Schrödinger Equation

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The time-dependent Schrödinger equation

$$i\hbar \frac{\partial |\Psi(\mathbf{r}, \mathbf{R}(t), \mathbf{s}, t)\rangle}{\partial t} = \hat{\mathcal{H}}^{tot}(\mathbf{r}, \mathbf{R}(t), \mathbf{s}) |\Psi(\mathbf{r}, \mathbf{R}(t), \mathbf{s}, t)\rangle \quad (1)$$

where the total Hamiltonian is given by

$$\hat{\mathcal{H}}^{tot}(\mathbf{r}, \mathbf{R}(t), \mathbf{s}) = \hat{\mathcal{H}}^0(\mathbf{r}, \mathbf{R}(t)) + \hat{\mathcal{H}}^{soc}(\mathbf{r}, \mathbf{R}(t), \mathbf{s}) \quad (2)$$

by expanding the wavefunction with a *basis set*  $\{|\psi_i\rangle\}$  or *representation*

$$|\Psi\rangle = \sum_i |\psi_i\rangle \langle \psi_i | \Psi \rangle = \sum_i c_i |\psi_i\rangle \quad (3)$$

and substituting eq (3) into eq (1), we have

$$\begin{aligned} \frac{\partial c_j(t)}{\partial t} &= - \sum_i \left[ i\hbar^{-1} \langle \psi_j | \hat{\mathcal{H}}^{tot} | \psi_i \rangle + \langle \psi_j | \frac{d}{dt} | \psi_i \rangle \right] c_i(t) \\ &= - \sum_i \left[ i\hbar^{-1} \langle \psi_j | \hat{\mathcal{H}}^0 | \psi_i \rangle + i\hbar^{-1} \langle \psi_j | \hat{\mathcal{H}}^{soc} | \psi_i \rangle + \langle \psi_j | \frac{d}{dt} | \psi_i \rangle \right] c_i(t) \\ &= - \sum_i \left( i\hbar^{-1} H_{ji}^0 + i\hbar^{-1} H_{ji}^{soc} + T_{ji} \right) c_i(t) \end{aligned} \quad (4)$$

# Spin-adiabatic Representation

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Expand the wavefunctions in the “*spin-adiabatic* representation”

$$|\Phi\rangle = \sum_n c_n |\psi_n\rangle = \sum_n c_n \begin{pmatrix} |\psi_n^\uparrow\rangle \\ |\psi_n^\downarrow\rangle \end{pmatrix}; \quad (\hat{\mathcal{H}}^0 + \hat{\mathcal{H}}^{soc}) |\psi_n\rangle = \varepsilon_n |\psi_n\rangle$$

and substitute eq. (10) into TDKS, we have

$$\frac{\partial}{\partial t} \begin{pmatrix} c_1 \\ c_2 \\ \vdots \\ c_n \end{pmatrix} = -i\hbar^{-1} \begin{pmatrix} \varepsilon_1 & \textcolor{blue}{T}_{12} & \dots & \textcolor{blue}{T}_{1n} \\ \textcolor{blue}{T}_{21} & \varepsilon_2 & \dots & \textcolor{blue}{T}_{2n} \\ \vdots & \vdots & \ddots & \vdots \\ \textcolor{blue}{T}_{n1} & \textcolor{blue}{T}_{n2} & \dots & \varepsilon_n \end{pmatrix} \begin{pmatrix} c_1 \\ c_2 \\ \vdots \\ c_n \end{pmatrix}$$

where

$$\textcolor{blue}{T}_{ij} = -i\hbar \left( \langle \psi_i^\uparrow | \langle \psi_i^\downarrow | \right) \frac{d}{dt} \begin{pmatrix} |\psi_j^\uparrow\rangle \\ |\psi_j^\downarrow\rangle \end{pmatrix}$$

- No spin up/down in this case.
- No SOC matrix elements needed, already included in  $\varepsilon_n$ .
- Straightforward with existing NAMD code.

# Spin-diabatic Representation

Expand the wavefunction in the “*spin-diabatic representation*”

$$|\Phi\rangle = \sum_{n\sigma} c_n^\sigma |\psi_n^\sigma\rangle = \sum_n \left[ c_n^\uparrow \begin{pmatrix} |\psi_n^\uparrow\rangle \\ 0 \end{pmatrix} + c_n^\downarrow \begin{pmatrix} 0 \\ |\psi_n^\downarrow\rangle \end{pmatrix} \right]; \quad \hat{\mathcal{H}}^0 |\psi_n^\sigma\rangle = \varepsilon_n^\sigma |\psi_n^\sigma\rangle$$

and substitute eq. (7) into TDKS, we have

$$\frac{\partial}{\partial t} \begin{pmatrix} c_1^\uparrow \\ c_2^\uparrow \\ \vdots \\ c_n^\uparrow \\ c_1^\downarrow \\ c_2^\downarrow \\ \vdots \\ c_n^\downarrow \end{pmatrix} = \frac{-i}{\hbar} \begin{pmatrix} \varepsilon_1^\uparrow + S_{11}^{\uparrow\uparrow} & T_{12}^{\uparrow\uparrow} + S_{12}^{\uparrow\uparrow} & \dots & T_{1n}^{\uparrow\uparrow} + S_{1n}^{\uparrow\uparrow} & | & S_{11}^{\uparrow\downarrow} & S_{12}^{\uparrow\downarrow} & \dots & S_{1n}^{\uparrow\downarrow} \\ T_{21}^{\uparrow\uparrow} + S_{21}^{\uparrow\uparrow} & \varepsilon_2^\uparrow + S_{22}^{\uparrow\uparrow} & \dots & T_{2n}^{\uparrow\uparrow} + S_{2n}^{\uparrow\uparrow} & | & S_{12}^{\uparrow\downarrow} & S_{22}^{\uparrow\downarrow} & \dots & S_{2n}^{\uparrow\downarrow} \\ \vdots & \vdots & \ddots & \vdots & | & \vdots & \vdots & \ddots & \vdots \\ T_{n1}^{\uparrow\uparrow} + S_{n1}^{\uparrow\uparrow} & T_{n2}^{\uparrow\uparrow} + S_{n2}^{\uparrow\uparrow} & \dots & \varepsilon_n^\uparrow + S_{nn}^{\uparrow\uparrow} & | & S_{n1}^{\uparrow\downarrow} & S_{n2}^{\uparrow\downarrow} & \dots & S_{nn}^{\uparrow\downarrow} \\ \hline S_{11}^{\downarrow\uparrow} & S_{12}^{\downarrow\uparrow} & \dots & S_{1n}^{\downarrow\uparrow} & | & \varepsilon_1^\downarrow + S_{11}^{\downarrow\downarrow} & T_{12}^{\downarrow\downarrow} + S_{12}^{\downarrow\downarrow} & \dots & T_{1n}^{\downarrow\downarrow} + S_{1n}^{\downarrow\downarrow} \\ S_{12}^{\downarrow\uparrow} & S_{22}^{\downarrow\uparrow} & \dots & S_{2n}^{\downarrow\uparrow} & | & T_{21}^{\downarrow\downarrow} + S_{21}^{\downarrow\downarrow} & \varepsilon_2^\downarrow + S_{22}^{\downarrow\downarrow} & \dots & T_{2n}^{\downarrow\downarrow} + S_{2n}^{\downarrow\downarrow} \\ \vdots & \vdots & \ddots & \vdots & | & \vdots & \vdots & \ddots & \vdots \\ S_{n1}^{\downarrow\uparrow} & S_{n2}^{\downarrow\uparrow} & \dots & S_{nn}^{\downarrow\uparrow} & | & T_{n1}^{\downarrow\downarrow} + S_{n1}^{\downarrow\downarrow} & T_{n2}^{\downarrow\downarrow} + S_{n2}^{\downarrow\downarrow} & \dots & \varepsilon_n^\downarrow + S_{nn}^{\downarrow\downarrow} \end{pmatrix} \begin{pmatrix} c_1^\uparrow \\ c_2^\uparrow \\ \vdots \\ c_n^\uparrow \\ c_1^\downarrow \\ c_2^\downarrow \\ \vdots \\ c_n^\downarrow \end{pmatrix}$$

where

$$S_{ij}^{\sigma_1 \sigma_2} = \langle \psi_i^{\sigma_1} | \hat{\mathcal{H}}^{soc} | \psi_j^{\sigma_2} \rangle \quad T_{ij}^{\sigma\sigma} = -i\hbar \langle \psi_i^\sigma | \frac{d}{dt} | \psi_j^\sigma \rangle$$

# Spin-orbit Couplings in DFT

Spin-orbit Hamiltonian and AE wavefunctions in PAW formalism

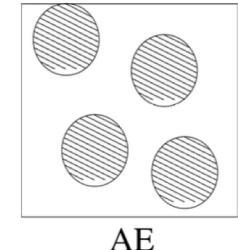
we have

$$\hat{\mathcal{H}}^{soc}(\mathbf{r}) = \frac{\hbar\vec{\sigma} \cdot \vec{\mathbf{p}} \times \nabla v_{KS}(\mathbf{r})}{4m^2c^2};$$

Due to the derivatives, SOC  
is dominated in core region.

$$|\psi_{n\sigma}\rangle = \sum_{i=\{l,m\},\alpha} \langle \tilde{\mathbf{p}}_{i\sigma} | \tilde{\psi}_{n\sigma} \rangle |\phi_{i\sigma}^\alpha\rangle$$

AE in PAW sphere



$$\langle \psi_{n_1\sigma_1} | \hat{\mathcal{H}}^{soc} | \psi_{n_2\sigma_2} \rangle = \sum_{\alpha,i,j} \langle \tilde{\psi}_{n_1\sigma_1} | \tilde{\mathbf{p}}_{i\sigma_1}^\alpha \rangle \langle \phi_{i\sigma_1}^\alpha | \hat{\mathcal{H}}^{soc} | \phi_{j\sigma_2}^\alpha \rangle \langle \tilde{\mathbf{p}}_{j\sigma_2}^\alpha | \tilde{\psi}_{n_2\sigma_2} \rangle$$

- The projector coefficients  $\langle \tilde{\mathbf{p}}_i | \tilde{\psi}_n \rangle$  is contained in NormalCAR, which can be easily output and parsed.

$$\langle \phi_{i\sigma_1}^\alpha | \hat{\mathcal{H}}^{soc} | \phi_{j\sigma_2}^\alpha \rangle = -\frac{1}{2m^2c^2} \langle Y_i^\alpha; \sigma_1 | \hat{\mathbf{S}} \cdot \hat{\mathbf{L}} | Y_j^\alpha; \sigma_2 \rangle \langle f_i^\alpha | \frac{1}{r} \frac{dv_{KS}^\alpha}{dr} | f_j^\alpha \rangle$$

- This term is *not as easily obtained*, where  $f_i^\alpha$  is the radial function,  $Y_i^\alpha$  is the spherical harmonics and KS potential  $v_{KS}^\alpha = v_{ext}^\alpha + v_{xc}^\alpha + v_{Har}^\alpha$  depend on the charge density of each atom site.

# OUTLINE

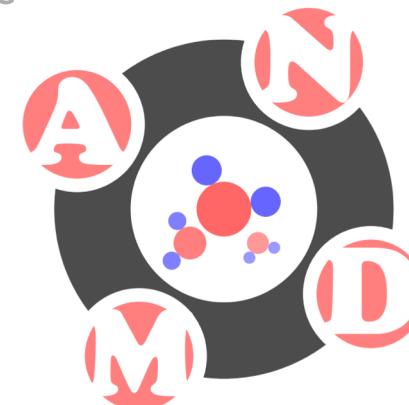
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## METHODOLOGY

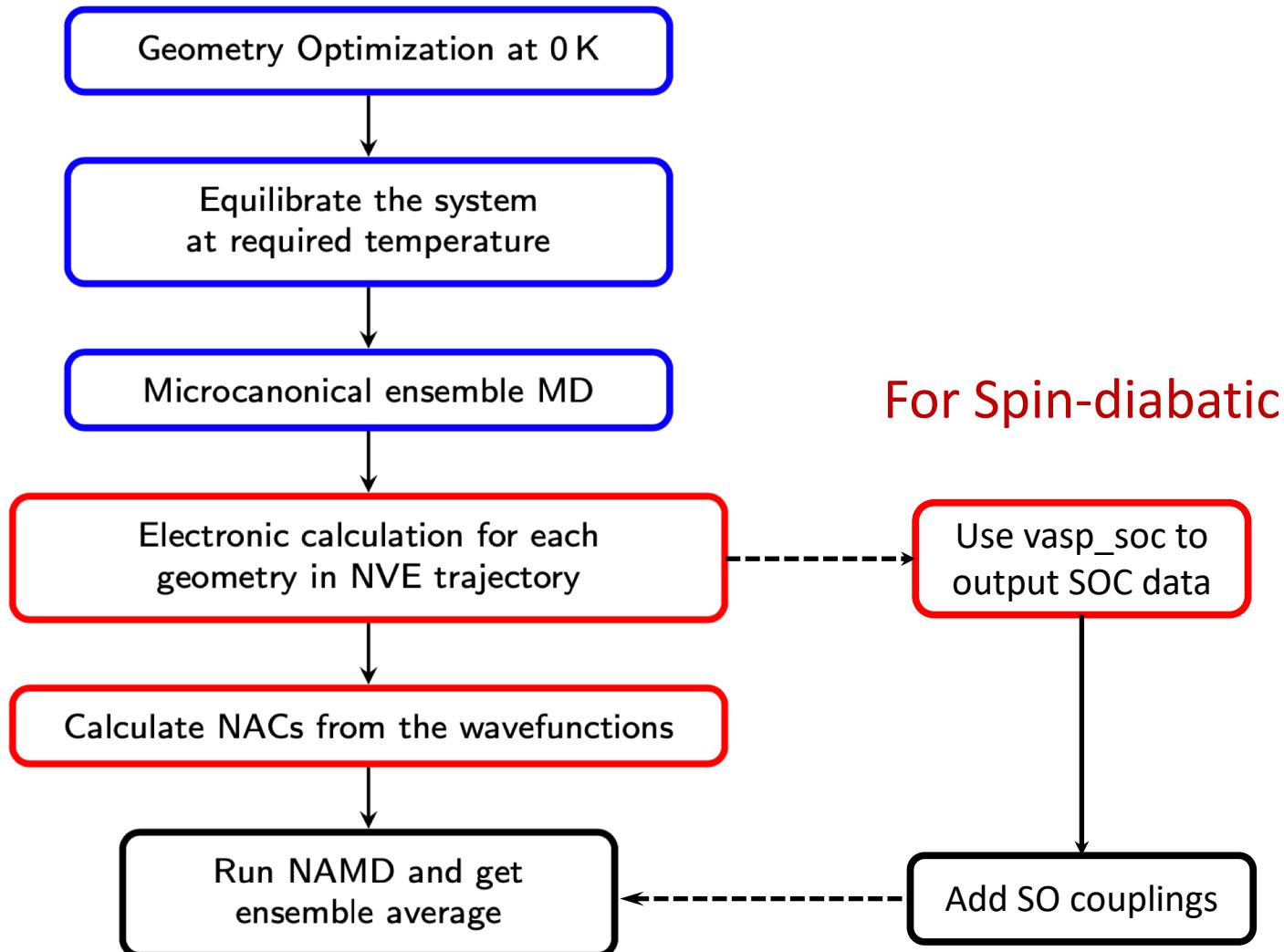
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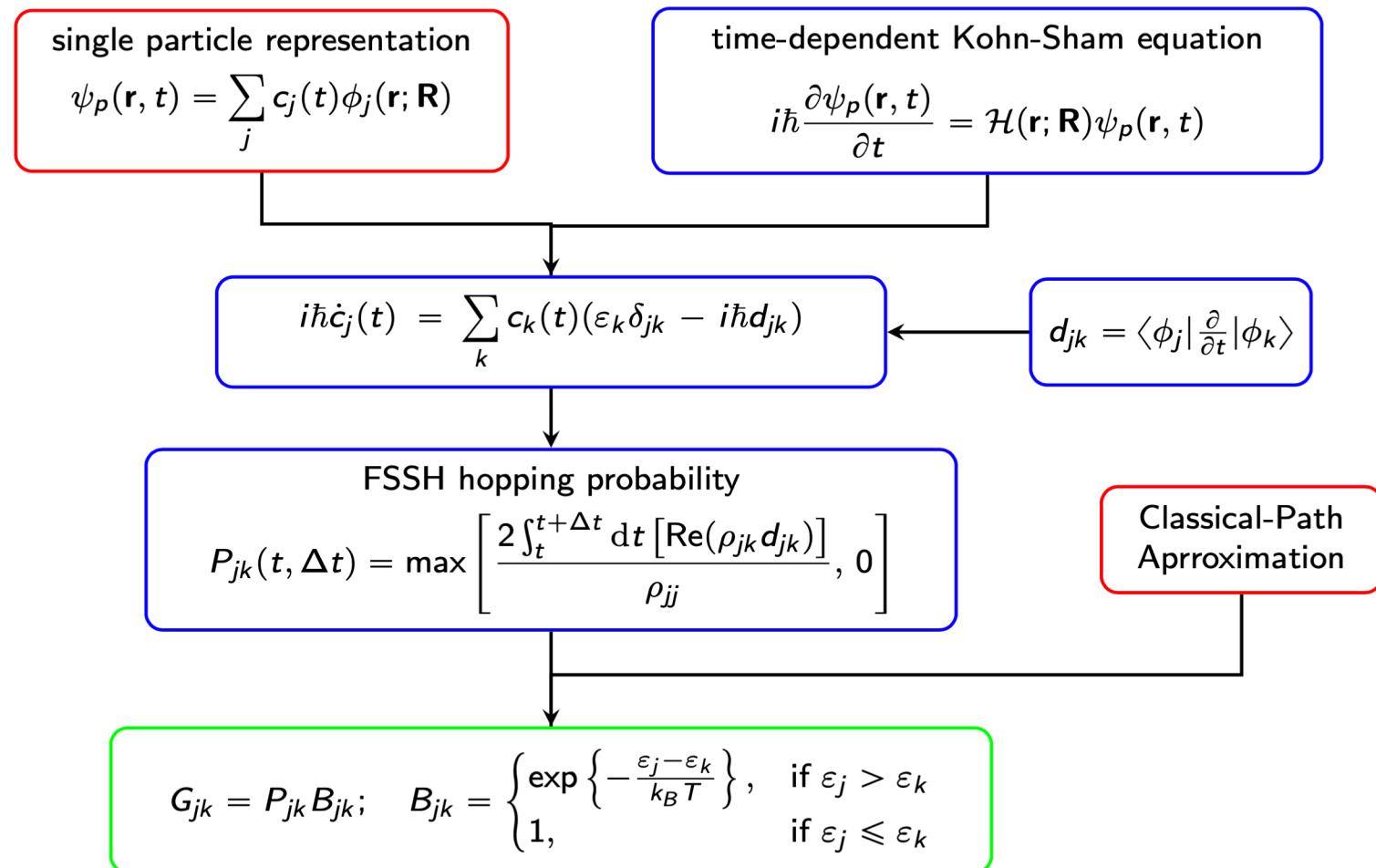


# Work Flow



# Surface Hopping with SOC

## Spin-adiabatic Representation



# Surface Hopping with SOC

## Spin-diabatic Representation

single particle representation

$$\psi_p(\mathbf{r}, t) = \sum_j c_j(t) \phi_j(\mathbf{r}; \mathbf{R})$$

time-dependent Kohn-Sham equation

$$i\hbar \frac{\partial \psi_p(\mathbf{r}, t)}{\partial t} = \mathcal{H}(\mathbf{r}; \mathbf{R}) \psi_p(\mathbf{r}, t)$$

$$\frac{\partial c_j(t)}{\partial t} = - \sum_i \left( i\hbar^{-1} H_{ji}^0 + i\hbar^{-1} H_{ji}^{soc} + T_{ji} \right) c_i(t)$$

$$d_{jk} = \langle \phi_j | \frac{\partial}{\partial t} | \phi_k \rangle$$

The hopping probility within FSSH

$$P_{j \rightarrow k}(t, \Delta t) = \max \left( - \frac{2\Delta t [\hbar^{-1} \text{Im}(c_j^* c_k (H_{jk}^0 + H_{jk}^{soc})) - \text{Re}(c_j^* c_k T_{jk})]}{c_j^* c_j}, 0 \right)$$

Classical-Path Aprroximation

$$G_{jk} = P_{jk} B_{jk}; \quad B_{jk} = \begin{cases} \exp \left\{ - \frac{\varepsilon_j - \varepsilon_k}{k_B T} \right\}, & \text{if } \varepsilon_j > \varepsilon_k \\ 1, & \text{if } \varepsilon_j \leq \varepsilon_k \end{cases}$$

# INCARs for Wavefunction Calculations

## Spin-adiabatic (vasp\_ncl)

```
1 Electronic Calculation for NAMD
2 SYSTEM = Ni
3
4 #Start Parameter
5   ISTART = 0
6   ICHARG = 2
7
8 #Electronic Relaxation
9   ALGO = Fast
10  NELM = 100
11  PREC = Normal
12  ISMEAR = 0
13  SIGMA = 0.02
14  NPAR = 4
15
16  #ISPIN = 2
17  MAGMOM = 0 0 1 0 0 1 0 0 1 0 0 1 0 0
18  LSORBIT = .TRUE
19  LMAXMIX = 4
20
21 #Ionic Relaxation
22  IBRION = -1
23
24 #Write Flags
25  NWRITE = 1
26  LREAL = Auto
27  LWAVE = .TRUE
28  LCHARG = .FALSE
29  LORBIT = 11
~
```

## Spin-diabatic (vasp\_soc)

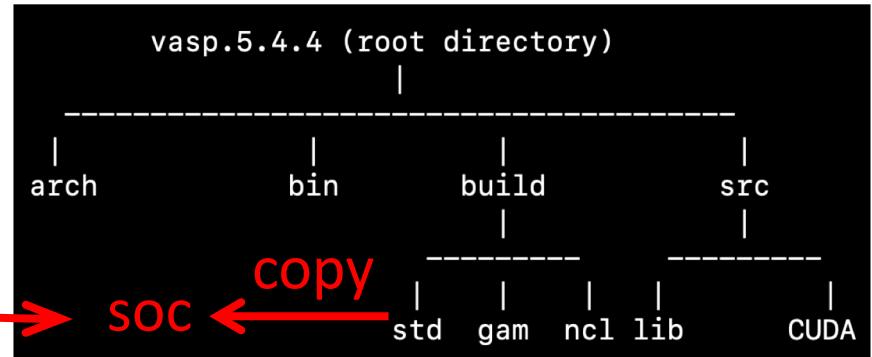
```
1 Electronic Calculation for NAMD
2 SYSTEM = Ni
3
4 #Start Parameter
5   ISTART = 0
6   ICHARG = 2
7
8 #Electronic Relaxation
9   ALGO = Fast
10  NELM = 100
11  PREC = Normal
12  ISMEAR = 0
13  SIGMA = 0.02
14  NPAR = 4
15
16  ISPIN = 2
17  MAGMOM = 32*1
18
19 #Ionic Relaxation
20  IBRION = -1
21
22 #Write Flags
23  NWRITE = 1
24  LREAL = Auto
25  LWAVE = .TRUE
26  LCHARG = .FALSE
27  LORBIT = 11
~
```

# VASP Patches

Patches for VASP 5.4.4.

```
zzf@SciNat:Patch_soc$ ls  
main.F  paw.F  relativistic.F  stm.F
```

*Replace original codes*



1. Copy “std” folder to “soc”.
2. Replace original codes by patches.
3. Complile!
4. Copy and rename “vasp” to “vasp\_soc”

NormalCAR    SocCar

```
zzf@SciNat:vasp.5.4.4$ ls  
arch  build  makefile.include  src  
bin   makefile  README  
zzf@SciNat:vasp.5.4.4$ cd build/  
zzf@SciNat:build$ ls  
gam  ncl  std  
zzf@SciNat:build$ cp -r std soc  
zzf@SciNat:build$ cp /data/zzf/Patch_soc/*.F soc/  
zzf@SciNat:build$ cd soc  
zzf@SciNat:soc$ make clean  
rm -f *.o *.mod *.f90  
zzf@SciNat:soc$ make
```

```
zzf@SciNat:soc$ ls vasp  
vasp  
zzf@SciNat:soc$ cp vasp ../../bin/vasp_soc  
zzf@SciNat:soc$ ls ../../bin/  
vasp_gam  vasp_ncl  vasp_soc  vasp_std
```

# namd\_soc Input Files

inp

```
1 &NAMD PARA
2   BMIN      = 343
3   BMAX      = 366
4   NBANDS    = 416
5
6   SOCTYPE   = 1 Spin-adiabatic
7
8   NSW       = 2000
9   POTIM     = 1
10  TEMP      = 300
11
12  NSAMPLE   = 10
13  NAMDTIME  = 1000
14  NELM      = 1000
15  NTRAJ     = 20000
16  LHOLE     = .F.
17
18  RUNDIR    = "../run"
19  LCPEXT    = .F.
20 /
~
```

```
1 &NAMD PARA
2   BMINU     = 168
3   BMAXU     = 183
4   BMIND     = 149
5   BMAXD     = 171
6   NBANDS    = 208
7
8   SOCTYPE   = 2 Spin-diabatic
9
10  NSW       = 2000
11  POTIM     = 1
12  TEMP      = 300
13
14  NSAMPLE   = 10
15  NAMDTIME  = 1000
16  NELM      = 1000
17  NTRAJ     = 20000
18  LHOLE     = .F.
19
20  RUNDIR    = "../run"
21  LCPEXT    = .F.
22 /
```

INICON

```
1 302 364
2 702 366
3 689 365
4 321 366
5 958 366
6 316 364
7 458 365
8 731 366
9 254 365
10 840 366
~
```

```
1 302 181 1
2 702 183 1
3 689 183 1
4 321 182 1
5 958 182 1
6 316 183 1
7 458 183 1
8 731 181 1
9 254 183 1
10 840 183 1
~
```

# Run NAMD with SOC

## Compile namd\_soc

```
zzf@SciNat:src$ ls
couplings.f90  lattice.f90  prec.f90      TimeProp.f90
fileio.f90     main.f90     soc.F90       wave.f90
hamil.f90      Makefile    SurfHop.f90
zzf@SciNat:src$ make clean
rm -f *.mod *.a namd
rm -f prec.o lattice.o wave.o fileio.o soc.o couplings.o ha
mil.o TimeProp.o SurfHop.o main.o namd_soc
zzf@SciNat:src$ make
gfortran -g -O2 -c prec.f90
gfortran -g -O2 -c lattice.f90
gfortran -g -O2 -c wave.f90
gfortran -g -O2 -c fileio.f90
gfortran -g -O2 -c soc.F90
gfortran -g -O2 -c couplings.f90
gfortran -g -O2 -c hamil.f90
gfortran -g -O2 -c TimeProp.f90
gfortran -g -O2 -c SurfHop.f90
gfortran -g -O2 -c main.f90
gfortran -g -O2 -o namd_soc prec.o lattice.o wave.o fileio.
o soc.o couplings.o hamil.o TimeProp.o SurfHop.o main.o
zzf@SciNat:src$ cp namd_soc ~/bin
zzf@SciNat:src$
```

## NAMD Results

```
zzf@SciNat:namd$ ls
COUPCAR  EIGTXT  INICON  inp  NATTXT  PSICT.4  SHPROP.4
zzf@SciNat:namd$ 
zzf@SciNat:namd$ ls
COUPCAR  EIGTXT  INICON  inp  NATTXT  PSICT.2  SHPROP.2  SOTXT
zzf@SciNat:namd$
```

```
zzf@SciNat:namd$ namd_soc
```

```
Hefei-NAMD (soc version 1.3.1, Dec 16, 2021)
```

```
../run/01/WAVECAR
../run/02/WAVECAR
../run/03/WAVECAR
../run/04/WAVECAR
../run/05/WAVECAR
../run/06/WAVECAR
../run/07/WAVECAR
../run/08/WAVECAR
../run/09/WAVECAR
../run/10/WAVECAR
../run/11/WAVECAR
../run/12/WAVECAR
../run/13/WAVECAR
../run/14/WAVECAR
../run/15/WAVECAR
../run/16/WAVECAR
../run/17/WAVECAR
../run/18/WAVECAR
../run/19/WAVECAR
../run/20/WAVECAR
```

```
BMIN = 316
BMAX = 354
NBANDS = 416
INIBAND = 354
SOCTYPE = 1
NSW = 21
POTIM = 1.0
TEMP = 300.0
NAMDTINI = 4
NAMDTIME = 15
NTRAJ = 20000
NELM = 1000
LHOLE = F
LSHP = T
LCPTXT = F
RUNDIR = ../run
```

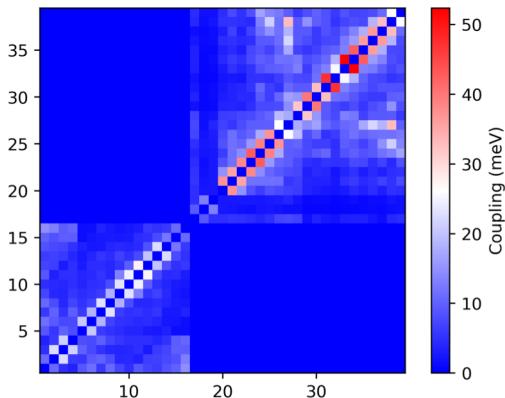
```
CPU Time [s]: 0.27
```

```
zzf@SciNat:namd$
```

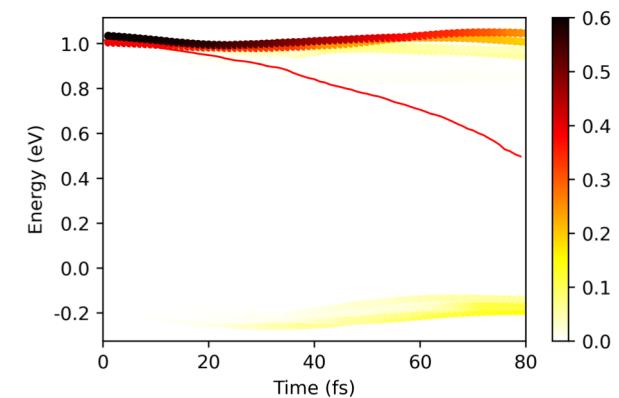
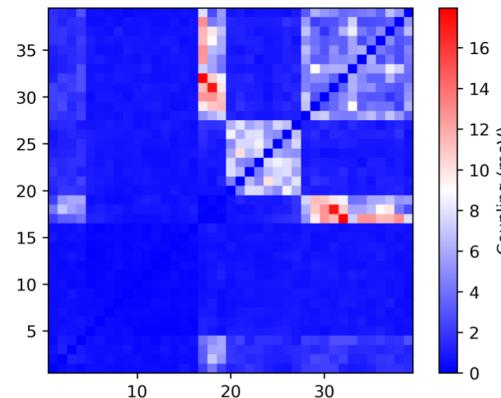
# Data Processing After NAMD

```
zzf@SciNat:namd$ python namdplt.py  
zzf@SciNat:namd$ ls *png  
COUPLE_NA.png  COUPLE_SO.png  NAMD.png
```

NAC



SOC



# More Information

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赵瑾老师主页

<http://home.ustc.edu.cn/~ljchen9>

Hefei-NAMD

<https://github.com/QijingZheng/Hefei-NAMD>

蔻享社区: 物理学/Hefei-NAMD板块

网址: <http://bbs.koushare.com>

往期培训

郑奇靖 | Hefei-NAMD Training Session 2020/12/03

<http://staff.ustc.edu.cn/~zqj/posts/Hefei-NAMD-Training/>

褚维斌 | Hefei-NAMD 使用的一些经验 2021/05/10

<https://www.koushare.com/video/videodetail/11720>

蒋翔 | Exciton Dynamics based on Hefei-NAMD 2021/05/10

<https://www.koushare.com/video/videodetail/11721>

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