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Nested regular polygon solutions of $2N$ -body problem[☆]

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Abstract

In this Letter we study the necessary conditions for the masses of the nested regular polygon solutions of the planar $2N$ -body problem. © 2001 Published by Elsevier Science B.V.

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1. Main results

This Letter uses the same notations as paper [5]. For $n \geq 2$ the equations of motion of the planar n -body problem [1,2,5,6] can be written in the form

$$\ddot{z}_k = - \sum_{j \neq k} m_j \frac{z_k - z_j}{|z_k - z_j|^3}, \quad (1.1)$$

where z_k is the complex coordinate of the k th mass m_k in an inertial coordinate system. In Eq. (1.1) and throughout this Letter, unless otherwise restricted, all indices and summations will range from 1 to N .

Let ρ_k denote the N complex k th roots of unity, i.e.,

$$\rho_k = \exp(2\pi i k / N). \quad (1.2)$$

This equation will also serve to define ρ_k for any number k . We assume mass m_k ($k = 1, \dots, N$) locate at the vertices ρ_k of a regular polygon inscribed on the unit circle, and $\bar{m}_k = b m_k$ ($b > 0$, $k = 1, \dots, N$) locate at

$$\bar{\rho}_k = a \rho_k \quad (a > 0, a \neq 1). \quad (1.3)$$

Then the center of masses $m_1, \dots, m_N, \bar{m}_1, \dots, \bar{m}_N$ is

$$z_0 = \sum_j (m_j \rho_j + \bar{m}_j \bar{\rho}_j) / M, \quad (1.4)$$

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where $M = \sum_j (m_j + \bar{m}_j)$. The functions describing their rotation about z_0 with angular velocity ω are then given by

$$z_k(t) = (\rho_k - z_0) \exp(i\omega t), \quad k = 1, \dots, N, \quad (1.5)$$

and

$$\bar{z}_k(t) = (a\rho_k - z_0) \exp(i\omega t), \quad k = 1, \dots, N. \quad (1.6)$$

Then the equations of motion of the planar $2N$ -body problem can be written in the form

$$\ddot{z}_k = \sum_{j \neq k} m_j \frac{z_j - z_k}{|z_j - z_k|^3} + \sum_j \bar{m}_j \frac{\bar{z}_j - z_k}{|\bar{z}_j - z_k|^3} \quad (1.7)$$

and

$$\ddot{\bar{z}}_k = \sum_j m_j \frac{z_j - \bar{z}_k}{|z_j - \bar{z}_k|^3} + \sum_{j \neq k} \bar{m}_j \frac{\bar{z}_j - \bar{z}_k}{|\bar{z}_j - \bar{z}_k|^3}. \quad (1.8)$$

Moeckel and Simo [4] proved the following result:

Theorem (Moeckel–Simo). *For every mass ratio b , there are exactly two planar central configurations consisting of two nested regular N -gons. For one of these, the ratio of the sizes of the two polygons is less than 1, and for the other it is greater than 1.*

In this Letter, we study the inverse problem of the theorem (Moeckel–Simo) and the following results are established.

Theorem 1. *If, for $N \geq 2$, the functions $z_k(t)$ and $\bar{z}_k(t)$ given by (1.5) and (1.6) are solutions of the $2N$ -body problem (1.7) and (1.8), it follows that $\omega^2 = M\gamma/N$, where*

$$\gamma = \frac{1}{1+b} \left(\sum_{j=1}^{N-1} \frac{1-\rho_j}{|1-\rho_j|^3} + \sum_{j=0}^{N-1} \frac{b(1-a\rho_j)}{|1-a\rho_j|^3} \right), \quad (1.9)$$

b and a has the following relationship:

$$b = \frac{(\sum_{j=0}^{N-1} (a-\rho_j)/|a-\rho_j|^3 - \sum_{j=1}^{N-1} a(1-\rho_j)/|1-\rho_j|^3)}{(\sum_{j=0}^{N-1} a(1-a\rho_j)/|1-a\rho_j|^3 - \sum_{j=1}^{N-1} a(1-\rho_j)/a^3|1-\rho_j|^3)}. \quad (1.10)$$

Theorem 2. *For $N \geq 2$, $m_k > 0$, and $\bar{m}_k = bm_k$, $b > 0$, the functions $z_k(t)$ and $\bar{z}_k(t)$ given by (1.5) and (1.6) with $\omega^2 = M\gamma/N$ and γ given by (1.9) are solutions of the $2N$ -body problem (1.7) and (1.8) if and only if $m_1 = m_2 = \dots = m_N$ and b is determined uniquely by (1.10).*

Corollary. *For $N = 2$ and $a > 1$ the functions $z_k(t)$ and $\bar{z}_k(t)$ given by (1.5) and (1.6) with*

$$\omega^2 = M\gamma/N \quad \text{and} \quad \gamma = \frac{1}{1+b} \left(\frac{1}{4} - \frac{4ab}{(a^2-1)^2} \right)$$

are solutions of the (2×2) -body problem (1.7) and (1.8) if and only if

$$m_1 = m_2 \quad \text{and} \quad b = \frac{a^7 - 2a^5 - 8a^4 + a^3 - 8a^2}{17a^4 - 2a^2 + 1}.$$

This is a collinear periodic solution for 4-body problem. When $0 < a < 1$,

$$\gamma = \frac{1}{1+b} \left(\frac{1}{4} + \frac{2b(a^2+1)}{(a^2-1)^2} \right) \quad \text{and} \quad b = \frac{a^7 - 2a^5 + 17a^3}{-8a^5 + a^4 - 8a^3 - 2a^2 + 1}.$$

2. The proof of the main results

For two nested regular polygons, we have defined

$$\rho_k = \exp(2\pi i k / N), \quad (2.1)$$

$$\bar{\rho}_k = a \exp(2\pi i k / N), \quad a \neq 1, \quad (2.2)$$

$$z_0 = \sum_j (m_j \rho_j + \bar{m}_j \bar{\rho}_j) / M, \quad (2.3)$$

where

$$M = \sum_j (m_j + \bar{m}_j), \quad (2.4)$$

$$z_k(t) = (\rho_k - z_0) \exp(i\omega t), \quad k = 1, \dots, N, \quad (2.5)$$

and

$$\bar{z}_k(t) = (a\rho_k - z_0) \exp(i\omega t), \quad k = 1, \dots, N. \quad (2.6)$$

Proof of Theorem 1. Direct substitution into the differential equations (1.7), (1.8) shows that the $z_k(t)$ and $\bar{z}_k(t)$ are the solution of (1.7) and (1.8) if and only if

$$(\rho_k - z_0) \omega^2 \exp(i\omega t) = \left(\sum_{j \neq k} m_j \frac{\rho_k - \rho_j}{|\rho_k - \rho_j|^3} + \sum_j \bar{m}_j \frac{\rho_k - \bar{\rho}_j}{|\rho_k - \bar{\rho}_j|^3} \right) \exp(i\omega t) \quad (2.7)$$

and

$$(\bar{\rho}_k - z_0) \omega^2 \exp(i\omega t) = \left(\sum_j m_j \frac{\bar{\rho}_k - \rho_j}{|\bar{\rho}_k - \rho_j|^3} + \sum_{j \neq k} \bar{m}_j \frac{\bar{\rho}_k - \bar{\rho}_j}{|\bar{\rho}_k - \bar{\rho}_j|^3} \right) \exp(i\omega t), \quad (2.8)$$

or if and only if

$$\sum_{j \neq k} m_j \left(\frac{1}{|\rho_k - \rho_j|^3} - \frac{\omega^2}{M} \right) (\rho_k - \rho_j) + \sum_j \bar{m}_j \left(\frac{1}{|\rho_k - \bar{\rho}_j|^3} - \frac{\omega^2}{M} \right) (\rho_k - \bar{\rho}_j) = 0 \quad (2.9)$$

and

$$\sum_j m_j \left(\frac{1}{|\bar{\rho}_k - \rho_j|^3} - \frac{\omega^2}{M} \right) (\bar{\rho}_k - \rho_j) + \sum_{j \neq k} \bar{m}_j \left(\frac{1}{|\bar{\rho}_k - \bar{\rho}_j|^3} - \frac{\omega^2}{M} \right) (\bar{\rho}_k - \bar{\rho}_j) = 0. \quad (2.10)$$

Multiplying both sides by ρ_{N-k} and noting that $|\rho_k - \rho_j| = |\rho_k| |1 - \rho_{j-k}| = |1 - \rho_{j-k}|$ and using $\bar{\rho}_k = a\rho_k$,

$$\sum_{j \neq k} m_j \left(\frac{1}{|1 - \rho_{j-k}|^3} - \frac{\omega^2}{M} \right) (1 - \rho_{j-k}) + \sum_j \bar{m}_j \left(\frac{1}{|1 - a\rho_{j-k}|^3} - \frac{\omega^2}{M} \right) (1 - a\rho_{j-k}) = 0 \quad (2.11)$$

and

$$\sum_j m_j \left(\frac{1}{|a - \rho_{j-k}|^3} - \frac{\omega^2}{M} \right) (a - \rho_{j-k}) + \sum_{j \neq k} \bar{m}_j \left(\frac{1}{|a - a\rho_{j-k}|^3} - \frac{\omega^2}{M} \right) (a - a\rho_{j-k}) = 0. \quad (2.12)$$

Notice that every step from (2.7)–(2.12) can be conversed respectively. Now we define the $N \times N$ circulant matrix [3] $C = [c_{k,j}]$, $A = [a_{k,j}]$, $B = [b_{k,j}]$, $D = [d_{k,j}]$ as follows:

$$c_{k,j} = 0, \quad \text{for } k = j, \quad c_{k,j} = \left(\frac{1}{|1 - \rho_{j-k}|^3} - \frac{\omega^2}{M} \right) (1 - \rho_{j-k}), \quad \text{for } k \neq j, \quad (2.13)$$

$$a_{k,j} = \left(\frac{1}{|1 - a\rho_{j-k}|^3} - \frac{\omega^2}{M} \right) (1 - a\rho_{j-k}), \quad (2.14)$$

$$b_{k,j} = \left(\frac{1}{|a - \rho_{j-k}|^3} - \frac{\omega^2}{M} \right) (a - \rho_{j-k}), \quad (2.15)$$

$$d_{k,j} = 0, \quad \text{for } k = j, \quad d_{k,j} = \left(\frac{1}{|a - a\rho_{j-k}|^3} - \frac{\omega^2}{M} \right) (a - a\rho_{j-k}), \quad \text{for } k \neq j. \quad (2.16)$$

Then (2.11) and (2.12) hold if and only if the matrix equation

$$\begin{bmatrix} C & A \\ B & D \end{bmatrix} \begin{pmatrix} m_1 \\ \vdots \\ m_N \\ \bar{m}_1 \\ \vdots \\ \bar{m}_N \end{pmatrix} = 0 \quad (2.17)$$

has a positive solution.

When $\bar{m}_j = bm_j$, $b > 0$, (2.17) is equivalent to

$$\begin{bmatrix} C & bA \\ B & bD \end{bmatrix} \begin{pmatrix} m_1 \\ \vdots \\ m_N \\ m_1 \\ \vdots \\ m_N \end{pmatrix} = 0 \quad (2.18)$$

$$\Leftrightarrow \begin{bmatrix} C + bA \\ B + bD \end{bmatrix} \begin{pmatrix} m_1 \\ \vdots \\ m_N \end{pmatrix} = 0 \quad (2.19)$$

$$\Leftrightarrow (C + bA) \begin{pmatrix} m_1 \\ \vdots \\ m_N \end{pmatrix} = 0 \quad (2.20)$$

and

$$(B + bD) \begin{pmatrix} m_1 \\ \vdots \\ m_N \end{pmatrix} = 0. \quad (2.21)$$

We notice that Eqs. (2.11) and (2.12) are equivalent to (2.20) and (2.21). According to (2.11), (2.12), (2.20) and (2.21), similar to the proof of Theorem 1 in [6], we have

$$\left(\sum_k m_k \right) \left[\sum_{j \neq N} \left(\frac{1}{|1 - \rho_j|^3} - \frac{\omega^2}{M} \right) (1 - \rho_j) + \sum_j b \left(\frac{1}{|1 - a\rho_j|^3} - \frac{\omega^2}{M} \right) (1 - a\rho_j) \right] = 0 \quad (2.22)$$

$$\Leftrightarrow \frac{\omega^2}{M} \sum_j [(1 - \rho_j) + b(1 - a\rho_j)] = \sum_{j \neq N} \frac{1 - \rho_j}{|1 - \rho_j|^3} + \sum_j \frac{b(1 - a\rho_j)}{|1 - a\rho_j|^3} \quad (2.23)$$

and

$$\left(\sum_k m_k \right) \left[\sum_j \left(\frac{1}{|a - \rho_j|^3} - \frac{\omega^2}{M} \right) (a - \rho_j) + \sum_{j \neq N} b \left(\frac{1}{|a - a\rho_j|^3} - \frac{\omega^2}{M} \right) (a - a\rho_j) \right] = 0 \quad (2.24)$$

$$\Leftrightarrow \frac{\omega^2}{M} \sum_j [(a - \rho_j) + b(a - a\rho_j)] = \sum_j \frac{a - \rho_j}{|a - \rho_j|^3} + \sum_{j \neq N} \frac{b(a - a\rho_j)}{|a - a\rho_j|^3}. \quad (2.25)$$

We note that

$$\sum_j (1 - \rho_j) = N, \quad (2.26)$$

$$\sum_j b(1 - a\rho_j) = Nb, \quad (2.27)$$

$$\sum_j (a - \rho_j) = Na, \quad (2.28)$$

$$\sum_j b(a - a\rho_j) = Nba, \quad (2.29)$$

then (2.23) and (2.25) are equivalent to the following equations, respectively:

$$\frac{\omega^2}{M} = \frac{1}{N(1+b)} \left[\sum_{j \neq N} \frac{1 - \rho_j}{|1 - \rho_j|^3} + \sum_j \frac{b(1 - a\rho_j)}{|1 - a\rho_j|^3} \right], \quad (2.30)$$

$$\frac{\omega^2}{M} = \frac{1}{Na(1+b)} \left[\sum_j \frac{a - \rho_j}{|a - \rho_j|^3} + \sum_{j \neq N} \frac{b(a - a\rho_j)}{|a - a\rho_j|^3} \right]. \quad (2.31)$$

So γ is determined by (1.9) and b is determined uniquely by (1.10). The proof is completed. \square

Proof of Theorem 2.

(1) Proof of the necessary conditions of Theorem 2: The proof hinges on showing that certain eigenvalues of the circulant matrix $C + bA$ are zero. This is accomplished by using the general formulas (2.32), (2.33) for the eigenvalues λ_k and the eigenvectors \vec{v}_k of a circulant matrix $C + bA = [c_{k,j} + ba_{k,j}]$:

$$\lambda_k = \sum_j (c_{1,j} + ba_{1,j}) \rho_{k-1}^{j-1}, \quad (2.32)$$

$$\vec{v}_k = (\rho_{k-1}, \rho_{k-1}^2, \dots, \rho_{k-1}^N)^T. \quad (2.33)$$

Since

$$(c_{1,1} + ba_{1,1}) = b \left(\frac{1}{|1 - a|^3} - \frac{\omega^2}{M} \right) (1 - a) \neq 0,$$

so $C + bA$ is not a zero matrix. According to the relations between the eigenvalues λ_k and the eigenvectors \vec{v}_k , one has

$$(C + bA)\vec{v}_k = \lambda_k \vec{v}_k, \quad 1 \leq k \leq N.$$

Then, in order to find the solution of (2.20), it is enough to find the zero eigenvalue λ_k with the positive real eigenvectors \vec{v}_k .

(i) N is an odd number. Only $\vec{v}_1 = (1, 1, \dots, 1)^T$ is the positive real eigenvector. At the same time

$$\lambda_1 = \sum_j (c_{1,j} + ba_{1,j}) \rho_{1-1}^{j-1} = \sum_j (c_{1,j} + ba_{1,j}) = 0$$

(see (2.22) in the proof of Theorem 1). That is, $(C + bA)\vec{v}_1 = 0$ and $\vec{m} = (m_1, m_1, \dots, m_1)^T$, $m_1 > 0$, is the unique solution of (2.20).

(ii) N is an even number. Only $\vec{v}_1 = (1, 1, \dots, 1)^T$ and $\vec{v}_{N/2+1} = (-1, 1, \dots, -1, 1)^T$ are real eigenvector, but only \vec{v}_1 is a positive real eigenvector. The corresponding eigenvalue is

$$\lambda_1 = \sum_j (c_{1,j} + ba_{1,j}) \rho_{1-1}^{j-1} = \sum_j (c_{1,j} + ba_{1,j}) = 0$$

(see (2.22) in the proof of Theorem 1). That is, $(C + bA)\vec{v}_1 = 0$ and $\vec{m} = (m_1, m_1, \dots, m_1)^T$, $m_1 > 0$, is the unique solution of (2.20).

Similar to the above proof, the matrix equation (2.21) have the same result as the matrix equation (2.20).

(2) Proof of sufficient conditions of Theorem 2: Assume $m_1 = m_2 = \dots = m_N > 0$, γ is determined by (1.9) and a, b are determined by (1.10), then $(m_1, \dots, m_1)^T$ is a solution of (2.21) and (2.22) or (1.7) and (1.8). Since

$$\begin{aligned} & \sum_{j \neq k} m_j \left(\frac{1}{|1 - \rho_{j-k}|^3} - \frac{\omega^2}{M} \right) (1 - \rho_{j-k}) + \sum_j b m_j \left(\frac{1}{|1 - a\rho_{j-k}|^3} - \frac{\omega^2}{M} \right) (1 - a\rho_{j-k}) \\ &= m_1 \left[\sum_{j \neq k} \left(\frac{1}{|1 - \rho_{j-k}|^3} - \frac{\omega^2}{M} \right) (1 - \rho_{j-k}) + \sum_j b \left(\frac{1}{|1 - a\rho_{j-k}|^3} - \frac{\omega^2}{M} \right) (1 - a\rho_{j-k}) \right] \\ &= m_1 \left[\sum_{j \neq N} \left(\frac{1}{|1 - \rho_j|^3} - \frac{\omega^2}{M} \right) (1 - \rho_j) + \sum_j b \left(\frac{1}{|1 - a\rho_j|^3} - \frac{\omega^2}{M} \right) (1 - a\rho_j) \right] = 0. \end{aligned} \quad (2.34)$$

Similar to the above proof, the same results for (2.22) are easily obtained. Thus the proof is completed. \square

The corollary is obvious from Theorem 2.

References

- [1] F. Diacu, Singularities of the N -Body Problem, Les Editions CRM, Montreal, 1992, p. 59.
- [2] B. Elmabsout, Celest. Mech. 41 (1988) 131.
- [3] M. Marcus, H. Minc, A Survey of Matrix Theory and Matrix Inequalities, Allyn and Bacon, Boston, MA, 1964.
- [4] R. Moeckel, C. Simo, SIAM J. Math. Anal. 26 (1995) 978.
- [5] L.M. Perko, E.L. Walter, Proc. Amer. Math. Soc. 94 (1985) 301.
- [6] Z. Xie, S. Zhang, Phys. Lett. A 277 (2000).