

# Poisson's Equation and Fast Method

# Poisson's Equation and its Fundamental Solution

$$\nabla^2 \phi = -\rho$$

$$\phi(x \rightarrow \infty) = 0$$

$$\phi(x) = \int_{\mathbb{R}^n} \frac{\rho(y)}{4\pi \|x-y\|_2} dy$$

$$\phi_j = \sum_{i=1}^N \frac{m_j}{4\pi R_{ij}}$$

# For Example: the Gravitational Problem

$$f(x) = \nabla \phi$$

$$f(x) = - \sum_{j=1, j \neq i}^N \nabla_{x_i} \left( \frac{\rho_j v_j}{4\pi \|x_i - x_j\|_2} \right)$$

$$f_i = - \sum_{j=1, j \neq i}^N \frac{\rho_j v_j (x_i - x_j)}{4\pi \|x_i - x_j\|_2^3}$$

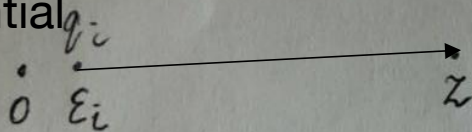


Given N particles and M evaluation position, direct computation requires  $O(NM)$  time!

# Introduction to Fast Summation

2D. Complex Number.  $\phi(z) = \text{Re}(\sum_j q_j \log(z - z_j))$

Consider a source and its potential

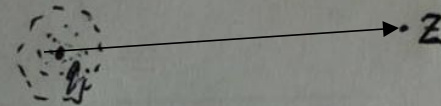


Apply Taylor expansion gives:

$$\begin{aligned}\phi(z) &= q_i \log(z - \epsilon_i) \\ &= q_i \log z - \sum_{k=1}^P \frac{q_i \epsilon_i^k}{k z^k}\end{aligned}$$

$\Rightarrow$

Multipole Expansion:



$$\begin{aligned}\phi(z) &= \sum_j q_j \log(z - z_j) \\ &= (\sum_j q_j) \log(z) - \sum_{k=1}^P \sum_j \frac{q_j z_j^k}{k z^k}\end{aligned}$$

$$Q = \sum_j q_j.$$

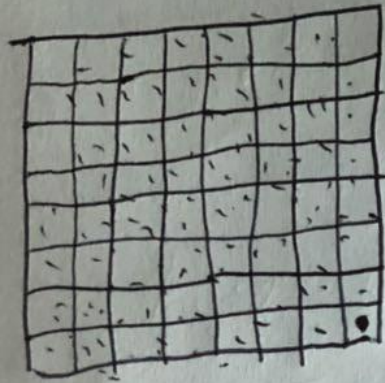
$$Q(k) = - \sum_j \frac{q_j z_j^k}{k}$$

$$\phi(z) = Q \log z + \sum_{k=1}^P \frac{Q_k}{z^k}$$

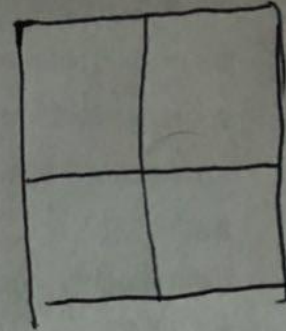
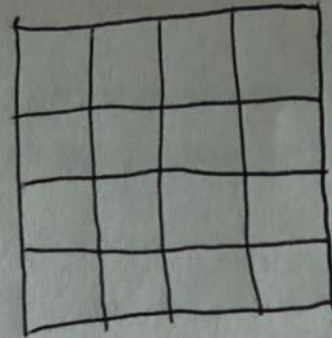


$N \log N$  Algorithm:

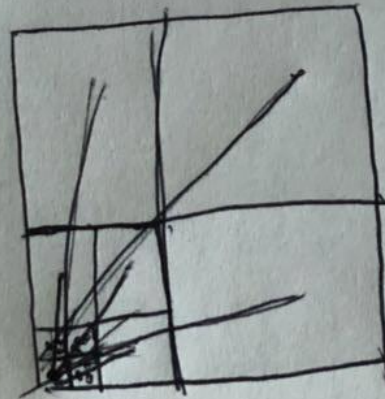
Tree Code



Compute  
 $Q_k$  for  
each level  $\rightarrow$

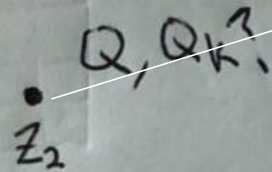
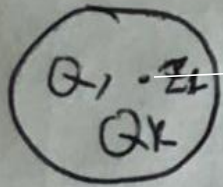


Eval:



Take one step further:

If we know M-Expansion at  $z_1$  ( $M_1 = \{z_1, Q, Q_k\}$ ),  
What is the M-Expansion at  $z_2$  ( $M_2 = \{z_2, Q_2, Q_{2,k}\}$ )?  
We want to obtain the coefficients from  $M_1$ , not  $q_i$ 's



$$\phi(z) = Q \log(z - z_2) + \sum_{k=1}^p \frac{b_k}{z^k}$$

Recall: 
$$\phi(z) = Q \log z + \sum_{k=1}^p \frac{Q_k}{z^k}$$

$b_k$  is a generalization of  $Q_k$ :

$$b_k = \underbrace{-\frac{Q(z_1 - z_2)^k}{k}}_{\text{Recall: } \frac{Q_j(z_j - c)^k}{k}} + \underbrace{\sum_{i=1}^k Q_i (z_1 - z_2)^{k-i} \binom{k-1}{i-1}}_{\text{rest of terms}}$$

Recall:  $\frac{Q_j(z_j - c)^k}{k}$

$$\phi(z) = Q \log(z - c) + \sum_{k=1}^p \frac{b_k}{(z - c)^k}$$



View Source as Multipole:

$$Q = q_i, \quad \text{~~Q_k = 0~~}$$

Reveal "Multipole Expansion"

$$\phi(z) = Q \log(z-c) + \sum_{k=1}^P \frac{\sum_j \frac{-q_j (z_j - c)^k}{k}}{(z-c)^k}$$

From "Multipoles":



compute  $b_k$  with "Rest of terms"

M2M Transform



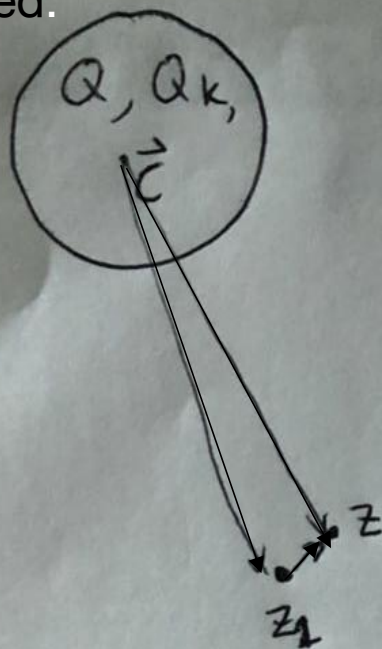
```
struct Multipole{
    vec2 center;
    complex q[p];
};
//source charge is a special Multipole,
//with center = charge pos, q0 = qi, q1...q4 =0
```

```
Multipole M2M(std::vector<Multipole> &qlist)
{
    Multipole res;
    res.center = weightedAverageof(qlist[i].center*qlist[i].q[0]);
    q[0] = sumof(qlist[i].q[0]);

    for(k=1:p){
        res.q[k]=0;
        for(j=0:qlist.size()) {
            res.q[k] += computeBk(qlist[i]);
        }
    }
    return res;
}
```

Question:

If we know M-Expansion at  $c_1$  ( $M_1 = \{c_1, Q, Q_k\}$ ),  
What is the polynomial at  $z_1$ , so that potentials at  
neighbor  $z$  can be evaluated.



M2L

Transform

$$\phi(z) = Q \log(z-c) + \sum_{k=1}^P \frac{b_k}{(z-c)^k}$$

$$= Q \log(z_1 - c + z - z_1) + \sum_{k=1}^P \frac{b_k}{(z_1 - c + z - z_1)^k}$$

$$= \underbrace{Q \log(z_1 - c) + \sum_{k=1}^P \frac{b_k}{(z_1 - c)^k}}_{\phi(z_1)}$$

+ H.O.T

$$\text{H.O.T.} = \sum_{l=1}^P b_l (z - z_1)^l$$

$$b_l = -\frac{Q}{l(c-z_1)^l} + \frac{1}{(c-z_1)^l} \sum_{k=1}^P \frac{b_k}{(z_1-c)^k} \binom{l+k-1}{k-1}$$

```

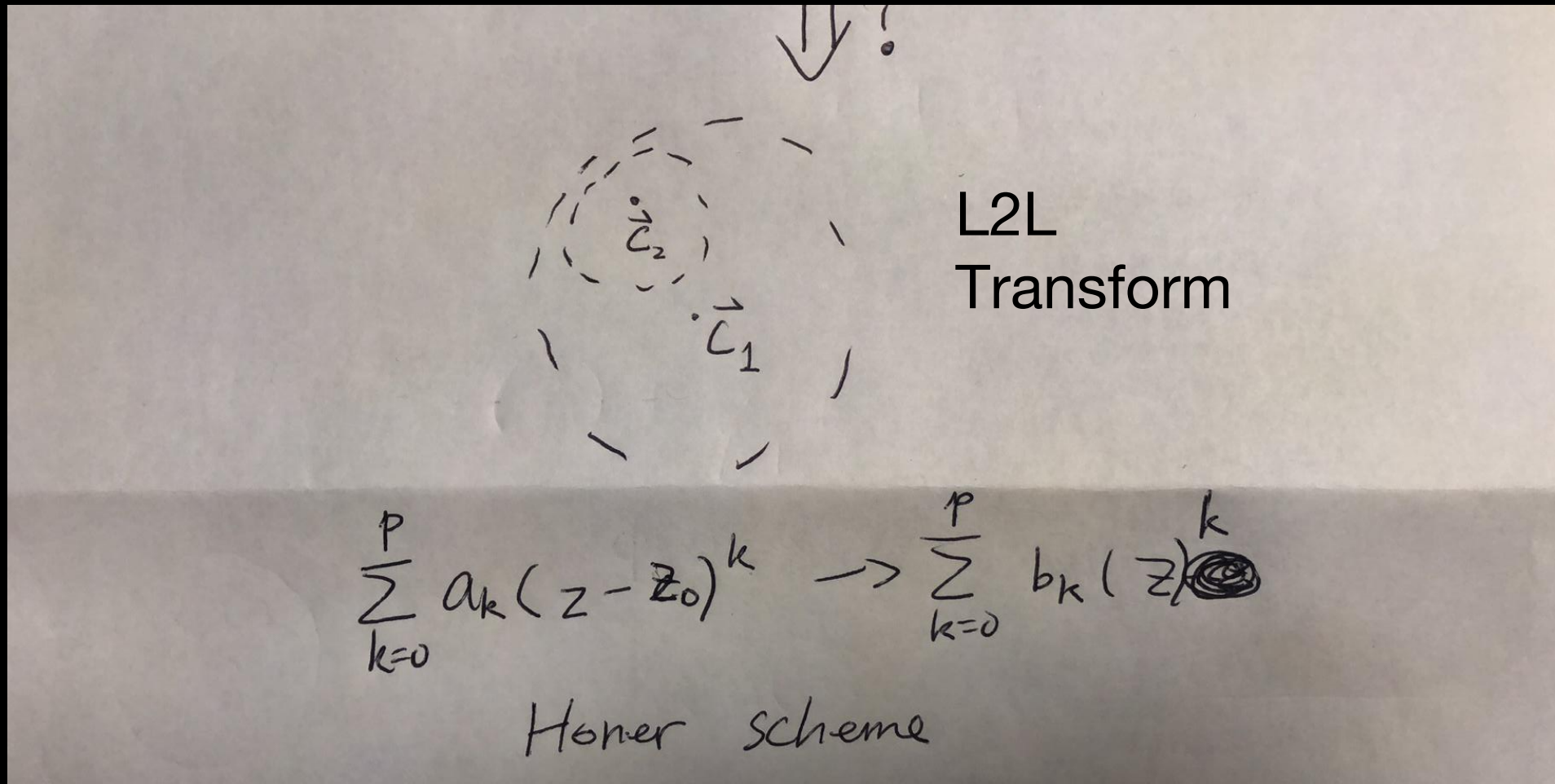
struct Localpole{
    vec2 center;
    complex b[p];
};

```

Question:

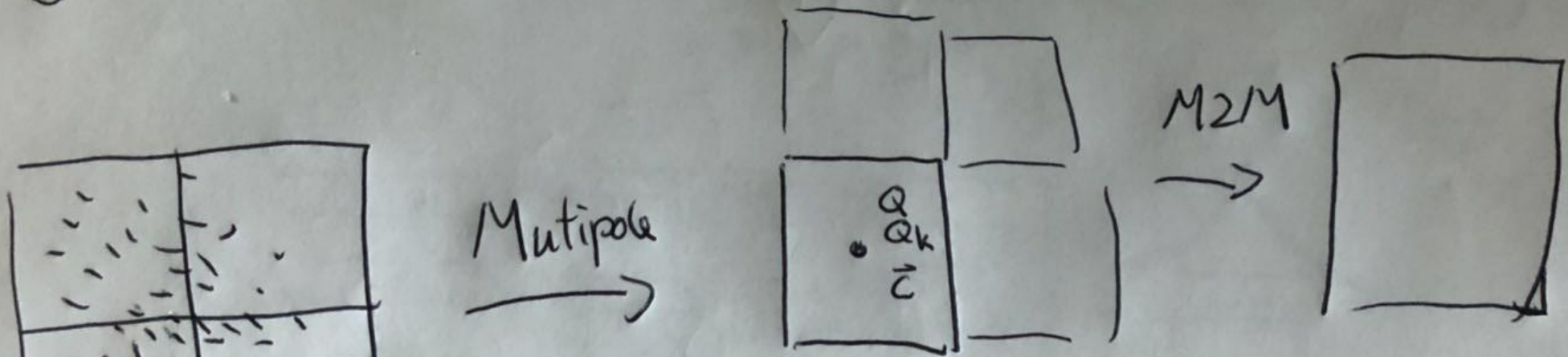
If we know L-Expansion at  $c_1$  ( $L_1 = \{c_1, B\}$ ),

What is the polynomial at  $c_2$ , so that potentials at neighbor  $z$  around  $c_2$  can be evaluated.



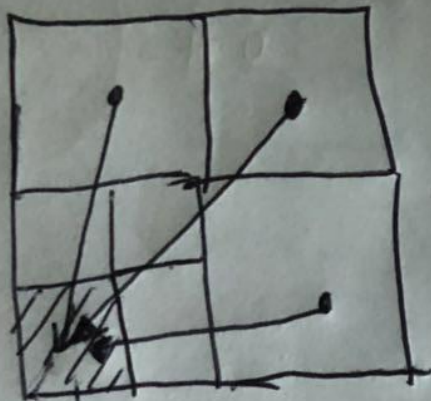
Multipole Expansion :      Coarsening  
Localpole Expansion :    Interpolation

$O(N)$  Algorithm:



$$\sum_l c^p \frac{N}{4^l} = O(c^p N)$$





M2L



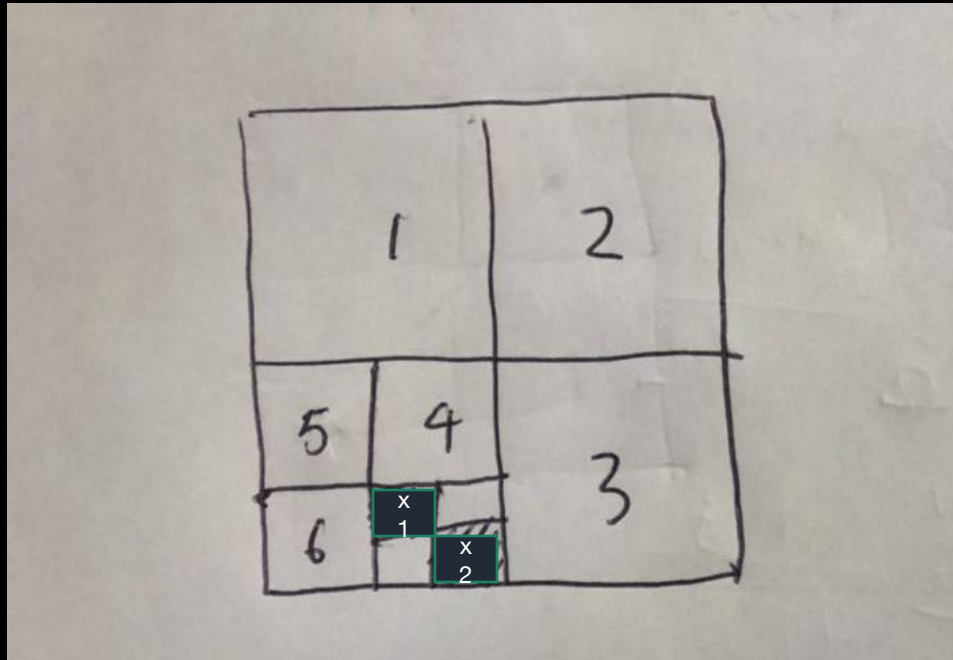
L2L + M2L

$$\sum_{\ell} c^p \frac{N}{4^{\ell}} = O(c^p N)$$

Tree code

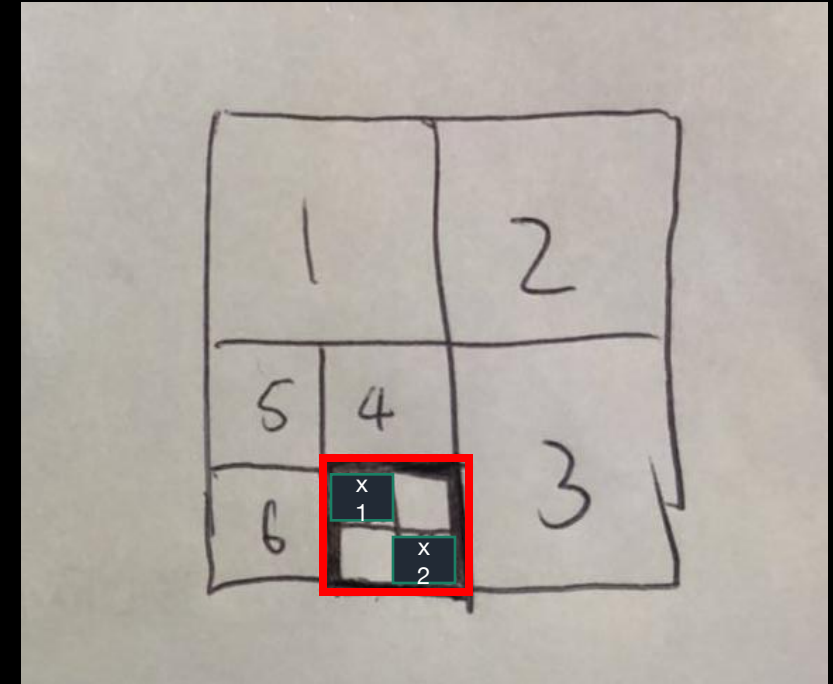
VS

FMM



$\Phi(x_1)$  = contribution from (node1, 2, 3, 4, 5, 6)

$\Phi(x_2)$  = contribution from (node1, 2, 3, 4, 5, 6)



= contribution from (node1, 2, 3)

$\Phi(x_1)$  = L2L from  + contribution from (4, 5, 6)

$\Phi(x_2)$  = L2L from  + contribution from (4, 5, 6)

In 3D, the algorithm can be obtained via:

- Taylor expansion of the Green's function in Cartesian system
- Taylor expansion of the Green's function in Spherical coordinates

For more details, checkout:

[https://math.nyu.edu/faculty/greengar/shortcourse\\_fmm.pdf](https://math.nyu.edu/faculty/greengar/shortcourse_fmm.pdf)

# Many, Many important applications

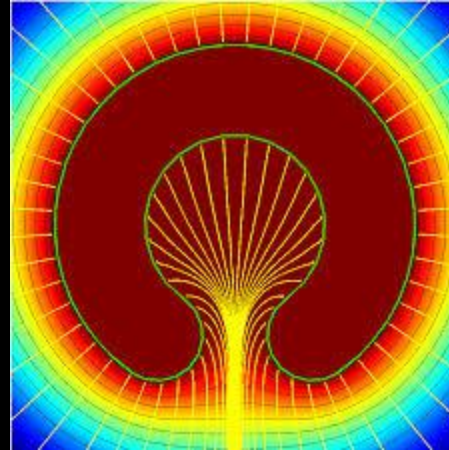
- Gravitational Force → Dark matter, cosmology...



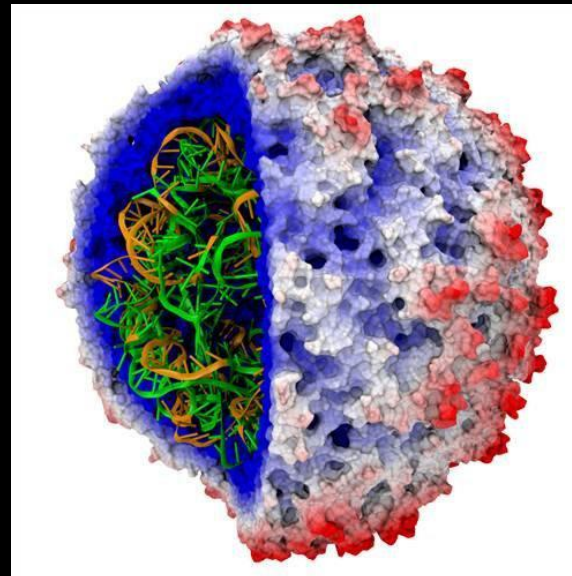


# Electrostatic

$$\nabla^2 \phi = \epsilon \rho$$
$$f = \nabla \phi$$



Major force  
between  
moleculars!!

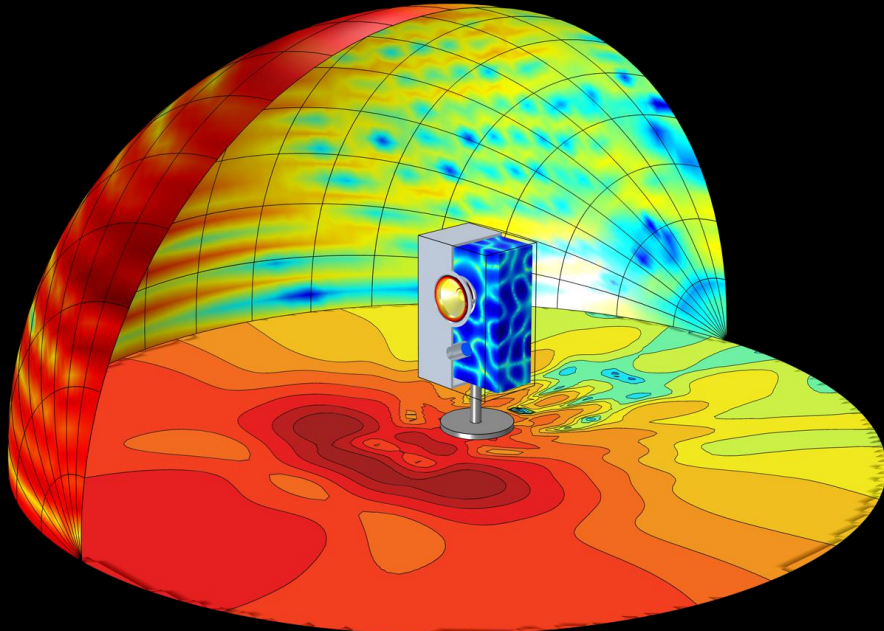


Cancer research  
Drug Design  
Virus Analysis

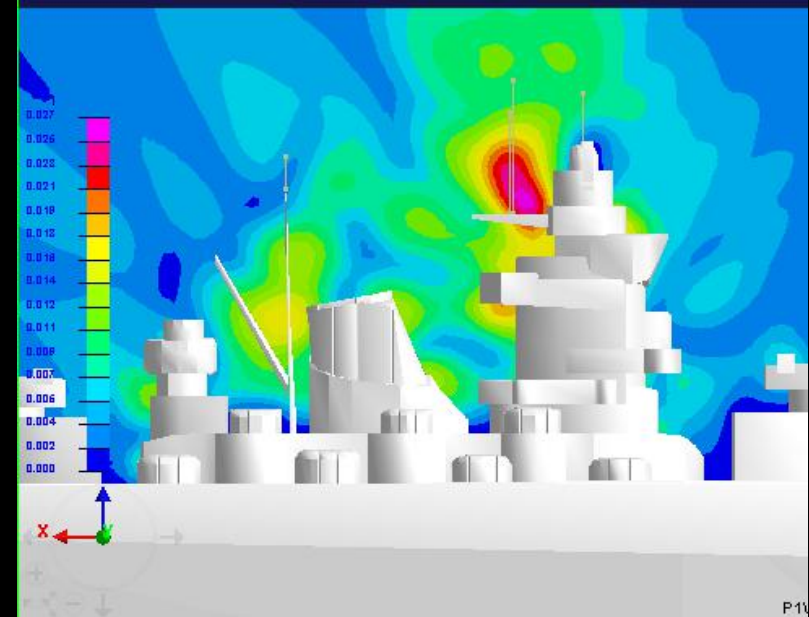
# Helmholtz

$$\nabla^2 \phi + k^2 \phi = -f$$

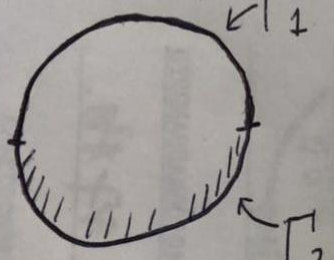
Acoustics



Electromagnetic



# Boundary Element Method



$\nabla^2 \phi = 0$   
 $q = \frac{\partial \phi}{\partial n} = f$  on  $\Gamma_2$   
 $u = \phi = g$  on  $\Gamma_1$   
 Solve for  $u$  on  $\Gamma_2$  and  $q$  on  $\Gamma_1$ , s.t.

$$\frac{1}{2} u_i + \sum_{j=1}^N u_j \vec{n}_j \cdot \nabla \frac{1}{r_{ij}} dA_j = \sum_{j=1}^N q_j \frac{1}{r_{ij}} dA_j$$

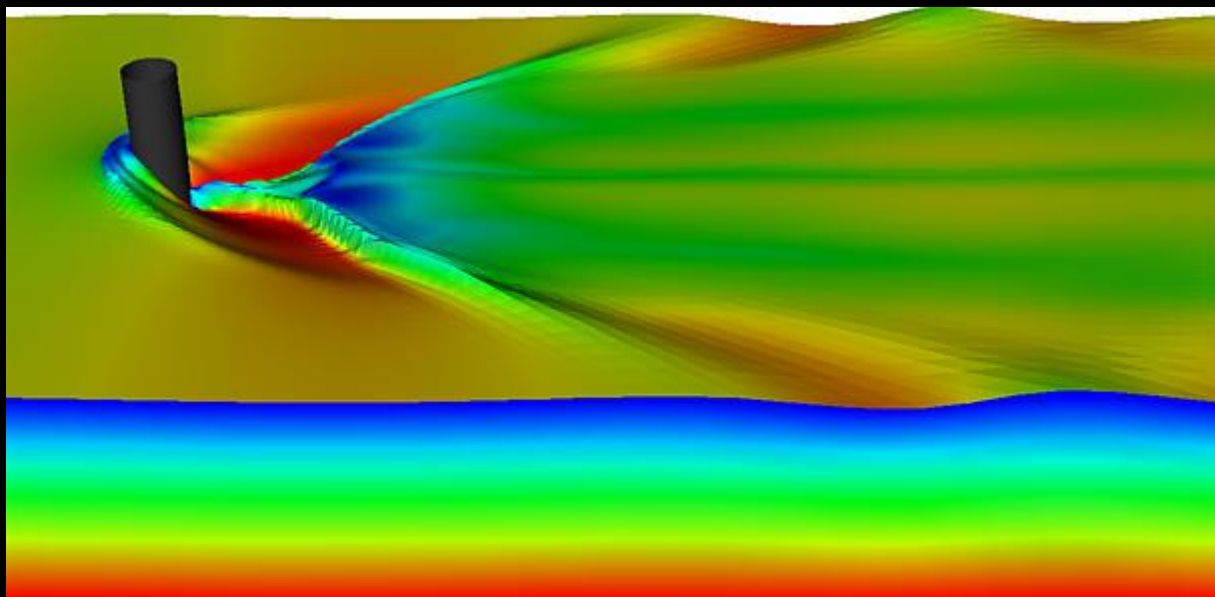
PDE  $\rightarrow$  Integral Equations

- Matrix is Dense!
- Condition number is Good!(usually converge in  $O(1)$  iterations)

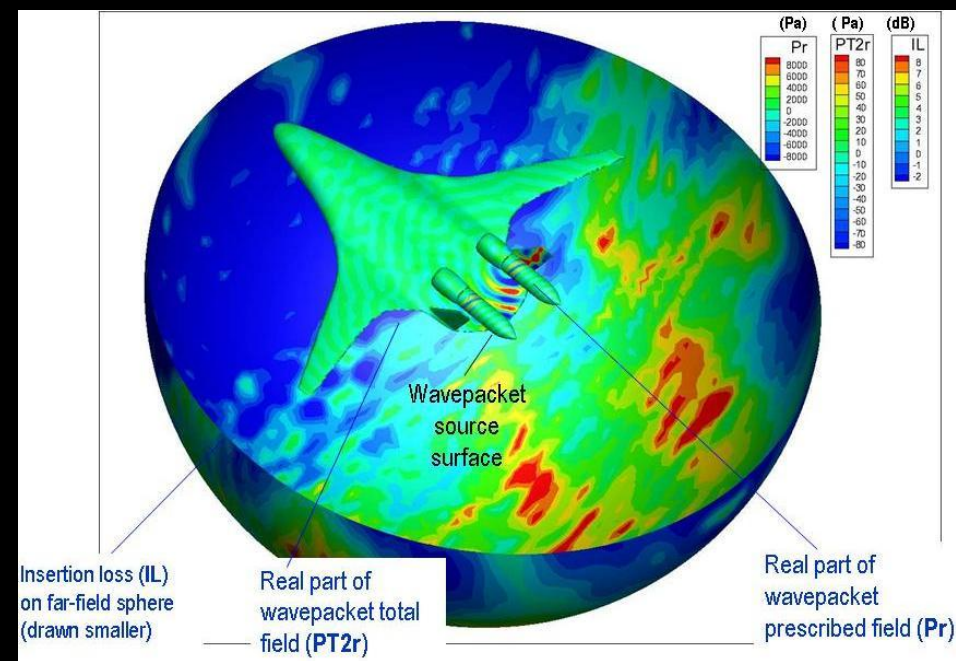
# Boundary Element Method

- In 2D
  - Full Domain is  $N^2$ , with Multigrid Methods,  $O(N^2)$  computation.
  - Boundary element has  $N$  elements, BiCGSTAB method converges in constant iteration, and each iteration took  $O(N^2)$  for Dense Matrix-vector!
- In 3D
  - Full Domain is  $N^3$ , with Multigrids,  $O(N^3)$
  - Boundary element has  $N^2$  elements, in total  $N^4$  computation!
- With FMM replacing the matrix vector multiplication operation
  - $O(N)$  in 2D
  - $O(N^2)$  in 3D.
  - Semi-Analytic!





Large scale, deep  
wave



Sound Pressure

# Vortex flow

$$u_i = \sum_{j=1, j \neq i}^N \frac{v_j \omega_j \times (x_i - x_j)}{4\pi \|x_i - x_j\|_2^3}$$



$$\begin{aligned}\nabla^2 \psi &= -\omega \\ u &= \nabla \times \psi\end{aligned}$$



# Type of summations

$$\nabla \phi(x) = \sum_{j=1}^N \frac{m_j \vec{r}_{ij}}{R_{ij}^3}$$

$$\text{B.I.E : } \sum_{j=1}^N \vec{n}_j \cdot \nabla \left( \frac{1}{R_{ij}} \right) \sigma_j A_j$$

$$\text{Biot-Savart } \sum_{j=1}^N \frac{\omega_j \times \vec{r}_{ij}}{R_{ij}^3} dV_j$$

# Compute them with same routine

- Given Routine:
  - $F = \text{computeGradPhi}(s)$ ; //returns gravity force using FMM from many source  $s$
- B.I.E:
  - Let  $s1[i] = n[i].x * s[i]$ ,  $s2[i] = n[i].y * s[i]$ ,  $s3[i] = n[i].z * s[i]$ ;
  - $F1 = \text{computeGradPhi}(s1)$ ,  $F2 = \text{computeGradPhi}(s2)$ ,  $F3 = \text{computeGradPhi}(s3)$
  - $\text{Res} = F1.x + F2.y + F3.z$ ;
- Biot-Savart:
  - Your homework



# Other fast summation methods:

- PPPM

# PPPM: Combining PDE form and summation forms

$$u_i = \sum_{j=1, j \neq i}^N \frac{v_j \omega_j \times (x_i - x_j)}{4\pi \|x_i - x_j\|_2^3}$$



$$\begin{aligned}\nabla^2 \psi &= -\omega \\ u &= \nabla \times \psi\end{aligned}$$

$$f_i = -\epsilon \sum_{j=1, j \neq i}^N \frac{\rho_j v_j (x_i - x_j)}{4\pi \|x_i - x_j\|_2^3}$$



$$\begin{aligned}\nabla^2 \phi &= \epsilon \rho \\ f &= \nabla \phi\end{aligned}$$

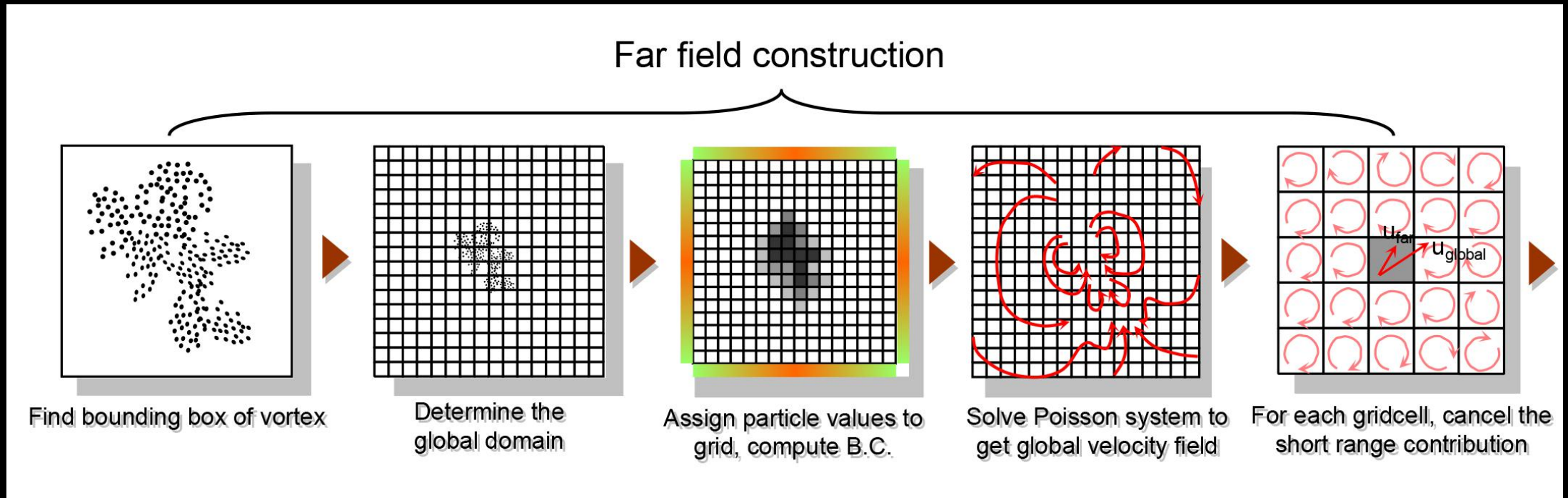
Direct summation for  
the turbulent part



Poisson's Equation  
for the smooth part

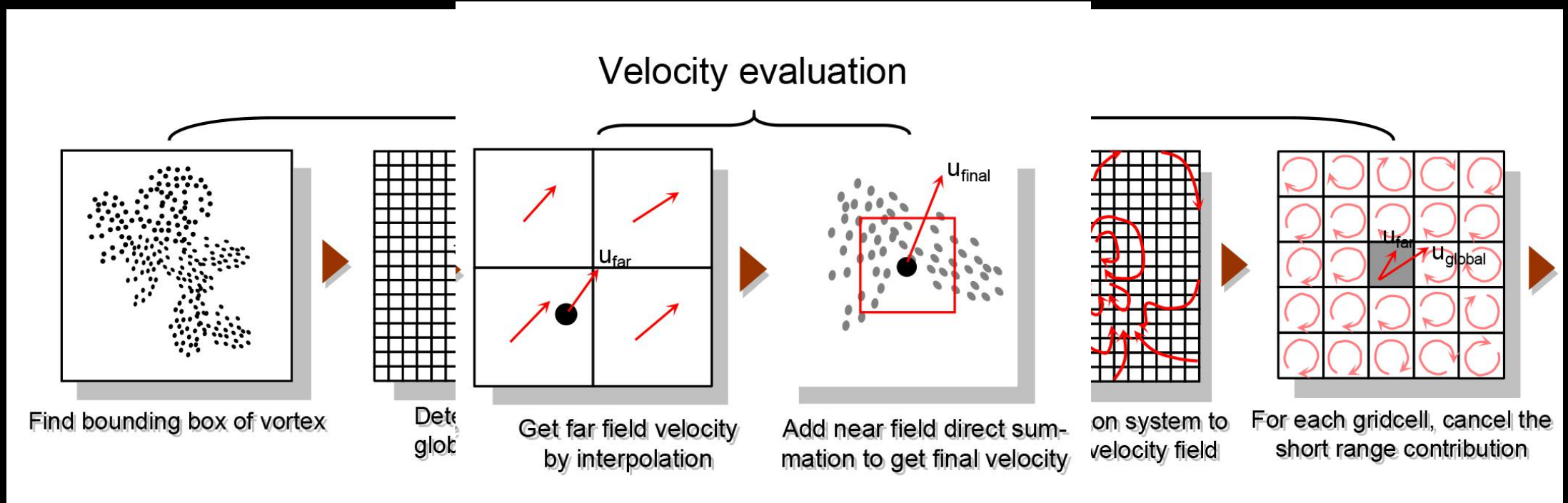
# PPPM

- Fast solution uses near-far decomposition to get acceleration. Can we do similar thing on a particle-mesh setup?

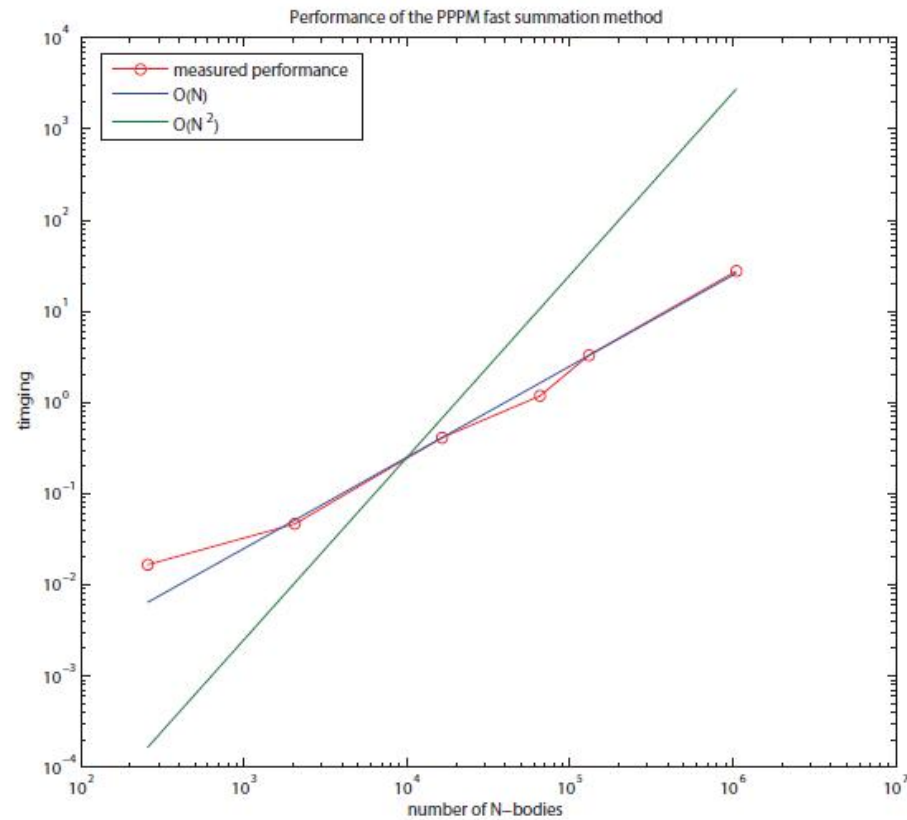


# PPPM

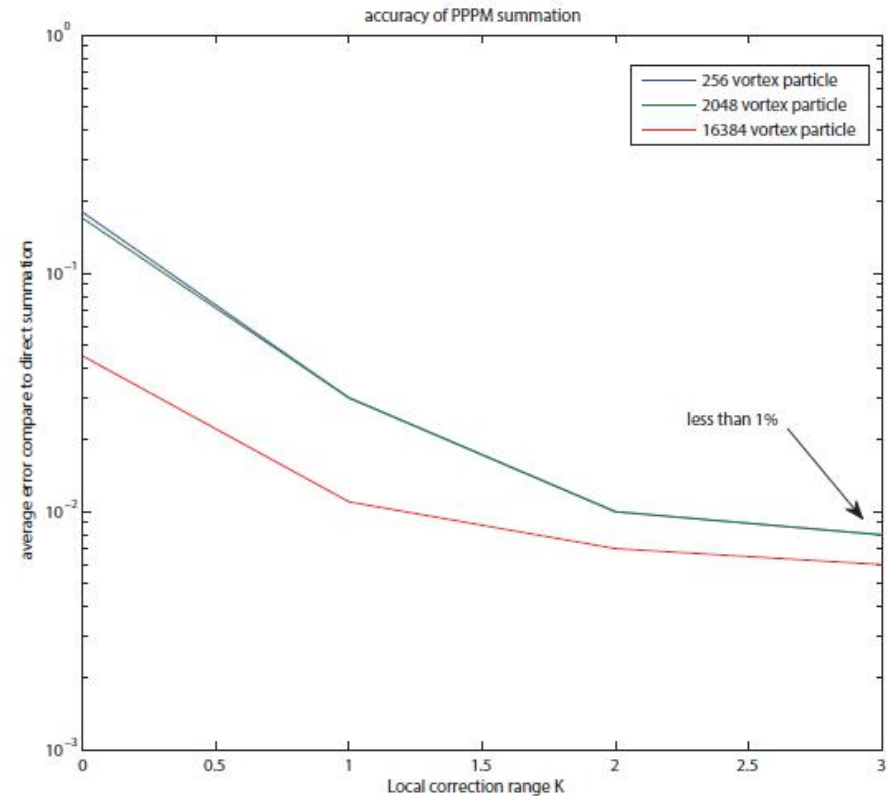
- Fast solution uses near-far decomposition to get acceleration. Can we do similar thing on a particle-mesh setup?



# PPPM



**Figure 4:** Performance of the PPPM fast summation. Computation time grows linearly with the number of computational elements.



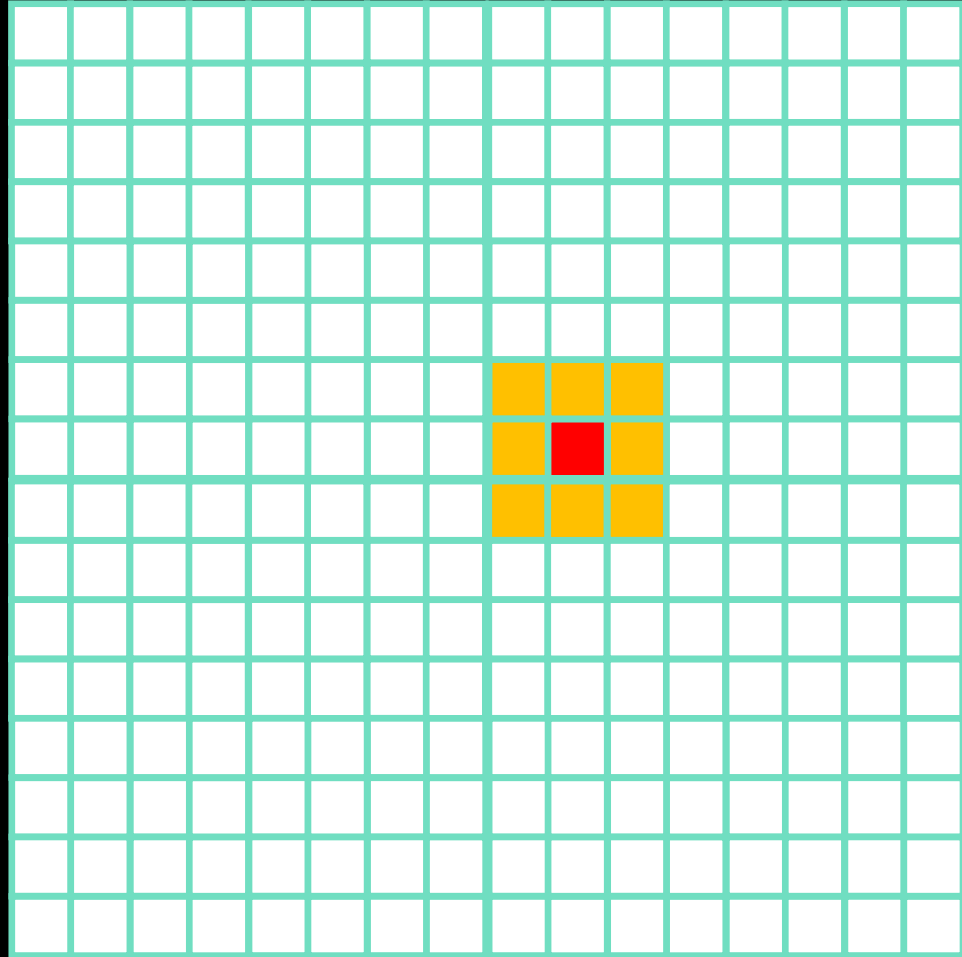
**Figure 5:** Accuracy statistics of the PPPM fast summation.



# Local Correction

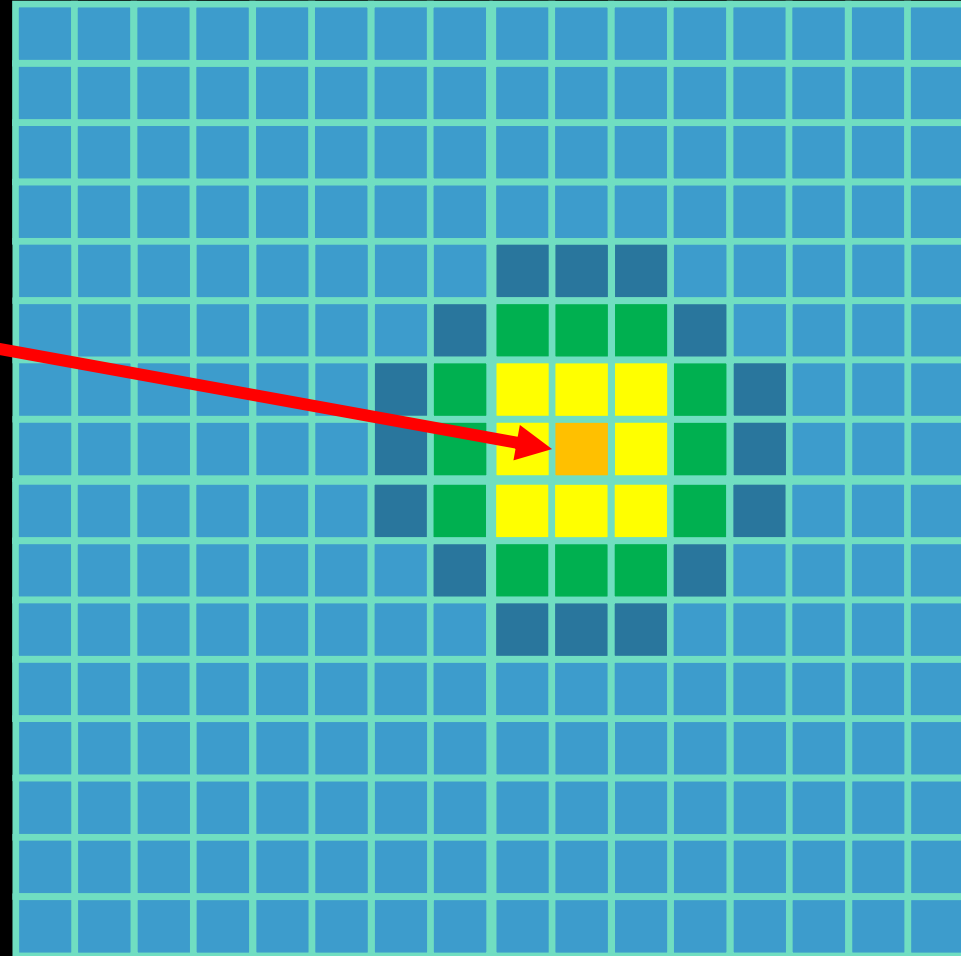
- In 3D, for a correction window of size  $K$  in each dimension, a local matrix of size  $K^3 \times K^3$  can be precomputed to cancel the local influence from grid.
- $T(N) = O(c K^6 N)$

# Local Correction



# Local Correction

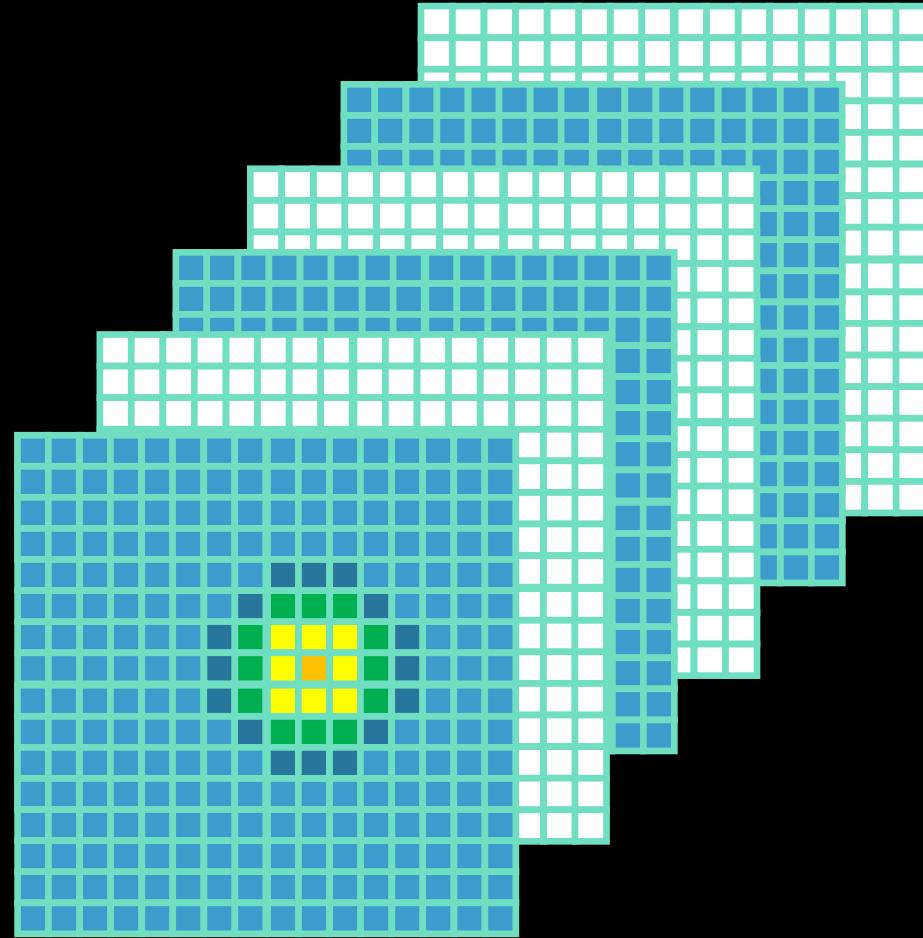
The influence made by  
neighbor cells.



# Local Correction

- *The matrix inverse reveals how the center cell's value depends linearly on its neighbors (including itself).*

$$s_c = \sum_{j \in \eta} a_j r_j$$



# PPPM in few lines

- `w_bar = particle_to_grid(w_p);`
- `dw = w_p - interpolate(w_bar);`
- `Psi = Poisson.Solve(w_bar);`
- `v_smooth = curl(Psi);`
- `v_p = interpolate(v_smooth) + nearSum(dw);`



# Summarize

- Fast Summation Methods
  - FMM
  - PPPM
- Equations solved by Fast Summation Methods:
  - Poisson's Equation
  - Laplace Equation
  - Helmholtz Equation
  - BEM
- Applications of Fast Summation Methods
  - Electrostatic::Molecular Dynamics::Cancer, drug design research  
Magnetics::Ship design
  - Acoustics::Urban planning, vehicle shape design, theatre design
  - Potential flow::aircraft, wave
  - Vortex method::turbulent flow