

# Lattice regularization of non-relativistic interacting fermions in 1+1 dimensions

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# Model Introduction

- The continuum Hamiltonian of a  $(1+1)$  system<sup>†</sup> reads:

$$H = \sum_{\sigma=1}^2 -\psi_{\sigma}^{\dagger} \frac{\partial_x^2}{2m_a} \psi_{\sigma} + \lambda \psi_1^{\dagger} \psi_2^{\dagger} \psi_2 \psi_1 + c \left( \psi_1^{\dagger} i \overleftrightarrow{\partial}_x \psi_2^{\dagger} \right) \left( \psi_2 i \overleftrightarrow{\partial}_x \psi_1 \right). \quad (1)$$

- $\psi_1$  and  $\psi_2$  are two species of non-relativistic fermion particles.
- $\lambda$  is a parameter for an on-site interaction.
- $c$  is a parameter for a pairwise local velocity-dependent interaction.

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<sup>†</sup>S.Y Yong and D. T. Son (2018)

# Lattice Regularization

- Regulating a system on a lattice, we want to study few-body systems non-perturbatively.

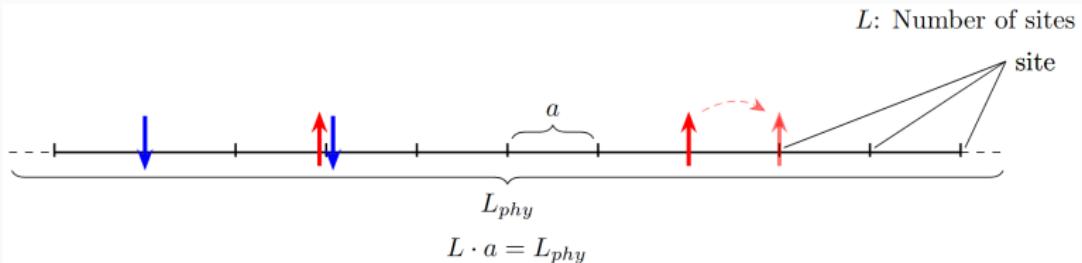


Figure 1: A graphical illustration of lattice theory.

- We are interested in whether we can recover the continuum theory at the continuum limit ( $a \rightarrow 0$  or  $L \rightarrow \infty$ ).

## System with Contact Interaction ( $c = 0$ )

- With only contact interaction, (also known as the Hubbard Model) the Hamiltonian is:

$$H_1 = H_0 + H_{\text{int}}^{(\lambda)}. \quad (2)$$

- $H_0 = \sum_{x,\sigma} \epsilon(a)(\psi_{\sigma,x}^\dagger \psi_{\sigma,x+\alpha} - 2\psi_{\sigma,x}^\dagger \psi_{\sigma,x} + \psi_{\sigma,x+\alpha}^\dagger \psi_{\sigma,x})$
- $H_{\text{int}}^{(\lambda)} = \epsilon(a)\lambda(a)\psi_{1,x}^\dagger \psi_{1,x} \psi_{2,x}^\dagger \psi_{2,x}$
- In the continuum limit,  $a\lambda(a)\epsilon(a)$  is constant.

# System with Contact Interaction ( $c = 0$ )

- We can show that, using exact diagonalization,  
 $\lim_{a \rightarrow 0} a\lambda(a)\epsilon(a) = -0.140966.$

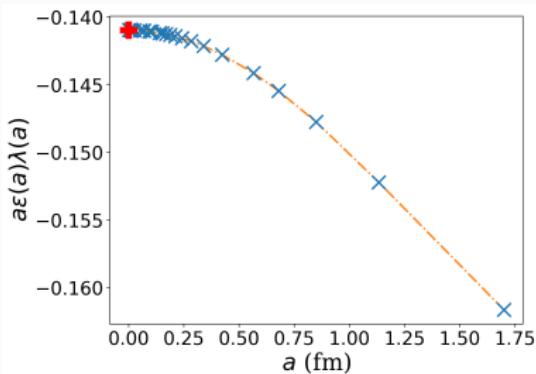


Figure 2: Interaction coefficient  $a\lambda(a)\epsilon(a)$  as a function of lattice spacing  $a$ .

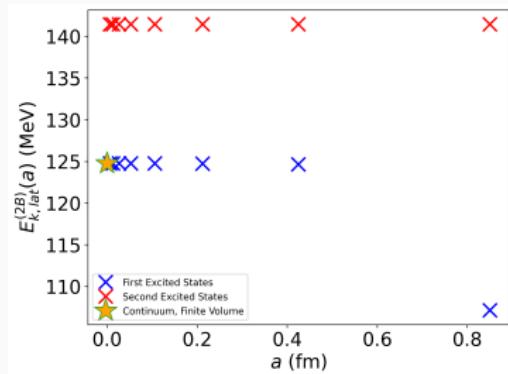
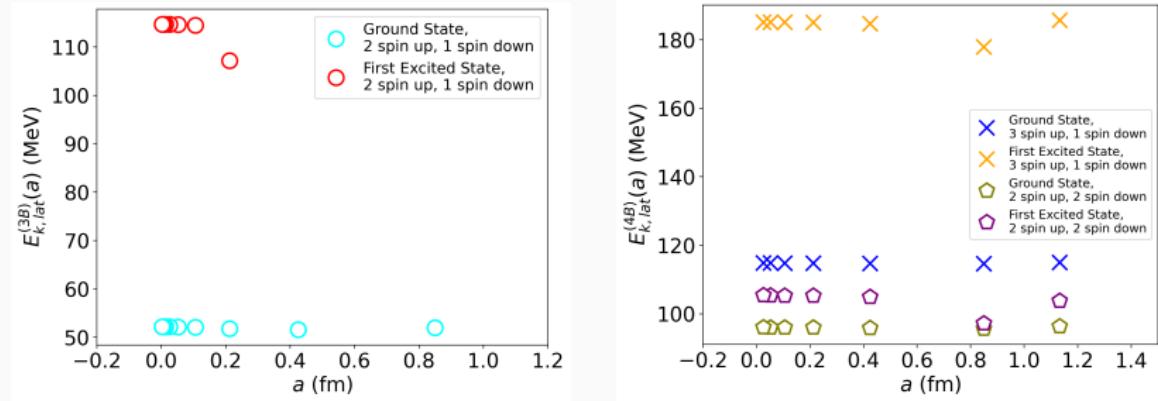


Figure 3: Comparison of energy levels generated by our method and the finite-volume calculation<sup>†</sup>.

<sup>†</sup>C. Körber, E. Berkowitz, and T. Luu (2019).

# System with Contact Interaction ( $c = 0$ )



**Figure 4:** Three- and four-body ground state and first excited state energy as functions of lattice spacing. They have well-defined continuum limits.

## System with Contact Interaction ( $c = 0$ )

- We present an infinite volume extrapolation method, fixing the lattice spacing  $a \leq 0.1\text{fm}$ , and increase site number,  $L$ .
- The interaction strength,  $U = -0.140966$  is used for delta potential calculation.

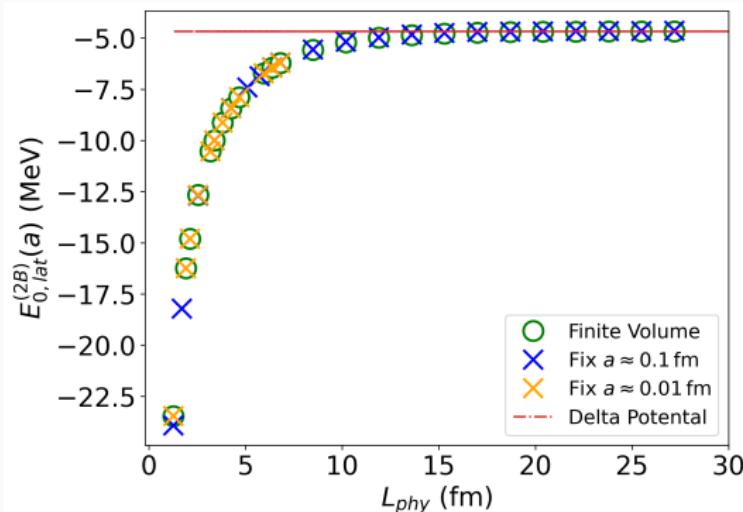


Figure 5: Comparison of different lattice spacing  $a$ .

# System with Momentum-Dependent Interaction ( $\lambda = 0$ )

- With only the momentum-dependent interaction, the Hamiltonian reads

$$H_2 = H_0 + H_{\text{int}}^{(c)}. \quad (3)$$

- $H_{\text{int}}^{(c)} = \epsilon c_2 \sum_x (\psi_{2,x+\alpha} \psi_{1,x} - \psi_{2,x} \psi_{1,x+\alpha})^\dagger (\psi_{2,x+\alpha} \psi_{1,x} - \psi_{2,x} \psi_{1,x+\alpha})$

- At the continuum limit,  
 $\lim_{a \rightarrow 0} c_2 = -1$ .

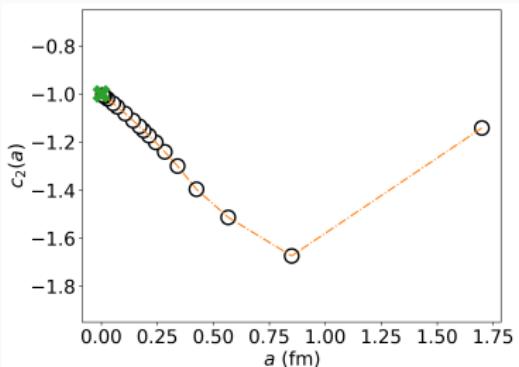
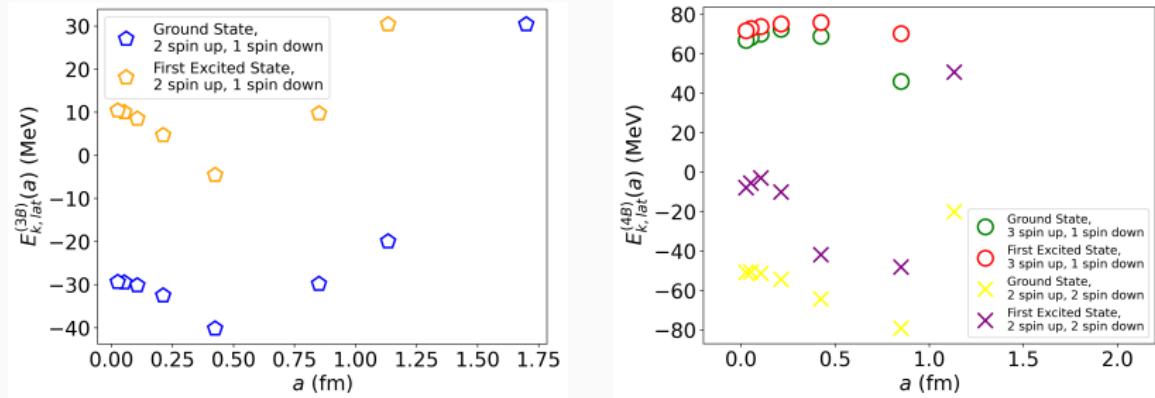


Figure 6: Plot of interaction strength  $c$  as a function of lattice spacing,  $a$ .

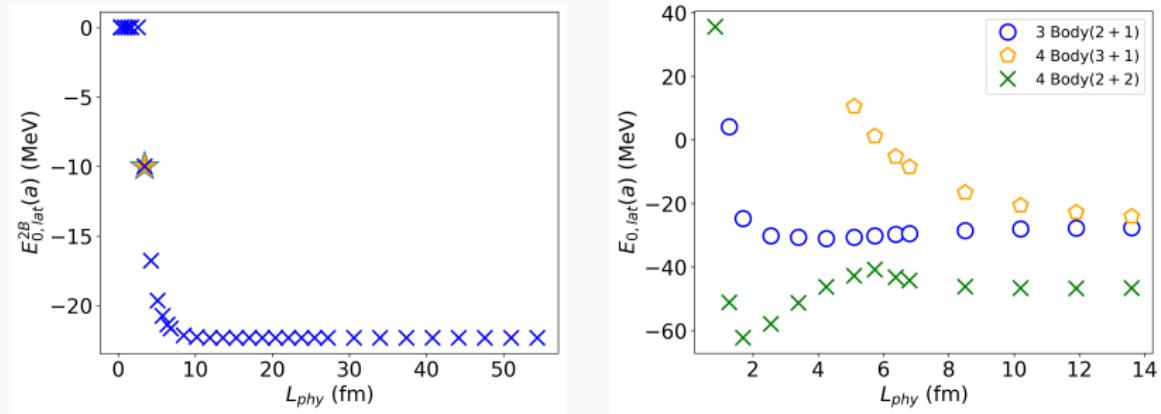
# System with Momentum-Dependent Interaction ( $\lambda = 0$ )

- Energy for three- and four-body systems can support further study of three-body force.



**Figure 7:** Three- and four-body ground state and first excited state energy as functions of lattice spacing.

# System with Momentum-Dependent Interaction ( $\lambda = 0$ )



**Figure 8:** Infinite volume extrapolation for two- (Left), three-, and four-body (Right) systems. at  $a \approx 0.1$  fm.  $E_{3B}/E_{2B} \approx 1.23$ .

- In the continuum effective theory calculation,  $E_{3B}/E_{2B} = 1.375$ .

## Summary and Outlook

- In this work, we study a lattice regularization of a continuum theory, and use exact diagonalization method to verify the consistency of two approaches at the continuum and infinite length limit.
- Further study the full Hamiltonian, including the renormalization of a three-body force and extend this method to many-body systems, using quantum Monte Carlo simulations.
- This method can be applied on more general mass-spin imbalanced system.

Thank you!

# Fixing Energy Scales

- The discretized version of a free Hamiltonian would look like

$$H_0 = -\frac{1}{2m_a\epsilon^2}(\psi_{\sigma,x}^\dagger\psi_{\sigma,x+\alpha} - 2\psi_{\sigma,x}^\dagger\psi_{\sigma,x} + \psi_{\sigma,x+\alpha}^\dagger\psi_{\sigma,x}) \quad (4)$$

- Notice that  $m_a$  here is not equal to the physical mass,  $m_{phy}$ , and depends on the lattice spacing  $a$ .
- Energy of a free Hamiltonian can be solved using exact diagonalization method. The energy scale  $\epsilon(a) = -\frac{1}{2m_a\epsilon^2}$  is fixed by comparing to the ground state energy of a one-dimensional infinite square well problem.

$$2\epsilon(a) \left[ 1 - \cos\left(\frac{2\pi}{L}\right) \right] = \frac{2\pi^2}{m_{phy}L_{phy}^2} \quad (5)$$

# System with Momentum-Dependent Interaction ( $\lambda = 0$ )

- With only contact interaction, the Hamiltonian reads

$$H_2 = H_0 + H_{\text{int}}^{(c)} \quad (6)$$

- $H_0 = \sum_{x,\sigma} \epsilon(a) (\psi_{\sigma,x}^\dagger \psi_{\sigma,x+\alpha} - 2\psi_{\sigma,x}^\dagger \psi_{\sigma,x} + \psi_{\sigma,x+\alpha}^\dagger \psi_{\sigma,x})$
- $H_{\text{int}}^{(c)} = \epsilon(a) c \sum_x (\psi_{2,x+\alpha} \psi_{1,x} - \psi_{2,x} \psi_{1,x+\alpha})^\dagger (\psi_{2,x+\alpha} \psi_{1,x} - \psi_{2,x} \psi_{1,x+\alpha})$
- We introduce a relative momentum  $q$ , and relative position,  $\delta$  to reduce to a one-body system.

$$|q, \delta\rangle = \frac{1}{\sqrt{L}} \sum_x \psi_{2,x}^\dagger \psi_{1,x+\delta}^\dagger e^{i(2\pi/L)qx} |0\rangle \quad (7)$$

- Define  $I_{\text{lat}}^{(c)} = 2\epsilon(a)\delta = 0(E - H_0)^{-1}\delta = 0$  that appears in the denominator of the first-order expansion of Green function.
- At pole,

$$\lim_{a \rightarrow 0} c_2(a) = 2/I_0^{(c)} = -2 / \int_0^{2\pi} \frac{dp}{2\pi} \frac{(1 - \cos 2p)}{1 - \cos p} = -1. \quad (8)$$