

# Fluorescence Emission Study on Plastic Scintillator EJ-200 After Irradiation

Alberto Belloni<sup>a</sup>, Sarah Eno<sup>a</sup>, Yongbin Feng<sup>a</sup>, Charles Hurlbut<sup>b</sup>, Aaron Hunt<sup>a</sup>,  
Geng-Yuan Jeng<sup>a</sup>, Zachary Thomas<sup>a</sup>, Yao Yao<sup>a</sup>, Zishuo Yang<sup>\*a</sup>

<sup>a</sup>*Department of Physics, University of Maryland, College Park, MD 20740, USA*

<sup>b</sup>*Eljen Technology, 1300 W. Broadway, Sweetwater, TX 79556, USA*

---

## Abstract

Fluorescence light output versus wavelengths for plastic scintillator EJ-200 with varying dopant concentrations is measured before and after irradiations of various doses and dose rates. Dose, dose rate, and temperature effects on the scintillator's dopants during irradiations are investigated.

*Keywords:* plastic scintillator, fluorescence, radiation damage, spectrophotometer

---

## 1. Introduction

Organic plastic scintillators made of polystyrene (PS) or polyvinyltoluene (PVT) substrate and wavelength-shifting dopants have long been popular in detectors used in particle physics, nuclear physics, radiation safety, and health physics applications due to their high light output, low cost, fast response, and versatility of physical construction. Prolonged exposure of plastic scintillator to ionizing radiation, however, can result in damage such that absorption in the substrate increases and the transfer efficiency of the initial excitation of the substrate to the dopants combined with a dopant's probability of radiative relaxation lessens. In this paper, we present fluorescence emission measurements before and after irradiations of various doses and dose rates for one type of plastic scintillator manufactured by Eljen Technology, EJ-200 (similar to BC-408 from Bicron) with different concentrations of primary and secondary dopants. Dose rate effects are of interest because material-testing is typically done at much higher dose rates than that which the scintillator will experience in use, due to reactor-time cost. Dose rate effects may be responsible for the factor-of-three-larger light loss with integrated luminosity in the Compact Muon Solenoid (CMS) endcap hadron calorimeter [1, 2] than expected based on high dose rate exposure using <sup>60</sup>Co sources [3, 4].

The effect of dose and dose rate on the radiation-induced increase of absorption have been the subject of many investigations [5]. The presence of oxygen plays an important role during and after irradiation. Studies show that the penetration depth of oxygen into the substrate depends on the dose rate: at

lower dose rates, oxygen penetrates more deeply [5, 6, 7]. Because of the importance of the interaction of oxygen with radicals produced in the substrate during irradiation, this can lead to dose rate effects. Two types of interactions occur in this process. The presence of oxygen increases the number of migration mechanisms for the radicals produced during irradiation. This allows the radicals that create color centers which absorb light to migrate, find other radicals, and reform chemical bonds that are not color centers [8]. However, oxygen can also interact with the radicals produced during irradiation in a way that forms additional color centers [9].

Studies on the effect of radiation damage to the light output are fewer than those on the absorption. Several studies indicate that damage is to the substrate or the transfer of substrate's emission to dopants, rather than on the dopants themselves. Some studies have indicated that while oxygen plays a beneficial role in regards to annealing of induced absorption effects after irradiation, it plays a detrimental role in regards to light output[10].

In order to understand the relative role of the destruction of dopants versus damage to the substrate for modern plastic scintillators, we have studied the light output for plastic scintillator EJ-200 with various dopant concentrations, for different total doses and dose rates.

## 2. Sample and Irradiation

EJ-200 uses PVT as a substrate. It has a light output that is 60% of anthracene and a maximum-emission wavelength of 425 nm (blue). Eljen Technology prepared samples with dimensions of 1x1x5 cm<sup>3</sup>. The samples were made with concentrations of primary and secondary dopants at 1.0 and 2.0 that of the nominal concentration (which has optimized light output before irradiation). 1P and 2P refer to nominal and double the nominal concentration of primary dopant, while 1X and 2X refer to nominal and double the nominal concentration of secondary dopant. The samples for irradiation were wrapped in Tyvek sheets and were in contact with air. Using a <sup>60</sup>Co source at the National Institute of Standard and Technology (NIST), some 1P and 2P samples were irradiated at rates of 9 krad/hr and 429 krad/hr for a total dose of 3 Mrad, at about 23 °C ambient temperature; One 1P and one 2P sample were irradiated at 9 krad/hr for a total dose of 4 Mrad, at about 23 °C ambient temperature; Two samples that previously received 3 Mrad at 9 krad/hr at NIST were irradiated again at 9 krad/hr for an additional dose of 4 Mrad; Two 1X samples were irradiated at 80 krad/hr for 6 Mrad, at different temperatures 23 °C and -30 °C. Using the LHC beamline irradiation facility CASTOR, where neutrons and gamma rays contribute the most to ionizing radiation, one 1X and one 2X sample were irradiated at 0.84 krad/hr for a total dose of 1.68 Mrad. Additional samples that were produced in the same batch were kept non-irradiated for calibration and comparison. Table 1 summarizes the samples irradiated in this study.

Table 1: Summary of irradiated EJ-200 samples. 1P and 2P refer to nominal and double the nominal concentration of primary dopant, while 1X and 2X refer to nominal and double the nominal concentration of secondary dopant.

Sample	Dose (Mrad)	Dose rate (krad/hr)	Irradiation type
1P	$2.95 \pm 0.05$	$9.48 \pm 0.16$	$^{60}\text{Co}$
2P	$2.95 \pm 0.05$	$9.48 \pm 0.16$	$^{60}\text{Co}$
1P	$2.95 \pm 0.05$	$429 \pm 7$	$^{60}\text{Co}$
2P	$2.95 \pm 0.05$	$429 \pm 7$	$^{60}\text{Co}$
1P	$4.00 \pm 0.07$	$9.20 \pm 0.16$	$^{60}\text{Co}$
2P	$4.00 \pm 0.07$	$9.20 \pm 0.16$	$^{60}\text{Co}$
1P	$(2.95+4.00) \pm 0.09$	$9.32 \pm 0.11$	$^{60}\text{Co}$
2P	$(2.95+4.00) \pm 0.09$	$9.32 \pm 0.11$	$^{60}\text{Co}$
1X	$1.68 \pm 0.17$	$0.84 \pm 0.17$	CASTOR
2X	$1.68 \pm 0.17$	$0.84 \pm 0.17$	CASTOR
1X	$5.82 \pm 0.10$	$80.3 \pm 1.4$	$^{60}\text{Co}$ , at $23^\circ\text{C}$
1X	$5.82 \pm 0.10$	$80.4 \pm 1.4$	$^{60}\text{Co}$ , at $-30^\circ\text{C}$

### 3. Instrumentation and Techniques

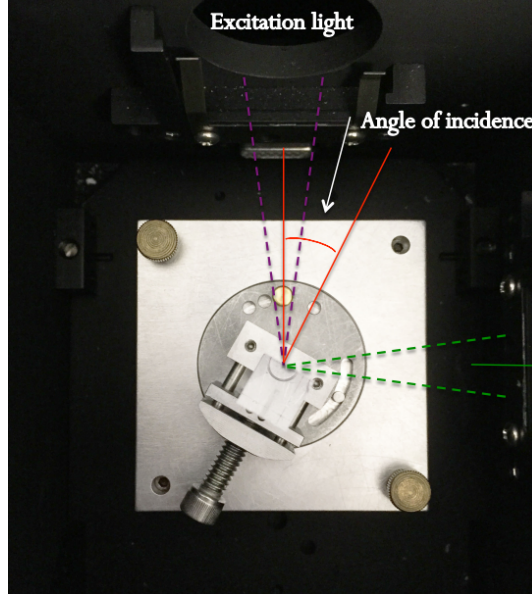
The fluorescence emission spectra were measured using a FluoroMax-Plus fluorescence spectrometer by Horiba Scientific. The spectrometer uses a Xenon arc lamp as a continuous light source. Light from the lamp is collected by a diamond-turned elliptical mirror and focused on the entrance slit of the excitation monochromator. The entrance and exit slits of the excitation monochromator are continuously adjustable, which controls the bandpass of the excitation beam. A toroidal mirror focuses the excitation beam onto the sample, while a small amount of the beam is split off by a beam-splitter to the reference silicon photodiode. Fluorescence emission produced by the sample is subsequently collected at  $90^\circ$  angle with respect to the excitation beam. Similar to the excitation monochromator, the emission monochromator selects the bandpass of emission light that is sent to a Hamamatsu R928 photomultiplier tube (PMT). After correcting the signals of the PMT and the photodiode by their quantum efficiencies, the fluorescence emission signal is calculated as the ratio of the PMT's signal over the photodiode's signal. For each fluorescence spectrum measurement, a fixed excitation wavelength is used, and the emission intensity at a range of wavelengths is measured [11].

In order to study the effect of radiation damage to primary and secondary dopants, three excitation wavelengths were chosen based on absorption and emission spectra provided by Eljen Technology [12]: 285 nm (exciting PVT and primary dopant), 350 nm (exciting mainly secondary dopant and non-negligible primary dopant), and 400 nm (exciting almost exclusively secondary dopant).

Because the measurements are very sensitive to slight position changes relative to the excitation beam, the sample holder was customized to ensure reproducibility. Initialization of the instrument includes three hours of warm-up time, during which the spectrophotometer takes blank measurements. The sam-

ple placement in the spectrometer adopts the front-surface configuration (shown in Figure 1), which gives the excitation beam a  $60^\circ$  incident angle and maximizes fluorescence emission at the surface. This minimizes the effect of self-absorption when measuring the fluorescence emission after irradiation. Continuous measurements after irradiation for up to 7 days have shown no significant change in fluorescence emission spectrum, while more than 15% recovery in the bulk transmission spectrum is seen (Appendix A).

Figure 1: Front-surface configuration for fluorescence spectrum measurements.

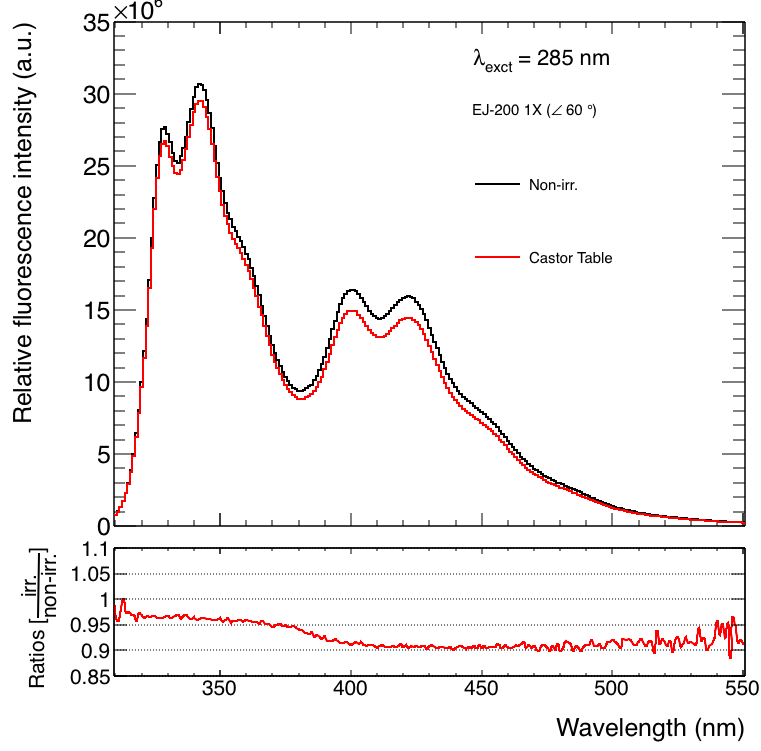


#### 4. Measurements

After irradiations, the samples showed coloration which partially disappeared after time due to the annealing of the substrate. The samples were repeatedly measured after irradiation until no additional annealing was seen. A typical fluorescence spectrum of EJ-200 is shown in Figure 2, after 3 Mrad at 9 krad/hr and after 3 Mrad at 429 krad/hr.

We calculate the light output ratio (R) using the integrated area of the irradiated sample's spectrum with respect to the corresponding non-irradiated one. Systematic uncertainties associated with the integrated area of each spectrum measurement are identified. Repeated measurements on the reference sample on different days, after turning off and on the instrument everyday, suggest the relative uncertainty associated with day-to-day fluctuation to be  $\pm 8\%$ . However, repeated measurements on the reference sample during the same day, when the spectrophotometer is continuously running, suggest the relative uncertainty

Figure 2: Comparison between fluorescence spectrum of an irradiated sample to that of the non-irradiated sample excited at 285 nm.



associated with same-day fluctuation to be  $\pm 1\%$ . This indicates that although initialization conditions of the instrument strongly affect photon-counting, the instrument is stable after initialization. Therefore the variation of intensity measured on different days can be described by a normalization factor as a function of day, by measuring the reference sample. After correcting the integrated areas of spectra with their normalization factors, the uncertainty associated with instrument is dominated by same-day fluctuation. Measurements on the reference sample when exciting its four different sides suggest the corresponding uncertainty to be  $\pm 1\%$ , due to sample imperfection (optical surface, dopant concentration, etc). The monochromators have an accuracy of  $\pm 0.5$  nm, which translates into a negligible uncertainty on the integrated area of a spectrum. Since PMT's photon-counting obeys Poisson distribution and the number of photons is large, the corresponding random uncertainty is negligible. Therefore, due to mainly two independent uncertainties (sample imperfection and instrument systematic uncertainty), the uncertainty on light output ratio is found to be equal to or less than  $\pm 2\%$ .

Table 2 shows light output ratios excited at three different wavelengths for

1P and 2P samples of the same dose (3 Mrad) at different dose rates (9 krad/hr and 429 krad/hr). Table 3 shows light output ratios excited at three different wavelengths for 1P and 2P samples of different doses at dose rate 9 krad/hr. Table 5 shows light output ratios excited at three different wavelengths for 1X samples of same dose (6 Mrad) and dose rate (80 krad/hr), but at two different ambient temperatures (23 °C and -30 °C) during irradiation. Table 4 shows light output ratios excited at three different wavelengths for 1X and 2X samples of the same dose (1.68 Mrad) at very low dose rate (0.84 krad/hr),

Table 2: Light output ratios after irradiations of the same total dose and at different dose rates. The uncertainties on the total dose and dose rate are less than 2%.

Sample	Dose (Mrad)	Dose rate (krad/hr)	$R_{285 \text{ nm}}$	$R_{350 \text{ nm}}$	$R_{400 \text{ nm}}$
1P	2.95	9.48	$0.95 \pm 0.02$	$0.95 \pm 0.02$	$0.92 \pm 0.02$
1P	2.95	429	$0.95 \pm 0.02$	$0.91 \pm 0.02$	$0.84 \pm 0.02$
2P	2.95	9.48	$0.95 \pm 0.02$	$0.95 \pm 0.02$	$0.92 \pm 0.02$
2P	2.95	429	$0.95 \pm 0.02$	$0.91 \pm 0.02$	$0.84 \pm 0.02$

Table 3: Light output ratios after irradiations of different total doses at dose rate 9 krad/hr. The uncertainties on the total dose and dose rate are less than 2%.

Sample	Dose (Mrad)	Dose rate (krad/hr)	$R_{285 \text{ nm}}$	$R_{350 \text{ nm}}$	$R_{400 \text{ nm}}$
1P	2.95	9.48	$0.95 \pm 0.02$	$0.95 \pm 0.02$	$0.92 \pm 0.02$
1P	4.00	9.20	$0.93 \pm 0.02$	$0.92 \pm 0.02$	$0.87 \pm 0.02$
1P	6.95	9.32	$0.90 \pm 0.02$	$0.88 \pm 0.02$	$0.81 \pm 0.02$
2P	2.95	9.48	$0.95 \pm 0.02$	$0.95 \pm 0.02$	$0.92 \pm 0.02$
2P	4.00	9.20	$0.94 \pm 0.02$	$0.93 \pm 0.02$	$0.90 \pm 0.02$
2P	6.95	9.32	$0.90 \pm 0.02$	$0.89 \pm 0.02$	$0.83 \pm 0.02$

Table 4: Light output ratios after irradiation of 1.68 Mrad at 0.84 krad/hr. Note that the uncertainty on the total dose is 10%, and the uncertainty on the dose rate is estimated to be 20%.

Sample	Dose (Mrad)	Dose rate (krad/hr)	$R_{285 \text{ nm}}$	$R_{350 \text{ nm}}$	$R_{400 \text{ nm}}$
1X	1.68	0.84	$0.94 \pm 0.02$	$0.94 \pm 0.02$	$0.88 \pm 0.02$
2X	1.68	0.84	$0.95 \pm 0.02$	$0.97 \pm 0.02$	$0.90 \pm 0.02$

## 5. Interpretation

Table 3 shows that radiation damage is not only to the base material, but also to the dopants of plastic scintillator. Measurements at 400 nm excitation

Table 5: Light output ratios after irradiations of the same total dose and dose rate at 23 °C and -30 °C. The uncertainties on the total dose and dose rate are less than 2%.

Sample	Dose (Mrad)	Dose rate (krad/hr)	$R_{285\text{ nm}}$	$R_{350\text{ nm}}$	$R_{400\text{ nm}}$
1X, 23 °C	5.82	80.3	$0.90 \pm 0.02$	$0.88 \pm 0.02$	$0.77 \pm 0.02$
1X, -30 °C	5.82	80.4	$0.90 \pm 0.02$	$0.84 \pm 0.02$	$0.70 \pm 0.01$

wavelength, which almost exclusively excites the secondary dopant, show that the secondary dopant is increasingly damaged as the total dose increases.

If the dose rate effect on light output is due to oxygen diffusion, then it is expected that the reduction in light output for the same dose would have a negative correlation with the dose rate, except for when the dose rate is low enough that oxygen permeates the entire sample. Comparing the light output ratios after irradiations of 3 Mrad dose at two different dose rates, 9 krad/hr and 429 krad/hr, we find no negative correlation between the reduction in light output and dose rate (Table 2). On the contrary, measurements at 400 nm excitation wavelength, which almost exclusively excites secondary dopant, show that the reduction in secondary dopant's light output is two times higher at 429 krad/hr than at 9 krad/hr. Measurements at 350 nm excitation wavelength, which excites largely secondary dopant and non-negligible primary dopant, confirm greater reduction at 429 krad/hr in secondary dopant's light output. Measurements at 285 nm excitation wavelength, which excites PVT and primary dopant, show no significant difference in light output ratios between two dose rates.

However, Table 4 and Table 3 show that the amount of light output reduction is comparable between a total dose of 1.68 Mrad at very low dose rate 0.84 krad/hr and a total dose of 3-4 Mrad at a much higher dose rate 9 krad/hr.

The diffusion rate of oxygen in PVT is expected to be higher at higher temperature. Temperature also affects molecules' average kinetic energy and therefore various energy transfer mechanisms in the scintillator. Studies have shown temperature effects on plastic scintillator in pre-irradiation treatment [13] and post-irradiation annealing [14]. However, there has been little study on the temperature effect during irradiation on plastic scintillators. Table 5 suggests that more secondary dopant is damaged during the irradiation at -30 °C than at 23 °C.

## 6. Conclusion

Using a fluorescence spectrophotometer, fluorescence light output versus wavelengths for plastic scintillator EJ-200 with varying dopant concentrations is measured before and after irradiation for various doses and dose rates. Dose, dose rate, and temperature effects on the scintillator's dopants are investigated after irradiations. We found that the plastic scintillator's secondary dopant is damaged by irradiation. This indicates that radiation damage is not only to the

substrate but also to the dopants of plastic scintillators. Measurements of the light output from surface excitation show a dose rate effect in radiation damage to the secondary dopant, that the higher dose rate 429 krad/hr causes more damage to the secondary dopant than the lower dose rate 9 krad/hr. However, a total dose of 1.68 Mrad at very low dose rate 0.84 krad/hr causes the similar amount of damage to the dopants as a total dose of 3-4 Mrad at higher dose rate 9 krad/hr, suggesting nontrivial mechanisms of radiation damage to the dopants. In particular, sub-krad/hr dose rate effect on scintillator light output is of interest to the upgrade of the CMS hadron calorimeter, and more comprehensive studies would help to design hadron calorimeters and to understand radiation damage to scintillator dopants. We observe that more secondary dopant is damaged during the irradiation at -30 °C than at 23 °C. This is also of interest to the upgrade of the CMS hadron calorimeter, with the possible scenario of cold environment, which requires a deeper understanding of the temperature effect on plastic scintillator during irradiation.

## Appendix A. Absorption and Emission Spectra while Annealing

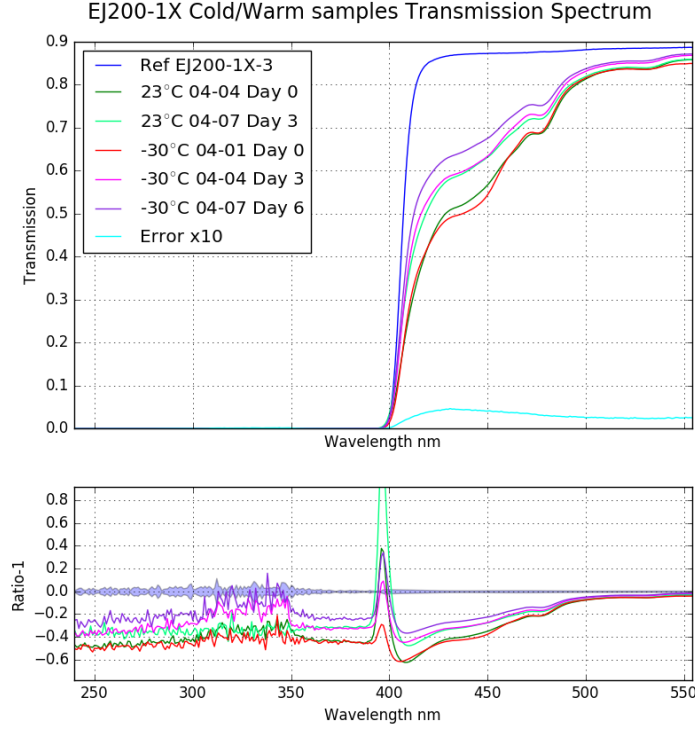
Comparison between the variations in the bulk-material absorption spectrum [15] and the fluorescence emission spectrum during the annealing process shows that significant change in the bulk absorption does not affect the fluorescence emission measurements.

## References

- [1] J. F. Butler, D. U. C. B.-L. I. Contardo, M. M. Klute, J. U. o. M. Mans, L. I.-B. Silvestris, Technical Proposal for the Phase-II Upgrade of the CMS Detector, Tech. Rep. CERN-LHCC-2015-010. LHCC-P-008, CERN, Geneva, upgrade Project Leader Deputies: Lucia Silvestris (INFN-Bari), Jeremy Mans (University of Minnesota) Additional contacts: Lucia.Silvestris@cern.ch, Jeremy.Mans@cern.ch (Jun 2015).  
URL <https://cds.cern.ch/record/2020886>
- [2] ECFA High Luminosity LHC Experiments Workshop: Physics and Technology Developments Summary submitted to ECFA. 96th Plenary ECFA meeting.  
URL <https://cds.cern.ch/record/1983664>
- [3] V. Hagopian, I. Daly, Radiation damage of fibers, AIP Conference Proceedings 450 (1) (1998) 53–61. doi:<http://dx.doi.org/10.1063/1.56958>.
- [4] A. Byon-Wagner, Radiation hardness test programs for the {SDC} calorimeter, Radiation Physics and Chemistry 41 (12) (1993) 263 – 271. doi:[http://dx.doi.org/10.1016/0969-806X\(93\)90064-2](http://dx.doi.org/10.1016/0969-806X(93)90064-2).
- [5] C. Zorn, Plastic and liquid organic scintillators, in: F. Sauli (Ed.), Instrumentation in High Energy Physics, 2nd Edition, World Scientific, 1993, Ch. 4, pp. 218–279. doi:10.1142/9789814360333\_0004.

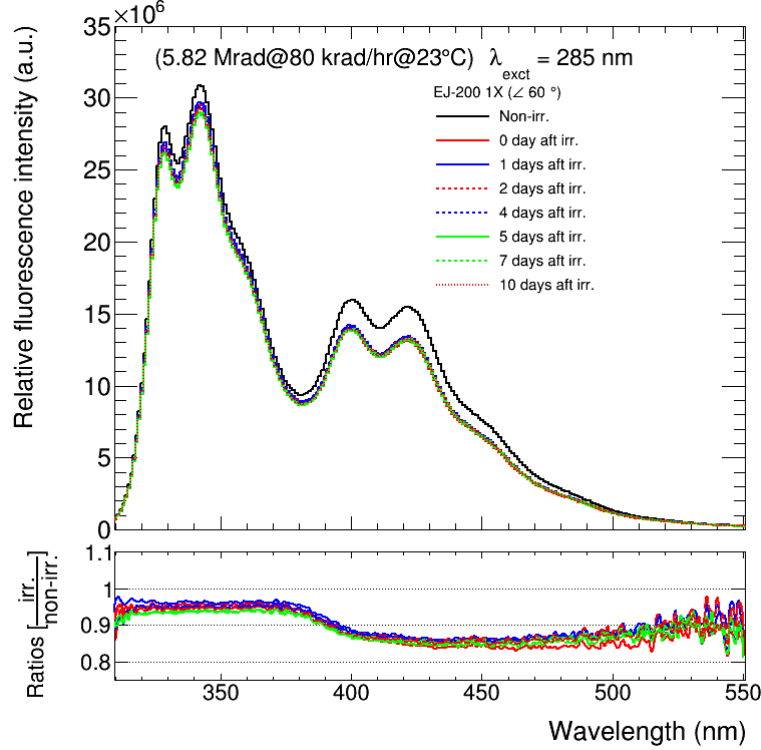


Figure A.3: Absorption spectra of irradiated samples (at 23 °C and -30 °C) repeatedly measured during the annealing process.



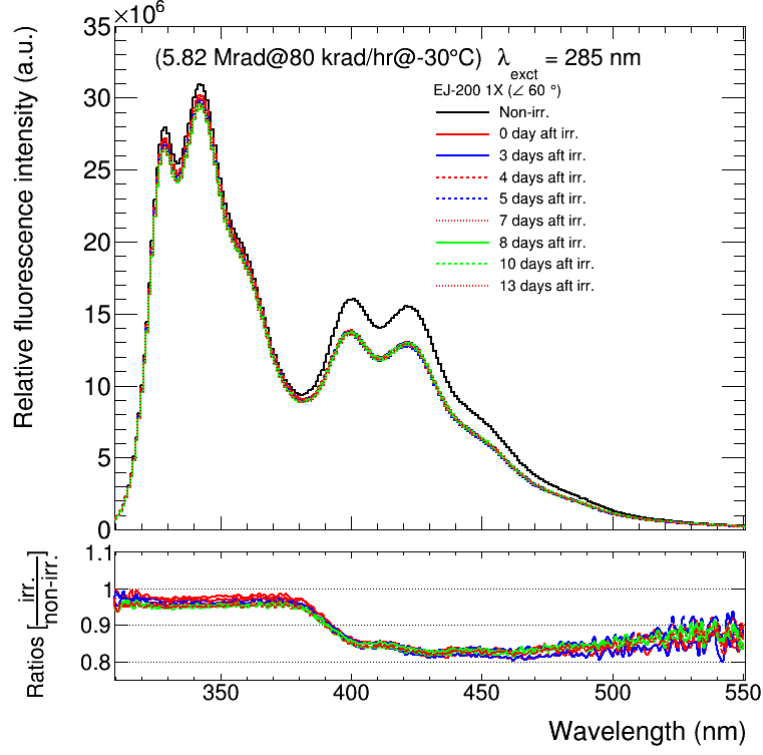
- [6] K. Wick, D. Paul, P. Schröder, V. Stieber, B. Bicken, Recovery and dose rate dependence of radiation damage in scintillators, wavelength shifters and light guides, Nuclear Instruments and Methods in Physics Research Section B: Beam Interactions with Materials and Atoms 61 (4) (1991) 472 – 486. doi:[http://dx.doi.org/10.1016/0168-583X\(91\)95325-8](http://dx.doi.org/10.1016/0168-583X(91)95325-8).
- [7] K. Wick, T. Zoufal, Unexpected behaviour of polystyrene-based scintillating fibers during irradiation at low doses and low dose rates, Nuclear Instruments and Methods in Physics Research Section B: Beam Interactions with Materials and Atoms 185 (14) (2001) 341 – 345. doi:[http://dx.doi.org/10.1016/S0168-583X\(01\)00776-5](http://dx.doi.org/10.1016/S0168-583X(01)00776-5).
- [8] B. Wulko, K. Wick, W. Busjan, A. Dannemann, U. Holm, Evidence for the creation of short-lived absorption centers in irradiated scintillators, Nuclear Instruments and Methods in Physics Research Section B: Beam Interactions with Materials and Atoms 95 (1) (1995) 141 – 143. doi:[http://dx.doi.org/10.1016/0168-583X\(94\)00435-8](http://dx.doi.org/10.1016/0168-583X(94)00435-8).

Figure A.4: Fluorescence spectra of the sample irradiated at 23°C repeatedly measured during the annealing process.



- [9] A. Bross, A. Pla-Dalmau, Radiation damage of plastic scintillators, Nuclear Science, IEEE Transactions on 39 (5) (1992) 1199–1204. doi:10.1109/23.173178.
- [10] E. Biagtan, E. Goldberg, R. Stephens, E. Valeroso, J. Harmon, Gamma dose and dose rate effects on scintillator light output, Nuclear Instruments and Methods in Physics Research Section B: Beam Interactions with Materials and Atoms 108 (12) (1996) 125 – 128. doi:http://dx.doi.org/10.1016/0168-583X(95)00874-8.
- [11] H. Scientific, FluoroMax-4 & FluoroMax-4P with USB Operation Manual, HORIBA Jobin Yvon Inc.
- [12] C. Hurlbut, Eljen Technology, private communication.
- [13] K. Johnson, H. Whitaker, Temperature treatment of plastic scintillator affects radiation hardness, Nuclear Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment 301 (2) (1991) 372 – 375. doi:http://dx.doi.org/10.1016/0168-9002(91)90480-E.

Figure A.5: Fluorescence spectra of the sample irradiated at  $-30^{\circ}\text{C}$  repeatedly measured during the annealing process.



- [14] J. B. Birks, The Theory and practice of scintillation counting, 1964.  
 URL <http://www.slac.stanford.edu/spires/find/books/www?cl=QCD928:B52>
- [15] K. Wong, A. Belloni, S. Eno, Absorption Study on Plastic Scintillators After Irradiation, in preparation.