

First we define known quantities

$$\begin{aligned}
 \Delta &:= r^2 - 2 M r + a^2 \\
 \Sigma &:= r^2 + a^2 \cos^2[\theta] \\
 J &:= C \text{DiracDelta}[r - r_0] \text{DiracDelta}\left[\theta - \frac{\pi}{2}\right] \\
 J_{m\uparrow} &:= \left(\sqrt{2} (r - i a \cos[\theta]) \right)^{-1} + i (r^2 + a^2) \sin[\theta] J \\
 J_n &:= - \frac{a \Delta}{\Sigma} \sin[\theta]^2 J
 \end{aligned}$$

Now for J_2 , we can ignore ∂_ϕ since J_m and J_n are independent of it

$$\begin{aligned}
 J_2 = & \frac{-\Delta}{2\sqrt{2} \Sigma (r - i a \cos \theta)^2} \left[\sqrt{2} \left(\frac{\partial}{\partial r} - \frac{a}{\Delta} \frac{\partial}{\partial \phi} + \frac{1}{r - i a \cos \theta} \right) (r - i a \cos \theta)^2 J_{\bar{m}} \right. \\
 & \left. + 2 \left(\frac{\partial}{\partial \theta} - \frac{i}{\sin \theta} \frac{\partial}{\partial \phi} + \frac{i a \sin \theta}{r - i a \cos \theta} \right) \frac{\Sigma (r - i a \cos \theta)}{\Delta} J_n \right] \quad (2.12)
 \end{aligned}$$

$$\begin{aligned}
 J_2 := & \frac{-\Delta}{2 \sqrt{2} \Sigma (r - i a \cos[\theta])^2} \\
 & \left(\sqrt{2} \left(\partial_r \# + \frac{1}{r - i a \cos[\theta]} \# \right) \&@((r - i a \cos[\theta])^2 J_{m\uparrow}) + 2 \left(\partial_\theta \# + \frac{i a \sin[\theta]}{r - i a \cos[\theta]} \# \right) \&@ \right. \\
 & \left. \left(\frac{\Sigma (r - i a \cos[\theta])}{\Delta} J_n \right) \right)
 \end{aligned}$$

Finally for the integration we need to first look at the definition of the SpinWeighted spherical harmonics.

Since all SpinWeighted spherical harmonics, here denoted ${}_s Y_{lm}(\theta, \phi) = Y[s, l, \theta, \phi]$ behave in the ϕ argument as $\text{Exp}[i m \phi]$ and since all other variables are independent of ϕ in the equation above We have

$\int Y[s, l, m, \theta, \phi] = 2 \pi \text{DiracDelta}[m, 0] Y[s, l, m, \theta, 0]$. Then we can integrate over θ .

the source term is given by

$${}^2 J_{lm}(r) = \int_0^{2\pi} \int_0^\pi \frac{(r - i a \cos \theta)^2}{(r_+ - r_-)^2} \Sigma J_2 {}_{-1} \bar{Y}_{lm} \sin \theta d\theta d\phi$$

In[]:=

We can split the integrand into two parts depending if we have $\text{DiracDelta}[\theta - \frac{\pi}{2}]$ or the derivative $\text{DiracDelta}'[\theta - \frac{\pi}{2}]$.

and we will use the standard properties

$$\int_{-\infty}^{\infty} f(x) \delta(x-a) dx = f(a)$$

$$\int f(x) \delta^{(n)}(x) dx \equiv - \int \frac{\partial f}{\partial x} \delta^{(n-1)}(x) dx.$$

First part

$$\begin{aligned} \text{In}[*]:= & \frac{(r - i a \cos[\theta])^2}{(rp - rm)^2} \Sigma J2 \text{Conjugate}[Y[-1, l, 0, \theta, 0]]; \\ & \% /. \text{DiracDelta}[-\pi + 2 \theta] \rightarrow 1 /. \text{DiracDelta}[-\pi + 2 \theta] \rightarrow 0; \\ & \text{FirstPart} = \% /. \theta \rightarrow \frac{\pi}{2} \\ \text{Out}[*]= & -\frac{1}{2 \sqrt{2} (-rm + rp)^2} (a^2 - 2 M r + r^2) \text{Conjugate}[Y[-1, l, 0, \frac{\pi}{2}, 0]] \\ & \left(-8 i a^2 C \text{DiracDelta}[r - r0] + \sqrt{2} \left(3 r \left(\frac{1}{\sqrt{2} r} + 2 i C (a^2 + r^2) \text{DiracDelta}[r - r0] \right) + \right. \right. \\ & \left. \left. r^2 \left(-\frac{1}{\sqrt{2} r^2} + 4 i C r \text{DiracDelta}[r - r0] + 2 i C (a^2 + r^2) \text{DiracDelta}'[r - r0] \right) \right) \right) \end{aligned}$$

Second part

$$\begin{aligned} \text{In}[*]:= & \frac{(r - i a \cos[\theta])^2}{(rp - rm)^2} \Sigma J2 \text{Conjugate}[Y[-1, l, 0, \theta, 0]]; \\ & \% /. \text{DiracDelta}[-\pi + 2 \theta] \rightarrow 0 /. \text{DiracDelta}[-\pi + 2 \theta] \rightarrow 1; \\ & \text{SecondPart} = -D[\%, \theta] /. \theta \rightarrow \frac{\pi}{2} /. \text{Conjugate}'[Y[-1, l, 0, \frac{\pi}{2}, 0]] \rightarrow 1 \\ \text{Out}[*]= & -\frac{2 i \sqrt{2} a^2 C (a^2 - 2 M r + r^2) \text{Conjugate}[Y[-1, l, 0, \frac{\pi}{2}, 0]] \text{DiracDelta}[r - r0]}{(-rm + rp)^2} + \\ & \frac{(a^2 - 2 M r + r^2) (2 - 8 a C r \text{DiracDelta}[r - r0]) Y^{(0,0,0,1,0)}[-1, l, 0, \frac{\pi}{2}, 0]}{2 \sqrt{2} (-rm + rp)^2} \end{aligned}$$

We get the solution

In[]:= **solution = FirstPart + SecondPart // Simplify**

$$\begin{aligned} \text{Out[]} = & - \frac{2 i \sqrt{2} a^2 C (a^2 - 2 M r + r^2) \text{Conjugate}[Y[-1, l, 0, \frac{\pi}{2}, 0]] \text{DiracDelta}[r - r0]}{(-r m + r p)^2} - \\ & \frac{1}{2 \sqrt{2} (-r m + r p)^2} (a^2 - 2 M r + r^2) \text{Conjugate}[Y[-1, l, 0, \frac{\pi}{2}, 0]] \\ & \left(-8 i a^2 C \text{DiracDelta}[r - r0] + \sqrt{2} \left(3 r \left(\frac{1}{\sqrt{2} r} + 2 i C (a^2 + r^2) \text{DiracDelta}[r - r0] \right) + \right. \right. \\ & \left. \left. r^2 \left(-\frac{1}{\sqrt{2} r^2} + 4 i C r \text{DiracDelta}[r - r0] + 2 i C (a^2 + r^2) \text{DiracDelta}'[r - r0] \right) \right) \right) + \\ & \frac{(a^2 - 2 M r + r^2) (2 - 8 a C r \text{DiracDelta}[r - r0]) Y^{(0,0,0,1,0)}[-1, l, 0, \frac{\pi}{2}, 0]}{2 \sqrt{2} (-r m + r p)^2} \end{aligned}$$

Finally, we can simplify the results by using identities for SpinWeightedSpherical Harmonics

HarmonicRules =

$$\begin{aligned} & \left\{ Y^{(0,0,0,1,0)}[s_-, l_-, m_-, \theta_-, \phi_-] \rightarrow s \cot[\theta] Y[s, l, m, \theta, \phi] - \sqrt{l + l^2 - s - s^2} Y[1 + s, l, m, \theta, \phi] - \right. \\ & \quad i \csc[\theta] Y^{(0,0,0,0,1)}[s, l, m, \theta, \phi], Y^{(0,0,0,0,1)}[s_-, l_-, m_-, \theta_-, \phi_-] \rightarrow i m Y[s, l, m, \theta, \phi], \\ & \quad \left. \text{Conjugate}[Y[s_-, l_-, m_-, \theta_-, \phi_-]] \rightarrow (-1)^{-m+s} Y[-s, l, -m, \theta, \phi] \right\}; \end{aligned}$$

$$\begin{aligned} \text{Out[]} = & \left\{ Y^{(0,0,0,1,0)}[s_-, l_-, m_-, \theta_-, \phi_-] \rightarrow \right. \\ & s \cot[\theta] Y[s, l, m, \theta, \phi] - \sqrt{l + l^2 - s - s^2} Y[1 + s, l, m, \theta, \phi] - i \csc[\theta] Y^{(0,0,0,0,1)}[s, l, m, \theta, \phi], \\ & Y^{(0,0,0,0,1)}[s_-, l_-, m_-, \theta_-, \phi_-] \rightarrow i m Y[s, l, m, \theta, \phi], \\ & \left. \text{Conjugate}[Y[s_-, l_-, m_-, \theta_-, \phi_-]] \rightarrow (-1)^{-m+s} Y[-s, l, -m, \theta, \phi] \right\} \end{aligned}$$

Applying these rules we get the final solution

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In[ ]:= solution /. HarmonicRules /. Y(0,0,0,0,1)[-1, l, 0,  $\frac{\pi}{2}$ , 0] → 0;

Collect[%, {DiracDelta[r - r0], DiracDelta[r - r0]}] /.

{Y[1, l, 0,  $\frac{\pi}{2}$ , 0] → -Conjugate[Y[-1, l, 0,  $\frac{\pi}{2}$ , 0]]};

Print["Solution == ", %]

Solution == -  $\frac{(a^2 - 2 M r + r^2) \text{Conjugate}[Y[-1, l, 0, \frac{\pi}{2}, 0]]}{\sqrt{2} (-r m + r p)^2}$  -  $\frac{\sqrt{l + l^2} (a^2 - 2 M r + r^2) Y[0, l, 0, \frac{\pi}{2}, 0]}{\sqrt{2} (-r m + r p)^2}$  +

DiracDelta[r - r0]  $\left( -\frac{2 i C r^3 (a^2 - 2 M r + r^2) \text{Conjugate}[Y[-1, l, 0, \frac{\pi}{2}, 0]]}{(-r m + r p)^2} - \right.$ 

 $\frac{3 i C r (a^2 + r^2) (a^2 - 2 M r + r^2) \text{Conjugate}[Y[-1, l, 0, \frac{\pi}{2}, 0]]}{(-r m + r p)^2} +$ 

 $\left. \frac{2 \sqrt{2} a C \sqrt{l + l^2} r (a^2 - 2 M r + r^2) Y[0, l, 0, \frac{\pi}{2}, 0]}{(-r m + r p)^2} \right)$  -

 $\frac{i C r^2 (a^2 + r^2) (a^2 - 2 M r + r^2) \text{Conjugate}[Y[-1, l, 0, \frac{\pi}{2}, 0]] \text{DiracDelta}[r - r0]}{(-r m + r p)^2}$ 

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We should further utilize identities for the Y function but we can already see that we must have made a mistake

Since the solution given in the article ($C = l \left(\frac{\Delta_0}{p_0} \right)$) and ($e = 0$)

$$\begin{aligned}
 {}^2J_{lm} = & - \frac{\Delta \delta_{m0}}{\sqrt{2} (r_+ - r_-)^2} \left[(M a e / \mathfrak{Q}_0) + \pi \mathfrak{J}(\Delta_0 / \mathfrak{Q}_0)^{1/2} \right] \\
 & \cdot \left[i(r_0^2 + a^2) {}_{-1}\bar{Y}_{l0} \left(\frac{\pi}{2}, 0 \right) \delta'(r - r_0) \right. \\
 & \left. + \left\{ i r_0 {}_{-1}\bar{Y}_{l0} \left(\frac{\pi}{2}, 0 \right) - a [l(l+1)]^{1/2} {}_0\bar{Y}_{l0} \left(\frac{\pi}{2}, 0 \right) \right\} \delta(r - r_0) \right]
 \end{aligned}$$

Is only quadratic in r, but our solution is at least quartic in r.