

First we define known quantities

$$\begin{aligned} \Delta &:= r^2 - 2 M r + a^2 \\ \Sigma &:= r^2 + a^2 \cos^2[\theta] \\ J &:= C \operatorname{DiracDelta}[r - r_0] \operatorname{DiracDelta}\left[\theta - \frac{\pi}{2}\right] \\ J_{m\bar{n}} &:= \left(\sqrt{2} (r - i a \cos[\theta]) \right)^{-1} i (r^2 + a^2) \sin[\theta] J \\ J_n &:= - \frac{a \Delta}{\Sigma} \sin[\theta]^2 J \end{aligned}$$

Now for J_2 , we can ignore ∂_ϕ since J_m and J_n are independent of it

$$\begin{aligned} J_2 = & \frac{-\Delta}{2\sqrt{2} \Sigma (r - i a \cos \theta)^2} \left[\sqrt{2} \left(\frac{\partial}{\partial r} - \frac{a}{\Delta} \frac{\partial}{\partial \phi} + \frac{1}{r - i a \cos \theta} \right) (r - i a \cos \theta)^2 J_{\bar{m}} \right. \\ & \left. + 2 \left(\frac{\partial}{\partial \theta} - \frac{i}{\sin \theta} \frac{\partial}{\partial \phi} + \frac{i a \sin \theta}{r - i a \cos \theta} \right) \frac{\Sigma (r - i a \cos \theta)}{\Delta} J_n \right] \end{aligned} \quad (2.12)$$

$$\begin{aligned} J_2 := & \frac{-\Delta}{2 \sqrt{2} \Sigma (r - i a \cos[\theta])^2} \\ & \left(\sqrt{2} \left(\partial_r \# + \frac{1}{r - i a \cos[\theta]} \# \right) \&@ \left((r - i a \cos[\theta])^2 J_{m\bar{n}} \right) + 2 \left(\partial_\theta \# + \frac{i a \sin[\theta]}{r - i a \cos[\theta]} \# \right) \&@ \right. \\ & \left. \left(\frac{\Sigma (r - i a \cos[\theta])}{\Delta} J_n \right) \right) \end{aligned}$$

Finally for the integration we need to first look at the definition of the SpinWeighted spherical harmonics.

Since all SpinWeighted spherical harmonics, here denoted ${}_s Y_{lm}(\theta, \phi) = Y[s, l, \theta, \phi]$ behave in the ϕ argument as $\exp[i m \phi]$ and since all other variables are independent of ϕ in the equation above We have

$\int Y[s, l, m, \theta, \phi] = 2 \pi \operatorname{DiracDelta}[m, 0] Y[s, l, m, \theta, 0]$. Then we can integrate over θ .

the source term is given by

$${}^2 J_{lm}(r) = \int_0^{2\pi} \int_0^\pi \frac{(r - i a \cos \theta)^2}{(r_+ - r_-)^2} \Sigma J_2 {}_{-1} \bar{Y}_{lm} \sin \theta d\theta d\phi$$

In[55]:=

We can split the integrand into two parts depending if we have $\operatorname{DiracDelta}\left[\theta - \frac{\pi}{2}\right]$ or the derivative $\operatorname{DiracDelta}'\left[\theta - \frac{\pi}{2}\right]$.

and we will use the standard properties

$$\int_{-\infty}^{\infty} f(x) \delta(x-a) dx = f(a)$$

$$\int f(x) \delta^{(n)}(x) dx \equiv - \int \frac{\partial f}{\partial x} \delta^{(n-1)}(x) dx.$$

First part

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In[56]:= (r - I a Cos[θ])^2
          (rp - rm)^2 Σ J2 Conjugate[Y[-1, l, 0, θ, 0]];
% /. DiracDelta[-π + 2 θ] → 1 /. DiracDelta[-π + 2 θ] → 0;
FirstPart = % /. θ → π/2
Out[58]:= - 1/(2 √2 (-rm + rp)^2) (a^2 - 2 M r + r^2) Conjugate[Y[-1, l, 0, π/2, 0]]
          (-8 I a^2 C DiracDelta[r - r0] + √2 (2 I √2 C r^2 DiracDelta[r - r0] +
          2 I √2 C (a^2 + r^2) DiracDelta[r - r0] + I √2 C r (a^2 + r^2) DiracDelta'[r - r0]))
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Second part

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In[59]:= (r - I a Cos[θ])^2
          (rp - rm)^2 Σ J2 Conjugate[Y[-1, l, 0, θ, 0]];
% /. DiracDelta[-π + 2 θ] → 0 /. DiracDelta[-π + 2 θ] → 1;
SecondPart = -D[%, θ] /. θ → π/2 /. Conjugate[Y[-1, l, 0, π/2, 0]] → 1
Out[61]:= - 2 I √2 a^2 C (a^2 - 2 M r + r^2) Conjugate[Y[-1, l, 0, π/2, 0]] DiracDelta[r - r0]
          (-rm + rp)^2
          2 √2 a C r (a^2 - 2 M r + r^2) DiracDelta[r - r0] Y^(0,0,0,1,0)[-1, l, 0, π/2, 0]
          (-rm + rp)^2
```

We get the solution

```
In[62]:= solution = FirstPart + SecondPart // Simplify
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$$\text{Out[62]} = -\frac{2i\sqrt{2}a^2C(a^2 - 2Mr + r^2)\text{Conjugate}[Y[-1, l, 0, \frac{\pi}{2}, 0]]\text{DiracDelta}[r - r0]}{(-rm + rp)^2} -$$

$$\frac{1}{2\sqrt{2}(-rm + rp)^2}(a^2 - 2Mr + r^2)\text{Conjugate}[Y[-1, l, 0, \frac{\pi}{2}, 0]]$$

$$(-8ia^2C\text{DiracDelta}[r - r0] + \sqrt{2}(2i\sqrt{2}Cr^2\text{DiracDelta}[r - r0] +$$

$$2i\sqrt{2}C(a^2 + r^2)\text{DiracDelta}[r - r0] + i\sqrt{2}Cr(a^2 + r^2)\text{DiracDelta}'[r - r0])) -$$

$$\frac{2\sqrt{2}aCr(a^2 - 2Mr + r^2)\text{DiracDelta}[r - r0]Y^{(0,0,0,1,0)}[-1, l, 0, \frac{\pi}{2}, 0]}{(-rm + rp)^2}$$

Finally, we can simplify the results by using identities for SpinWeightedSpherical Harmonics

```
In[63]:= HarmonicRules =
```

$$\left\{ \begin{aligned} &Y^{(0,0,0,1,0)}[s_-, l_-, m_-, \theta_-, \phi_-] \rightarrow s \cot[\theta] Y[s, l, m, \theta, \phi] - \sqrt{l + l^2 - s - s^2} Y[1 + s, l, m, \theta, \phi] - \\ &i \csc[\theta] Y^{(0,0,0,0,1)}[s, l, m, \theta, \phi], Y^{(0,0,0,0,1)}[s_-, l_-, m_-, \theta_-, \phi_-] \rightarrow i m Y[s, l, m, \theta, \phi], \\ &\text{Conjugate}[Y[s_-, l_-, m_-, \theta_-, \phi_-] \rightarrow (-1)^{-m+s} Y[-s, l, -m, \theta, \phi] \end{aligned} \right\};$$

Applying these rules we get the final solution

```
In[64]:= solution /. HarmonicRules /. Y^{(0,0,0,0,1)}[-1, l, 0, \frac{\pi}{2}, 0] -> 0;
```

```
% /. {Y[1, l, 0, \frac{\pi}{2}, 0] -> -Conjugate[Y[-1, l, 0, \frac{\pi}{2}, 0]]} // Simplify;
```

```
Print["Solution == ", %]
```

$$\text{Solution} == -\frac{2i\sqrt{2}a^2C(a^2 - 2Mr + r^2)\text{Conjugate}[Y[-1, l, 0, \frac{\pi}{2}, 0]]\text{DiracDelta}[r - r0]}{(-rm + rp)^2} +$$

$$\frac{2\sqrt{2}aC\sqrt{l + l^2}r(a^2 - 2Mr + r^2)\text{DiracDelta}[r - r0]Y[0, l, 0, \frac{\pi}{2}, 0]}{(-rm + rp)^2} -$$

$$\frac{1}{2\sqrt{2}(-rm + rp)^2}(a^2 - 2Mr + r^2)\text{Conjugate}[Y[-1, l, 0, \frac{\pi}{2}, 0]]$$

$$(-8ia^2C\text{DiracDelta}[r - r0] + \sqrt{2}(2i\sqrt{2}Cr^2\text{DiracDelta}[r - r0] + 2i\sqrt{2}$$

$$C(a^2 + r^2)\text{DiracDelta}[r - r0] + i\sqrt{2}Cr(a^2 + r^2)\text{DiracDelta}'[r - r0]))$$

We should further utilize identities for the Y function but we can already see that we must have made a mistake

Since the solution given in the article ($C = I(\frac{\Delta 0}{\rho 0})$) and ($e = 0$)

$$\begin{aligned}
{}^2J_{lm} = & -\frac{\Delta\delta_{m0}}{\sqrt{2}(r_+-r_-)^2} [(Mae/\mathfrak{A}_0) + \pi \mathfrak{J}(\Delta_0/\mathfrak{A}_0)^{1/2}] \\
& \cdot \left[i(r_0^2 + a^2) {}_{-1}\overline{Y}_{l0} \left(\frac{\pi}{2}, 0 \right) \delta'(r - r_0) \right. \\
& \left. + \left\{ i r_0 {}_{-1}\overline{Y}_{l0} \left(\frac{\pi}{2}, 0 \right) - a [l(l+1)]^{1/2} {}_0\overline{Y}_{l0} \left(\frac{\pi}{2}, 0 \right) \right\} \delta(r - r_0) \right]
\end{aligned}$$

Is only quadratic in r, but our solution is at least quartic in r.