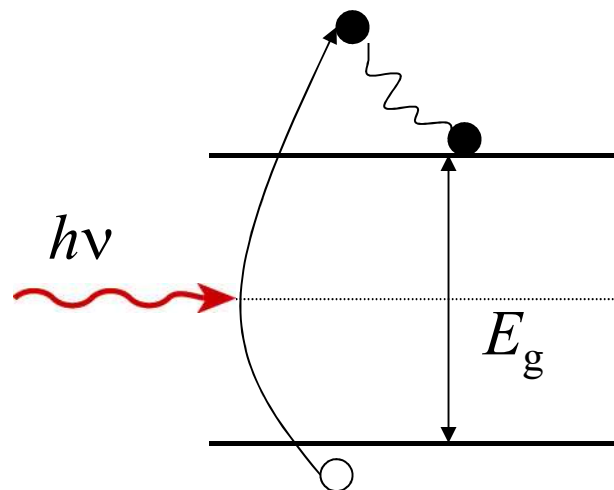


## Chapter 3-4. R-G statistics

- R-G statistics is the **mathematical characterization of R-G processes**
- An important generation process in device operation is **photo-generation**

If the photon energy ( $h\nu$ ) is greater than the band gap energy, then the light will be absorbed thereby creating electron-hole pairs



## Photo-generation

The intensity of monochromatic light that passes through a material is given by:  $I = I_0 \exp(-\alpha x)$  where  $I_0$  is the light intensity *just* inside the material at  $x = 0$ , and  $\alpha$  is the **absorption coefficient**. Note that  $\alpha$  is material dependent and is a strong function of  $\lambda$ .

Since photo-generation creates equal #s of holes and electrons, and each photon creates one e-h pair, we can write:

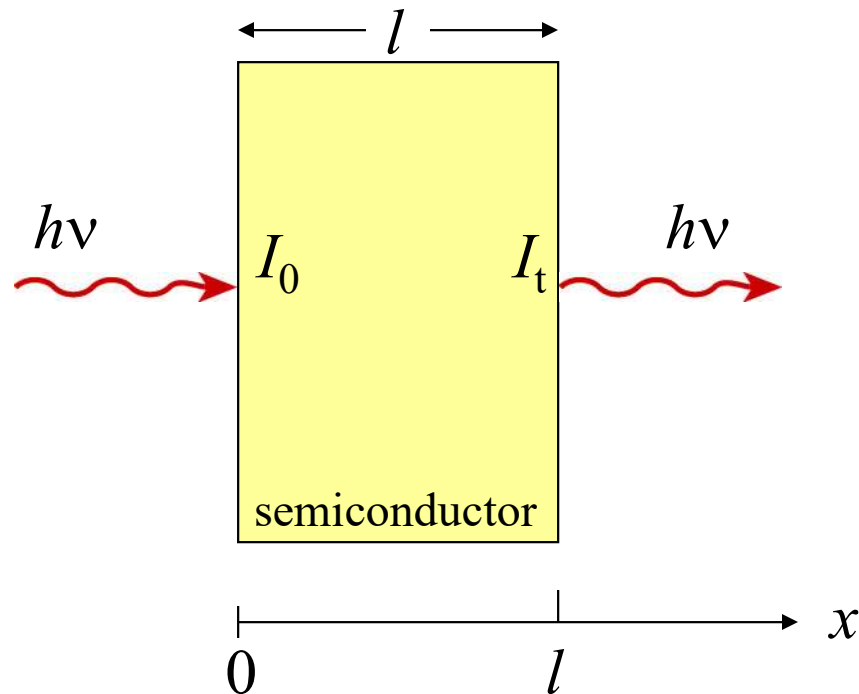
$$\frac{\partial n}{\partial t} \Big|_{\text{light}} = \frac{\partial p}{\partial t} \Big|_{\text{light}} = G_L(x, \lambda) = G_{L0} e^{-\alpha x}$$

where  $G_{L0}$  is the photo-generation rate [ $\# / (\text{cm}^3 \text{ s})$ ] at  $x = 0$

Question: What happens if the energy of photons is less than the band gap energy?

## Light absorption and transmittance

Consider a slab of semiconductor of thickness  $l$ .



$$I_t = I_0 \exp(-\alpha l)$$

where  $I_0$  is light intensity at  $x = 0$  and  $I_t$  is light intensity at  $x = l$ .

# Absorption coefficient vs. wavelength in semiconductors

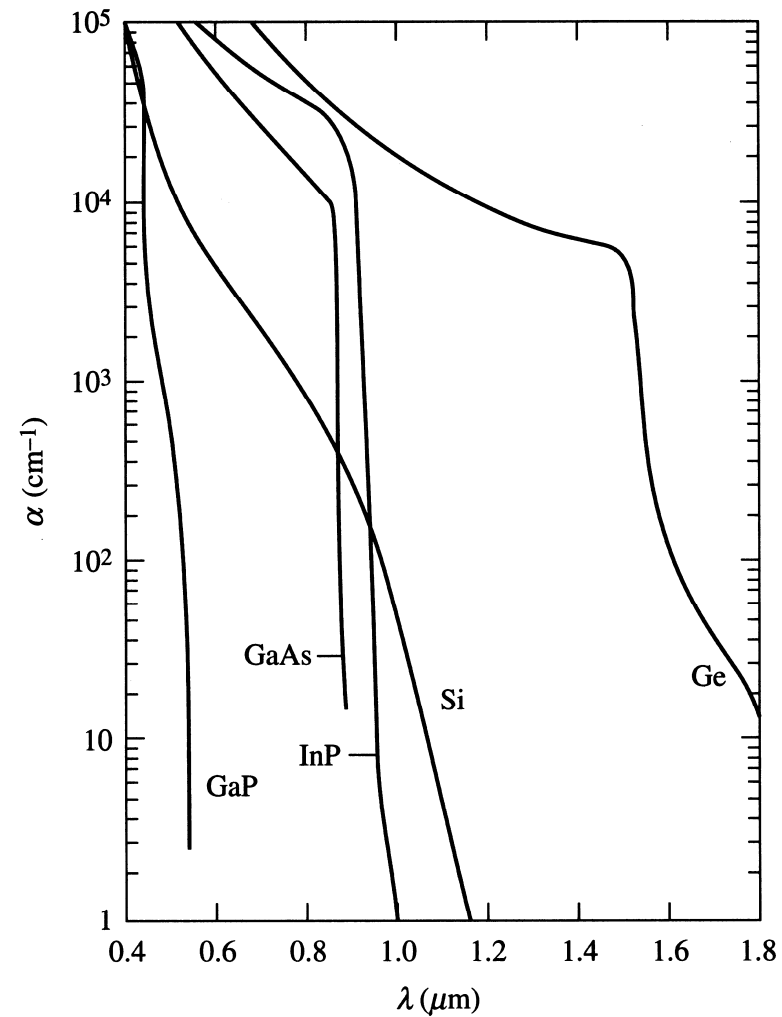
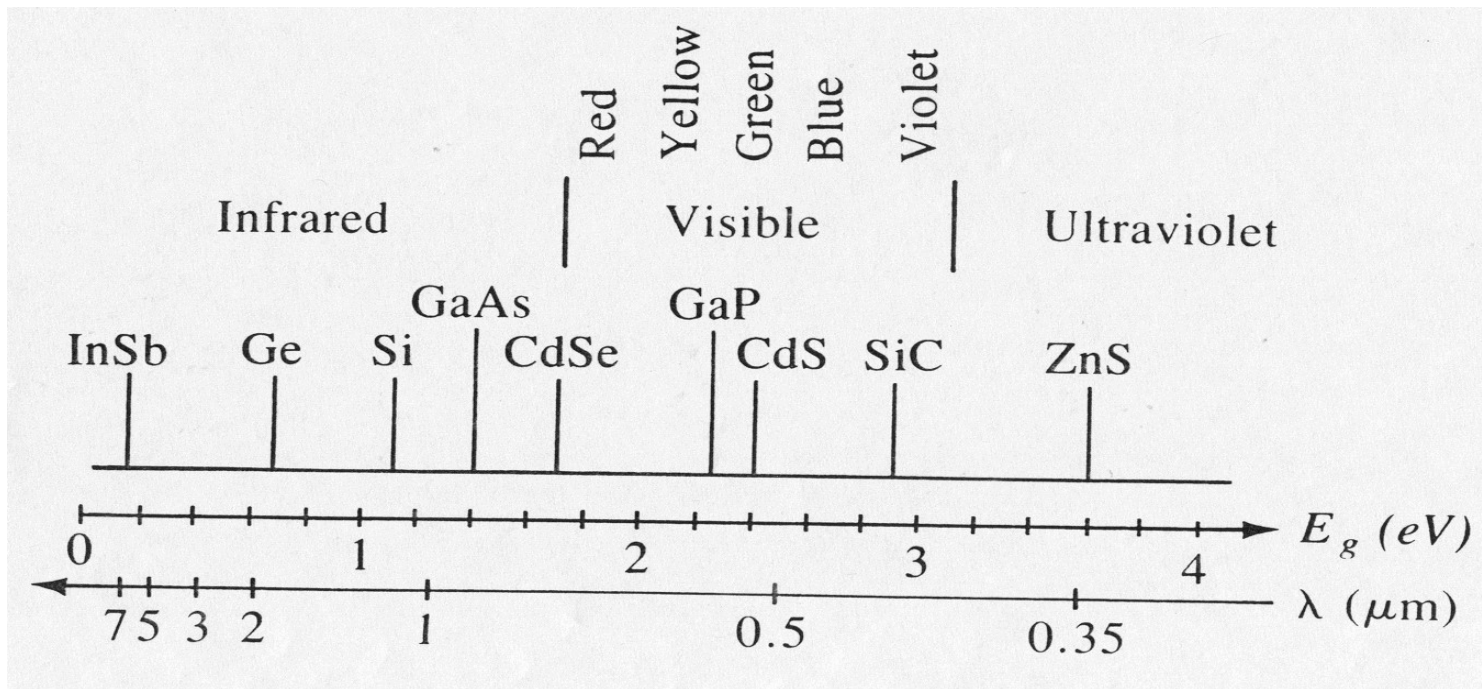


Figure 3.20

# Bandgaps of common semiconductors



# Indirect thermal recombination-generation

$n_0, p_0$  - under thermal equilibrium

$n, p$  - under arbitrary conditions, functions of  $t$

$$\left. \begin{aligned} \Delta n &= n - n_0 \\ \Delta p &= p - p_0 \end{aligned} \right\} \begin{array}{l} \Delta n \text{ and } \Delta p \text{ are deviations in carrier concentrations} \\ \text{from their equilibrium values. } \Delta n \text{ and } \Delta p \text{ can be} \\ \text{both positive or negative.} \end{array}$$

$N_t$  is the # of R-G centers/cm<sup>3</sup>

**Low-level injection** condition is assumed.

- Change in the majority carrier concentration is negligible. For example, in n-type material,  $\Delta p \ll n_0$ ;  $n \approx n_0$ .

## R-G statistics

Consider n-type silicon under perturbation:

We look at only minority carriers, and in this case, holes.

In general,

$$\frac{\partial p}{\partial t} = \frac{\partial p}{\partial t} \Big|_{\text{R}} + \frac{\partial p}{\partial t} \Big|_{\text{G}} + G_{\text{L}}(x, \lambda)$$

rate of	(loss)	(gain)	external
hole	due to	due to	such as
build up	recomb.	generation	light

$\frac{\partial p}{\partial t} \Big|_{\text{R}}$  should be proportional to  $p$  and  $N_{\text{t}}$ . Why?

$$\frac{\partial p}{\partial t} \Big|_{\text{R}} = -C_{\text{p}} N_{\text{t}} p$$

## R-G statistics (continued)

Under thermal equilibrium,  $G_L = 0$ ; and  $dp/dt = 0$

$$\longrightarrow \left. \frac{\partial p}{\partial t} \right|_G = - \left. \frac{\partial p}{\partial t} \right|_R = C_p N_t p_0$$

So, under arbitrary conditions, when  $G_L = 0$ ,

$$\begin{aligned} \frac{\partial p}{\partial t} &= -C_p N_t p + C_p N_t p_0 \\ &= -\frac{\Delta p}{\tau_p} \quad \text{defining} \quad \tau_p = \frac{1}{C_p N_t} \end{aligned}$$

$$\text{Since } \frac{\partial p}{\partial t} = \frac{\partial}{\partial t} (p_0 + \Delta p) = \frac{\partial \Delta p}{\partial t}$$



## R-G statistics (continued)

We can write:

$$\frac{\partial \Delta p}{\partial t} = -\frac{\Delta p}{\tau_p}$$

For holes in n-type

Similarly,

$$\frac{\partial \Delta n}{\partial t} = -\frac{\Delta n}{\tau_n}$$

For electrons in p-type

$\tau_p$  (or  $\tau_n$ ) is called “**minority carrier lifetime**” indicating the average time an excess minority carrier will survive in a sea of majority carriers.

Minority carrier lifetime is an important material parameter. Depends strongly on the concentration of deep-level of impurities, crystalline quality etc.. Varies strongly from material to material. Varies from a few ns to few ms in silicon depending on the quality!

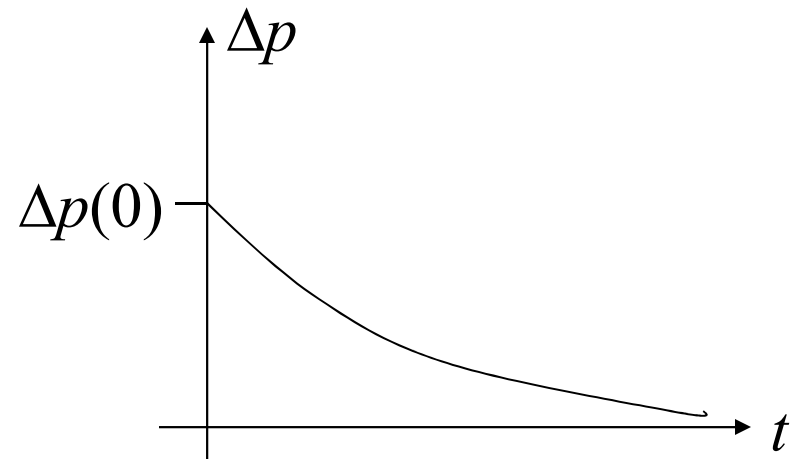
What happens when the perturbation is removed at  $t = 0$ ?

$$\frac{\partial \Delta p}{\partial t} = -\frac{\Delta p}{\tau_p}$$

For holes in an n-type semiconductor

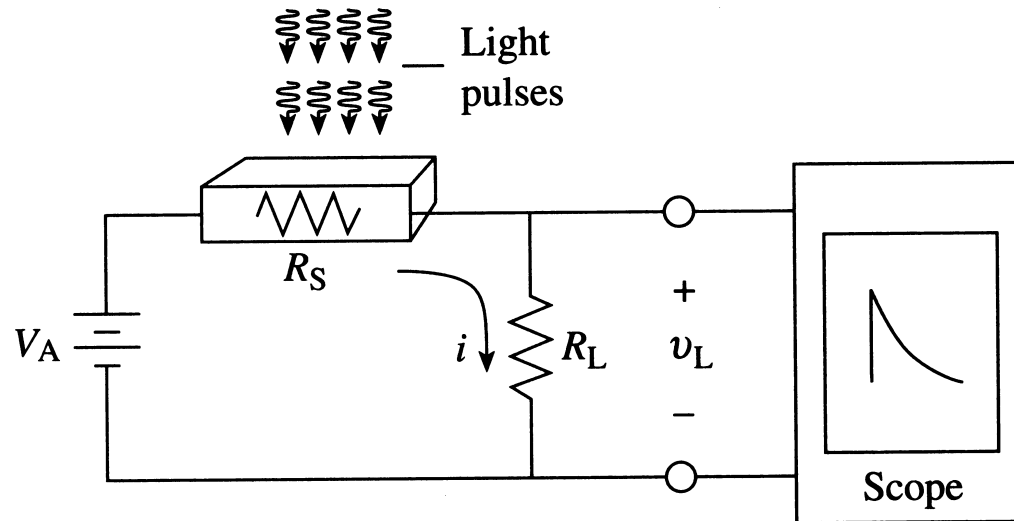
The solution for  $t > 0$  is:

$$\Delta p = \Delta p(0) \exp(-t / \tau_p)$$



The excess carrier concentration exponentially decays to zero when the external perturbation is removed. This fact is used to [measure lifetimes](#) using [photo-conductivity decay technique](#). See Sect. 3.3.4.

# Photoconductivity decay measurement



**Figure 3.22**

# Photoconductivity decay measurement system

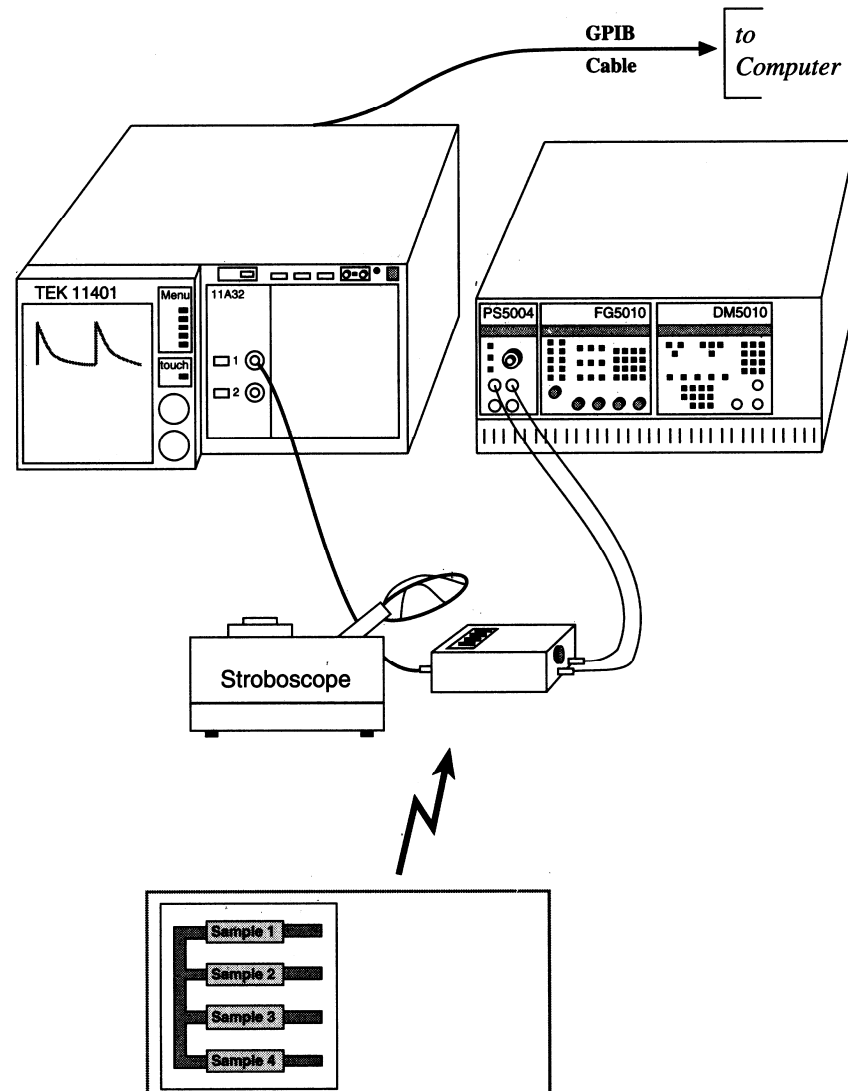
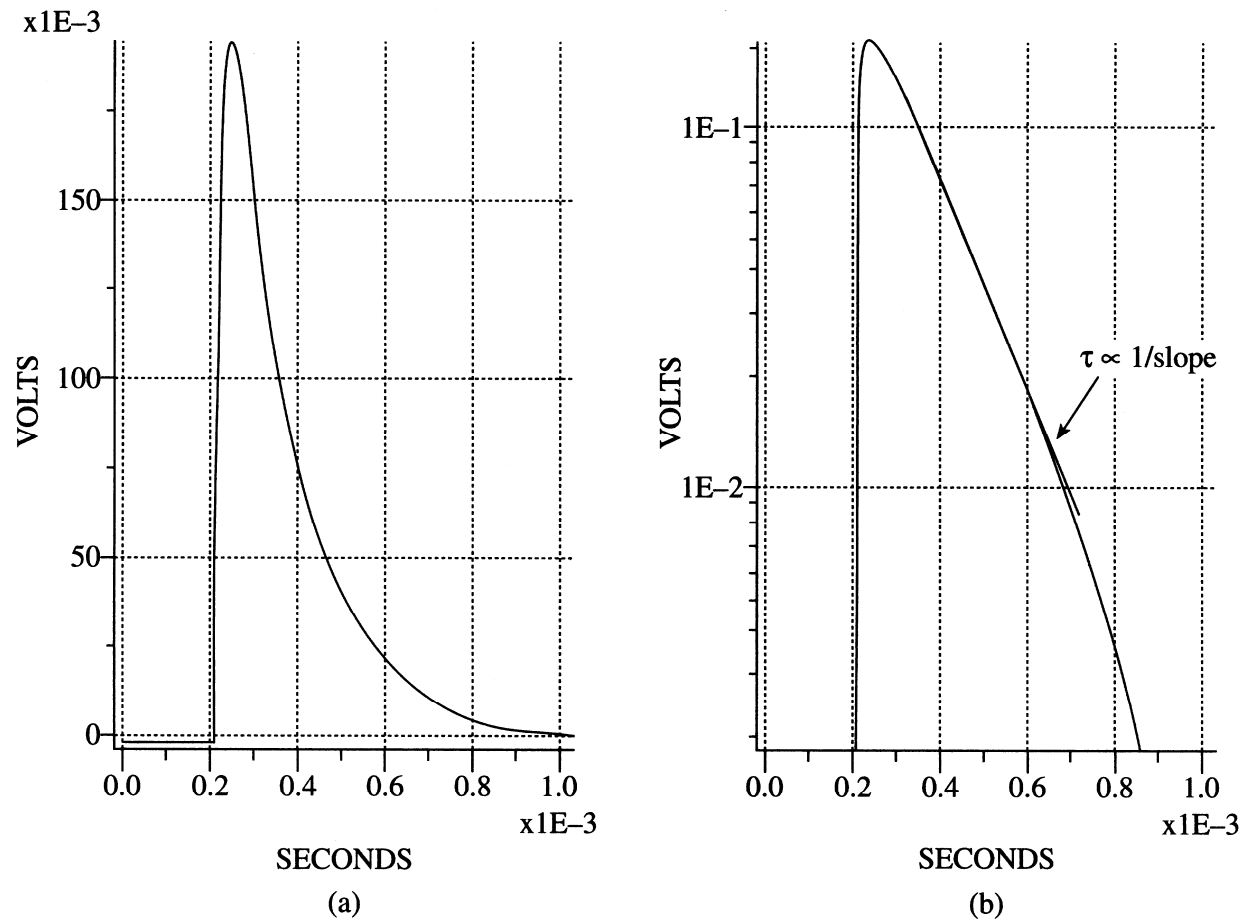


Figure 3.23

# Photoconductivity transient response



**Figure 3.24**