

Chapter 5-1. PN-junction electrostatics

In this chapter you will learn about pn junction electrostatics:
Charge density, electric field and electrostatic potential existing
inside the diode under equilibrium and steady state conditions.

You will also learn about:

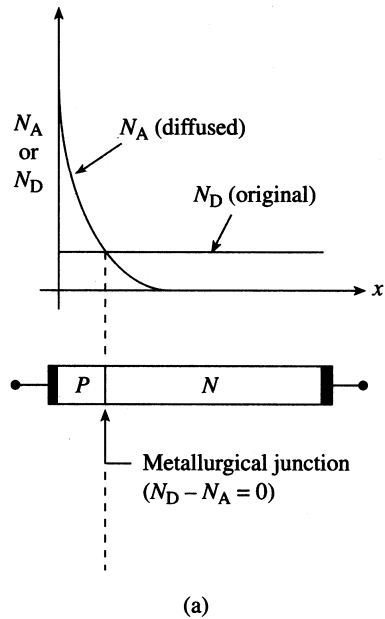
Poisson's Equation

Built-In Potential

Depletion Approximation

Step-Junction Solution

PN-junction fabrication



PN-junctions are created by several processes including:

1. Diffusion
2. Ion-implantation
3. Epitaxial deposition

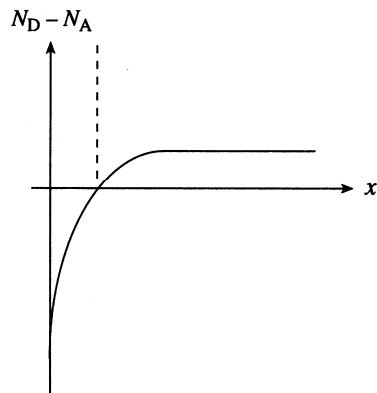
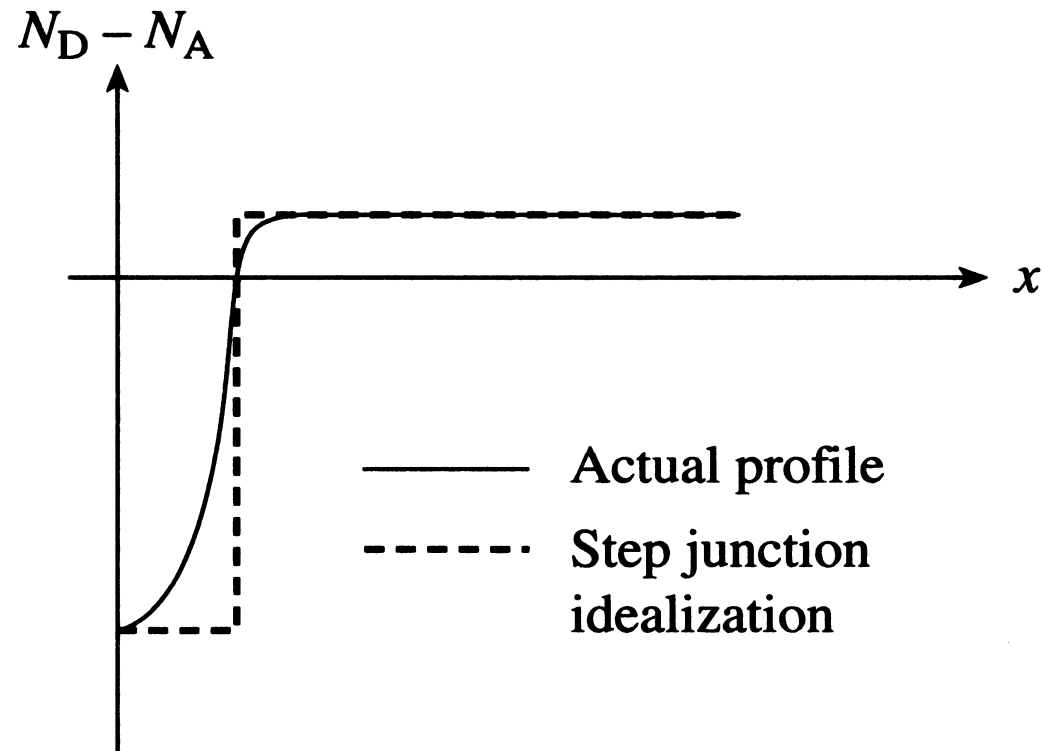


Figure 5.1 (b)

Each process results in different doping profiles

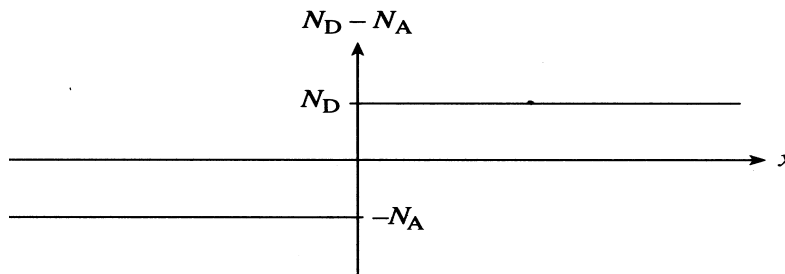
Ideal step-junction doping profile



(a)

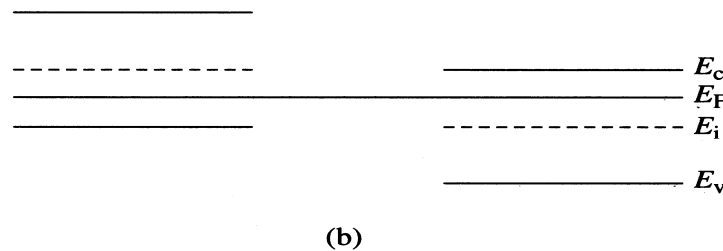
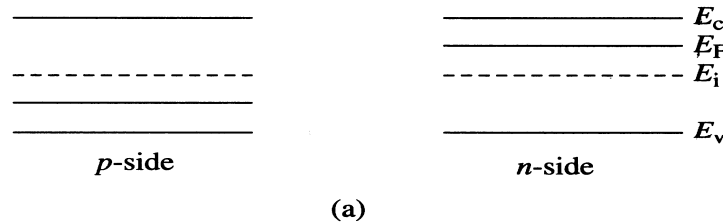
Figure 5.2

Equilibrium energy band diagram for the pn junction



$$n = n_i \exp\left(\frac{E_F - E_i}{kT}\right)$$

$$p = n_i \exp\left(\frac{E_i - E_F}{kT}\right)$$



E_F = same everywhere
under equilibrium

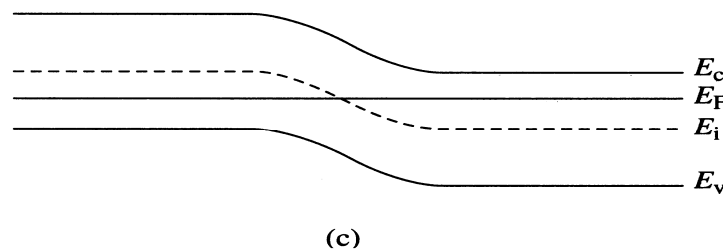
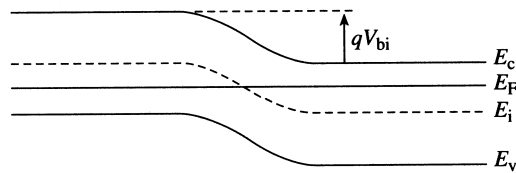


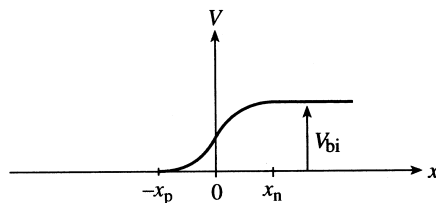
Figure 5.3

Join the two sides of the
band by a smooth curve.

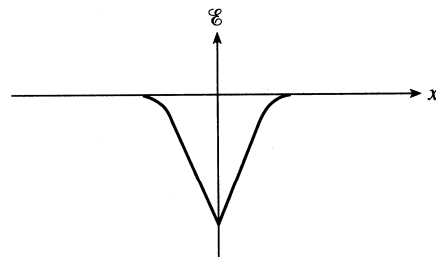
Electrostatic variables for the equilibrium pn junction



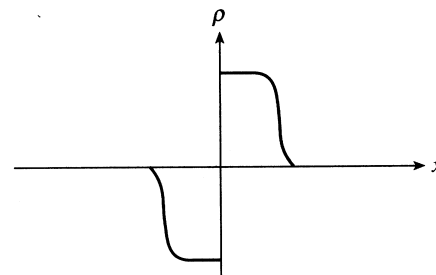
(a)



(b)



(c)



(d)

Figure 5.4

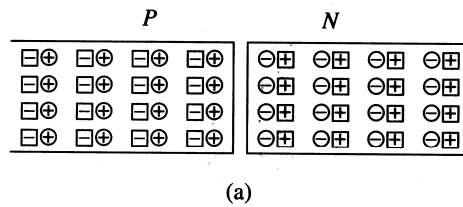
Potential, $V = - (1/q) (E_C - E_{\text{ref}})$. So, potential difference between the two sides (also called built-in voltage, V_{bi}) is equal to $-(1/q)(\Delta E_C)$.

$$V = -\frac{1}{q} (E_C - E_{\text{ref}})$$

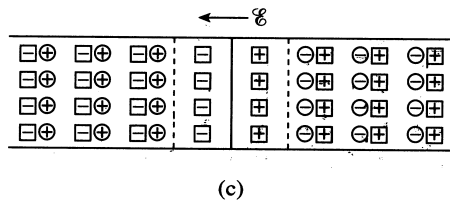
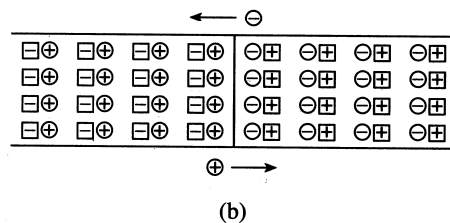
$$E = \frac{1}{q} \frac{dE_C}{dx} = \frac{1}{q} \frac{dE_i}{dx}$$

$$\left. \begin{aligned} \frac{dE}{dx} &= \frac{\rho}{\epsilon} \end{aligned} \right\} \begin{aligned} \rho &= \text{charge density} \\ \epsilon &= K_s \epsilon_0 \end{aligned}$$

Conceptual pn-junction formation



p and n type regions
before junction formation



Holes and electrons will diffuse
towards opposite directions,
uncovering ionized dopant atoms.
This will build up an electric field
which will prevent further
movement of carriers.

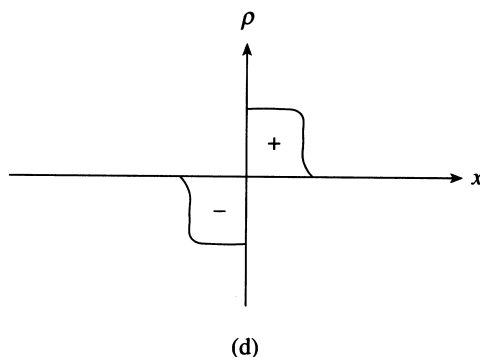


Figure 5.5

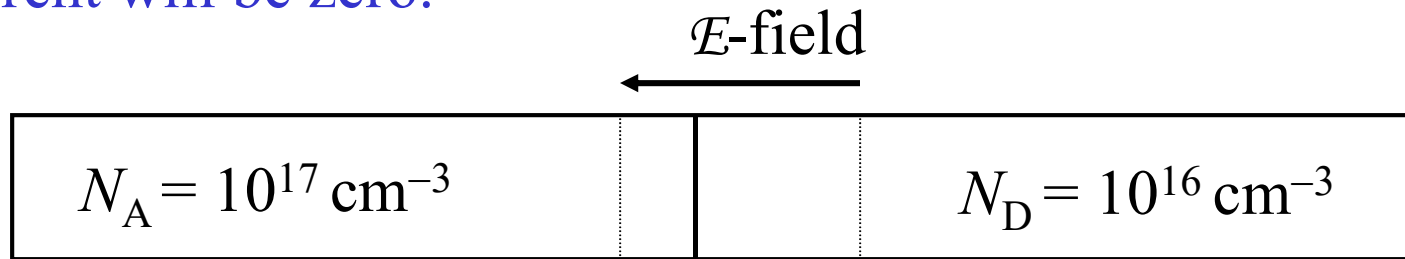
The built-in potential, V_{bi}

When the junction is formed, electrons from the n-side and holes from the p-side will diffuse leaving behind charged dopant atoms. Remember that the dopant atoms cannot move! Electrons will leave behind positively charged donor atoms and holes will leave behind negatively charged acceptor atoms.

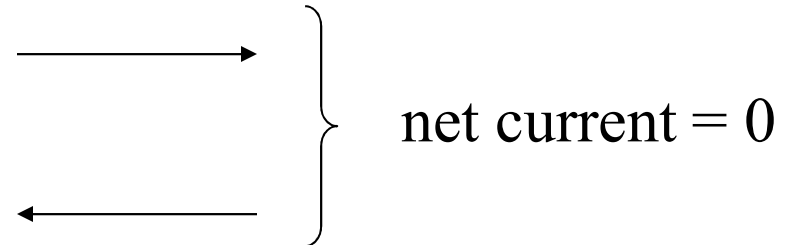
The net result is the build up of an electric field from the positively charged atoms to the negatively charged atoms, i.e., from the n-side to p-side. When steady state condition is reached after the formation of junction (how long this takes?) the net electric field (or the built in potential) will prevent further diffusion of electrons and holes. In other words, there will be **drift and diffusion currents such that net electron and hole currents will be zero.**

Equilibrium conditions

Under equilibrium conditions, the net electron current and hole current will be zero.

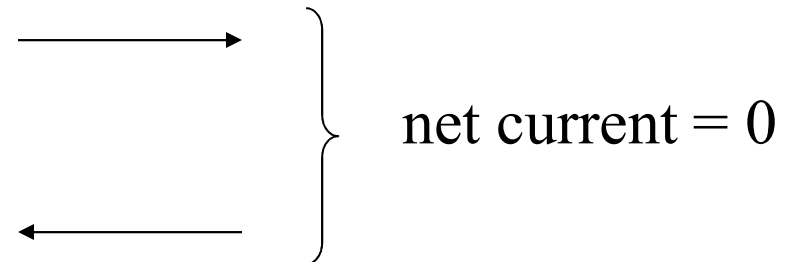


hole diffusion current



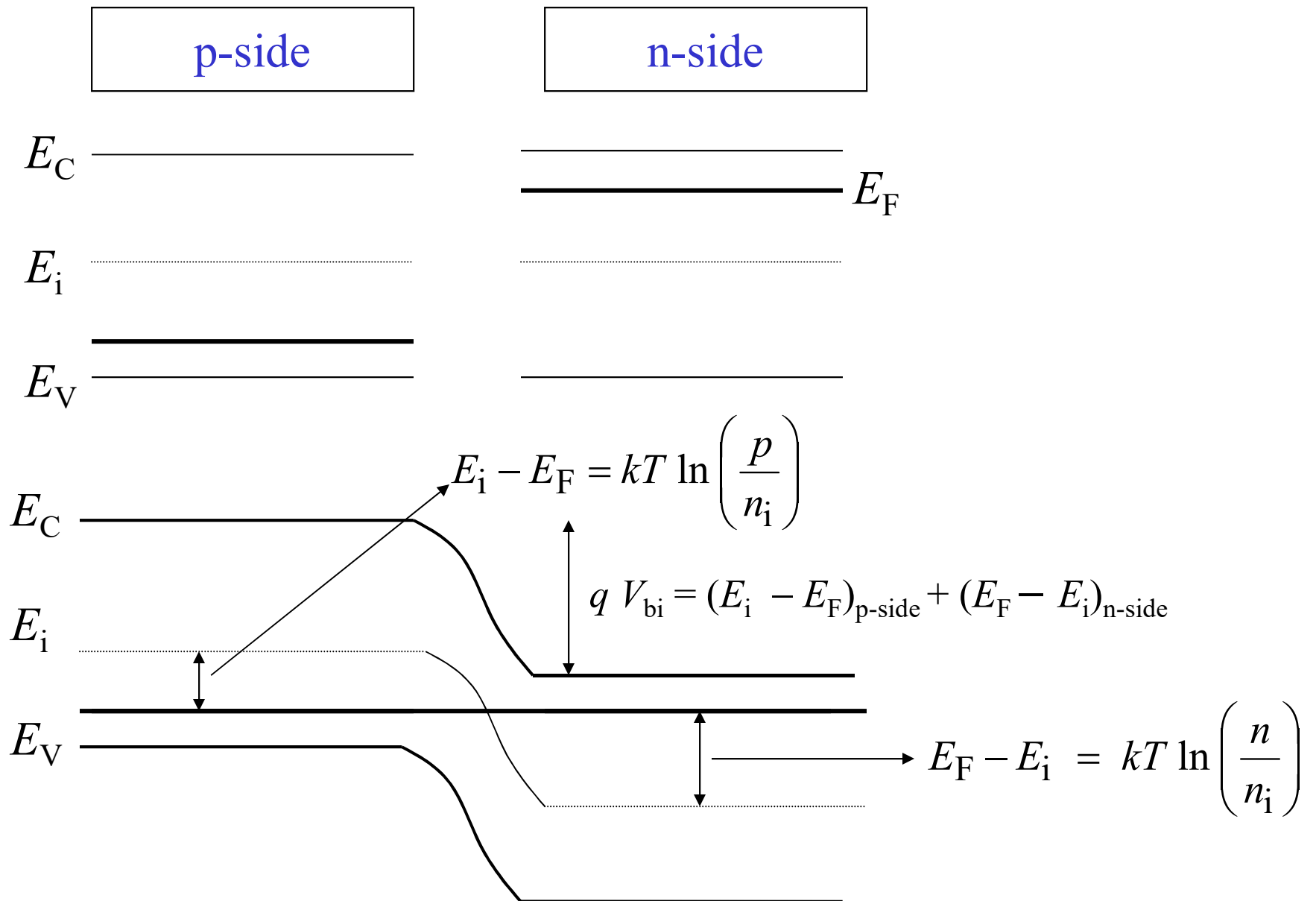
hole drift current

electron diffusion current
opposite to electron flux



electron drift current
opposite to electron flux

The built-in potential, V_{bi}



The built-in potential, V_{bi}

The built-in potential, V_{bi} , measured in Volts, is numerically equal to the “shift” in the bands expressed in eV.

$$\begin{aligned} V_{bi} &= (1/q) \{ (E_i - E_F)_{p\text{-side}} + (E_F - E_i)_{n\text{-side}} \} \\ &= \frac{kT}{q} \ln \left(\frac{p}{n_i} \right) + \frac{kT}{q} \ln \left(\frac{n}{n_i} \right) \\ &= \frac{kT}{q} \ln \left(\frac{p_p n_n}{n_i^2} \right) \end{aligned}$$

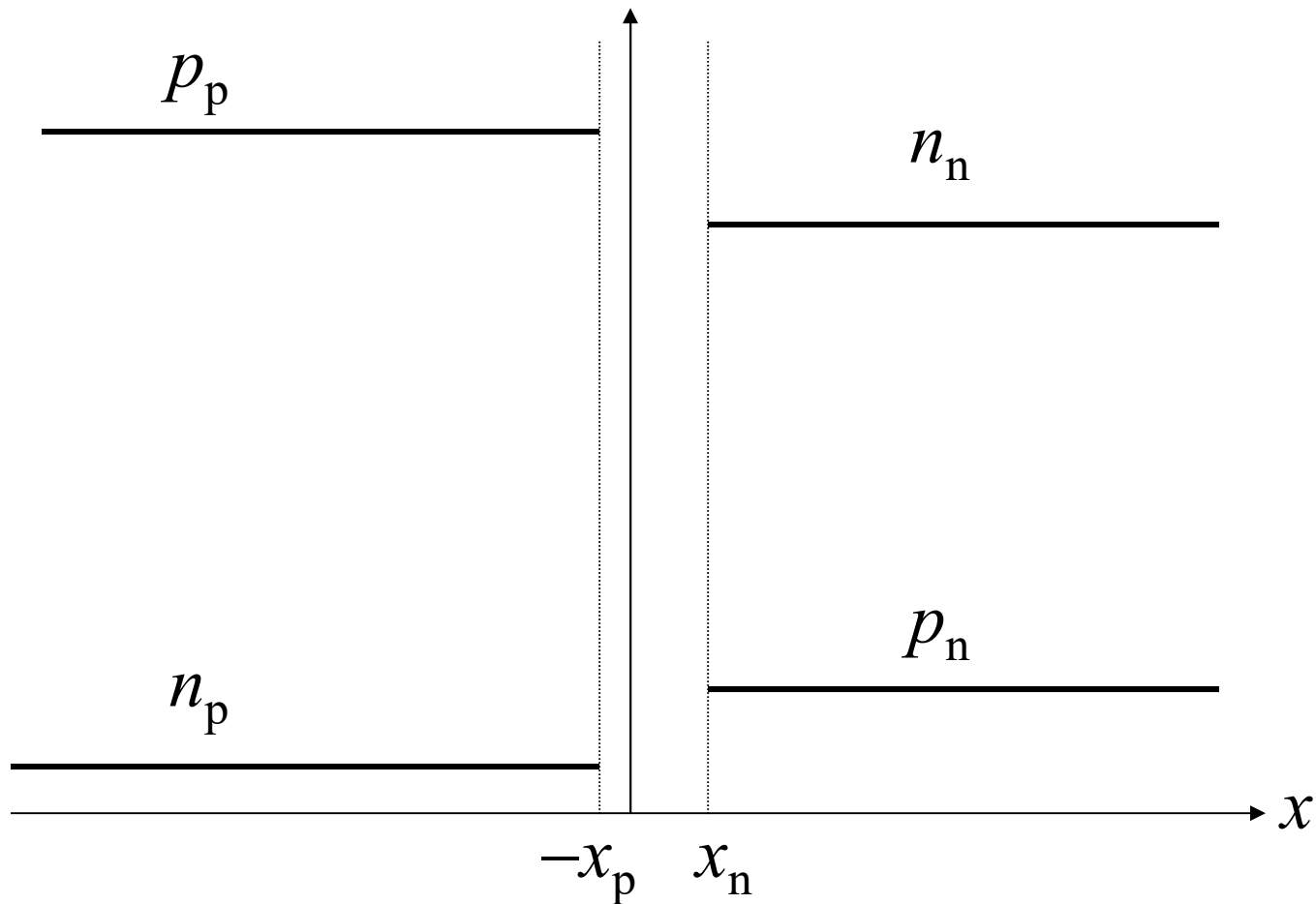
where p_p = hole – concentration on p – side

and n_n = electron – concentration on n – side

An interesting fact: $\frac{p_p}{p_n} = \frac{n_n}{n_p} = \exp \left(\frac{q V_{bi}}{kT} \right)$

Majority and minority carrier concentrations

p-side	N_A		N_D	n-side
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Built-in potential as a function of doping concentration for an abrupt p^+n or n^+p junction

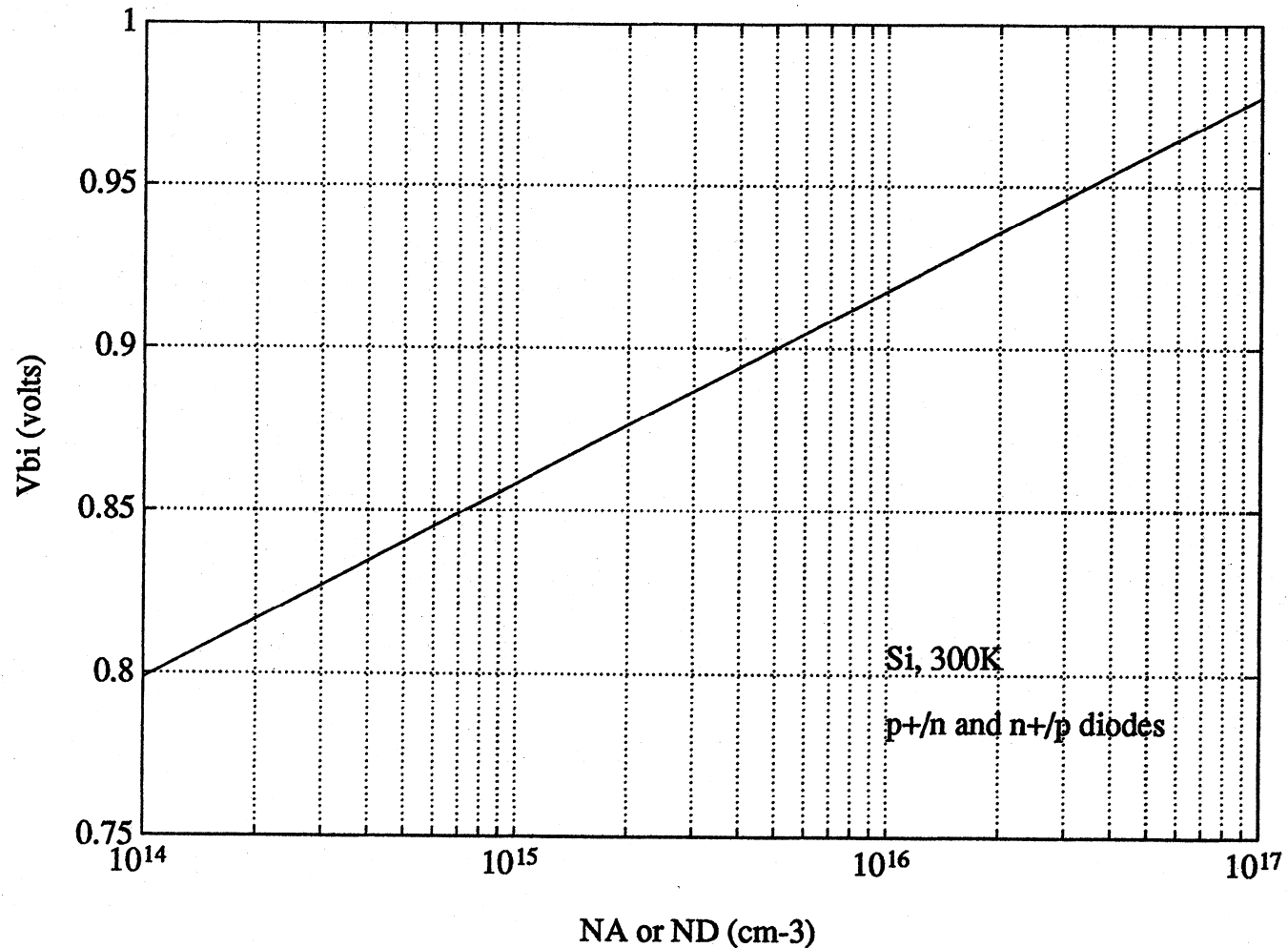
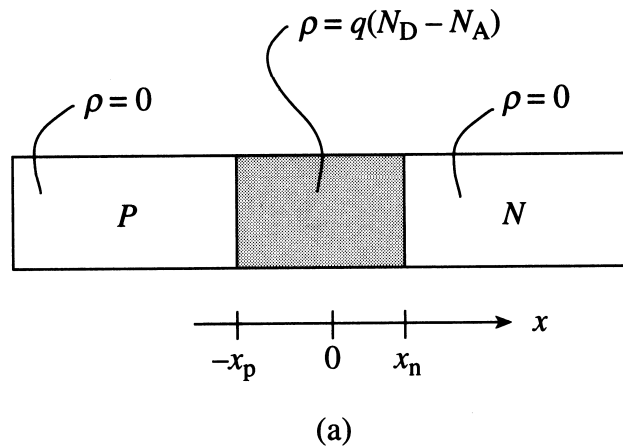


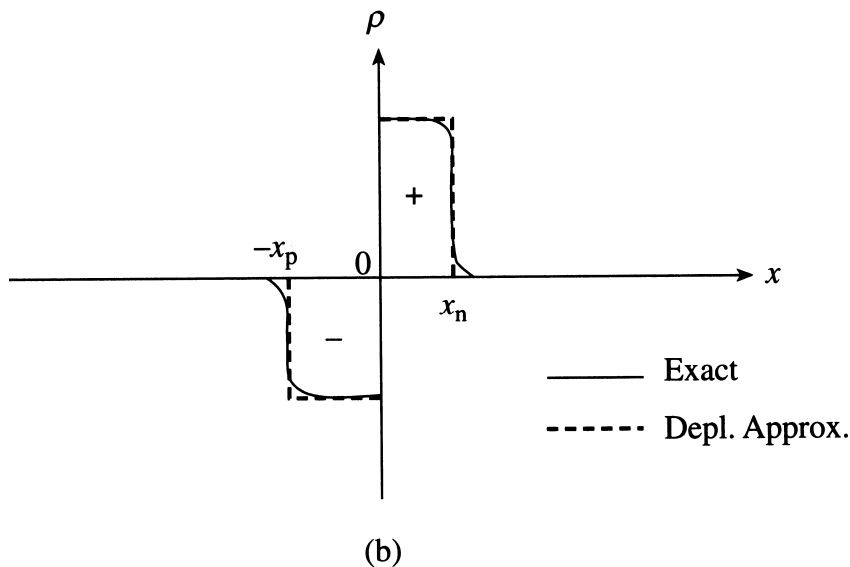
Figure E5.1

Depletion approximation



$$\frac{d\mathcal{E}}{dx} = \frac{\rho}{K_s \epsilon_0} \quad \text{Poisson equation}$$

$$= \begin{cases} \frac{q}{K_s \epsilon_0} (N_D - N_A) & \text{for } -x_p \leq x \leq x_n \\ 0 & \text{everywhere else} \end{cases}$$



We *assume* that the free carrier concentration inside the depletion region is zero.

Figure 5.6

Example 1

A p-n junction is formed in Si with the following parameters.
Calculate the built-in voltage, V_{bi} .

$N_D = 10^{16} \text{ cm}^{-3}$	$N_A = 10^{17} \text{ cm}^{-3}$
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Calculate majority carrier concentration in n-side and p-side.
Assume $n_n = N_D = 10^{16} \text{ cm}^{-3}$ and $p_p = N_A = 10^{17} \text{ cm}^{-3}$.

$$V_{bi} = \frac{kT}{q} \ln \left(\frac{p_p n_n}{n_i^2} \right) = \frac{kT}{q} \ln \left(\frac{N_A N_D}{n_i^2} \right)$$

Plug in the numerical values to calculate V_{bi} .

Example 2

A pn junction is formed in Si with the following parameters.
Calculate the built-in voltage, V_{bi} .

$N_D = 2 \times 10^{16} \text{ cm}^{-3}$ $N_A = 1 \times 10^{16} \text{ cm}^{-3}$	$N_A = 3 \times 10^{17} \text{ cm}^{-3}$ $N_D = 2 \times 10^{17} \text{ cm}^{-3}$
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Calculate majority carrier concentration in n-side and p-side.

$$n_n = \text{“effective } N_D \text{”} = 10^{16} \text{ cm}^{-3}; p_p = \text{“effective } N_A \text{”} = 10^{17} \text{ cm}^{-3}$$

$$V_{bi} = \frac{kT}{q} \ln \left(\frac{p_p n_n}{n_i^2} \right) = \frac{kT}{q} \ln \left(\frac{N_A N_D}{n_i^2} \right) \quad \left. \vphantom{\frac{kT}{q} \ln \left(\frac{N_A N_D}{n_i^2} \right)} \right\} \begin{array}{l} \text{Here } N_A \text{ and } N_D \\ \text{are “effective” or} \\ \text{net values.} \end{array}$$

Plug in the numerical values to calculate V_{bi} .