Search for a Bloch point using the Mean-field model

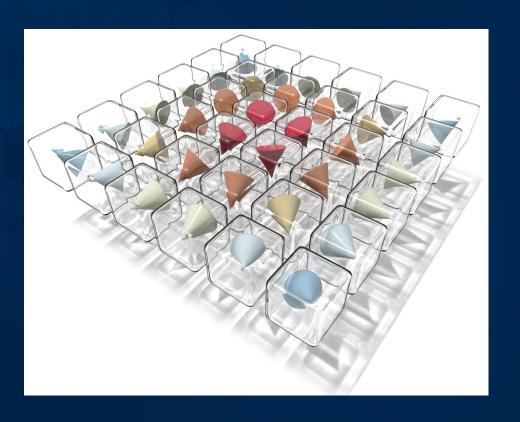
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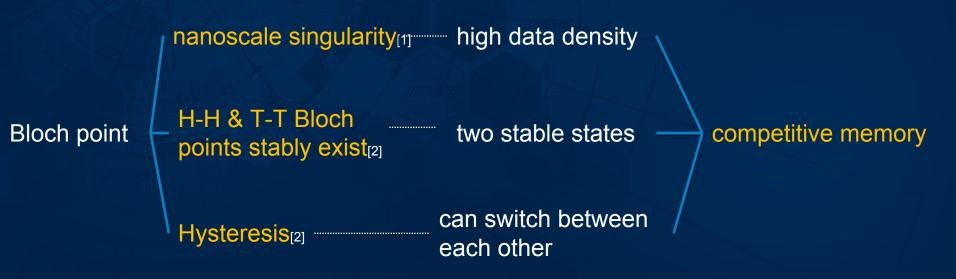
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PART 0 - MOTIVATION

WHY BLOCH POINT?



^[1] W Do ring. Point singularities in micromagnetism. Journal of Applied Physics, 39(2):1006–1007, 1968.

^[2] Beg, Marijan, et al. "Stable and manipulable Bloch point." Scientific reports 9.1 (2019): 1-8.

PART 0 - MOTIVATION

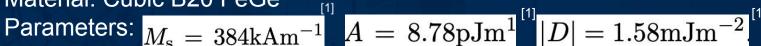
WHY SIMULATIONS?

- > More efficient
- > Easy to use
- Get results quickly

. . .

PART 1 - OBJECT OF STUDY

- Material: Cubic B20 FeGe



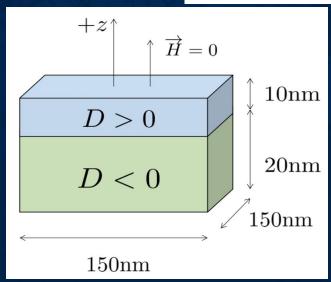
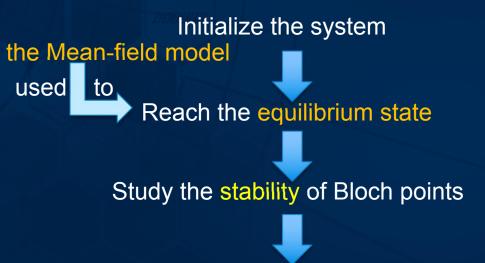


Fig1 The geometry of the studied system



Study the hysteretic behaviour of Bloch points

PART 2 - MEAN-FIELD MODEL

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Algorithm 1 Relax the system using the Mean-field model
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Require: temperature T, max iterations max_{iters} , changing rate λ

Require: Initialize the system with a uniform magnetisation configuration in +z or -z direction.

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iter \leftarrow 0
```

 $\mathbf{M}_{\mathsf{old}} \leftarrow \mathbf{M}_{\mathsf{ini}}$

do

Calculate effective field: $\mathbf{H}_{\mathrm{eff}} = \frac{2A}{\mu_{\mathrm{o}} \mathbf{M}} \nabla^2 \mathbf{m} - \frac{2D}{\mu_{\mathrm{o}} \mathbf{M}} \nabla \times \mathbf{m} + \mathbf{H}$ Part A

Update the length: $\mathbf{M}_{\mathsf{len}} = M_{\mathrm{s}} \cdot \mathcal{L} \left(eta \mu_0 \left| \mathbf{H}_{\mathrm{eff}}
ight|
ight) \cdot rac{\mathbf{H}_{\mathrm{eff}}}{\left| \mathbf{H}_{\mathrm{eff}}
ight|}$

where $\mathcal{L}(x)=\coth x-\frac{1}{x}$ is the Langevin function and the parameter $\beta=\frac{1}{\mathrm{K_b}T}$ introduces temperature to the algorithm.

Update the direction: $\mathbf{M}_{\mathsf{dir}} = \mathbf{M}_{\mathsf{old}} + \lambda \cdot (\mathbf{M}_{\mathsf{len}} - \mathbf{M}_{\mathsf{old}})$

Part B

Compute update: $\mathbf{M}_{\mathsf{new}} = |\mathbf{M}_{\mathsf{len}}| \cdot \frac{\mathbf{M}_{\mathsf{dir}}}{|\mathbf{M}_{\mathsf{dir}}|}$

Apply update: $M_{old} = M_{new}$

Update the counter of iterations: $iter \leftarrow iter + 1$

Calculate new effective field: $\mathbf{H}_{eff_{new}} = \frac{2A}{\mu_0 M_s} \nabla^2 \mathbf{m}_{new} - \frac{2D}{\mu_0 M_s} \nabla \times \mathbf{m}_{new} + \mathbf{H}$

while $iter \geq max_{iters}$ or $\mathbf{m}_{new} \times \mathbf{H}_{eff_{new}} = 0$

Part C

PART 2 - MEAN-FIELD MODEL

A. EFFECTIVE FIELD

Definition of effective field

$$\mathbf{H}_{ ext{eff}} = -rac{1}{\mu_0 ext{M}_{ ext{s}}} rac{\delta E(\mathbf{m})}{\delta \mathbf{m}}$$

Effective field points in the opposite direction of the system energy gradient.

Calculation of effective field

$$\mathbf{H}_{\mathrm{eff}} = egin{array}{c} \frac{2\mathrm{A}}{\mu_0 \mathrm{M}_{\mathrm{s}}} \nabla^2 \mathbf{m} & -\frac{2\mathrm{D}}{\mu_0 \mathrm{M}_{\mathrm{s}}} \nabla \times \mathbf{m} & + & \mathbf{H} \end{array}$$
 Exchange DMI Zeeman

PART 2 - MEAN-FIELD MODEL

B. Update Magnetization

- Calculate effective field: $\mathbf{H}_{\mathrm{eff}} = \frac{2\mathrm{A}}{\mu_0 \mathrm{M_s}} \nabla^2 \mathbf{m} \frac{2\mathrm{D}}{\mu_0 \mathrm{M_s}} \nabla \times \mathbf{m} + \mathbf{H}$
- Update the length: $\mathbf{M}_{\mathsf{len}} = M_{\mathsf{s}} \cdot \mathcal{L} \left(eta \mu_0 \left| \mathbf{H}_{\mathsf{eff}} \right| \right) \cdot rac{\ddot{\mathbf{H}}_{\mathsf{eff}}}{\left| \mathbf{H}_{\mathsf{eff}} \right|}$
- Update the direction: $\mathbf{M}_{\sf dir} = \mathbf{M}_{\sf old} + \lambda \cdot (\mathbf{M}_{\sf len} \mathbf{M}_{\sf old})$
- Compute update: $\mathbf{M}_{\mathsf{new}} = |\mathbf{M}_{\mathsf{len}}| \cdot rac{\mathbf{M}_{\mathsf{dir}}}{|\mathbf{M}_{\mathsf{dir}}|}$

Fig2: The Geometric schematic of updating magnetization

PART 2 - MEAN-FIELD MODEL

C. EQUILIBRIUM STATE

Energy minimization can be reprsent by the Brown's condition:

$$\mathbf{M} imes rac{\delta E}{\delta \mathbf{M}} = 0$$

The ideal stopping criteria of the simulation:

$$\mathbf{M} \times \mathbf{H}_{\text{eff}} = 0$$

PART 2 - MEAN-FIELD MODEL

C. STOP CRITERIA

- 1. Stop when arriving at the maximum iterations $iter > max_{iters}$
 - > represents the maximum computational cost we can accept
- 2. Stop when the magnetization and its effective field are in the same direction
 - If the algorithm convergent before the max iteration, we can save computing resources. $\mathbf{M} \times \mathbf{H}_{eff} = 0$
- 3. Stop according to the physical meaning of the Bloch point $107^{\circ} \le \theta_{\text{interface}} \le 180^{\circ}$
 - ➤ The Mean-field model cannot convergent when there is a singularity in the system, because the magnetization is discontinuous and not differentiable
 - when the angle of the magnetization at the interface is greater than 107 degrees.

PART 2 - MEAN-FIELD MODEL

ADVANTAGE

- High efficiency
 - along the defined direction (effective field)
 - > update all lattice at once
 - Introduce the influence of temperature
- > Show the detailed structures of Bloch points

RESULT: STABILITY

Two types of Bloch points appear at the interface, when the system is relaxed to the equilibrium state using the Mean-field model

PART 3 - RESULTS

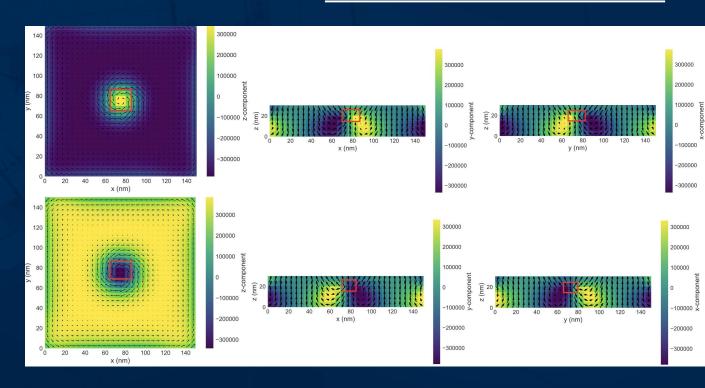


Fig3: The structures of two types of Bloch points

RESULT: STABILITY

The magnetization at the centre of the interface offset each other, resulting in a stable Head-to-Head Bloch point.

PART 3 - RESULTS

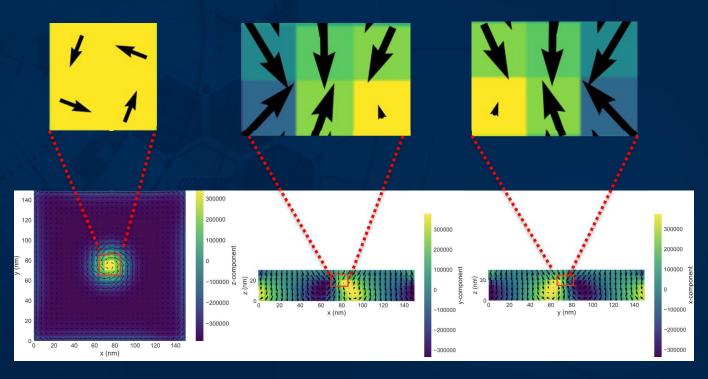


Fig4: The detailed structure of Head-to-Head Bloch points

PART 3 - RESULTS

RESULT: HYSTERESIS

- > H-H to T-T:
 - > increase to 1T
 - decrease to -0.1T
- > T-T to H-H:
 - decrease to -1T
 - > increase to 0.1T

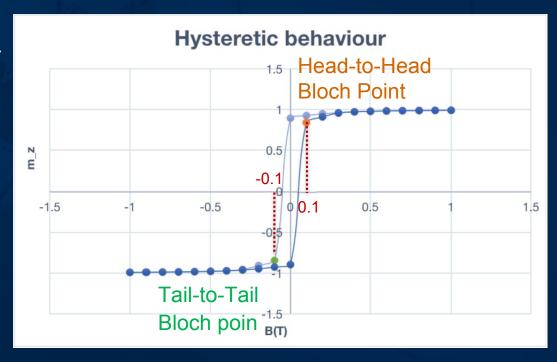


Fig5: The hysteretic behaviour of the system

PART 4 - CONCLUSION

- This project implements the Mean-field model
- Using the Mean-field model, the stability and hysteretic behaviour of Bloch points are studied.
 - Stability: the detailed structures of two Bloch points
 - Hysteresis: the two can be switched between when changing the external magnetic field
- ➤ In summary, this project demonstrates that Bloch point is a very promising topological quasiparticle for data storage.

Q&A

Thanks for listening!