

# Chapter 1

## Neural Encoding I: Firing Rates and Spike Statistics

### 1.1 1.1

**Principle 1.1** (Conservation of angular momentum). The rate of change of angular momentum of a system is equal to the net torque acting on the system, i.e.,

$$\frac{d\mathbf{L}}{dt} = \boldsymbol{\tau}, \quad (1.1)$$

where  $\boldsymbol{\tau}$  is the torque of all external forces on the system about any chosen axis, and  $d\mathbf{L}/dt$  is the rate of change of angular momentum of the system about the same axis.

### 1.2 1.2

**Remark 1.1.** Many physical laws are cumbersome when written in coordinate form but become more compact and attractive looking when written in tensorial form. For example, the incompressible Navier-Stokes equations in cylindrical coordinates are

$$\begin{aligned} \rho \left( \frac{Dv_r}{Dt} - \frac{v_\theta^2}{r} \right) &= \rho f_r - \frac{\partial p}{\partial r} + \mu \left( \Delta v_r - \frac{v_r}{r^2} - \frac{2}{r^2} \frac{\partial v_\theta}{\partial \theta} \right), \\ \rho \left( \frac{Dv_\theta}{Dt} + \frac{v_r v_\theta}{r} \right) &= \rho f_\theta - \frac{1}{r} \frac{\partial p}{\partial \theta} + \mu \left( \Delta v_\theta + \frac{2}{r^2} \frac{\partial v_r}{\partial \theta} - \frac{v_\theta}{r^2} \right), \\ \rho \frac{Dv_z}{Dt} &= \rho f_z - \frac{\partial p}{\partial z} + \mu \Delta v_z, \end{aligned}$$

where

$$\Delta = \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2} + \frac{\partial^2}{\partial z^2},$$

and

$$\frac{D}{Dt} = \frac{\partial}{\partial t} + v_r \frac{\partial}{\partial r} + \frac{v_\theta}{r} \frac{\partial}{\partial \theta} + v_z \frac{\partial}{\partial z}.$$

### 1.3 1.3

**Proposition 1.2.** The mass of fluid in a region  $W$  at time  $t$  is

$$m(W, t) = \int_W \rho(\mathbf{x}, t) dV, \quad (1.2)$$

where  $dV$  is the area element in the plane or the volume element in space.

### 1.4 1.4

**Assumption 1.3.** From now on, assume that

$$\text{force on } S \text{ per unit area} = -p(\mathbf{x}, t) \mathbf{n} + \mathbf{n} \cdot \boldsymbol{\sigma}(\mathbf{x}, t), \quad (1.3)$$

where  $\boldsymbol{\sigma}$  is the (*deviatoric*) *stress tensor* and  $\mathbf{n}$  is the unit outward normal of  $S$ .

## Chapter 2

# Neural Encoding II: Reverse Correlation and Visual Receptive Fields

### 2.1 2.1

**Principle 2.1** (Conservation of angular momentum). The rate of change of angular momentum of a system is equal to the net torque acting on the system, i.e.,

$$\frac{d\mathbf{L}}{dt} = \boldsymbol{\tau}, \quad (2.1)$$

where  $\boldsymbol{\tau}$  is the torque of all external forces on the system about any chosen axis, and  $d\mathbf{L}/dt$  is the rate of change of angular momentum of the system about the same axis.

### 2.2 2.2

**Remark 2.1.** Many physical laws are cumbersome when written in coordinate form but become more compact and attractive looking when written in tensorial form. For example, the incompressible Navier-Stokes equations in cylindrical coordinates are

$$\begin{aligned} \rho \left( \frac{Dv_r}{Dt} - \frac{v_\theta^2}{r} \right) &= \rho f_r - \frac{\partial p}{\partial r} + \mu \left( \Delta v_r - \frac{v_r}{r^2} - \frac{2}{r^2} \frac{\partial v_\theta}{\partial \theta} \right), \\ \rho \left( \frac{Dv_\theta}{Dt} + \frac{v_r v_\theta}{r} \right) &= \rho f_\theta - \frac{1}{r} \frac{\partial p}{\partial \theta} + \mu \left( \Delta v_\theta + \frac{2}{r^2} \frac{\partial v_r}{\partial \theta} - \frac{v_\theta}{r^2} \right), \\ \rho \frac{Dv_z}{Dt} &= \rho f_z - \frac{\partial p}{\partial z} + \mu \Delta v_z, \end{aligned}$$

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and

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### 2.3 2.3

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# Chapter 3

## Neural Decoding

### 3.1 3.1

**Principle 3.1** (Conservation of angular momentum). The rate of change of angular momentum of a system is equal to the net torque acting on the system, i.e.,

$$\frac{d\mathbf{L}}{dt} = \boldsymbol{\tau}, \quad (3.1)$$

where  $\boldsymbol{\tau}$  is the torque of all external forces on the system about any chosen axis, and  $d\mathbf{L}/dt$  is the rate of change of angular momentum of the system about the same axis.

### 3.2 3.2

**Remark 3.1.** Many physical laws are cumbersome when written in coordinate form but become more compact and attractive looking when written in tensorial form. For example, the incompressible Navier-Stokes equations in cylindrical coordinates are

$$\begin{aligned} \rho \left( \frac{Dv_r}{Dt} - \frac{v_\theta^2}{r} \right) &= \rho f_r - \frac{\partial p}{\partial r} + \mu \left( \Delta v_r - \frac{v_r}{r^2} - \frac{2}{r^2} \frac{\partial v_\theta}{\partial \theta} \right), \\ \rho \left( \frac{Dv_\theta}{Dt} + \frac{v_r v_\theta}{r} \right) &= \rho f_\theta - \frac{1}{r} \frac{\partial p}{\partial \theta} + \mu \left( \Delta v_\theta + \frac{2}{r^2} \frac{\partial v_r}{\partial \theta} - \frac{v_\theta}{r^2} \right), \\ \rho \frac{Dv_z}{Dt} &= \rho f_z - \frac{\partial p}{\partial z} + \mu \Delta v_z, \end{aligned}$$

where

$$\Delta = \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2} + \frac{\partial^2}{\partial z^2},$$

and

$$\frac{D}{Dt} = \frac{\partial}{\partial t} + v_r \frac{\partial}{\partial r} + \frac{v_\theta}{r} \frac{\partial}{\partial \theta} + v_z \frac{\partial}{\partial z}.$$

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