

# Phase Transitions in the Hard Spheres Model

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## Motivation

Why does a collection of water molecules suddenly freeze or boil when we slightly change its temperature? From the perspective of individual water molecules, very little seems to change from 99°C to 101°C. But macroscopically, we see a huge qualitative difference. More generally, even very simple mathematical models of collections of particles are observed to undergo dramatic *phase transitions* as a controlling parameter, such as temperature or density, is varied.

The most natural transition to study is from a gaseous phase, defined as one in which there are no long-distance correlations between the positions of particles, to a more viscous phase in which there are. Taking density,  $\rho$  as our controlling parameter, one tries to find the critical density,  $\rho_c$ , where  $\rho < \rho_c$  is a gas and  $\rho > \rho_c$  is not.

From a computer science perspective, if one wants to sample random configurations for this model it is important to know the running time of some algorithm that generates these random configurations. A very simple Markov chain Monte Carlo (MCMC) algorithm to sample random configurations in the model is the *Single-Sphere Global-Move dynamics* (described on page 2). It is believed that there is some  $\rho_c$  where for  $\rho < \rho_c$  the running time of the algorithm scales polynomially, and for  $\rho > \rho_c$  it scales exponentially.

There is some evidence that the answers to these two questions are related. Empirical evidence has shown that the algorithm is faster when the configuration is gaseous and slower when it is crystalline. But the values for these theoretical run-times are not known.

## Proposed Work

One such model for a collection of molecules is the *Hard Spheres Model*. In the model, there is a box with volume  $V$  and  $N$  non-overlapping spheres of volume  $v_s$  contained inside. The percentage of the volume that the spheres take up within the box is represented by  $\rho$ , that is,

$$\rho = \frac{Nv_s}{V}.$$

This  $\rho$  is also known as the density. Numerical results (see, e.g., [1][2]) show that the critical density for a phase transition to a crystalline configuration is  $\rho_c \approx 0.7$ .

## Single-Sphere Global-Move Dynamics

A dynamics for generating random configurations in the Hard Spheres Model is the *Single-Sphere Global-Move dynamics*. In this Markov chain, a sphere and location within the box are

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both picked uniformly at random. If the picked sphere can be placed at the picked location without overlapping another sphere, then it is moved there. Otherwise the move is rejected and the sphere remains in its original location. By a path-coupling argument (see, e.g., [3]) Kannan, Mahoney, and Montenegro [4] proved that  $\rho_c \geq 1/16$  for the running time transitioning from polynomial to exponential in two dimensions. More generally, for  $d$  dimensions  $\rho_c \geq 2^{-(d+1)}$ .

The basic idea for the coupling argument is that there are two configurations of spheres,  $X$  and  $Y$ , that differ only in the position of a single sphere. Then one wants to find the greatest value of  $\rho$  below which the Hamming distance between  $X$  and  $Y$  is expected to decrease. The typical Hamming metric for the coupling argument counts the number of spheres whose positions differ under  $X$  and  $Y$ . In other words, each sphere that disagrees contributes 1 to the sum.

Hayes and Moore [5] designed a modified Hamming metric that depends continuously on the difference between the two sphere positions. So a sphere whose two positions are very close contributes a small fraction less than 1. With this they improved the bound of the critical density from  $\rho_c \geq 0.125$  to  $\rho_c \geq 0.154483$  in the two dimensional setting.

## Event Chain Dynamics

In the *Event Chain dynamics* [6], a sphere and direction are both chosen uniformly at random. The sphere attempts to move in the picked direction a distance  $\delta$ . If the sphere collides with another sphere after only moving a distance  $\epsilon < \delta$ , then the sphere stops at the point of collision and the collided sphere attempts to move in the same picked direction for a distance  $\delta - \epsilon$ . This process is repeated as many times as needed.

The hope is that Event Chain dynamics will still converge quickly at higher densities where the rejected move rate of Single-Sphere Global-Move dynamics is high. The intuition is that Event Chain dynamics never rejects a move. By empirical observations the Event Chain dynamics appears to have a short run-time, but there are no theoretical guarantees; where as Single-Sphere Global-Move dynamics has theoretical guarantees but performs much slower at high densities.

## Open Questions

Proving some rigorous lower bound for Event Chain dynamics is a major open problem in the field. The dynamics is being used to produce random configurations but without any proofs about its running time. The only bearing for using the dynamics is because it appears to work well. With some accurate guarantees, there would be validity in using this dynamics over another.

In a slightly different model, there is an alternative parameter of interest, the *fugacity*  $\lambda$ , that controls the probability of a sphere entering or leaving the box. The number of spheres  $N$  within the box does not remain constant, which means  $\rho$  does not either. An interesting question to ask is given some fixed  $\lambda$  value, what would  $\rho$  converge to? If these two parameters have a strong relation, is there some dynamics involving  $\lambda$  that would allow short running times at high densities? The motivation for looking into this varied model is that problematic configurations may become easier to mix if one of the spheres has the probability to be removed; also the theoretical guarantees may be easier to analyze.

## References

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