

# Exploiting Non-Slip Wall Contacts to Position Two Particles Using The Same Control Input

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**Abstract**—Steered particles offer a method for targeted therapy, interventions, and drug delivery in regions inaccessible by large robots. Magnetic actuation has the benefits of requiring no tethers, being able to operate from a distance, and in some cases allows imaging for feedback (e.g. MRI). This paper investigates particle control with uniform magnetic gradients (the same force is applied everywhere in the workspace). Given three orthogonal magnetic fields, steering one particle in 3D is trivial. Adding additional particles to steer makes the system underactuated because there are more states than control inputs. However, the walls of in vivo and artificial environments often have surface roughness such that the particles do not move unless actuation pulls them away from the wall. In previous works, we showed that the individual 2D position of two particles is controllable in a square workspace with non-slip wall contact. Because in vivo environments are usually not square, this paper extends the previous work to convex workspaces including circles and 3D positioning. This paper also implements the algorithms using a hardware setup inspired by the gastrointestinal tract.

## I. INTRODUCTION

Particle swarms propelled by a uniform field, where each particle receives the same control input, are common in applied mathematics, biology, and computer graphics.

The small size of these robots makes it difficult to perform onboard computation. Instead, these robots are often controlled by a broadcast signal. The tiny robots themselves are often just rigid bodies, and it may be more accurate to define the *system*, consisting of particles, a uniform control field, and sensing, as the robot. Such systems are severely underactuated, having 2 degrees of freedom in the shared control input, but  $2n$  degrees of freedom for the  $n$ -particle swarm. Techniques are needed that can handle this underactuation. In previous work, we showed that the 2D position of each particle in such a swarm is controllable if the workspace contains a single obstacle the size of one particle [1].

Positioning is a foundational capability for a robotic system, e.g. placement of brachytherapy seeds. However, requiring a single, small, rigid obstacle suspended in the middle of the workspace is often an unreasonable constraint, especially in 3D. This paper relaxes that constraint, and provides position control algorithms that only require non-slip wall contacts. We assume that particles in contact with the boundaries have zero velocity if the uniform control input pushes the particle into the wall.

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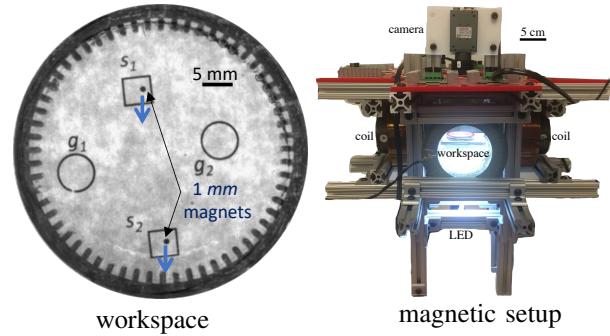


Fig. 1. Workspace and magnetic setup for an experiment of positioning particles that receive the same control inputs, but cannot move while a control input pushes them into a boundary.

The paper is arranged as follows. After a review of recent related work in Sec. II, Sec. III-A introduces a model for boundary interaction. We provide a shortest-path algorithm to arbitrarily position two particles in Sec. IV. Sec. V describes implementations of the algorithms in simulation and Sec. VI describes hardware experiments, as shown in Fig. 1. We end with directions for future research in Sec. VII.

This paper is elaboration of preliminary work in a conference paper [2] which considered only square workspaces. This work extends the analysis to convex workspaces and 3D positioning. This paper also implements the algorithms using a hardware setup inspired by the anatomy of the gastrointestinal tract.

## II. RELATED WORK

Controlling the *shape*, or relative positions, of a swarm of robots is a key ability for a range of applications. Correspondingly, it has been studied from a control-theoretic perspective in both centralized and decentralized approaches. For examples of each, see the centralized virtual leaders in [3], and the gradient-based decentralized controllers using control-Lyapunov functions in [4]. However, these approaches assume a level of intelligence and autonomy in individual robots that exceeds the capabilities of many systems, including current micro- and nano-robots. Current micro- and nano-robots, such as those in [5]–[7] lack onboard computation.

Instead, this paper focuses on centralized techniques that apply the same control input to both particles. Precision control requires breaking the symmetry caused by the uniform input. Symmetry can be broken using particles that respond differently to the uniform control signal, either through

agent-agent reactions [8], or engineered inhomogeneity [9]–[11]. This work assumes a uniform control with homogenous particles, as in [12], and breaks the control symmetry using obstacles in the workspace. The magnetic gradients of MRI scanners are uniform in the workspace, [13].

Alternative techniques rely on non-uniform inputs, such as artificial force-fields. Applications have included techniques to design shear forces for sensorless manipulation of a single object by [14]. [15] demonstrated a collection of 2D force fields generated by six degree-of-freedom vibration inputs to a rigid plate. These force fields, including shear forces, could be used as a set of primitives for motion control to steer the formation of multiple objects. However unlike the uniform control model in this paper, their control was multi-modal and position-dependent.

Much recent work has focused on exploiting inhomogeneities in the magnetic field to control multiple micro particles using gradient-based pulling [16], [17]. Unfortunately, using large-scale external magnetic fields makes it challenging to independently control more than one micro-robot unless the distance between the electromagnetic coils is at the same length scales as the robot workspace [16]–[18]. In contrast this paper requires only a controllable constant gradient in orthogonal directions to position the particles.

### III. THEORY

If a control input causes the particles to collide with obstacles at different times, inverting the control input does not undo the action. Due to this lack of time-reversibility, techniques that require a bidirectional graph, e.g. PRM and RRT\* are not suitable [19], [20]. Instead, this paper employs a graph search. This section starts with a boundary interaction model in subsection III-A.

Our algorithms rely on holding one particle stationary by pushing it into the boundary while moving the other particle. In subsections III-B and III-C we provide shortest-path results for two representative workspaces, squares and circles.

#### A. Boundary Interaction Model

In the absence of obstacles, uniform inputs move a swarm identically. Independent control requires breaking this symmetry. The following sections examine using non-slip boundary contacts to break the symmetry caused by uniform inputs.

If the  $i^{\text{th}}$  particle has position  $\mathbf{x}_i(t)$  and velocity  $\dot{\mathbf{x}}_i(t)$ , we assume the following system model:

$$\dot{\mathbf{x}}_i(t) = \mathbf{u}(t) + F(\mathbf{x}_i(t), \mathbf{u}(t)), \quad i \in [1, n]. \quad (1)$$

$$F(\mathbf{x}_i(t), \mathbf{u}(t)) = \begin{cases} -\mathbf{u}(t) & \mathbf{x}_i(t) \in \text{boundary and} \\ & \mathbf{N}(\text{boundary}(\mathbf{x}_i(t))) \cdot \mathbf{u}(t) \leq 0 \\ 0 & \text{else.} \end{cases}$$

Here  $F(\mathbf{x}_i(t), \mathbf{u}(t))$  is the frictional force provided by the boundary, and  $\mathbf{N}(\text{boundary}(\mathbf{x}_i(t)))$  is the normal to the boundary at position  $\mathbf{x}_i(t)$ .

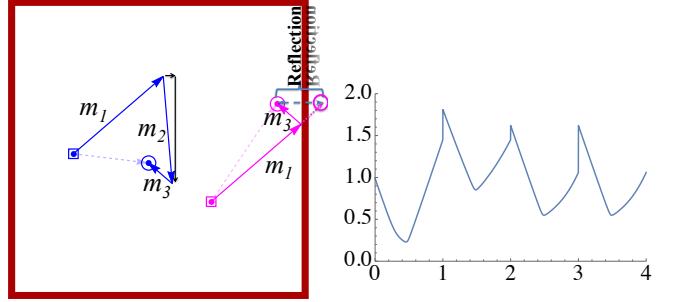


Fig. 2. The shortest ?three-move path that reconfigures two particles has the property that the incident angle equals the reflected angle.

These system dynamics represent particle swarms in low-Reynolds number environments, where viscosity dominates inertial forces and so velocity is proportional to input force [21]. In this regime, the input force command  $\mathbf{u}(t)$  controls the velocity of the particles. The same model can be generalized to particles moved by fluid flow where the vector direction of fluid flow  $\mathbf{u}(t)$  controls the velocity of particles, or for a swarm of particles that move at a constant speed in a direction specified by a uniform input  $\mathbf{u}(t)$  [22]. As in our model, fluid flowing in a pipe has zero velocity along the boundary. Similar mechanical systems exist at larger scales, e.g. all tumblers of a combination lock move uniformly unless obstructed by an obstacle. Our control problem is to design the control inputs  $\mathbf{u}(t)$  to make all  $n$  particles achieve a task.

#### B. Example: Shortest Path in a Square Workspace

If the goal configuration cannot be reached in one move but can be reached in three moves, the shortest path has a simple solution. The first move,  $m_1$ , makes one particle hit a wall,  $m_2$  adjusts the relative spacing error to zero, and  $m_3$  takes the particles to their final positions.  $m_2$  cannot be shortened, so optimization depends on choosing the location where the particle hits the wall. Since the shortest distance between two points is a straight line, reflecting the goal position across the boundary wall and plotting a straight line gives the optimal hit location, as shown in Fig. 2. There are four walls, and four candidate solutions.

#### C. Shortest Path in Unit Disk that Intersects Circumference

The shortest path between two points in the unit disk that intersects the circumference is composed of two straight line segments, as shown in Fig. 3. The problem can be simplified by choosing the coordinate system carefully. We define the  $x$  axis along the position of the starting point:  $S = (s, 0)$ , and define the point of intersection by the angle  $\theta$  from the  $x$  axis:  $P = (\cos \theta, \sin \theta)$ . Define the final point  $E$  by a radius  $e$  and angle  $\beta$ :  $E = e(\cos \beta, \sin \beta)$ . Then the length of the two line segments is

$$\sqrt{(s - \cos \theta)^2 + \sin^2 \theta} + \sqrt{(e \sin \beta - \cos \theta)^2 + (e \sin \beta - \sin \theta)^2}, \quad (2)$$

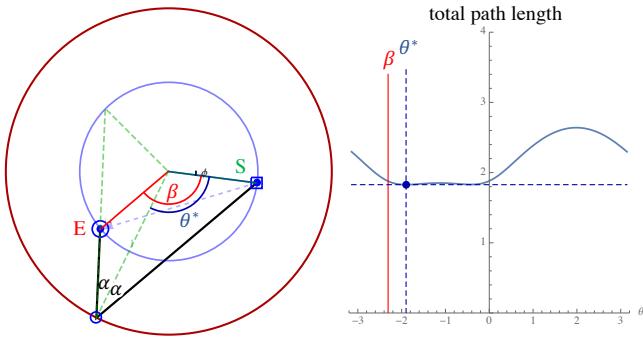


Fig. 3. The shortest path between two points (blue square) to (blue ellipse) in the unit disk that intersects the circumference. The path length as a function of intersection point,  $(\cos \theta, \sin \theta)$  is shown at right. See [23].

which is minimized by choosing an appropriate  $\theta$  value. This equation can be simplified to

$$\sqrt{1 + s^2 - 2s \cos \theta} + \sqrt{1 + e^2 - 2e \cos(\beta - \theta)}. \quad (3)$$

The length of the two line segments as a function of  $\theta$  is drawn in the right plot of Fig. 3. There are several simple solutions. If  $s$  is 1 or  $e$  is 0 or  $\beta$  is 0, the optimal angle  $\theta^*$  is 0. If  $e$  is 1 or  $s$  is 0, the optimal angle is  $\beta$ . Label the origin  $O$ . The optimal solution shows that the angle  $\angle OPS$  (from the origin to  $P$  to  $S$ ) is the same as the angle  $\angle OPE$  (from the origin to  $P$  to  $E$ ). We name these angles  $\alpha$ . This can be proved by drawing an ellipse whose foci are  $S$  and  $E$ . When the ellipse is tangent to the circle, the point of tangency is exactly  $P$ . Since the distance from the origin to  $P$  is always 1, we can set up three equalities using the law of sines: From triangle  $OSP$ :  $\frac{\sin \alpha}{s} = \frac{\sin(\alpha+\theta)}{1} = \frac{\sin \theta}{\|SP\|}$ , and from triangle  $OEP$ :  $\frac{\sin \alpha}{e} = \frac{\sin(\beta-\theta)}{\|EP\|}$ . If we mirror the point  $S$  about the  $\theta$  axis and label this point  $C$ , from triangle  $CEO$ :  $\frac{\sin(\alpha+\theta)}{e} = \frac{\sin(2\theta-\beta)}{\|CE\|}$ .

Simplifying this system of equations results in  $s = e \csc \theta (s \sin(2\theta - \beta) + \sin(\beta - \theta))$ . Solving this last equation results in a quartic solution that has a closed-form solution with four roots, each of which can be either a clockwise or a counterclockwise rotation  $\theta$ , depending on the sign of  $\beta$ , with  $-\pi \leq \beta \leq \pi$ . We evaluate each and select the solution that results in the shortest length path. This optimal path satisfies the law of reflection off the unit circle, with angle of incidence equal to angle of reflection. For an interactive Mathematica demonstration of this shortest path, see [23].

#### IV. POSITION CONTROL OF TWO PARTICLES USING BOUNDARY INTERACTION

This section presents algorithms that use non-slip contacts with walls to arbitrarily position two particles in a convex workspace. Our previous work used a square workspace [2]. Alg. 1 can now handle any convex workspace, including the special limit case of a circular workspace. In the last subsection we present techniques to control 3D positioning of two particles.

Workspaces are 2D convex polygons with no internal obstacles. Assume two particles are initialized at  $s_1$  and  $s_2$  with corresponding goal destinations  $g_1$  and  $g_2$ . Denote the current positions of the particles  $p_1$  and  $p_2$

##### A. Two Particle Path Planning

The configuration space for two particles is a four dimensional manifold. Translating both particles the same amount is a trivial operation, but changing the relative positions requires boundary interaction. For this reason, our algorithms use the two dimensional  $\Delta$  configuration space, defined as the difference in position of the first particle from the second particle:  $\Delta p = p_2 - p_1$ . Values  $.x$  and  $.y$  denote the  $x$  and  $y$  coordinates, i.e.,  $p_1.x$  and  $p_1.y$  denote the  $x$  and  $y$  locations of  $p_1$ . Alg. 1 assigns a uniform control input at every instance. The goal is to move the particles within  $\epsilon$  of the goal positions using a shared control input where  $\epsilon$  is an arbitrary small number. We do this by first moving them within  $\epsilon$  of the correct relative position and then translating the particles to the goal. The relative position is  $\|\Delta g - \Delta p\| = \|(g_2 - g_1) - (p_2 - p_1)\|$ .

The  $\Delta$  configuration space is a set of all possible  $\Delta p$  values.  $\Delta$  configuration spaces for a representative set of workspaces are shown in Fig. 4. The *reachable set* is the locus of points in the  $\Delta$  configuration space corresponding to any two-move sequence where the first move brings one particle into contact with the boundary, and the second move translates the second particle without moving the first. Fig. 7 shows the starting and ending relative distance as  $\Delta s$  and  $\Delta g$  in the  $\Delta$  configuration space. The next subsections give procedures to compute the reachable set.

Alg. 1 first computes the reachable set. If the goal relative position is in the reachable set, we move particles to achieve relative position. If it is not in the reachable set, we move particles to achieve the closest point on the reachable set from  $\Delta g$ . Achieving a  $\Delta$  configuration requires two-moves, the first to push until one particle touches a wall, and the second to adjust the relative spacing. Once the correct *relative* position has been achieved, a final translation delivers both particles to their goal destinations. Otherwise, we iterate until we reach the goal.

##### B. Convex Polygonal Workspaces: Reachable Set

Fig. 4 shows different workspaces and their representative  $\Delta$  configuration spaces.

If a particle is touching a wall, the other particle can move freely in the reachable set as shown in Fig. 5. Alg. 2 computes the reachable set for any convex workspace. Fig. 6 illustrates the procedure to construct the reachable set generated by collisions with the  $i^{\text{th}}$  side. If one particle hits side  $i$  before the other ( $\overline{s_1 s_2} \nparallel \overline{p_i p_{i+1}}$ ), the reachable set is defined by a polygon, constructed in lines 2-13 of Alg. 2. The union of these polygons for all  $n$  sides is the reachable set of  $\Delta$  configurations.

##### C. Circular Workspaces: Reachable Set

To compute the reachable set for a circular workspace, first we consider all possible first contact locations. The set of

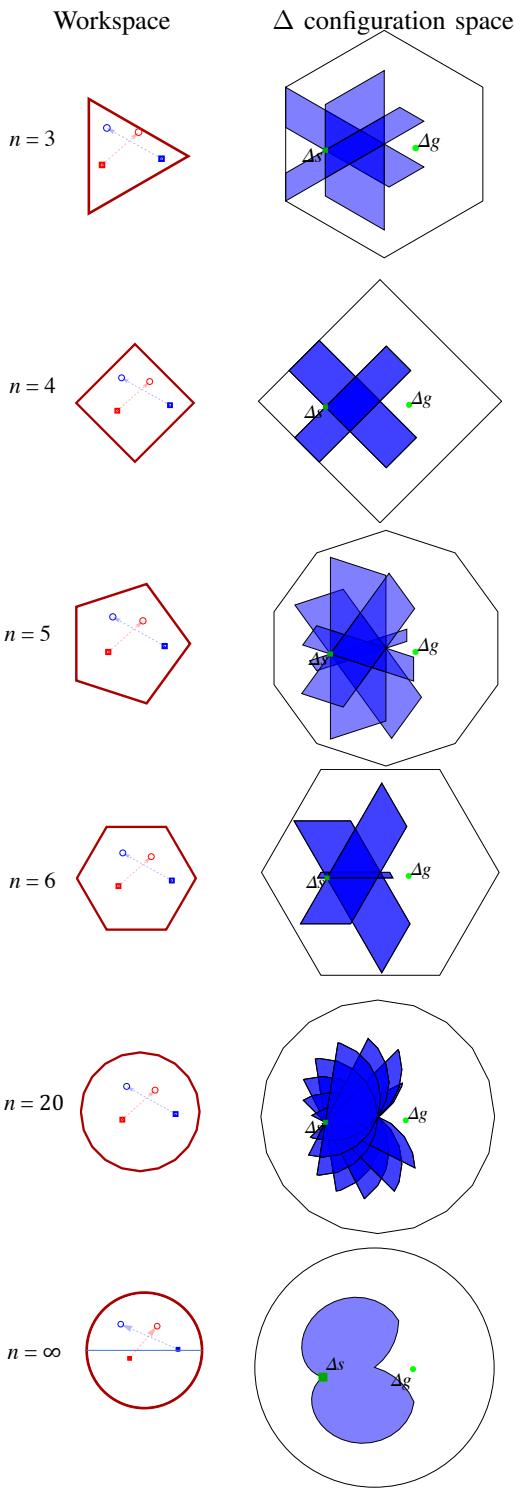


Fig. 4. Reachable sets are drawn with transparent blue polygons. A polygon with  $n$  sides has  $n$  reachable sets, but if  $n$  is even and the polygon is regular, half the reachable sets overlap.

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**Algorithm 1** 2-PARTICLEPATHPLANNING( $s_1, s_2, g_1, g_2, P, \epsilon$ )

**Require:** knowledge of starting  $(s_1, s_2)$  and goal  $(g_1, g_2)$   
positions of two particles.  $P$  is a description of the  
workspace.  $\epsilon$  is a positive error bound.

- 1:  $(p_1, p_2) \leftarrow (s_1, s_2)$   $\triangleright p_1, p_2$  are current positions
- 2: moves  $\leftarrow \{\}$
- 3:  $\Delta p \leftarrow p_2 - p_1$
- 4:  $\Delta g \leftarrow g_2 - g_1$
- 5: **while**  $||\Delta p - \Delta g|| > \epsilon$  **do**
- 6:      $R_{SET} \leftarrow$  Compute reachable set  $\triangleright$  use Alg. 2 or 3
- 7:      $\Delta g_c \leftarrow$  nearest point in  $R_{SET}$  to  $\Delta g$
- 8:      $m \leftarrow$  move-to-wall corresponding to  $\Delta g_c$
- 9:     moves  $\leftarrow$  Append  $m$  to moves
- 10:     $(p_1, p_2) \leftarrow$  ApplyMove  $m$  to  $(p_1, p_2)$
- 11:     $\Delta p \leftarrow p_2 - p_1$
- 12: **end while**
- 13: moves  $\leftarrow$  Append  $g_2 - p_2$  to moves  $\triangleright$  translate to goal
- 14: **return** moves

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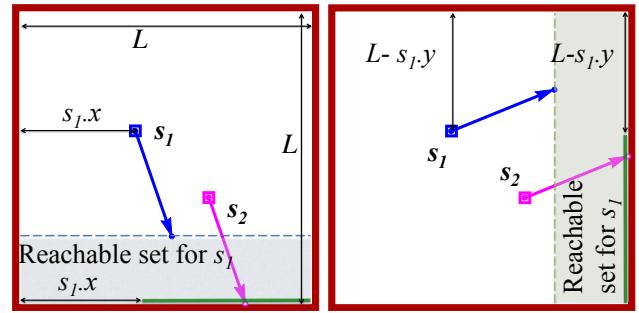


Fig. 5. Boundary interaction is used to change the relative positions of the particles. Each particle gets the same control input. (left) If particle 2 hits the bottom wall before particle 1 reaches a wall, particle 2 can reach anywhere along the green line, and particle 1 can move to anywhere in the shaded area. (right) Similarly, if particle 2 hits the right wall before particle 1 reaches a wall, particle 2 can reach anywhere along the green line, and then particle 1 can move to anywhere in the shaded area.

boundary points that a particle can touch before the other particle touches is an arc of angle  $2(\pi - \frac{\arcsin d_{12}}{r})$ , where  $d_{12} = ||s_1 - s_2||$  and  $r$  is the radius of the workspace. We define the angle between two particles as  $\theta = \arctan(\frac{p_1.x - p_2.x}{p_1.y - p_2.y})$ .

A circle has an infinite number of sides, thus infinite reachable sets. However, the reachable set can be parameterized by the angle of first contact location  $\psi$ , as shown in Fig. 7 where

$$\psi \in [\psi_{\min}, \psi_{\max}] = \theta + \left[ \frac{\sin^{-1} d}{2r} - \frac{\pi}{2}, \frac{\pi}{2} - \frac{\sin^{-1} d}{2r} \right]. \quad (4)$$

The possible first contact locations are on an arc with interior angle  $\gamma$  parameterized by  $\psi$ :

$$\gamma(\psi) = \cos^{-1} \left( 1 - \frac{d_{\perp}(\psi)}{r} \right), \text{ where:} \quad (5)$$

$$d_{\perp}(\psi) = 2||s_1.p_{\psi}(\psi) - s_2.p_{\psi}(\psi)||, \quad (6)$$

$$p_{\psi}(\psi) = r[\cos(\psi), \sin(\psi)]. \quad (7)$$

Reachable sets with  $\pi$  difference in  $\psi$  value are equivalent in the  $\Delta$  configuration space, so we can plan in this space and

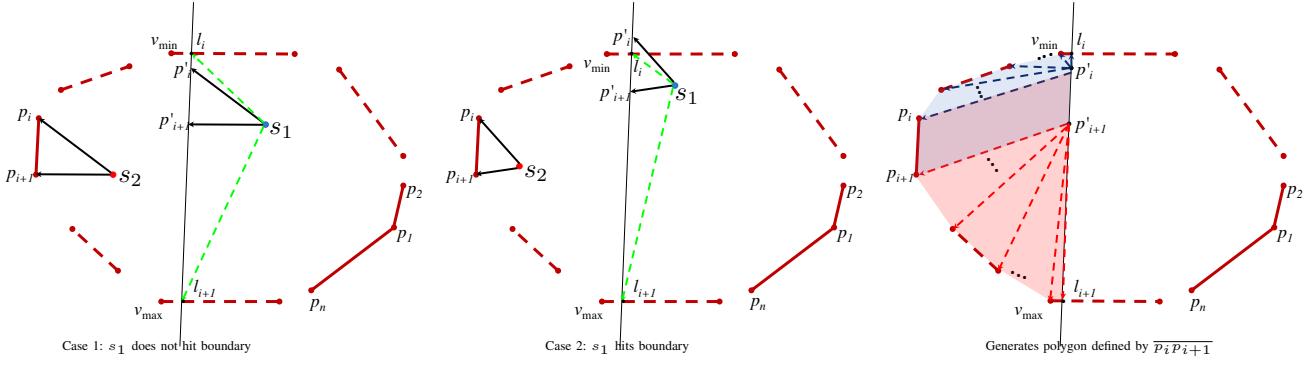


Fig. 6. Steps to generate the reachable set when one particle collides with edge  $i, i + 1$  of a convex polygonal workspace.

**Algorithm 2** REACHABLESETPOLYGON( $s_1, s_2, g_1, g_2, P$ )

**Require:** knowledge of starting  $(s_1, s_2)$  and goal  $(g_1, g_2)$  positions of two particles.  $P$  is a list of the vertices of a convex polygon.

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1:  $R_{\text{SET}} \leftarrow \{\}$ 
2: for  $p_i$  in  $P$  do
3:    $p'_i \leftarrow s_1 + s_2 - p_i$ 
4:    $p'_{i+1} \leftarrow s_1 + s_2 - p_{i+1}$ 
5:    $L \leftarrow \overline{p_i p_{i+1}}$             $\triangleright$  line  $(p'_i, p'_{i+1})$ 
6:    $l_i, l_{i+1} \leftarrow$  intersections of  $L$  and polygon  $P$ 
7:   if  $p'_i$  not inside polygon  $P$  then
8:      $p'_i \leftarrow l_i$ 
9:   end if
10:  if  $p'_{i+1}$  not inside polygon  $P$  then
11:     $p'_{i+1} \leftarrow l_{i+1}$ 
12:  end if
13:   $D \leftarrow s_2 - s_1 - ([l_i, v_{\min}, \dots, p_i] - p'_i,$ 
    $[p_{i+1}, p_{i+2}, \dots, v_{\max}, l_{i+1}] - p'_{i+1})$ 
14:   $R_{\text{SET}} \leftarrow$  Append polygon  $D$  to  $R_{\text{SET}}$ 
15: end for
16: Return  $R_{\text{SET}}$ 
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choose between the two options to immobilize the particle closest to a wall. The reachable  $\Delta$  configuration set for any first contact point defined by  $\psi$  is the area under a chord from angle  $\psi - \frac{\gamma(\psi)}{2}$  to  $\psi + \frac{\gamma(\psi)}{2}$ , for a circle of radius  $r$  centered at  $c = r(\cos(\psi - \pi), \sin(\psi - \pi))$ . One such chord is drawn in red in Fig. 7.

The equations for the four lines outlining the union of reachable  $\Delta$  configuration sets are as follows:

$$\begin{aligned}
l_1 &= r \left( (\cos \psi_{\min} - \cos(\gamma + \psi_{\min})) \right. & (8) \\
&\quad \left. + (\sin \psi_{\min} - \sin(\gamma + \psi_{\min})) \right) & 0 < \gamma < \gamma(\psi_{\min}), \\
l_2 &= r \left( (\cos \psi_{\max} - \cos(\gamma + \psi_{\max})) \right. \\
&\quad \left. + (\sin \psi_{\max} - \sin(\gamma + \psi_{\max})) \right) & \gamma(\psi_{\max}) < \gamma < 0, \\
l_3 &= r \left( (\cos \psi - \cos(\psi + \gamma(\psi))) \right. \\
&\quad \left. + (\sin \psi - \sin(\psi + \gamma(\psi))) \right) & \psi_{\min} < \psi < \psi_{\max},
\end{aligned}$$

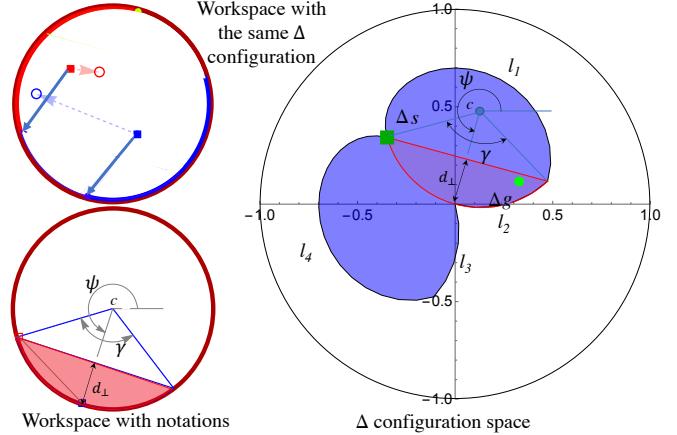


Fig. 7. Left top: The possible first contact points for the blue and red particles are shown with blue and red arcs. Left bottom: Once the blue particle touches the wall (blue square) the other particle (red region) can go anywhere in the reachable set (red region). Right: The  $\Delta$  configuration space for the corresponded starting positions of the particles is shown. The possible reachable sets before contact are shown in the  $\Delta$  configuration space as a blue region. Once the blue particle contacts the boundary as shown, the reachable  $\Delta$  configuration is the red set.

$$\begin{aligned}
l_4 &= r \left( (\cos \psi_{\min} - \cos(\psi - \gamma(\psi))) \right. \\
&\quad \left. + (\sin \psi_{\min} - \sin(\psi - \gamma(\psi))) \right) & \psi_{\min} < \psi < \psi_{\max}.
\end{aligned}$$

We combine these boundaries to compute the reachable set summarized in Alg. 3. Next, find a  $\psi$  that would enable us to reach  $\Delta g_c$ , the nearest point in the reachable set to  $\Delta g$ . We first check if  $\Delta g_c$  is in the  $\Delta$  configuration space chords defined by either  $\psi_{\min}$  or  $\psi_{\max}$  using:

$$\begin{aligned}
&(\Delta g_c.x - c.x)^2 + (\Delta g_c.y - c.y)^2 > r^2 \text{ and} \\
&(c.x - \Delta g_c.x) \cos \psi + (c.y - \Delta g_c.y) \sin \psi > r \cos \gamma.
\end{aligned} \tag{9}$$

If  $\Delta g_c$  is not in either chord, we draw a line from  $\Delta g_c$  to the current relative position,  $\Delta p$ . This line is a chord of the circle centered at  $c$ . The  $\psi$  to this chord obeys:

$$\psi = \tan^{-1} \left( \frac{\Delta p.x - \Delta g_c.x}{\Delta p.y - \Delta g_c.y} \right). \tag{10}$$

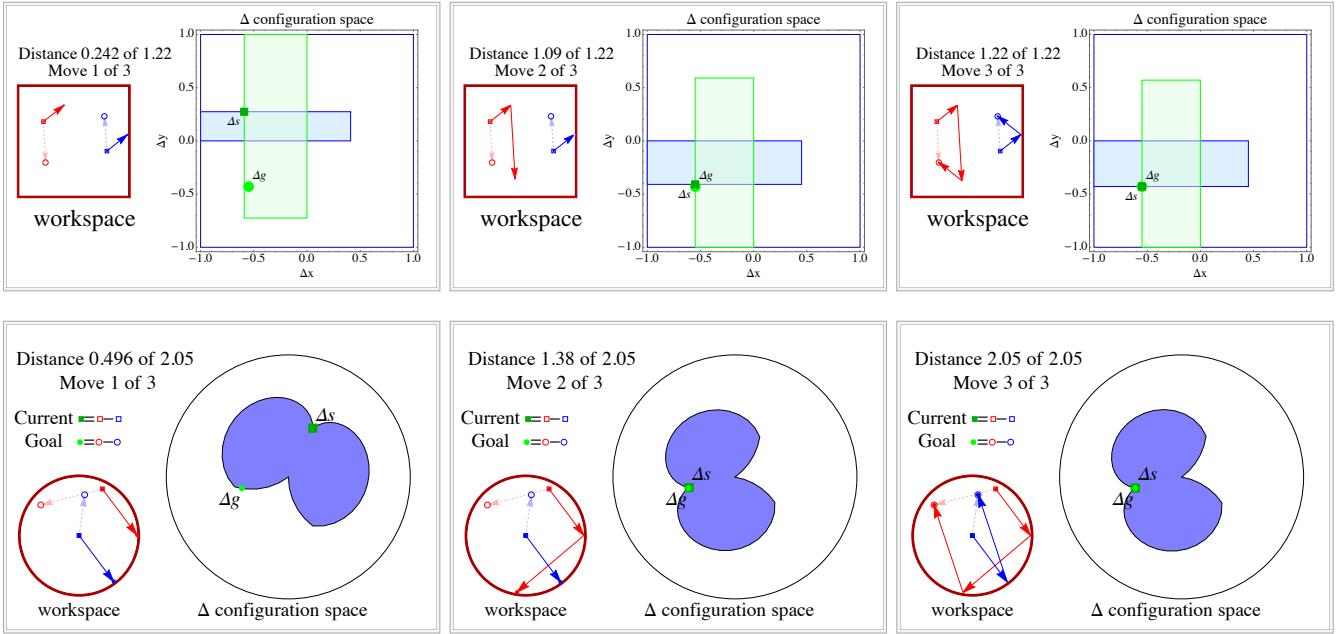


Fig. 8. Top row shows a polygonal workspace with its reachable sets. Bottom row, left circle shows the workspace. Right shows the  $\Delta$  configuration space and the reachable set that is shown in red is representative of the point we need to go to get to the goal relative distance in one move.

The particles achieve  $\Delta g_c$  in two moves. The first move causes one particle to touch the wall at  $p_\psi$ , (7). The second move achieves the required relative position.

### Algorithm 3 REACHABLESETCIRCLE( $s_1, s_2, g_1, g_2$ )

**Require:** knowledge of starting ( $s_1, s_2$ ) and goal ( $g_1, g_2$ ) positions of two particles.

- 1: Calculate  $p_\psi$  ▷ use (7)
- 2: Calculate  $\gamma$  ▷ use (5)
- 3: Calculate  $l_1, l_2, l_3, l_4$  ▷ use (8)
- 4: Return the union of  $(l_1, l_2, l_3, l_4)$

#### D. 3D workspaces: Cylinders and Prisms

Extending path planning to 3D is possible only if the two particles do not initially have the same  $x$  and  $y$  positions. For ease of analysis, we assume the workspace boundaries extend in the  $\pm z$  direction to form either right cylinders or right prisms. If the 3D projection is at a different angle, redefine the 2D workspace as a region perpendicular to the projection. First, we move the closest particle to the boundary, which prevents its  $z$  coordinate from changing. We next apply actuation in either the  $\pm z$  direction to achieve the desired  $\Delta z$ . Then the particles are actuated away from the boundary and to the appropriate  $z$  positions. Path planning continues using Alg. 1 to position the particles to the desired  $x$  and  $y$  positions. As an example, consider Fig. 9 which shows a cylindrical workspace. The blue particle starts in the blue disk and the red particle starts in the red disk. The two candidate shortest-length paths that touch the wall are shown with parallel arrows. Each arrow will cause one of the particles to touch the wall, enabling the other particle's reachable set to move freely in the  $z$  axis to achieve the required relative

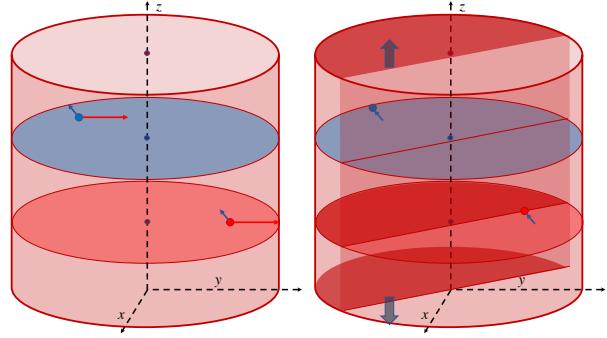


Fig. 9. Illustration on how boundary contacts enable 3D positioning. Once one particle contacts a boundary, the other particle's reachable set is a prism formed by extending the 2D reachable set in the  $z$ -direction.

position. This can be extended to other 3D workspaces if the workspace can be locally approximated as a 3D prism or cylinder. Other workspaces may be better handled by other path planners, such as [1], which used collisions with protrusions of the workspace to rearrange particles.

## V. SIMULATION

Algorithm 1 was implemented in Mathematica using particles with zero radius.

The contour plots in Fig. 10 left shows the length of the path for given  $s_1, s_2, g_1$  with  $g_2$  ranging over all the workspace. Fig. 10 right shows the total distance of the path. This plot clearly shows the nonlinear nature of the path planning. The hardest point to achieve is the when the goals have  $\pi$  difference and are very close to the boundary. Fig. 12 shows the same concepts in a square workspace. Fig. 12 top row considers the particles are interchangeable

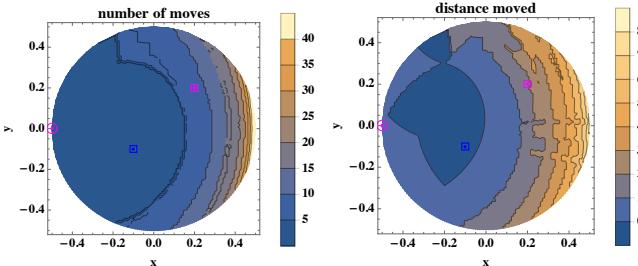


Fig. 10. Plots show performance with one goal on the boundary.

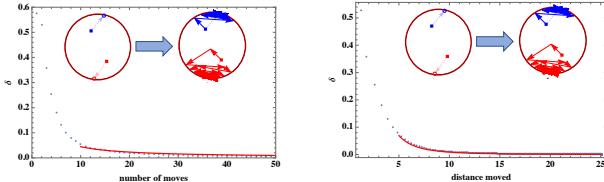


Fig. 11. Particle swap poles as their goal is the worst case. It can never be achieved but can be asymptotically reached. Plots show decreasing error as the number of moves grows.

for three arbitrary starting and goal positions for the particles. Using interchangeable particles makes the path length significantly smaller. The middle and bottom row shows the same configurations with unique particles.

The plots in Fig. 11 show the exponentially increasing number of moves and distance when the accuracy of reaching to the goal ( $\delta$ ) is getting to zero when the goal positions have  $\pi$  difference with each on the boundaries.

## VI. EXPERIMENTAL RESULTS

To demonstrate Alg. 1 experimentally, we performed several tests. Each used the same magnetic setup shown in Fig. 1. Two different intestine models were employed, the first a 3D-printed cross-section representation of a small intestine, and the second a cross-section of a bovine stomach.

### A. Magnetic Manipulation Setup

The magnetic manipulation system has two pairs of electromagnetic coils each with iron cores at their centers, and arranged orthogonal to each other. The iron core at the center of each coil concentrates the magnetic field towards the workspace. An Arduino and four SyRen regenerative motor drivers were used for control inputs to the coils. Finally, a FOculus F0134SB 659 x 494 pixel camera was attached to the top of the system focusing on the workspace which was back-lit by a 15 W LED light strip.

To obtain experimental data, the test samples which comprised of the phantom intestine model and the bovine cross section, were placed in laser cut acrylic discs and then immersed in corn syrup. Corn syrup was used to have the best viscosity(12000 cP) for the experiments. The velocities of particles immersed in it were damped enough to control their movements with ease. Spherical 1 mm magnets (super-magnetman #SP0100-50) were used as our particles.

todo: image of the magnetic setup with scale bar

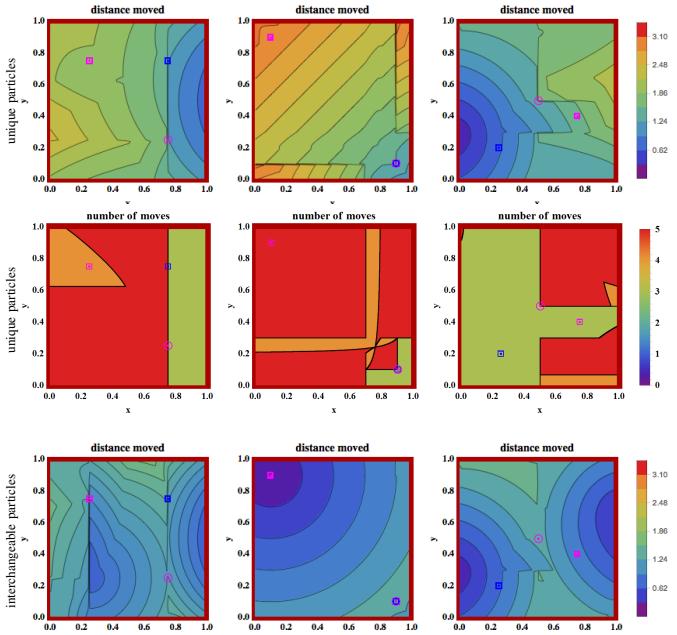


Fig. 12. Starting positions of particles 1 and 2 and goal position of particle 2 are fixed, and  $\epsilon = 0.001$ . The top row of contour plots show the distance if robot 1's goal position is varied in  $x$  and  $y$ . The bottom row shows the number of moves required for the same configurations.

### B. Intestine Phantom Model

The intestine phantom model was used as the first test field of the project and was made to mimic the geometry of an intestine and its villi. The model consists of a circular ring with an outer diameter of 50 mm, an inner diameter of 46 mm, and 60 2 mm long protrusions on its inner surface cut out of 6 mm thick acrylic to model the geometry of intestinal villi. Fig. 13 shows an experiment with starting and ending positions drawn with marker on the workspace.

### C. Bovine Stomach Cross-section

Strips of cow stomach approximately 5mm thick were cut and placed in neutrally buffered formalin for 24 hours for fixation. After fixation, each sample was transferred to 70% ethanol for storage. For the experiments, a slice of fixed intestine was attached to the acrylic disc with cyanoacrylate (superglue) and then submerged in corn syrup. A drawback of fixing the tissue samples before experiments is that they tended to shrivel and dry up a few minutes after being removed from the 70% ethanol.

## VII. CONCLUSION AND FUTURE WORK

This paper presented techniques for controlling the positions of two particles using uniform inputs and boundary friction forces. The paper provided algorithms for precise position control. The algorithms relied on calculating reachable sets in a 2D delta configuration space. Extending Alg. 1 to 3D is straightforward, but increases the complexity. Additionally, this paper assumed friction was sufficient to completely stop particles in contact with the boundary. The algorithms would require retooling to handle small friction

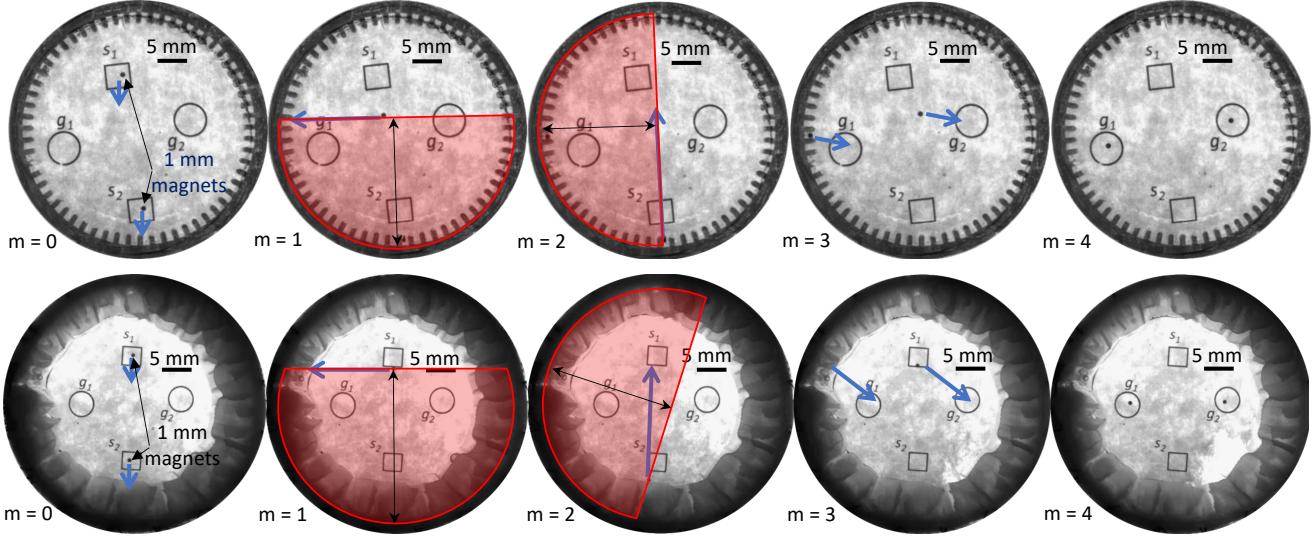


Fig. 13. Frames showing particle positions before and after control inputs. Top row: small intestine phantom. Bottom row: cow's stomach tissue.

coefficients. Hardware experiments illustrated the algorithms in ex vivo and in artificial workspaces that mimic the geometry of biological tissue

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