

PHYSICS 1 : COMPENDIUM

AAKASH JOG

1. NON-LINEAR MOTION

$$\begin{aligned}\vec{r} &= r\hat{r} \\ \vec{v} &= \dot{\vec{r}} = \dot{r}\hat{r} + r\dot{\theta}\hat{\theta} \\ \vec{a} &= \ddot{\vec{r}} = \left(\ddot{r} - r(\dot{\theta})^2\right)\hat{r} + \left(2\dot{r}\dot{\theta} + r\ddot{\theta}\right)\hat{\theta}\end{aligned}$$

2. CONSERVATIVE FORCES

$$\begin{aligned}\text{curl } \vec{F} &\doteq \vec{\nabla} \times \vec{F} \\ &= \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ F_x & F_y & F_z \end{vmatrix}\end{aligned}$$

If \vec{F} is conservative, $\text{curl } \vec{F} = 0$.

3. VARIABLE MASS SYSTEMS

If a rocket is releasing gasses with velocity u with respect to it,

$$\frac{dp}{dt} = m \frac{dv}{dt} + \frac{dm}{dt} u$$

4. CENTRES OF MASS

Solid hemisphere	$\left(\frac{3}{8}\right) R$
Hollow hemisphere	$\left(\frac{1}{2}\right) R$
Solid cone (from vertex)	$\left(\frac{3}{4}\right) h$
Hollow cone (from vertex)	$\left(\frac{2}{3}\right) h$

5. MOMENTS OF INERTIA

$$I = \int r^2 dm$$

Ring (\perp to plane)	mR^2
Disk (\perp to plane)	$\left(\frac{1}{2}\right) mR^2$
Solid sphere	$\left(\frac{2}{5}\right) mR^2$
Hollow sphere	$\left(\frac{2}{3}\right) mR^2$
Rod (centre)	$\left(\frac{1}{12}\right) ml^2$
Rod (end)	$\left(\frac{1}{3}\right) ml^2$
Cone (axis of symmetry)	$\left(\frac{3}{10}\right) mR^2$

6. ACCELERATING SYSTEMS

$$F_{\text{centrifugal}} = -m\vec{\omega} \times (\vec{\omega} \times \vec{r})$$

$$F_{\text{coriolis}} = -2m\vec{\omega} \times \vec{v}$$

7. OSCILLATIONS

7.1. Simple Oscillations.

$$\ddot{x} = -\omega^2 x$$

$$\omega_{\text{physical pendulum}} = \sqrt{\frac{d_{\text{axis, COM}} mg}{I_{\text{axis}}}}$$

7.2. Damped Oscillations.

$$\ddot{x} + \frac{\beta}{m}\dot{x} + \omega_0^2 x = 0$$

$$\text{Strong damping} \quad \frac{\beta}{2m} > \omega_0$$

$$\text{Critical damping} \quad \frac{\beta}{2m} = \omega_0$$

$$\text{Weak damping} \quad \frac{\beta}{2m} < \omega_0$$

Oscillations occur in case of weak damping.

$$\text{Let } \omega_1 = \sqrt{\omega_0^2 - \left(\frac{\beta}{2m}\right)^2}$$

For weak damping,

$$x = e^{-\beta/2m \cdot t} \left(\tilde{A} \cos \omega_1 t + \tilde{B} \sin \omega_1 t \right)$$

For critical damping,

$$x = e^{-\beta/2m \cdot t} \left(\tilde{A} + \tilde{B}t \right)$$

In case of strong damping,

$$x = \tilde{A}e^{(-\beta/2m + \sqrt{-\omega_1^2})t} + \tilde{B}e^{(-\beta/2m - \sqrt{-\omega_1^2})t}$$

7.3. Forced Oscillations.

$$m\ddot{x} + \frac{k}{m}x = \frac{F_0}{m} \cos \omega t$$

$$\therefore \ddot{x} + \frac{k}{m}x = \frac{F_0}{m} \cos \omega t$$

Therefore, solving

$$x = A \cos \omega_0 t + B \sin \omega_0 t + \frac{F_0}{k - m\omega^2} \cos \omega t$$

$$\therefore \dot{x} = \omega_0(-A \sin \omega_0 t + B \cos \omega_0 t) - \frac{F_0}{k - m\omega^2} \omega \sin \omega t$$

Substituting initial conditions,

$$\begin{aligned} \frac{F_0}{m} \\ x = \frac{\frac{F_0}{m}}{k - \omega^2} (-\cos \omega_0 t + \cos \omega t) \end{aligned}$$

$$\text{Let } \frac{F_0}{m} = f_0$$

$$\therefore x = -\frac{2f_0}{\omega_0^2 - \omega^2} \sin\left(\frac{\omega t - \omega_0 t}{2}\right) \sin\left(\frac{\omega t + \omega_0 t}{2}\right)$$

$$\omega - \omega_0 = \Delta\omega \text{ and } \omega + \omega_0 \approx 2\omega_0$$

$$\therefore x \approx \frac{2f_0}{\Delta\omega \cdot 2\omega_0} \sin\left(\frac{\Delta\omega}{2} t\right) \cdot \sin(\omega_0 t)$$