

Quantum and Solid State Physics

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Part I

Quantum Physics

1 Lecturer Information

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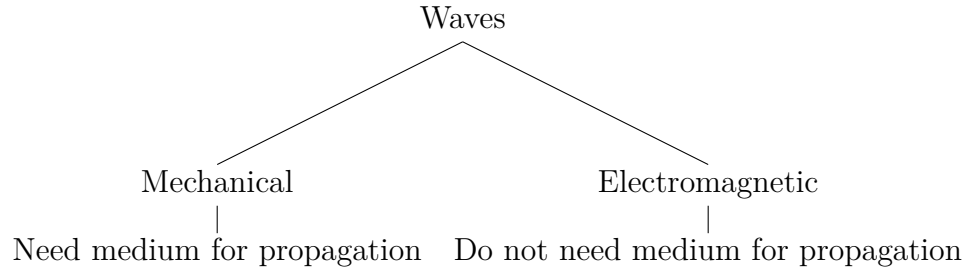
2 Required Reading

1. Griffiths, D. Introduction to quantum mechanics

3 Additional Reading

1. Tang: Fundamentals of quantum mechanics, Cambridge press.
2. Miller, Quantum mechanics for scientists and engineers.

4 Waves



4.1 1D Wave Equation

Definition 1 (1D wave equation). The equation

$$\frac{\partial^2 \psi}{\partial t^2} = v^2 \frac{\partial^2 \psi}{\partial x^2}$$

where ψ is a function of x and t , and v is the velocity of the wave, is called a 1D wave equation.

5 Harmonic Waves

Definition 2 (Harmonic waves). If a wave satisfies the equation

$$\psi(x, t) = A \cos(kx - \omega t + \varphi)$$

it is called a harmonic wave.

A is called the amplitude of the wave.

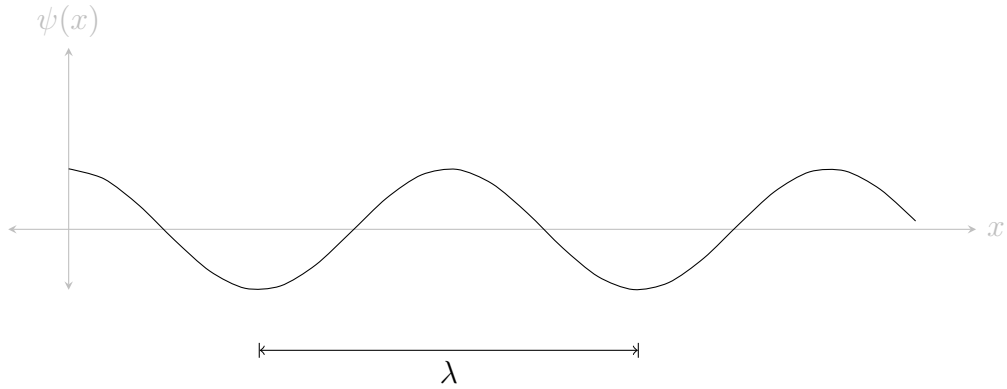
k is called the wave number, or spatial frequency of the wave.

ω is called the angular frequency of the wave.

Definition 3 (Wavelength). For a harmonic wave, a number λ , such that

$$\psi(x) = \psi(x + \lambda)$$

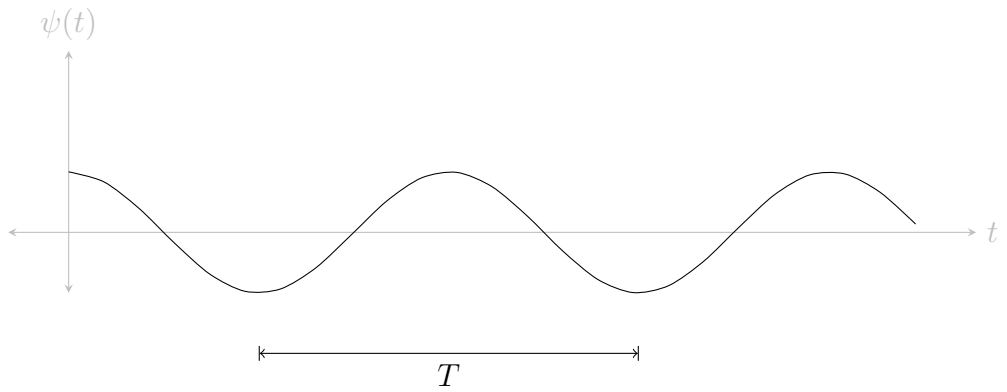
is called the wavelength of the wave.



Definition 4 (Time period). For a harmonic wave, a number T , such that

$$\psi(t) = \psi(t + T)$$

is called the time period of the wave.



Theorem 1.

$$k = \frac{2\pi}{\lambda}$$

where k is the wave number, and λ is the wavelength.

Proof. If $t = 0$,

$$\psi(x) = A \cos(kx)$$

By the definition of wavelength,

$$\psi(x) = \psi(x + \lambda)$$

$$\therefore A \cos(kx) = A \cos(k(x + \lambda))$$

$$\therefore k\lambda = 2\pi$$

$$\therefore k = \frac{2\pi}{\lambda}$$

□

Theorem 2.

$$\omega = \frac{2\pi}{T}$$

where ω is the angular frequency, and T is the time period.

Proof. If $x = 0$,

$$\psi(t) = A \cos(\omega t)$$

By the definition of time period,

$$\begin{aligned}\psi(t) &= \psi(t + T) \\ \therefore A \cos(\omega t) &= A \cos(\omega(t + T)) \\ \therefore \omega T &= 2\pi \\ \therefore \omega &= \frac{2\pi}{T}\end{aligned}$$

□

5.1 Complex Representation of Waves

Let

$$\tilde{\psi} = Ae^{i(kx - \omega t + \varphi)}$$

Then,

$$\psi = \Re\{\tilde{\psi}\}$$

5.2 Interference of Waves

Theorem 3. *Wave equations are linear, i.e. if ψ_1 and ψ_2 are solutions to the equation, then $\psi_1 + \psi_2$ is also a solution to the equation.*

5.2.1 Interference of Waves with a Phase Difference

Let

$$\begin{aligned}\psi_1 &= A \cos(kx - \omega t + \varphi) \\ \psi_2 &= A \cos(kx - \omega t)\end{aligned}$$

Therefore,

$$\begin{aligned}
 \psi_3 &= \psi_1 + \psi_2 \\
 &= A \cos(kx - \omega t + \varphi) + A \cos(kx - \omega t) \\
 &= 2A \cos\left(\frac{\varphi}{2}\right) \cos\left(kx + \omega t + \frac{\varphi}{2}\right)
 \end{aligned}$$

$\because \cos a + \cos b = 2 \cos\left(\frac{a+b}{2}\right) \cos\left(\frac{a-b}{2}\right)$

Therefore, the resultant wave is a wave with amplitude $2A \cos\left(\frac{\varphi}{2}\right)$ and phase $\frac{\varphi}{2}$.

6 Young's Double Slit Experiment (1801)

This experiment provided substantial proof that light behaves like a wave.

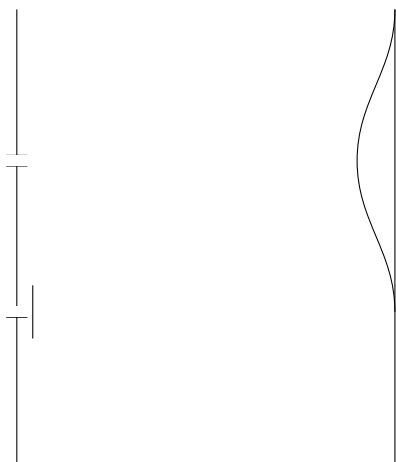


Figure 1: Intensity of light with only first slit open

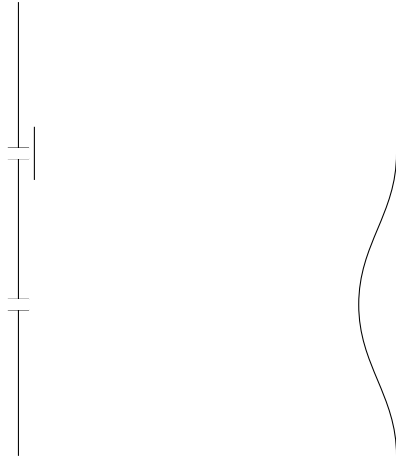


Figure 2: Intensity of light with only second slit open

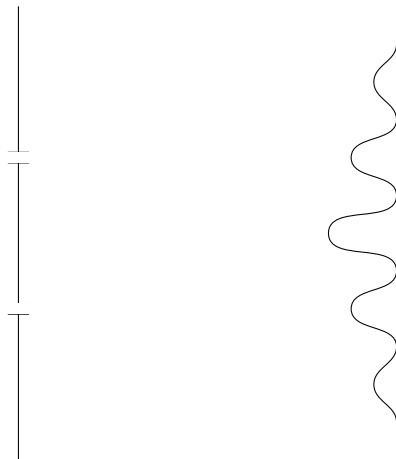


Figure 3: Intensity of light with both slits open

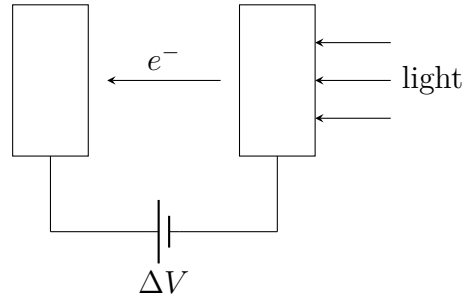
6.1 YDSE with Classical Particles

If the double slit experiment is performed with classical particles, instead of waves, the intensities add up. There is no fringe pattern, as observed in the experiment with waves.

7 The Photoelectric Effect

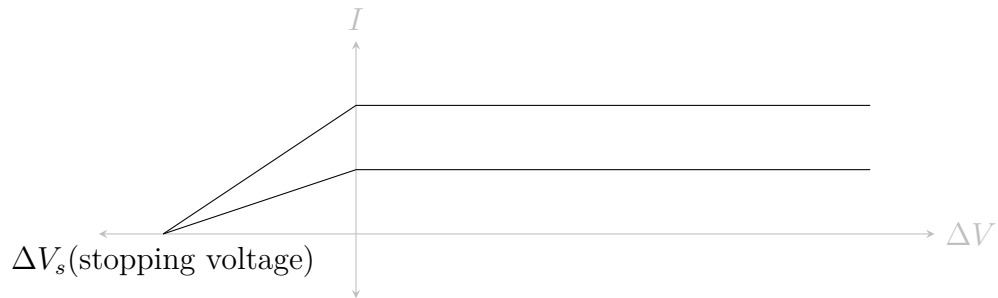
The first experiment in which the photoelectric effect was observed was performed by Hertz in 1887.

Two metallic plates, acting as electrodes were arranged as shown. They were connected to a voltage source ΔV , as shown.



The results observed were as shown.

1. The relationship between ΔV and the current in the wire was observed to be as shown.



The conclusions were as follows.

- (a) If the light intensity is constant, a specific amount of electrons is emitted. Therefore, the current is constant, and independent of ΔV .
- (b) If $\Delta V \gg 0$, all electrons emitted reached the other plate, and hence contributed to the current.
If $\Delta V < 0$, some electrons were unable to reach the other plate, and hence did not contribute to the current.
- (c) ΔV_s is not dependent on the intensity of the light.

As the energy of an electron is conserved,

$$E_{K_i} + E_{P_i} = E_{K_f} + E_{P_f}$$

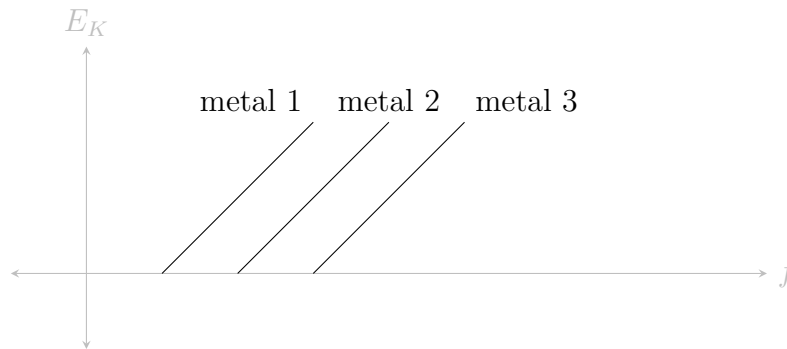
Therefore, if the electron barely reaches the other plate, i.e. if the voltage is ΔV_s

$$E_{K_i} + 0 = (-e)(-\Delta V_s)$$

$$\therefore e\Delta V_s = E_K$$

Therefore, as ΔV_s is independent of the intensity of light, the kinetic energy of the emitted electrons is also independent of the intensity of light.

2. The relationship between the kinetic energy of the emitted electrons, and the frequency of the incident light was observed to be as shown.



The conclusions were as follows.

- (a) There is a cutoff frequency, i.e. a frequency below which no electrons are emitted.
- (b) The kinetic energy of the emitted electrons is linearly dependent on the frequency of light.

These conclusions were inconsistent with the accepted notion of light being a wave.

7.1 Einstein's Explanation of the Photoelectric Effect (1905)

According to Einstein's explanation, light is a stream of particles, called photons. Each photon has energy equal to hf , where h is Planck's constant, and f is the frequency of the light, which is in fact a property of the wave nature of the light. This theory can explain the conclusions of Hertz's

experiment, which could not be explained by classical theories.

According to the explanation, each material has a property called the work function (W). The fact that there exists a cutoff voltage is justified due to this energy barrier. For an electron to be emitted, it needs to be provided energy to overcome this barrier. The cutoff frequency is such that all energy in a photon of this frequency to be used to overcome the work function.

Therefore,

$$hf_{\text{cutoff}} = W$$

Also, as each photon provides all its energy to a single electron, increasing the intensity of light just increases the number of electrons emitted, but does not increase the kinetic energy of the emitted electrons.

8 Quantum Particles

Definition 5 (Momentum of a quantum particle). The momentum of a quantum particle is defined as

$$p = \frac{E}{c}$$

where E is the energy of the particle.

Theorem 4 (Einstein Equation).

$$E = mc^2 + pc$$

Definition 6. The reduced Planck's constant is defined as

$$\hbar = \frac{h}{2\pi}$$

where h is Planck's constant.

Theorem 5.

$$\begin{aligned} E &= \hbar\omega \\ p &= \hbar k \end{aligned}$$

Proof.

$$\begin{aligned} E &= hf \\ &= h \frac{\omega}{2\pi} \\ &= \hbar\omega \end{aligned}$$

$$\begin{aligned} p &= \frac{E}{c} \\ &= \frac{h}{\lambda} \\ &= \frac{h}{2\pi} k \\ &= \hbar k \end{aligned}$$

□

8.1 de Broglie Wavelength

According to de Broglie's theory, particles have waves associated with them. Therefore, according to this theory,

$$\begin{aligned} \lambda &= \frac{h}{p} \\ &= \frac{h}{mv} \\ f &= \frac{E}{h} \end{aligned}$$

where m is the mass of the particle and v is its velocity.

Therefore, according to this theory, particles must exhibit wave-like behaviour. Hence, if the double slit experiment is performed with particles, the pattern observed must be similar to the fringe pattern observed with waves.

If the double slit experiment is performed with a single electron emitted at a time, over a long period of time, a fringe-like pattern, made up of dots corresponding to single electrons, is observed. This is consistent with de Broglie's theory.

Exercise 1.

Find the de Broglie wavelength of an electron moving at $10^7 \frac{\text{m}}{\text{s}}$.

Solution 1.

$$\begin{aligned}
\lambda &= \frac{h}{mv} \\
&= \frac{6.63 \times 10^{-34} \text{ J s}}{(9.11 \times 10^{-31} \text{ kg}) \left(10^7 \frac{\text{m}}{\text{s}}\right)} \\
&= 7.27 \times 10^{-11} \text{ m} \\
&= 72 \text{ pm}
\end{aligned}$$

Exercise 2.

Find the de Broglie wavelength of a rock of 50 g, thrown with a speed of $40 \frac{\text{m}}{\text{s}}$.

Solution 2.

$$\begin{aligned}
\lambda &= \frac{h}{mv} \\
&= \frac{6.63 \times 10^{-34}}{(50 \times 10^{-3}) (40)} \\
&= 3.3 \times 10^{-34} \text{ m}
\end{aligned}$$

8.2 Impact of Observation on the Result of Experiments

If 100% of the electrons are observed to determine which slit they pass through, the pattern observed is exactly like the pattern observed with classical particles.

If only some electrons are observed, say 70%, then those 70% electrons behave like classical particles, and the rest behave like a wave. Therefore, the effective result is a superposition both of these pattern.

Therefore, the outcome of the experiment is affected by the fact that the particles are being observed.

The reason behind this is that the act of observation involves interacting with the particles, usually in the form of the particle being hit by photons, which are necessary for the observer to make the observations.

Classical particles are not affected by this factor on such a large scale, as their mass is much larger than that of quantum particles.

9 Schrödinger Equation

Theorem 6 (Schrödinger Equation).

$$i\hbar \frac{\partial \psi(x, t)}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi(x, t)}{\partial x^2} + V(x)\psi(x, t)$$

Proof. Let the wave be given by

$$\psi(x, t) = Ae^{i(kx - \omega t)}$$

Differentiating with respect to time,

$$\begin{aligned} \frac{\partial \psi(x, t)}{\partial t} &= -i\omega \psi(x, t) \\ \therefore \omega &= \frac{E}{\hbar} &= -i \frac{E}{\hbar} \psi(x, t) \\ &= -\frac{i}{\hbar} \left(\frac{p^2}{2m} + E_P \right) \psi(x, t) \\ &= -\frac{i}{\hbar} \left(\frac{p^2}{2m} + V(x) \right) \psi(x, t) \end{aligned} \tag{1}$$

where $E_P = V(x)$ is
the potential energy of
the particle

Differentiating with respect to x ,

$$\begin{aligned} \frac{\partial \psi(x, t)}{\partial x} &= ik\psi(x, t) \\ \therefore \frac{\partial^2 \psi(x, t)}{\partial x^2} &= (ik)(ik)\psi(x, t) \\ &= -k^2\psi(x, t) \\ &= -\frac{p^2}{\hbar^2}\psi(x, t) \\ \therefore -\hbar^2 \frac{\partial^2 \psi(x, t)}{\partial x^2} &= p^2\psi(x, t) \\ \therefore -\frac{\hbar^2}{2m} \frac{\partial^2 \psi(x, t)}{\partial x^2} &= \frac{p^2}{2m}\psi(x, t) \end{aligned} \tag{2}$$

Therefore, solving the above equations simultaneously,

$$\begin{aligned} \frac{\partial \psi(x, t)}{\partial t} &= -\frac{i}{\hbar} \left(-\frac{\hbar^2}{2m} \frac{\partial^2 \psi(x, t)}{\partial x^2} + V(x) \right) \psi(x, t) \\ \therefore i\hbar \frac{\partial \psi(x, t)}{\partial t} &= \left(-\frac{\hbar^2}{2m} \frac{\partial^2 \psi(x, t)}{\partial x^2} + V(x) \right) \psi(x, t) \end{aligned}$$

□

10 Basic Postulates of Quantum Mechanics

1. $\psi(x, t)$ describes the configuration of the system.
2. The probability of finding the particle between x and $x + dx$ at time t is

$$|\psi(x, t)|^2 = \psi(x, t)\psi^*(x, t)$$

where ψ^* is the complex conjugate of ψ ,

3. If $\psi_1(x, t)$ and $\psi_2(x, t)$ are solutions to the Schrödinger equation, then $\alpha\psi_1(x, t) + \beta\psi_2(x, t)$ is also a solution to the Schrödinger equation.

11 Properties of $\psi(x, t)$

1.
$$\int_{-\infty}^{\infty} |\psi(x, t)|^2 dx = 1$$

$|\psi(x, t)|^2$ is the probability of the particle being in the neighbourhood of x . Therefore, the equation is obvious, as the probability of the particle being in the entire universe must be 1.

2. $|\psi(x, t)|$ is a single valued function.
This is obvious, as the probability of the particle existing at a point in space must have exactly one value.

12 Conservation of Normalization

Definition 7. If

$$\int_{-\infty}^{\infty} |\psi(x, t)|^2 dx = 1$$

Then the equation is said to be normalized.

Theorem 7. *The Schrödinger equation conserves normalization, i.e.,*

$$\frac{d}{dt} \int_{-\infty}^{\infty} |\psi(x, t)|^2 dx = 0$$

Proof.

$$\begin{aligned}\frac{d}{dt} \int_{-\infty}^{\infty} |\psi(x, t)|^2 dx &= \int_{-\infty}^{\infty} \frac{\partial}{\partial t} (\psi(x, t) \psi^*(x, t)) dx \\ &= \int_{-\infty}^{\infty} \frac{\partial \psi}{\partial t} \psi^* + \psi \frac{\partial \psi^*}{\partial t} dx\end{aligned}$$

From the Schrödinger Equation,

$$\begin{aligned}\frac{\partial \psi}{\partial t} &= \frac{i\hbar}{2m} \frac{\partial^2 \psi}{\partial x^2} - \frac{i}{\hbar} V \psi \\ \therefore \frac{\partial \psi^*}{\partial t} &= -\frac{i\hbar}{2m} \frac{\partial^2 \psi^*}{\partial x^2} + \frac{i}{\hbar} V \psi^*\end{aligned}$$

Therefore,

$$\begin{aligned}\frac{d}{dt} \int_{-\infty}^{\infty} |\psi(x, t)|^2 dx &= \int_{-\infty}^{\infty} \left(\frac{i\hbar}{2m} \frac{\partial^2 \psi}{\partial x^2} \psi^* - \frac{i\hbar}{2m} \frac{\partial^2 \psi^*}{\partial x^2} \psi \right) dx \\ &= \int_{-\infty}^{\infty} \left(\frac{i\hbar}{2m} \frac{\partial}{\partial x} \left(\psi^* \frac{\partial \psi}{\partial x} - \psi \frac{\partial \psi^*}{\partial x} \right) \right) dx \\ &= \frac{i\hbar}{2m} \left(\psi^* \frac{\partial \psi}{\partial x} - \psi \frac{\partial \psi^*}{\partial x} \right) \Big|_{-\infty}^{\infty}\end{aligned}$$

Therefore,

$$\begin{aligned}\frac{d}{dt} \int_{-\infty}^{\infty} |\psi(x, t)|^2 dx &= 0 \\ \iff \frac{i\hbar}{2m} \left(\psi^* \frac{\partial \psi}{\partial x} - \psi \frac{\partial \psi^*}{\partial x} \right) \Big|_{-\infty}^{\infty} &= 0\end{aligned}$$

Therefore, the normalization is conserved if the wave function tapers off at infinity. \square

13 Measurements and Operators

13.1 The Collapse of the Wave Function

Consider a probability distribution of finding a particle at x , given by the wave function $\psi(x, t)$. The position of the particle is actually measured, and

hence found to be at some position $x = A$, at time t . Therefore, after a small time interval dt after t , it can be predicted with certainty that the particle is around $x = A$. Hence, the original wave function is no longer valid, and there is a new wave function which describes the particle. This phenomenon is called the collapse of the wave function.

13.2 Momentum

Definition 8 (Observable). A physical parameter which describes a particle, such as position, momentum, energy, velocity, etc., is called an observable.

Definition 9 (Expected value). The average value of an observable, measured with N particles, with the same wave function $\psi(x, t)$, is called the expected value of the observable. It is denoted as $\langle \alpha \rangle$, where α is the observable.

Theorem 8 (Ehrenfest's Theorem). *Expected values obey classical laws, i.e., for example,*

$$\begin{aligned}\langle v \rangle &= \frac{d}{dt} \langle x \rangle \\ \langle p \rangle &= m \frac{d}{dt} \langle x \rangle\end{aligned}$$

Theorem 9.

$$\langle x \rangle = \int_{-\infty}^{\infty} \psi^* x \psi \, dx$$

Proof. By definition,

$$\begin{aligned}\langle x \rangle &= \int_{-\infty}^{\infty} x |\psi(x, t)|^2 \, dx \\ &= \int_{-\infty}^{\infty} x \psi(x, t) \psi^*(x, t) \, dx \\ &= \int_{-\infty}^{\infty} \psi^* x \psi \, dx\end{aligned}$$

□

Theorem 10.

$$\langle p \rangle = \int_{-\infty}^{\infty} \psi^* \left(-i\hbar \frac{\partial}{\partial x} \right) \psi \, dx$$

Proof. By Ehrenfest's Theorem,

$$\begin{aligned} \langle p \rangle &= m \frac{d}{dt} \langle x \rangle \\ &= m \frac{d}{dt} \int_{-\infty}^{\infty} x |\psi(x, t)|^2 \, dx \\ &= m \int_{-\infty}^{\infty} x \frac{\partial}{\partial t} (\psi \psi^*) \, dx \\ &= m \frac{i\hbar}{2m} \int_{-\infty}^{\infty} x \frac{\partial}{\partial x} \left(\psi^* \frac{\partial \psi}{\partial t} + \psi \frac{\partial \psi^*}{\partial t} \right) \, dx \\ &= \frac{i\hbar}{2} \left(x \left(\psi^* \frac{\partial \psi}{\partial x} - \psi \frac{\partial \psi^*}{\partial x} \right) \Big|_{-\infty}^{\infty} - \int_{-\infty}^{\infty} \left(\psi^* \frac{\partial \psi}{\partial x} - \psi \frac{\partial \psi^*}{\partial x} \right) \, dx \right) \\ &= -\frac{i\hbar}{2} \int_{-\infty}^{\infty} \left(\psi^* \frac{\partial \psi}{\partial x} - \psi \frac{\partial \psi^*}{\partial x} \right) \, dx \\ &= -\frac{i\hbar}{2} 2 \int_{-\infty}^{\infty} \frac{\partial \psi}{\partial x} \psi^* \, dx \\ &= -i\hbar \int_{-\infty}^{\infty} \frac{\partial \psi}{\partial x} \psi^* \, dx \\ &= \int_{-\infty}^{\infty} \psi^* (-i\hbar) \frac{\partial}{\partial x} \psi \, dx \end{aligned}$$

□

13.3 Operators

Definition 10 (Operator). A function which acts on a function is called an operator. There exists an operator associated with every observable.

Theorem 11. Every operator can be written as a function of \hat{x} , \hat{p} , $Q(\hat{x}, \hat{p})$.
In general,

$$\langle Q \rangle = \int_{-\infty}^{\infty} \psi^* \hat{Q} \psi \, dx$$

where

$$\hat{x} = x$$

$$\hat{p} = -i\hbar \frac{\partial}{\partial x}$$

Theorem 12.

$$\begin{aligned} E_k &= \frac{1}{2}mv^2 \\ &= \frac{p^2}{2m} \end{aligned}$$

Therefore,

$$\hat{E}_k = \frac{\hat{p}^2}{2m}$$

Therefore,

$$\begin{aligned} \langle E_k \rangle &= \int_{-\infty}^{\infty} \psi^* \left(\frac{\hat{p}^2}{2m} \right) \psi \, dx \\ &= \frac{1}{2m} \int_{-\infty}^{\infty} \psi^* (-\hbar^2) \frac{\partial^2}{\partial x^2} \psi \, dx \end{aligned}$$

Exercise 3.

A particle with mass m has wave function

$$\psi(x, t) = Ae^{-a\left(\frac{mx^2}{\hbar} + it\right)}$$

where a and A are constants.

1. Find A .
2. Find $V(x)$ which gives $\psi(x, t)$ as a solution of the Schrödinger Equation.

3. Find the expected values of

- (a) x .
- (b) x^2 .
- (c) p .
- (d) p^2 .

Solution 3.

1.

$$\begin{aligned}
 1 &= \int_{-\infty}^{\infty} |\psi(x, t)|^2 dx \\
 &= \int_{-\infty}^{\infty} \psi(x, t) \psi^*(x, t) dx \\
 &= \int_{-\infty}^{\infty} A e^{-a\left(\frac{mx^2}{\hbar} + it\right)} A e^{-a\left(\frac{mx^2}{\hbar} - it\right)} dx \\
 &= A^2 \int_{-\infty}^{\infty} e^{-2a\frac{mx^2}{\hbar}} dx \\
 &= 2A^2 \int_0^{\infty} e^{-2a\frac{mx^2}{\hbar}} dx \\
 &= 2A^2 \left(\frac{1}{2} \left(\frac{\pi \hbar}{2am} \right)^{\frac{1}{2}} \right) \\
 &= A^2 \left(\frac{\pi \hbar}{2am} \right)^{\frac{1}{2}}
 \end{aligned}$$

As the function is even, the integral from $-\infty$ to ∞ is twice that of the integral from 0 to ∞ .

$$\int_{-\infty}^{\infty} e^{-bx^2} dx = \frac{1}{2} \left(\frac{\pi}{b} \right)^{\frac{1}{2}}$$

Therefore,

$$A = \left(\frac{2am}{\pi \hbar} \right)^{\frac{1}{4}}$$

2.

$$\begin{aligned}\psi(x, t) &= Ae^{-a\left(\frac{mx^2}{\hbar}+it\right)} \\ &= \left(\frac{2am}{\pi\hbar}\right)^{\frac{1}{4}} e^{-a\left(\frac{mx^2}{\hbar}+it\right)}\end{aligned}$$

Therefore,

$$\begin{aligned}\frac{\partial\psi}{\partial t} &= (-ia)\psi(x, t) \\ \frac{\partial\psi}{\partial x} &= -\frac{am}{\hbar}2x\psi(x, t) \\ \therefore \frac{\partial^2\psi}{\partial x^2} &= -\frac{2am}{\hbar} \frac{\partial}{\partial x} (x\psi(x, t)) \\ &= -\frac{2am}{\hbar} \left(\psi(x, t) - \frac{am}{\hbar}2x^2\psi(x, t) \right)\end{aligned}$$

Therefore, substituting in the Schrödinger Equation,

$$\begin{aligned}i\hbar \frac{\partial\psi}{\partial t} &= \left(-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + V(x) \right) \psi \\ \therefore a\hbar\psi &= a\hbar \left(\psi - \frac{am}{\hbar}2x^2\psi \right) + V(x)\psi \\ \therefore V(x) &= 2a^2mx^2\end{aligned}$$

3. (a)

$$\begin{aligned}\langle x \rangle &= \int_{-\infty}^{\infty} \psi^* x \psi \, dx \\ &= \int_{-\infty}^{\infty} x |\psi(x, t)|^2 \, dx \\ &= \int_{-\infty}^{\infty} x A^2 e^{-2a\left(\frac{mx^2}{\hbar}\right)} \, dx\end{aligned}$$

Therefore, as this function is odd, its integral over a symmetric interval is zero.

Therefore,

$$\langle x \rangle = 0$$

(b)

$$\begin{aligned} \text{As } \int_{-\infty}^{\infty} x^2 e^{-bx^2} dx &= \frac{1}{2b} \sqrt{\frac{\pi}{b}} \\ \langle x^2 \rangle &= \int_{-\infty}^{\infty} x^2 A^2 e^{-2a\left(\frac{mx^2}{\hbar}\right)} dx \\ &= \frac{\hbar}{4am} \end{aligned}$$

(c)

$$\begin{aligned} \langle p \rangle &= \int_{-\infty}^{\infty} \psi^* \left(-i\hbar \frac{\partial}{\partial x} \right) \psi dx \\ &= m \frac{d\langle x \rangle}{dt} \\ &= 0 \end{aligned}$$

(d)

$$\begin{aligned} \langle p^2 \rangle &= \int_{-\infty}^{\infty} \psi^* (-\hbar^2) \frac{\partial^2 \psi}{\partial x^2} dx \\ &= -\hbar^2 \int_{-\infty}^{\infty} \psi^* \left(\left(\frac{2am}{\hbar} x \right)^2 - \frac{2am}{\hbar} \right) \psi dx \\ &= -\hbar^2 \int_{-\infty}^{\infty} A^2 e^{-2a\frac{mx^2}{\hbar}} \left(\left(\frac{2am}{\hbar} x \right)^2 - \frac{2am}{\hbar} \right) dx \\ &= am\hbar \end{aligned}$$

13.4 Eigenfunctions and Eigenvalues of Operators

Definition 11 (Eigenvalues and eigenfunctions). Let \hat{O} be an operator. The scalar α_n such that

$$\hat{O}(f_n) = \alpha_n f_n$$

is called the eigenvalue corresponding to the eigenfunction f_n .

Definition 12. If the wave function of a particle is given by

$$\psi(x) = Ae^{ikx}$$

the particle is said to be in a state with definite momentum.

Exercise 4.

Find the eigenvalue of the momentum operator for a particle with definite momentum.

Solution 4.

$$\begin{aligned}\hat{p}\psi(x) &= \alpha\psi(x) \\ \therefore -i\hbar \frac{d\psi(x)}{dx} &= \alpha\psi(x) \\ \therefore \psi(x) &= Ae^{-\frac{\alpha}{\hbar}x}\end{aligned}$$

As

$$\psi(x) = Ae^{ikx}$$

Therefore,

$$\begin{aligned}\frac{\alpha}{\hbar} &= k \\ \therefore \alpha &= \hbar k \\ &= p\end{aligned}$$

Therefore,

$$\begin{aligned}\hat{p}\psi(x) &= \alpha\psi(x) \\ &= p\psi(x)\end{aligned}$$

Therefore, the eigenvalue of \hat{p} is p .

Definition 13 (Dirac delta function). The Dirac delta function is defined to be

$$\delta(x - a) = \begin{cases} \infty & ; \quad x = a \\ 0 & ; \quad x \neq a \end{cases}$$

Theorem 13.

$$\int_{-\infty}^{\infty} \delta(x - a) dx = 1$$

Theorem 14.

$$\delta(x - a) = \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{ip(x-a)} dp$$

Theorem 15.

$$f(a) = \int_{-\infty}^{\infty} \delta(x - a) f(x) \, dx$$

13.5 Dirac Delta Function

Theorem 16. *The Dirac delta function is even, i.e.*

$$\delta(-x) = \delta(x)$$

13.6 Fourier Transform

Definition 14 (Fourier transform). The Fourier transform of $f(x)$ is defined as

$$F(k) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x) e^{-ikx} \, dx$$

$$f(k) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} F(x) e^{ikx} \, dx$$

Exercise 5.

Find the Fourier transform of $f(x) = 1$.

Solution 5.

$$F(k) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x) e^{-ikx} \, dx$$

$$= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{-ikx} \, dx$$

$$= \sqrt{2\pi} \delta(-k)$$

$$= \sqrt{2\pi} \delta(k)$$

As $\delta(x - a) =$
 $\frac{1}{2\pi} \int_{-\infty}^{\infty} e^{ip(x-a)} \, dp.$
 As $\delta(x)$ is even.

Exercise 6.

Find the Fourier transform of $f(x) = e^{iax}$.

Solution 6.

$$\begin{aligned}
 F(k) &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x) e^{-ikx} \, dx \\
 &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{iak} e^{-ikx} \, dx \\
 &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{-i(k-a)x} \, dx \\
 &= \sqrt{2\pi} \delta(-(k-a)) \\
 &= \sqrt{2\pi} \delta(k-a)
 \end{aligned}$$

As $\delta(x-a) =$
 $\frac{1}{2\pi} \int_{-\infty}^{\infty} e^{ip(x-a)} \, dp$.
 As $\delta(x)$ is even.

14 Heisenberg's Uncertainty Principle

Heisenberg's Uncertainty Principle states that

$$\sigma_x \sigma_p \geq \frac{\hbar}{2}$$

where

$$\sigma_x = \sqrt{\langle x^2 \rangle - \langle x \rangle^2}$$

is the standard deviation of the position of the particle, and

$$\sigma_p = \sqrt{\langle p^2 \rangle - \langle p \rangle^2}$$

is the standard deviation of the momentum of the particle.

14.1 Attempting to Measure both the Position and Momentum of an Electron



Consider a screen with a single slit kept in the path of a beam of electrons as shown.

Let the momentum of the electrons be a known value p_0 .

Let the thickness of the slit be dy .

Consider the electrons which pass through the slit.

As the thickness of the slit is dy , the uncertainty in the position of the electrons, which pass through the slit, is dy .

As the probability wave of the electron passes through the slit, it diffracts. Hence, some of the electrons which pass through the slit continue on the original trajectory, and some of them change their trajectory and also move in the y -axis. Hence, there is an uncertainty in the momentum of the electrons. Let the angle that a particular electron deviates with be θ . Therefore, the uncertainty in the momentum of the electron is $p_0 \sin \theta$.

Therefore,

$$\begin{aligned}\sigma_x \sigma_p &= p_0 \sin \theta \, dy \\ &= p_0 \frac{\lambda}{dy} \, dy \\ &= p_0 \lambda \\ &= p_0 \frac{h}{p_0} \\ &= h \\ &\approx \hbar\end{aligned}$$

15 Definite Momentum States

Theorem 17. $|\tilde{\psi}(k)|^2 dk$ is the probability of the momentum of a particle being between $k\hbar$ and $(k + dk)\hbar$.

Consider the wave function

$$\psi(x) = Ae^{ik_1x}$$

Therefore, using Fourier transformation,

$$\tilde{\psi}(k) = A\sqrt{2\pi}\delta(k + k_1)$$

Therefore,

$$|\psi(x)|^2 = A^2$$

$$|\tilde{\psi}(k)|^2 = \delta(k + k_1)$$

Therefore, the probability distributions are as shown in Figure 4, and Figure 5.



Figure 4: Probability distribution of x for a definite momentum state

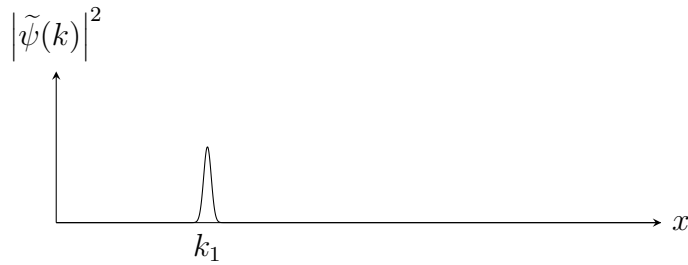


Figure 5: Probability distribution of k for a definite momentum state

Hence, σ_x is high, and σ_p is low.

16 Definite Position States

Consider the wave function

$$\psi(x) = \delta(x - x_1)$$

Therefore, using Fourier transformation,

$$\tilde{\psi}(k) = \frac{e^{ikx_1}}{\sqrt{2\pi}}$$

Therefore,

$$|\psi(x)| = \delta(x - x_1)$$

$$|\tilde{\psi}(k)| = \frac{1}{\sqrt{2\pi}}$$

Therefore, the probability distributions are as shown in Figure 6, and Figure 7.



Figure 6: Probability distribution of x for a definite position state

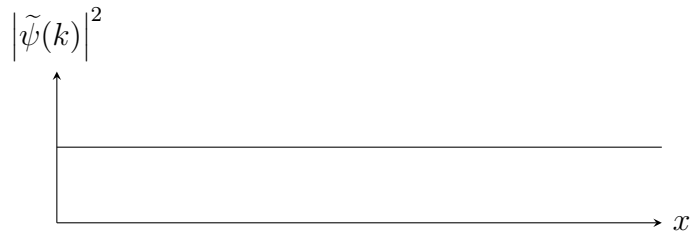


Figure 7: Probability distribution of k for a definite position state

Hence, σ_x is low, and σ_p is high.

Theorem 18. *The eigenstates or eigenfunctions of the momentum operator form a basis. Therefore, if*

$$\hat{A}\varphi_a = a\varphi_a$$

for discrete variables,

$$\psi(x) = \sum_a c_a \varphi_a(x)$$

for continuous variables,

$$\psi(x) = \int c_a \varphi_a(x) da$$

Proof.

$$\psi(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \tilde{\psi}(k) e^{ikx} dk$$

Therefore, any value of $\psi(x)$ can be written in terms of the corresponding $\tilde{\psi}(k)$ and e^{ikx} . Hence, the eigenstates of the momentum operator form a basis. \square

Theorem 19. *If*

$$\hat{A}\varphi_a = a\varphi_a$$

then, for discrete variables,

$$\psi(x) = \sum_a c_a \varphi_a(x)$$

then, for continuous variables,

$$\psi(x) = \int c_a \varphi_a(x) da$$

Then, the probability of the observable A having value a is

$$P(A = a) = |c_a|^2$$

17 Description of Physical State by Eigenstates

Let the particle be in a state of definite momentum.
Therefore,

$$\psi(x) = Ae^{ik_1x}$$

Therefore, the total probability over the entire space is

$$\begin{aligned}\int_{-\infty}^{\infty} |\psi|^2 dx &= \int_{-\infty}^{\infty} Ae^{ik_1x} A^* e^{ik_2x} dx \\ &= \int_{-\infty}^{\infty} |A|^2 dx \\ &\rightarrow \infty\end{aligned}$$

Therefore, as this integral cannot be finite, irrespective of A , this state does not represent a physical state.

Let

$$\begin{aligned}\varphi_n &= Ae^{ik_nx} \\ \varphi_m &= Ae^{ik_mx}\end{aligned}$$

be eigenfunctions of \hat{p} .

Therefore, as they are eigenfunctions, they are orthonormal to each other.

Therefore,

$$\begin{aligned}\int_{-\infty}^{\infty} \varphi_n^* \varphi_m dx &= |A|^2 \int_{-\infty}^{\infty} e^{i(k_m - k_n)x} dx \\ &= |A|^2 \cdot 2\pi \delta(k_m - k_n)\end{aligned}$$

As the eigenfunctions are orthonormal, and not just orthogonal to each other, the coefficient of the delta function must be 1. Therefore,

$$\begin{aligned}1 &= |A|^2 \cdot 2\pi \\ \therefore |A|^2 &= \frac{1}{2\pi} \\ \therefore |A| &= \frac{1}{\sqrt{2\pi}}\end{aligned}$$

Hence, a single eigenstate cannot describe the physical state of the particle, but if more than one eigenstates are known, then together they can describe the physical state of the particle

18 The Commutator

Definition 15 (Commutator). The commutation of two functions \hat{A} and \hat{B} is defined to be

$$\begin{aligned} [\hat{A}, \hat{B}] &= \hat{A}\hat{B} - \hat{B}\hat{A} \\ \therefore [\hat{A}, \hat{B}] f(x) &= \hat{A}(\hat{B}(f(x))) - \hat{B}(\hat{A}(f(x))) \end{aligned}$$

The operator which gives the commutation of two functions is called the commutator.

Theorem 20. *If*

$$[\hat{A}, \hat{B}] = 0$$

then \hat{A} and \hat{B} have some eigenfunctions in common.

Proof. Let a be an eigenvalue of \hat{A} , and let φ_a be an eigenfunction of \hat{A} . Therefore,

$$\hat{A}\varphi_a = a\varphi_a$$

Therefore, applying \hat{B} on the LHS, as the commutation is zero,

$$\hat{B}\hat{A}\varphi_a = \hat{A}\hat{B}\varphi_a$$

Therefore, applying \hat{B} on the RHS,

$$\hat{B}a\varphi_a = a\hat{B}\varphi_a$$

Therefore,

$$\hat{A}(\hat{B}\varphi_a) = a(\hat{B}\varphi_a)$$

This equation is of the form

$$\hat{A}\varphi = a\varphi$$

Therefore, $\hat{B}\varphi_a$ is also an eigenfunction of \hat{A} with eigenvalue a .

Hence, assuming there is no degeneracy in eigenvalues, i.e. every eigenvalue has a single corresponding eigenfunction, \hat{B} cannot change φ_a , except by a constant.

Therefore,

$$\hat{B}\varphi_a = b\varphi_a$$

Therefore, φ_a is an eigenfunction of \hat{B} corresponding to the eigenvalue b .

Hence there is a common eigenfunction for \hat{A} and \hat{B} . \square

Exercise 7.

Find the commutation of \hat{x} and \hat{p} .

Solution 7.

$$\begin{aligned}
 [\hat{x}, \hat{p}] f(x) &= \hat{x}\hat{p}f(x) - \hat{p}\hat{x}f(x) \\
 &= (x) \left(-i\hbar \frac{d}{dx} f(x) \right) + i\hbar \frac{d}{dx} (xf(x)) \\
 &= -i\hbar x f'(x) + i\hbar f(x) + i\hbar x f'(x) \\
 &= i\hbar f(x)
 \end{aligned}$$

Therefore,

$$[\hat{x}, \hat{p}] = i\hbar \hat{I}$$

where \hat{I} is the identity operator.

Exercise 8.

Can a wave function be an eigenfunction of both \hat{x} and \hat{p} ?

Solution 8.

$$\begin{aligned}
 \hat{x}f(x) &= xf(x) \\
 \hat{p}f(x) &= -i\hbar \frac{\partial}{\partial x} f(x) \\
 &= \hbar k f(x)
 \end{aligned}$$

Therefore, as the two equations cannot be solved simultaneously, the two operators cannot have a common eigenfunction.

Exercise 9.

Can any wave function be written as a superposition of momentum eigenstates?

Solution 9.

Therefore, as

$$f(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} F(k) e^{ikx} dk$$

every $f(x)$ can be written in terms of e^{ikx} .

Hence, any wave function can be written as a superposition of momentum eigenstates.

Exercise 10.

It is impossible to demonstrate the wave-like behaviour of macroscopic objects as λ is too small. True or false?

Solution 10.

$$\lambda = \frac{h}{p}$$

Therefore, as the p for macroscopic objects is very large, the λ is very small. Hence, the statement is true.

Exercise 11.

If

$$\langle x \rangle = 0$$

then σ_p is infinite.
True or false?

Solution 11.

The standard deviation of the position, σ_x , is not known, and cannot be calculated using the given data. Therefore, as nothing is known about σ_x , nothing can be said about σ_p . Hence, the statement is false.

19 Schrödinger Equation in 1D for Separable Wave Functions

Let the wave function $\Psi(x, t)$ be separable, i.e. let it be expressible in the form

$$\Psi(x, t) = \psi(x)\varphi(t)$$

Let the potential V be independent of time, i.e. let it be a function of x only. Therefore, the Schrödinger equation is

$$\begin{aligned}
i\hbar \frac{\partial \Psi}{\partial t} &= -\frac{\hbar^2}{2m} \frac{\partial^2 \Psi}{\partial x^2} + V\Psi \\
&= \left(-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + V(x) \right) \Psi \\
\therefore \left(i\hbar \frac{\partial}{\partial t} \right) \Psi &= \left(-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + V(x) \right) (\psi(x)\varphi(t)) \\
\therefore i\hbar (\psi(x)\varphi'(t)) &= -\frac{\hbar^2}{2m} \psi''(x)\varphi(t) + V(x)\psi(x)\varphi(t) \\
\therefore i\hbar \frac{\varphi'(t)}{\varphi(t)} &= -\frac{\hbar^2}{2m} \frac{\psi''(x)}{\psi(x)} + V(x)
\end{aligned}$$

For both sides of the equation to be satisfied, they need to be constant. Let this constant be E .

Therefore,

$$\begin{aligned}
\frac{\varphi'(t)}{\varphi(t)} &= \frac{E}{i\hbar} \\
&= -i \frac{E}{\hbar} \\
\therefore \varphi(t) &= c_1 e^{-\frac{iEt}{\hbar}}
\end{aligned}$$

The RHS of the equation is

$$-\frac{\hbar^2}{2m} \psi''(x) + V(x)\psi(x) = E\psi(x)$$

Definition 16. The operator

$$\hat{H} = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + V(x)$$

is called the Hamiltonian or the energy operator.

Therefore,

$$\hat{H}\psi(x) = E\psi(x)$$

Therefore, the E is the eigenvalue of \hat{H} corresponding to $\psi(x)$. Therefore,

$$\begin{aligned}
\Psi(x, t) &= \psi(x)\varphi(t) \\
&= \sum_n c_n \psi_n(x) e^{-\frac{iE_n t}{\hbar}}
\end{aligned}$$

Therefore, $\psi_n(x)$ is an eigenfunction of \hat{H} with eigenvalue E_n .

Theorem 21.

$$\hat{H} = \frac{\hat{p}^2}{2m} + V(\hat{x})$$

20 Stationary States

Definition 17. A state with definite energy, $\psi(x)e^{-\frac{iEt}{\hbar}}$, is called a stationary state.

Theorem 22. Any state can be expressed as a superposition of stationary states.

$$\begin{aligned}\Psi(x, t) &= \psi(x)\varphi(t) \\ &= \sum_n c_n \psi_n(x) e^{-\frac{iE_n t}{\hbar}}\end{aligned}$$

Theorem 23. For a particle in a stationary state, the probability of finding the particle at some x is independent of time.

Theorem 24. For a particle in a stationary state, the expected value of x is independent of time.

Theorem 25. For a particle in a stationary state, the expected value of p is independent of time.

Theorem 26. For a particle in a stationary state, the expected value of any operator Q is independent of time.

Theorem 27. A definite state is a state of definite energy.

Definition 18 (Kronecker delta function). The Kronecker delta function is defined as

$$\delta_{ij} = \begin{cases} 1 & ; \quad i = j \\ 0 & ; \quad i \neq j \end{cases}$$

Theorem 28. The coefficients in

$$\Psi(x) = \sum_{m=1}^{\infty} c_m \psi_m(x)$$

are given by

$$c_m = \int_{-\infty}^{\infty} \psi_m^*(x) \Psi(x) dx$$

Proof.

$$\begin{aligned}
\int_{-\infty}^{\infty} \psi_m^*(x) \Psi(x) \, dx &= \int_{-\infty}^{\infty} \psi_m^*(x) \sum_{n=1}^{\infty} c_n \psi_n(x) \, dx \\
&= \sum_{n=1}^{\infty} \int_{-\infty}^{\infty} c_n \psi_m^*(x) \psi_n(x) \, dx \\
&= \sum_{n=1}^{\infty} c_n \delta_{nm} \\
&= c_m
\end{aligned}$$

As $\psi_m^*(x)$ and $\psi_n(x)$ are eigenfunctions of the Hamiltonian operator, their product is either 0 or 1 depending on the combination of m and n . If $m = n$, then the product is 1, otherwise the product is 0. Therefore, the product can be expressed as the Kronecker delta function.

□

Theorem 29. *For a general state*

$$\psi(x) = \sum_{n=1}^{\infty} c_n \psi_n(x)$$

the expected value of the Hamiltonian operator is

$$\langle H \rangle = \sum_{n=1}^{\infty} P(E = E_n) E_n$$

Proof.

$$\begin{aligned}
\langle H \rangle &= \int_{-\infty}^{\infty} \psi^*(x) \hat{H} \psi(x) \, dx \\
&= \int_{-\infty}^{\infty} \left(\sum_{n=1}^{\infty} c_n^* \psi_n^*(x) \right) \hat{H} \left(\sum_{m=1}^{\infty} c_m \psi_m(x) \right) \, dx \\
&= \int_{-\infty}^{\infty} \sum_{n=1}^{\infty} c_n^* \psi_n^*(x) \sum_{m=1}^{\infty} c_m E_m \psi_m(x) \, dx \\
&= \sum_{n=1}^{\infty} \sum_{m=1}^{\infty} c_n^* c_m E_m \int_{-\infty}^{\infty} \psi_n(x) \psi_m(x) \, dx \\
&= \sum_{n=1}^{\infty} \sum_{m=1}^{\infty} c_n^* c_m E_n \delta_{nm} \\
&= \sum_{n=1}^{\infty} c_n^* c_n E_n \\
&= \sum_{n=1}^{\infty} |c_n|^2 E_n \\
&= \sum_{n=1}^{\infty} P(E = E_n) E_n
\end{aligned}$$

□

21 Infinite Potential Wells

Consider an infinite square well, from $x = 0$ to $x = a$, with the potential given by

$$V(x) = \begin{cases} \infty & ; \quad x < 0 \\ 0 & ; \quad 0 < x < a \\ \infty & ; \quad a < x \end{cases}$$

Consider a particle with mass m , under this potential. Therefore,

$$\hat{H}\psi(x) = E\psi(x)$$

Therefore,

$$-\frac{\hbar^2}{2m} \frac{d^2\psi(x)}{dx^2} + V(x)\psi(x) = E\psi(x)$$

For $x < 0$,

$$V(x) = \infty$$

Therefore,

$$\psi(x) = 0$$

For $a < x$,

$$V(x) = \infty$$

Therefore,

$$\psi(x) = 0$$

For $0 < x < a$,

$$V(x) = 0$$

Therefore,

$$\begin{aligned} -\frac{\hbar^2}{2m} \frac{d^2\psi(x)}{dx^2} &= E\psi(x) \\ \therefore \frac{d^2\psi(x)}{dx^2} &= -\frac{2mE}{\hbar^2}\psi(x) \end{aligned}$$

Let

$$k = \sqrt{\frac{2mE}{\hbar^2}}$$

Therefore, for $0 < x < a$,

$$\psi(x) = A \sin(kx) + B \cos(kx)$$

Therefore, at the boundaries,

$$\psi(x=0) = 0$$

$$\psi(x=a) = 0$$

Therefore, for $\psi(0) = 0$,

$$\begin{aligned} A \sin(0) + B \cos(0) &= 0 \\ \therefore B &= 0 \end{aligned}$$

Therefore, for $\psi(a) = 0$,

$$\begin{aligned} A \sin(ka) &= 0 \\ \therefore \sin ka &= 0 \end{aligned}$$

Therefore, for $n \in \mathbb{N}$,

$$ka = n\pi$$

Therefore,

$$k_n = \frac{n\pi}{a}$$

Therefore,

$$\begin{aligned} k &= \sqrt{\frac{2mE}{\hbar^2}} \\ \therefore \frac{n\pi}{a} &= \sqrt{\frac{2mE}{\hbar^2}} \\ \therefore E &= \frac{\pi^2 n^2 \hbar^2}{2ma^2} \end{aligned}$$

Therefore,

$$\begin{aligned} \psi_n(x) &= A \sin(k_n x) \\ &= A \sin\left(\frac{n\pi}{a} x\right) \end{aligned}$$

Therefore, normalizing the wave function,

$$A = \sqrt{\frac{2}{a}}$$

Therefore,

$$\psi_n(x) = \sqrt{\frac{2}{a}} \sin\left(\frac{n\pi}{a} x\right)$$

If A is zero, the wave function will not be normalizable. Therefore, A cannot be zero. Hence, $\sin(ka)$ has to be zero.

Therefore,

$$\begin{aligned}\psi(x) &= \sum_{n=1}^{\infty} c_n \psi_n(x) \\ &= \sum_{n=1}^{\infty} c_n \sqrt{\frac{2}{a}} \sin\left(\frac{n\pi}{a}x\right)\end{aligned}$$

Therefore, the probability distribution of x is as shown in Figure 8 and Figure 9.

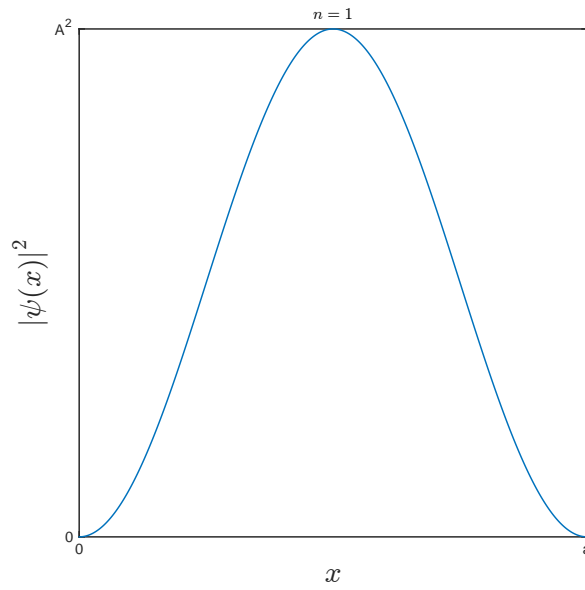


Figure 8: Probability distribution of x , with $n = 1$, in an infinite square well

Therefore, as these states are stationary

$$\Psi(x, t) = \sum_{n=1}^{\infty} c_n \sqrt{\frac{2}{a}} \sin\left(\frac{n\pi}{a}x\right) e^{-\frac{iE_n t}{\hbar}}$$

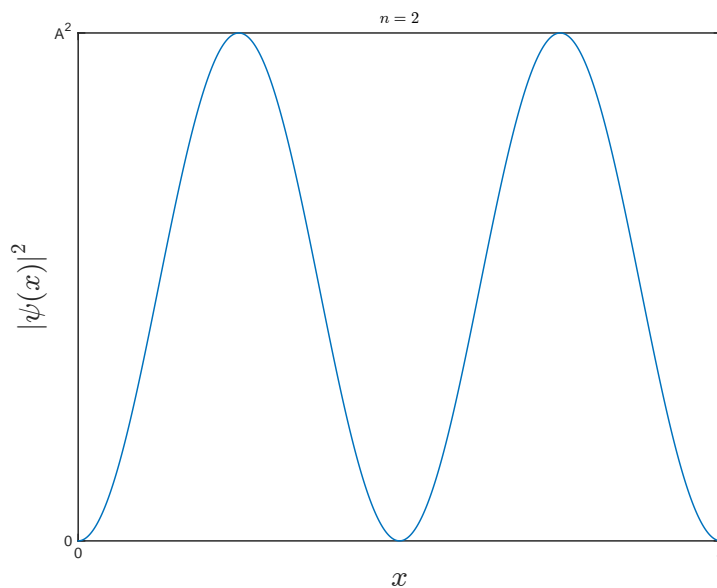


Figure 9: Probability distribution of x , with $n = 2$, in an infinite square well

Exercise 12.

The wave function of a particle of mass m in an infinite square well from $x = 0$ to $x = a$, at time $t = 0$ is given by

$$\psi(x, 0) = Ax(a - x)$$

for $0 < x < a$.

1. Find the general wave function $\psi(x, t)$.
2. What is the probability of measuring the following energies at time $t = 0$?

(a) $E_1 = \frac{\pi^2 \hbar^2}{2ma^2}$

(b) $E_2 = \frac{4\pi^2 \hbar^2}{2ma^2}$

(c) $E_3 = \frac{9\pi^2 \hbar^2}{2ma^2}$

Solution 12.

1.

$$1 = \int_0^a A^2 x^2 (a - x)^2 dx$$

Therefore, solving,

$$A = \sqrt{\frac{30}{a^5}}$$

Therefore,

$$Ax(a - x) = \sum_{n=1}^{\infty} c_n \sqrt{\frac{2}{a}} \sin\left(\frac{n\pi}{a}x\right)$$

Therefore,

$$\begin{aligned} c_n &= \int_0^a \sqrt{\frac{2}{a}} \sin\left(\frac{n\pi}{a}x\right) Ax(a - x) dx \\ &= \sqrt{\frac{2}{a}} Aa \int_0^a \sin\left(\frac{n\pi}{a}x\right) dx - \sqrt{\frac{2}{a}} A \int_0^a \sin\left(\frac{n\pi}{a}x\right) x^2 dx \end{aligned}$$

Therefore,

$$c_n = \begin{cases} 0 & ; \quad n \text{ is even} \\ \frac{8\sqrt{15}}{n^3\pi^3} & ; \quad n \text{ is odd} \end{cases}$$

Therefore,

$$\psi(x, t) = \sum_{n=1,3,5,\dots} \sqrt{\frac{2}{a}} \frac{8\sqrt{15}}{n^3\pi^3} \sin\left(\frac{n\pi}{a}x\right) e^{-\frac{iE_n t}{\hbar}}$$

where

$$E_n = \frac{n^2\pi^2\hbar^2}{2ma^2}$$

2. (a)

$$\begin{aligned} P(E = E_1) &= |c_1|^2 \\ &= \left(\frac{8\sqrt{15}}{\pi^3}\right)^2 \\ &= 0.9986 \end{aligned}$$

(b)

$$\begin{aligned} P(E = E_2) &= |c_2|^2 \\ &= 0 \end{aligned}$$

For even n , c_n is zero.

(c)

$$\begin{aligned} P(E = E_3) &= |c_3|^2 \\ &= \left(\frac{8\sqrt{15}}{27\pi^3} \right)^2 \\ &= 0.0014 \end{aligned}$$

Exercise 13.

A particle with mass m in an infinite square well, has initial wave function

$$\psi(x, 0) = A (\psi_1(x) + \psi_2(x))$$

1. Find A .
2. Find the general wave function $\psi(x, t)$.
3. What is the probability of measuring the following energies at time $t = 0$?

(a) $E_1 = \frac{\pi^2 \hbar^2}{2ma^2}$

(b) $E_2 = \frac{4\pi^2 \hbar^2}{2ma^2}$

4. Find the wave function if the energy is measured to be E_1 .
5. Find the expectation value of x .

Solution 13.

1. As the square of the coefficients represents the probability of a particular value,

$$\begin{aligned} 1 &= \sum_{n=1}^{\infty} |c_n|^2 \\ &= 2A^2 \\ \therefore \frac{1}{\sqrt{2}} &= A \end{aligned}$$

2.

$$\begin{aligned}\psi(x, t) &= A \left(\psi_1(x) e^{-\frac{iE_1 t}{\hbar}} + \psi_2(x) e^{-\frac{iE_2 t}{\hbar}} \right) \\ &= \frac{1}{\sqrt{2}} \left(\sqrt{\frac{2}{a}} \sin\left(\frac{\pi}{a}x\right) e^{-\frac{iE_1 t}{\hbar}} + \sqrt{\frac{2}{a}} \sin\left(\frac{2\pi}{a}x\right) e^{-\frac{iE_2 t}{\hbar}} \right)\end{aligned}$$

where

$$\begin{aligned}E_1 &= \frac{\pi^2 \hbar^2}{2ma^2} \\ E_2 &= \frac{4\pi^2 \hbar^2}{2ma^2}\end{aligned}$$

3. (a)

$$\begin{aligned}P(E = E_1) &= |c_1|^2 \\ &= \left| \frac{1}{\sqrt{2}} e^{-\frac{iE_1 t}{\hbar}} \right|^2 \\ &= \frac{1}{2}\end{aligned}$$

(b)

$$\begin{aligned}P(E = E_2) &= |c_2|^2 \\ &= \left| \frac{1}{\sqrt{2}} e^{-\frac{iE_2 t}{\hbar}} \right|^2 \\ &= \frac{1}{2}\end{aligned}$$

4. If E is measured to be E_1 , then the wave function collapsed to

$$\psi(x) = \psi_1(x) e^{-\frac{iE_1 t}{\hbar}}$$

5. At time t ,

$$\begin{aligned}\langle x \rangle_t &= \int_0^a x |\psi(x, t)|^2 dx \\ &= A^2 \int_0^a x \left(\psi_1^*(x) e^{\frac{iE_1 t}{\hbar}} + \psi_2^*(x) e^{\frac{iE_2 t}{\hbar}} \right) \left(\psi_1(x) e^{-\frac{iE_1 t}{\hbar}} + \psi_2(x) e^{-\frac{iE_2 t}{\hbar}} \right) dx \\ &\neq \langle x \rangle_{t=0}\end{aligned}$$

22 Free Particles

For a free particle,

$$V(x) = 0$$

Therefore,

$$\begin{aligned}\hat{H}\psi(x) &= E\psi(x) \\ \therefore -\frac{\hbar^2}{2m} \frac{d^2\psi(x)}{dx^2} &= E\psi(x) \\ \therefore \frac{d^2\psi(x)}{dx^2} &= -\frac{2Em\psi(x)}{\hbar^2} \\ &= -k^2\psi(x)\end{aligned}$$

where

$$k = \sqrt{\frac{2mE}{\hbar^2}}$$

Therefore,

$$\psi(x) = Ae^{ikx}$$

However, this wave function is not normalizable. Hence, a free particle cannot exist in a stationary state.

The general solution to the time-dependent Schrödinger equation is a linear combination of separable solutions, i.e.

$$\Psi(x, t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \varphi(k) e^{i\left(kx - \frac{\hbar k^2}{2m}t\right)} dk$$

This function can be normalized, but it carries a multiple ks , and hence multiple energies and speeds.

Definition 19 (Wave packet). A solution of the form

$$\Psi(x, t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \varphi(k) e^{i\left(kx - \frac{\hbar k^2}{2m}t\right)} dk$$

is called a wave packet.

Therefore, for $t = 0$,

$$\Psi(x, 0) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \varphi(k) e^{ikx} dk$$

Therefore,

$$\varphi(k) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \Psi(x, 0) e^{-ikx} dx$$

Exercise 14.

A free particle, initially localized as

$$\psi(x, 0) = \begin{cases} A & ; \quad x \in (-a, a) \\ 0 & ; \quad x \notin (-a, a) \end{cases}$$

where A and a are real and positive.
Find $\psi(x, t)$.

Solution 14.

$$\begin{aligned} 1 &= \int_{-\infty}^{\infty} |\psi(x, 0)|^2 dx \\ &= \int_{-a}^a A^2 dx \\ &= 2aA^2 \end{aligned}$$

Therefore,

$$A = \frac{1}{\sqrt{2a}}$$

Therefore, taking the inverse Fourier transform,

$$\begin{aligned}
\tilde{\psi}(k) &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \psi(x, 0) e^{-ikx} dx \\
&= \frac{1}{\sqrt{2\pi}} \frac{1}{\sqrt{2a}} \int_{-a}^a e^{-ikx} dx \\
&= \frac{1}{\sqrt{2\pi}} \frac{1}{\sqrt{2a}} \int_{-a}^a e^{-ikx} dx \\
&= \frac{1}{\sqrt{4\pi a}} \frac{1}{-ik} e^{-ikx} \Big|_{-a}^a \\
&= \frac{1}{\sqrt{4\pi a}} \frac{1}{-ik} (e^{-ika} - e^{ika}) \\
&= \frac{1}{\sqrt{\pi a}} \frac{1}{k} \frac{e^{ika} - e^{-ika}}{2i} \\
&= \frac{1}{k\sqrt{\pi a}} \sin(ka)
\end{aligned}$$

Therefore,

$$\begin{aligned}
\psi(x, t) &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \tilde{\psi}(k) e^{ikx} e^{-\frac{iEt}{\hbar}} dx \\
&= \frac{1}{\sqrt{2\pi^2 a}} \int_{-\infty}^{\infty} \frac{\sin(ka)}{k} e^{ikx} e^{-\frac{ik^2 \hbar}{2m} t} dk
\end{aligned}$$

If a is very small,

$$\begin{aligned}
\tilde{\psi}(k) &= \frac{\sin(ka)}{k\sqrt{\pi a}} \\
&= \frac{ka}{k\sqrt{\pi a}} \\
&= \sqrt{\frac{a}{\pi}}
\end{aligned}$$

If a is very large,

$$\begin{aligned}\tilde{\psi}(k) &= \frac{\sin(ka)}{k\sqrt{\pi a}} \\ &= \sqrt{\frac{a}{\pi}} \frac{\sin(ka)}{ka} \\ &= \sqrt{\frac{a}{\pi}} \text{sinc}(ka)\end{aligned}$$

23 Finite Potential Wells

Definition 20 (Bound state). An eigenfunction of \hat{H} , which satisfies

$$\begin{aligned}\lim_{x \rightarrow \infty} \psi(x) &= 0 \\ \lim_{x \rightarrow -\infty} \psi(x) &= 0\end{aligned}$$

is called a bound state.

Definition 21 (Scattering state). An eigenfunction of \hat{H} , which satisfies

$$\begin{aligned}\lim_{x \rightarrow \infty} \psi(x) &\neq 0 \\ \lim_{x \rightarrow -\infty} \psi(x) &\neq 0\end{aligned}$$

is called a scattering state.

Theorem 30. *A particle in a finite or infinite well is in a bound state.*

Theorem 31. *A free particle is in a scattering state.*

Consider a finite square well such that

$$V(x) = \begin{cases} -V_0 & ; \quad x \in (-a, a) \\ 0 & ; \quad x \notin (-a, a) \end{cases}$$

Consider a particle with energy $-V_0 < E < 0$.

For $x < -a$,

$$V = 0$$

Therefore,

$$\frac{d^2\psi(x)}{dx^2} = -\frac{2mE}{\hbar^2}\psi(x)$$

Let

$$k = \sqrt{-\frac{2mE}{\hbar^2}}$$

Therefore if $E < 0$, k is real.

Therefore,

$$\frac{d^2\psi(x)}{dx^2} = k^2\psi(x)$$

Therefore,

$$\psi(x) = Ae^{kx} + Be^{-kx}$$

For the wave function to be normalizable, $\lim_{x \rightarrow -\infty} \psi(x)$ must be zero.

Therefore,

$$\begin{aligned} 0 &= \lim_{x \rightarrow -\infty} \psi(x) \\ &= \lim_{x \rightarrow -\infty} Ae^{kx} + Be^{-kx} \end{aligned}$$

Therefore B must be zero.

Therefore,

$$\psi(x) = Ae^{kx}$$

Similarly, for $x > a$,

$$\psi(x) = Ce^{-kx}$$

where

$$k = \sqrt{-\frac{2mE}{\hbar^2}}$$

For $-a < x < a$,

$$\begin{aligned} \frac{d^2\psi}{dx^2} &= -\frac{2m}{\hbar} (E - V(x)) \psi(x) \\ &= -\frac{2m}{\hbar} (E + V_0) \psi(x) \end{aligned}$$

Let

$$\beta = \sqrt{\frac{2m(E + V_0)}{\hbar^2}}$$

Therefore,

$$\frac{d^2\psi(x)}{dx^2} = -\beta^2\psi(x)$$

Therefore,

$$\psi(x) = D \cos(\beta x) + F \sin(\beta x)$$

Therefore,

$$\psi(x) = \begin{cases} Ae^{kx} & ; \quad x < -a \\ D \cos(\beta x) + F \sin(\beta x) & ; \quad -a < x < a \\ Ce^{-kx} & ; \quad a < x \end{cases}$$

where

$$k = \sqrt{-\frac{2mE}{\hbar^2}}$$

$$\beta = \sqrt{\frac{2m(E + V_0)}{\hbar^2}}$$

Theorem 32. *If $V(x)$ is an even function, then the eigenfunctions of \hat{H} are either odd or even.*

23.1 Even Wave Functions for Finite Potential Wells

If $\psi(x)$ is even,

$$\psi(x) = \begin{cases} Ae^{kx} & ; \quad x < -a \\ D \cos(\beta x) & ; \quad -a < x < a \\ Ce^{-kx} & ; \quad a < x \end{cases}$$

As $|\psi(x)|^2$ represents probability, $\psi(x)$ must always be continuous.

As the well is finite, the jump in the expression of $V(x)$ is also finite. Hence, $\psi'(x)$ must always be continuous.

Therefore, as $\psi(x)$ is always continuous,

$$Ae^{-ka} = D \cos(-\beta a)$$

Therefore, as $\psi'(x)$ is always continuous,

$$kAe^{-ka} = \beta D \cos(\beta a)$$

Similarly, for $x = a$,

$$\begin{aligned} Ce^{-ka} &= D \cos(\beta a) \\ -kCe^{-ka} &= -\beta D \sin(\beta a) \end{aligned}$$

Therefore,

$$A = C$$

Therefore,

$$k = \beta \tan(\beta a)$$

where

$$\begin{aligned} k &= \sqrt{\frac{-2mE}{\hbar^2}} \\ \beta &= \sqrt{\frac{2m(E + V_0)}{\hbar^2}} \end{aligned}$$

Let

$$\begin{aligned} z &= \beta a \\ y &= ka \end{aligned}$$

Therefore,

$$\begin{aligned} k &= \beta \tan(\beta a) \\ \therefore ka &= \beta a \tan(\beta a) \\ \therefore y &= z \tan(z) \end{aligned}$$

Also,

$$\begin{aligned} \beta^2 + k^2 &= \frac{2mV_0}{\hbar^2} \\ \therefore z^2 + y^2 &= a^2 (\beta^2 + k^2) \\ &= \frac{2mV_0 a^2}{\hbar^2} \end{aligned}$$

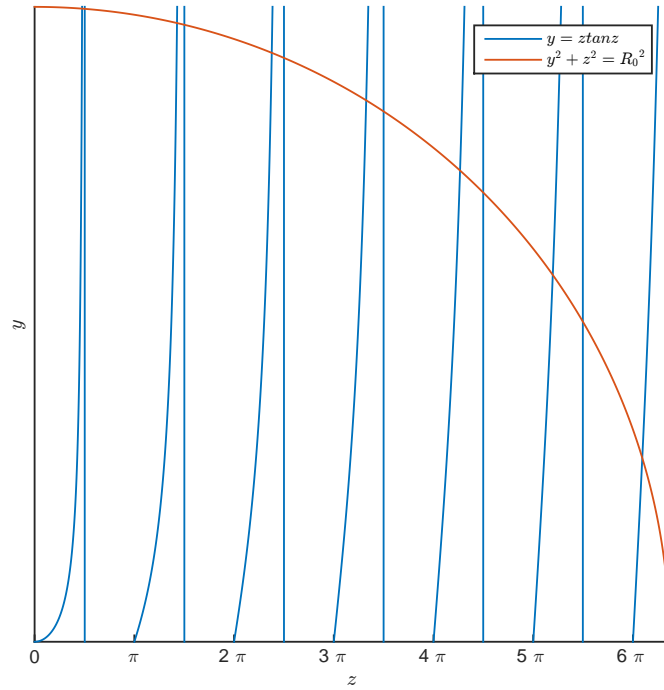
Let

$$R_0^2 = \frac{2mV_0a^2}{\hbar^2}$$

Therefore,

$$z^2 + y^2 = R_0^2$$

As a , k , and B are real and positive, z and y also must be real and positive. Therefore,



As the two graphs must intersect at least once, the equation has at least one solution.

If $R_0 < \pi$, the graphs intersect exactly once. Hence, if $R_0 < \pi$, there exists a single solution.

If $R_0 \rightarrow \infty$, there are infinitely many solutions, $z = \frac{n\pi}{2}$, where n is odd. Hence,

$$z_n = \frac{n\pi}{2}$$

$$\therefore \beta_n = \frac{n\pi}{2a}$$

Exercise 15.

The potential is given by

$$V(x) = \begin{cases} \infty & ; \quad x < 0 \\ 0 & ; \quad 0 < x < L \\ V_0 & ; \quad L < x \end{cases}$$

Solve for the bound states for $E < V_0$.

Solution 15.

For $x < 0$,

$$\psi(x) = 0$$

For $0 < x < L$,

$$\psi(x) = A \sin(kx) + B \cos(kx)$$

where

$$k = \sqrt{\frac{2mE}{\hbar^2}}$$

For $x > L$,

$$\psi(x) = Ce^{-\alpha x}$$

where

$$\alpha = \sqrt{\frac{2m(V_0 - E)}{\hbar^2}}$$

For ψ to be continuous at $x = L$, and ψ and ψ' to be continuous at $x = 0$,

$$B = 0$$

$$A \sin(kL) = Ce^{-\alpha L}$$

$$kA \cos(kL) = -\alpha Ce^{-\alpha L}$$

Therefore,

$$k \cot(kL) = -\alpha$$

$$\therefore kL \cot(kL) = -\alpha L$$

Let

$$z = kL$$

$$y = \alpha L$$

Therefore,

$$z \cot(z) = -y$$

Also,

$$\begin{aligned}\alpha^2 + k^2 &= \frac{2mV_0}{\hbar^2} \\ \therefore \alpha^2 L^2 + k^2 L^2 &= \frac{2mV_0 L^2}{\hbar^2} \\ \therefore y^2 + z^2 &= \frac{2mV_0 L^2}{\hbar^2}\end{aligned}$$

23.2 Odd Wave Functions for Finite Potential Wells

If $\psi(x)$ is odd,

$$\psi(x) = \begin{cases} Ae^{kx} & ; \quad x < -a \\ F \sin(\beta x) & ; \quad -a < x < a \\ Ce^{-kx} & ; \quad a < x \end{cases}$$

As $|\psi(x)|^2$ represents probability, $\psi(x)$ must always be continuous.

As the well is finite, the jump in the expression of $V(x)$ is also finite. Hence, $\psi'(x)$ must always be continuous.

Therefore, as $\psi(x)$ is always continuous,

$$Ae^{-ka} = F \sin(-\beta a)$$

Therefore, as $\psi'(x)$ is always continuous,

$$kAe^{-ka} = \beta F \cos(-\beta a)$$

Similarly, for $x = a$,

$$\begin{aligned}Ce^{-ka} &= F \sin(\beta a) \\ -kCe^{-ka} &= \beta F \cos(\beta a)\end{aligned}$$

Therefore,

$$A = -C$$

Therefore,

$$-k = \beta \cot(\beta a)$$

where

$$k = \sqrt{\frac{-2mE}{\hbar^2}}$$
$$\beta = \sqrt{\frac{2m(E + V_0)}{\hbar^2}}$$

Let

$$z = \beta a$$

$$y = ka$$

Therefore,

$$k = -\beta \cot(\beta a)$$
$$\therefore ka = -\beta a \cot(\beta a)$$
$$\therefore y = -z \cot(z)$$

Also,

$$\beta^2 + k^2 = \frac{2mV_0}{\hbar^2}$$
$$\therefore z^2 + y^2 = a^2 (\beta^2 + k^2)$$
$$= \frac{2mV_0 a^2}{\hbar^2}$$

Let

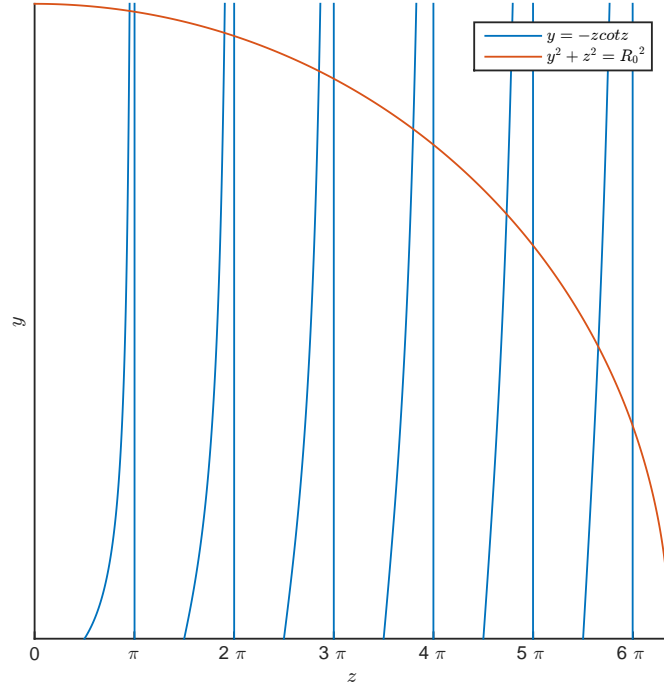
$$R_0^2 = \frac{2mV_0 a^2}{\hbar^2}$$

Therefore,

$$z^2 + y^2 = R_0^2$$

As a , k , and B are real and positive, z and y also must be real and positive.

Therefore,



If $R_0 < \frac{\pi}{2}$, the two graphs do not intersect. Therefore, there is no solution. If $R_0 \rightarrow \infty$, there are infinitely many solution, $z = n\pi$, where n is positive.

$$z_n = n\pi$$

$$\therefore \beta_n = \frac{n\pi}{a}$$

24 $\delta(x)$ Potential

Consider a potential

$$V(x) = -V_0\delta(x)$$

Therefore, for $x < 0$,

$$\begin{aligned} \frac{d^2\psi}{dx^2} &= -\frac{2m}{\hbar^2} (E - V(x)) \psi \\ &= -\frac{2mE}{\hbar^2} \psi \end{aligned}$$

Let

$$k = \sqrt{-\frac{2mE}{\hbar^2}}$$

Therefore, for $E < 0$, k is real.

Therefore,

$$\frac{d^2\psi}{dx^2} = k^2\psi$$

Therefore,

$$\psi(x) = Ae^{kx} + Be^{-kx}$$

If $x \rightarrow -\infty$, as the function is normalizable, $\psi = 0$.

Therefore,

$$B = 0$$

Similarly, for $x > 0$,

$$\psi(x) = Ce^{kx} + De^{-kx}$$

If $x \rightarrow \infty$, as the function is normalizable, $\psi = 0$.

Therefore,

$$C = 0$$

Therefore,

$$\psi(x) = \begin{cases} Ae^{kx} & ; \quad x < 0 \\ Ce^{-kx} & ; \quad x > 0 \end{cases}$$

Therefore, as $\psi(x)$ represents probability, it is continuous at $x = 0$.

Therefore,

$$A = C$$

Therefore,

$$\psi(x) = \begin{cases} Ae^{kx} & ; \quad x < 0 \\ Ae^{-kx} & ; \quad x > 0 \end{cases}$$

Therefore,

$$\psi'(x) = \begin{cases} Ake^{kx} & ; \quad x < 0 \\ -Ake^{-kx} & ; \quad x > 0 \end{cases}$$

By the time independent Schrödinger equation,

$$\psi''(x) = -\frac{2m}{\hbar^2} (E + V_0\delta(x)) \psi$$

Therefore,

$$\begin{aligned} \int_{-\varepsilon}^{\varepsilon} \psi''(x) dx &= \int_{-\varepsilon}^{\varepsilon} -\frac{2m}{\hbar^2} (E + V_0\delta(x)) \psi(x) dx \\ \therefore \psi'(\varepsilon) - \psi'(-\varepsilon) &= -\frac{2m}{\hbar^2} \int_{-\varepsilon}^{\varepsilon} E\psi(x) dx - \frac{2m}{\hbar^2} \int_{-\varepsilon}^{\varepsilon} V_0\delta(x)\psi(x) dx \\ &= -\frac{2m}{\hbar^2} (E\psi(0) \cdot 2\varepsilon) - \frac{2m}{\hbar^2} V_0\psi(0) \\ &= -\frac{2m}{\hbar^2} V_0\psi(0) \\ \therefore -kA - kA &= -\frac{2m}{\hbar^2} V_0A \\ \therefore k &= \frac{mV_0}{\hbar^2} \end{aligned}$$

Therefore,

$$\begin{aligned} k &= \sqrt{-\frac{2mE}{\hbar^2}} \\ \therefore E &= -\frac{k^2\hbar^2}{2m} \\ &= -\frac{mV_0^2}{2\hbar^2} \end{aligned}$$

Normalizing,

$$\begin{aligned} A &= \sqrt{k} \\ &= \sqrt{\frac{mV_0}{\hbar^2}} \end{aligned}$$

25 Tunneling

Consider a potential

$$V(x) = \begin{cases} 0 & ; \quad x < 0 \\ V_0 & ; \quad x > 0 \end{cases}$$

Therefore, for $x < 0$,

$$\begin{aligned} \frac{d^2\psi(x)}{dx^2} &= -\frac{2m}{\hbar^2} (E - V(x)) \psi(x) \\ &= -\frac{2m}{\hbar^2} E \psi(x) \end{aligned}$$

Let

$$k_1 = \frac{2mE}{\hbar^2}$$

Therefore,

$$\psi'' = -k_1^2 \psi$$

Therefore, for $x > 0$,

$$\begin{aligned} \frac{d^2\psi(x)}{dx^2} &= -\frac{2m}{\hbar^2} (E - V(x)) \psi(x) \\ &= -\frac{2m}{\hbar^2} (E - V_0) \psi(x) \end{aligned}$$

Let

$$k_2 = -\frac{2m(E - V_0)}{\hbar^2}$$

Therefore,

$$\psi'' = k_2^2 \psi$$

Therefore,

$$\psi(x) = \begin{cases} Ae^{ik_1x} + Be^{-ik_1x} & ; \quad x < 0 \\ Ce^{k_2x} + De^{-k_2x} & ; \quad x > 0 \end{cases}$$

For the function to be normalizable, for $x \rightarrow \infty$, $\psi = 0$.

Therefore,

$$C = 0$$

Therefore,

$$\psi(x) = \begin{cases} Ae^{ik_1x} + Be^{-ik_1x} & ; \quad x < 0 \\ De^{-k_2x} & ; \quad x > 0 \end{cases}$$

25.1 δ Barrier

Consider a potential

$$V(x) = V_0\delta(x)$$

Consider a single particle approaching the δ barrier from the left. Therefore, solving the time independent Schrödinger equation,

$$\psi(x) = \begin{cases} Ae^{ikx} + Be^{-ikx} & ; \quad x < 0 \\ Ce^{ikx} + De^{-ikx} & ; \quad 0 < x \end{cases}$$

where

$$k = \sqrt{\frac{2mE}{\hbar^2}}$$

Therefore, Ae^{ikx} represents the incident wave, Be^{-ikx} represents the reflected wave, and Ce^{ikx} represents the transmitted wave.

As a single particle approaches the barrier from the left, the wave represented by De^{-ikx} does not exist. Hence,

$$D = 0$$

Therefore,

$$\psi(x) = \begin{cases} Ae^{ikx} + Be^{-ikx} & ; \quad x < 0 \\ Ce^{ikx} & ; \quad 0 < x \end{cases}$$

As $\psi(x)$ is continuous at $x = 0$,

$$A + B = C$$

Therefore, solving,

$$\psi'(x) = \begin{cases} ik(Ae^{ikx} - Be^{-ikx}) & ; \quad x < 0 \\ ikCe^{ikx} & ; \quad 0 < x \end{cases}$$

As the barrier is a Dirac Delta function,

$$\psi'(\varepsilon) - \psi'(-\varepsilon) = \frac{2mv_0}{\hbar^2}\psi(0)$$

Therefore,

$$(ikC) - (ikA - ikB) = \frac{2mv_0}{\hbar^2}\psi(0)$$

Let

$$\beta = \frac{mv_0}{\hbar^2 k}$$

Therefore,

$$\frac{B}{A} = \frac{-i\beta}{1 + i\beta}$$

$$\frac{C}{A} = \frac{1}{1 + i\beta}$$

Similarly, if

$$V(x) = -V_0\delta(x)$$

then,

$$\psi'(\varepsilon) - \psi'(-\varepsilon) = -\frac{2mv_0}{\hbar^2}\psi(0)$$

Therefore,

$$(ikC) - (ikA - ikB) = -\frac{2mv_0}{\hbar^2}\psi(0)$$

Let

$$\beta = \frac{mv_0}{\hbar^2 k}$$

Therefore,

$$\frac{B}{A} = \frac{i\beta}{1 + i\beta}$$

$$\frac{C}{A} = \frac{1}{1 - i\beta}$$

25.2 Double δ Barrier

Consider a potential

$$V(x) = -\alpha (\delta(x + a) + \delta(x - a))$$

Therefore,

$$\psi(x) = \begin{cases} Ae^{kx} & ; \quad x < -a \\ Be^{kx} + Ce^{-kx} & ; \quad -a < x < a \\ De^{-kx} & ; \quad a < x \end{cases}$$

where

$$k = \sqrt{-\frac{2mE}{\hbar^2}}$$

Therefore,

$$\begin{aligned} \psi(-a + \varepsilon) &= \psi(-a - \varepsilon) \\ \psi(a + \varepsilon) &= \psi(a - \varepsilon) \\ \psi'(-a + \varepsilon) - \psi'(-a - \varepsilon) &= -\frac{2m\alpha}{\hbar^2}\psi(-a) \\ \psi'(a + \varepsilon) - \psi'(a - \varepsilon) &= -\frac{2m\alpha}{\hbar^2}\psi(a) \end{aligned}$$

As $V(x)$ is even, the wave function $\psi(x)$ can be split into odd and even parts, i.e. $\psi_{\text{even}}(x)$ and $\psi_{\text{odd}}(x)$

Therefore, for $\psi_{\text{even}}(x)$,

$$\begin{aligned} A &= D \\ B &= C \end{aligned}$$

Therefore,

$$\psi_{\text{even}}(x) = \begin{cases} Ae^{kx} & ; \quad x < -a \\ B(e^{kx} + e^{-kx}) & ; \quad -a < x < a \\ Ae^{-kx} & ; \quad a < x \end{cases}$$

Therefore,

$$\psi'(x) = \begin{cases} kAe^{kx} & ; \quad x < -a \\ kB(e^{kx} - e^{-kx}) & ; \quad -a < x < a \\ -kAe^{-kx} & ; \quad a < x \end{cases}$$

Therefore,

$$\begin{aligned}
\psi(-a + \varepsilon) &= \psi(-a - \varepsilon) \\
\therefore Ae^{-ka} &= B(e^{-ka} + e^{ka}) \\
\psi'(-a + \varepsilon) - \psi'(-a - \varepsilon) &= -\frac{2m\alpha}{\hbar^2}\psi(-a) \\
\therefore kB(e^{-ka} - e^{ka}) - kAe^{-ka} &= -\frac{2m\alpha}{\hbar^2}Ae^{-ka}
\end{aligned}$$

Therefore, solving,

$$\begin{aligned}
A &= B(e^{2ka} + 1) \\
B(e^{2ka} - 1) &= A\left(\frac{2m\alpha}{\hbar^2 k} - 1\right)
\end{aligned}$$

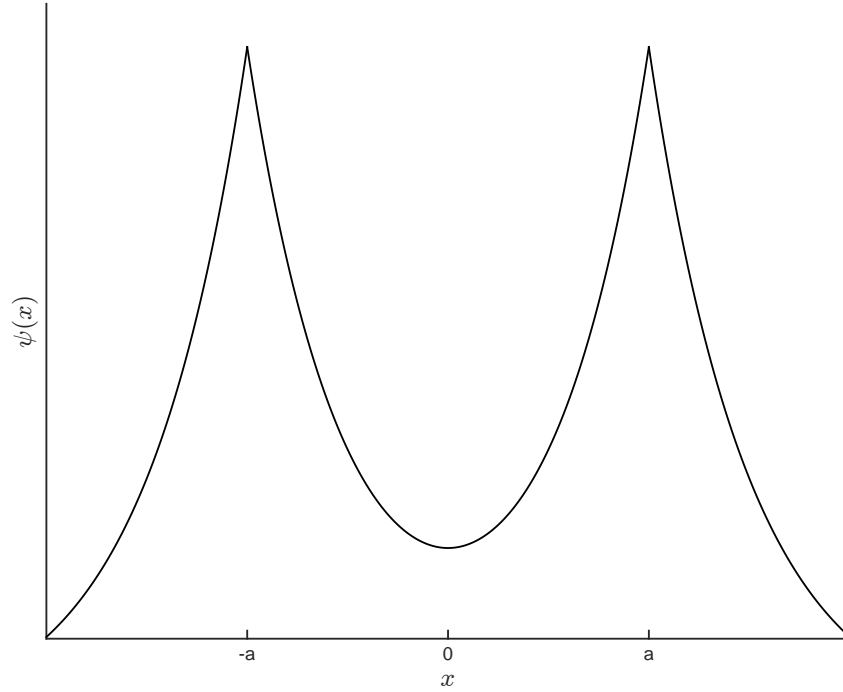
Let

$$\begin{aligned}
z &= 2ka \\
c &= \frac{\hbar^2}{2am\alpha}
\end{aligned}$$

Therefore,

$$\begin{aligned}
e^{-2ka} &= \frac{\hbar^2 k}{m\alpha} - 1 \\
\therefore e^{-z} &= cz - 1
\end{aligned}$$

Therefore, this equation has exactly one solution.



Therefore, for $\psi_{\text{odd}}(x)$,

$$A = -D$$

$$B = -C$$

Therefore,

$$\psi_{\text{odd}}(x) = \begin{cases} -Ae^{kx} & ; \quad x < -a \\ B(e^{kx} - e^{-kx}) & ; \quad -a < x < a \\ Ae^{-kx} & ; \quad a < x \end{cases}$$

26 Probability Current

Definition 22 (Probability current). The rate at which the probability distribution passes through the position x is defined to be the probability current. It is denoted as

$$J(x, t) = \frac{i\hbar}{2m} \left(\psi \frac{\partial \psi^*}{\partial x} - \psi^* \frac{\partial \psi}{\partial x} \right)$$

26.1 δ Barrier

Consider a potential

$$V(x) = V_0\delta(x)$$

Therefore, the wave function is

$$\psi(x) = \begin{cases} Ae^{ikx} + Be^{-ikx} & ; \quad x < 0 \\ Ce^{ikx} & ; \quad x > 0 \end{cases}$$

Let the incident, reflected, and transmitted wave functions be

$$\psi_I = Ae^{ikx}$$

$$\psi_R = Be^{-ikx}$$

$$\psi_T = Ce^{ikx}$$

Therefore,

$$\psi_I(x) = Ae^{ikx}$$

$$\psi_I(x, t) = Ae^{ikx} e^{-\frac{iEt}{\hbar}}$$

$$\psi_I^*(x, t) = A^* e^{-ikx} e^{\frac{iEt}{\hbar}}$$

Therefore,

$$\frac{\partial \psi_I}{\partial x} = ikAe^{ikx} e^{-\frac{iEt}{\hbar}}$$

$$\frac{\partial \psi_I^*}{\partial x} = -ikA^* e^{-ikx} e^{\frac{iEt}{\hbar}}$$

Therefore, the probability current for the incident wave is

$$J(x, t) = \frac{i\hbar}{2m} \left(\psi \frac{\partial \psi^*}{\partial x} - \psi^* \frac{\partial \psi}{\partial x} \right)$$

$$\begin{aligned} J_I(x, t) &= \frac{i\hbar}{2m} \left(\psi_I \frac{\partial \psi_I^*}{\partial x} - \psi_I^* \frac{\partial \psi_I}{\partial x} \right) \\ &= \frac{i\hbar}{2m} (-ik|A|^2 - ik|A|^2) \\ &= \frac{\hbar|A|^2 k}{m} \end{aligned}$$

Similarly, the probability current for the reflected wave is

$$\begin{aligned}
J(x, t) &= \frac{i\hbar}{2m} \left(\psi \frac{\partial \psi^*}{\partial x} - \psi^* \frac{\partial \psi}{\partial x} \right) \\
J_R(x, t) &= \frac{i\hbar}{2m} \left(\psi_R \frac{\partial \psi_R^*}{\partial x} - \psi_R^* \frac{\partial \psi_R}{\partial x} \right) \\
&= \frac{i\hbar}{2m} (ik|B|^2 + ik|B|^2) \\
&= -\frac{\hbar|B|^2 k}{m}
\end{aligned}$$

Similarly, the probability current for the incident wave is

$$\begin{aligned}
J(x, t) &= \frac{i\hbar}{2m} \left(\psi \frac{\partial \psi^*}{\partial x} - \psi^* \frac{\partial \psi}{\partial x} \right) \\
J_T(x, t) &= \frac{i\hbar}{2m} \left(\psi_T \frac{\partial \psi_T^*}{\partial x} - \psi_T^* \frac{\partial \psi_T}{\partial x} \right) \\
&= \frac{i\hbar}{2m} (-ik|C|^2 - ik|C|^2) \\
&= \frac{\hbar|C|^2 k}{m}
\end{aligned}$$

Let the transmission and reflection coefficients be T and R respectively. Therefore, substituting the boundary conditions,

$$\begin{aligned}
T &= \frac{|J_T|}{|J_I|} \\
&= \frac{|C|^2}{|A|^2} \\
&= \frac{1}{1 + \frac{mV_0^2}{2\hbar^2 E}} \\
R &= \frac{|J_R|}{|J_I|} \\
&= \frac{|B|^2}{|A|^2} \\
&= \frac{1}{1 + \frac{2\hbar^2 E}{mV_0^2}}
\end{aligned}$$

Therefore,

$$T + R = 1$$

This is consistent with the expectations, as the total energy must be conserved. The fact that $T \neq 0$ shows that even though the barrier exists, the particle does not bounce back, but instead tunnels through the barrier. A classical particle would not exhibit such a behaviour.

26.2 Finite Barrier

Consider a potential

$$V(x) = \begin{cases} 0 & ; \quad x < 0 \\ V_0 & ; \quad 0 < x < L \\ 0 & ; \quad L < x \end{cases}$$

Therefore, this potential can be considered to be

$$V(x) = V_1(x) - V_2(x)$$

where

$$V_1(x) = \begin{cases} 0 & ; \quad x < 0 \\ V_0 & ; \quad x > 0 \end{cases}$$

$$V_2(x) = \begin{cases} 0 & ; \quad x < L \\ V_0 & ; \quad x > L \end{cases}$$

Consider a single particle approaching the barrier from the left, Hence, due to tunnelling, a quantum particle can go through this barrier, and the wave function is

$$\psi(x) = \begin{cases} Ae^{ik_1x} + Be^{-ik_1x} & ; \quad x < 0 \\ Ce^{k_2x} + De^{-k_2x} & ; \quad 0 < x < L \\ Fe^{ik_1x} + Ge^{-ik_1x} & ; \quad L < x \end{cases}$$

where

$$k_1 = \sqrt{\frac{2mE}{\hbar^2}}$$

$$k_2 = \sqrt{-\frac{2m(E - V_0)}{\hbar^2}}$$

Therefore, solving as for the δ barrier,

$$\begin{aligned}
 T &= \frac{|J_T|}{|J_I|} \\
 &= \frac{|F|^2}{|A|^2} \\
 &= \frac{4k_1^2 k_2^2}{(k_1^2 + k_2^2) \sinh^2(k_2 L) + 4k_1^2 k_2^2} \\
 R &= \frac{|J_R|}{|J_I|} \\
 &= \frac{|B|^2}{|A|^2} \\
 &= \frac{(k_1^2 + k_2^2) \sinh^2(k_2 L)}{(k_1^2 + k_2^2) \sinh^2(k_2 L) + 4k_1^2 k_2^2}
 \end{aligned}$$

27 Quantum Harmonic Oscillator

Consider a potential

$$V(x) = \frac{1}{2}m\omega^2 x^2$$

Therefore,

$$\hat{H} = \frac{\hat{p}^2}{2m} + \frac{1}{2}m\omega^2 \hat{x}^2$$

and

$$\hat{H}\psi(x) = E\psi(x)$$

Therefore,

$$\begin{aligned}
 (m\omega\hat{x} + i\hat{p})(m\omega\hat{x} - i\hat{p}) &= (m\omega\hat{x})^2 - im\omega\hat{x}\hat{p} + im\omega\hat{p}\hat{x} + \hat{p}^2 \\
 &= (m\omega\hat{x})^2 + im\omega[\hat{p}, \hat{x}] + \hat{p}^2 \\
 &= (m\omega\hat{x})^2 + im\omega(-i\hbar) + \hat{p}^2 \\
 &= (m\omega\hat{x})^2 + m\omega\hbar + \hat{p}^2 \\
 \therefore \hat{p}^2 + (m\omega\hat{x})^2 &= (m\omega\hat{x} + i\hat{p})(m\omega\hat{x} - i\hat{p}) - m\omega\hbar
 \end{aligned}$$

$[\hat{p}, \hat{x}]$ is the
commutation of \hat{p} and
 \hat{x} .

Therefore,

$$\begin{aligned}
\hat{H} &= \frac{1}{2m} (\hat{p} + m^2 \omega^2 \hat{x}^2) \\
&= \frac{1}{2m} (\hat{p} + (m\omega \hat{x})^2) \\
&= \frac{1}{2m} ((m\omega \hat{x} + i\hat{p})(m\omega \hat{x} - i\hat{p}) - m\omega \hbar) \\
&= \hbar\omega \left(\frac{1}{\sqrt{2m\hbar\omega}} (m\omega \hat{x} + i\hat{p}) \frac{1}{\sqrt{2m\hbar\omega}} (m\omega \hat{x} - i\hat{p}) - \frac{1}{2} \right)
\end{aligned}$$

Definition 23 (Lowering and Raising Operators). The raising operator \hat{a}_+ and the lowering operator \hat{a}_- are defined as

$$\begin{aligned}
\hat{a}_+ &= \frac{1}{\sqrt{2m\hbar\omega}} (m\omega \hat{x} - i\hat{p}) \\
\hat{a}_- &= \frac{1}{\sqrt{2m\hbar\omega}} (m\omega \hat{x} + i\hat{p})
\end{aligned}$$

Therefore,

$$\hat{H} = \hbar\omega \left(\hat{a}_- \hat{a}_+ - \frac{1}{2} \right)$$

The commutation of \hat{a}_- and \hat{a}_+ is

$$\begin{aligned}
[\hat{a}_-, \hat{a}_+] &= \frac{1}{2m\hbar\omega} [m\omega \hat{x} + i\hat{p}, m\omega \hat{x} - i\hat{p}] \\
&= \frac{1}{2m\hbar\omega} ((m\omega \hat{x} + i\hat{p})(m\omega \hat{x} - i\hat{p}) - (m\omega \hat{x} - i\hat{p})(m\omega \hat{x} + i\hat{p})) \\
&= \frac{1}{2m\hbar\omega} ((m\omega \hat{x})^2 + \hat{p}^2 + m\omega \hbar - (m\omega \hat{x})^2 - \hat{p}^2 - (im\omega \hat{x}\hat{p} - im\omega \hat{p}\hat{x})) \\
&= \frac{1}{2m\hbar\omega} ((m\omega \hat{x})^2 + \hat{p}^2 + m\omega \hbar - (m\omega \hat{x})^2 - \hat{p}^2 - im\omega [\hat{x}, \hat{p}]) \\
&= \frac{1}{2m\hbar\omega} ((m\omega \hat{x})^2 + \hat{p}^2 + m\omega \hbar - (m\omega \hat{x})^2 - \hat{p}^2 - im\omega(i\hbar)) \\
&= \frac{1}{2m\hbar\omega} ((m\omega \hat{x})^2 + \hat{p}^2 + m\omega \hbar - (m\omega \hat{x})^2 - \hat{p}^2 + m\omega \hbar) \\
&= \frac{1}{2m\hbar\omega} (m\omega \hbar + m\omega \hbar) \\
&= 1
\end{aligned}$$

Therefore,

$$\begin{aligned}
[\hat{a}_-, \hat{a}_+] &= 1 \\
[\hat{a}_+, \hat{a}_-] &= -1
\end{aligned}$$

The commutation of \hat{H} and \hat{a}_+ is

$$\begin{aligned}
[\hat{H}, \hat{a}_+] &= \left[\hbar\omega \left(\hat{a}_- \hat{a}_+ - \frac{1}{2} \right), \hat{a}_+ \right] \\
&= \hbar\omega \left(\hat{a}_- \hat{a}_+ - \frac{1}{2} \right) \hat{a}_+ - \hbar\omega \hat{a}_+ \left(\hat{a}_- \hat{a}_+ - \frac{1}{2} \right) \\
&= \hbar\omega \left(\hat{a}_- \hat{a}_+ \hat{a}_+ - \frac{1}{2} \hat{a}_+ - \hat{a}_+ \hat{a}_- \hat{a}_+ + \frac{1}{2} \hat{a}_+ \right) \\
&= \hbar\omega [\hat{a}_-, \hat{a}_+] \hat{a}_+ \\
&= \hbar\omega \hat{a}_+
\end{aligned}$$

Similarly,

$$[\hat{H}, \hat{a}_-] = -\hbar\omega \hat{a}_-$$

27.1 Effect of \hat{a}_\pm on $\psi(x)$

Let E be an eigenvalue of \hat{H} corresponding to the eigenfunction $\psi(x)$. Therefore,

$$\hat{H}\psi(x) = E\psi(x)$$

Therefore,

$$\begin{aligned}
H\hat{a}_+\psi(x) &= (\hbar\omega \hat{a}_+ + \hat{a}_+ \hat{H}) \psi(x) \\
&= \hbar\omega \hat{a}_+ \psi(x) + E \hat{a}_+ \psi(x) \\
&= (E + \hbar\omega) \hat{a}_+ \psi(x)
\end{aligned}$$

Therefore, $\hat{a}_+\psi(x)$ is also an eigenfunction of \hat{H} with eigenvalue $(E + \hbar\omega)$. Similarly, $\hat{a}_-\psi(x)$ is also an eigenfunction of \hat{H} with eigenvalue $(E - \hbar\omega)$.

Hence, as the difference between any two consecutive eigenvalues of H is $\hbar\omega$, all eigenvalues are evenly spaced. Therefore, the energy levels form a ladder as in Figure 10. Hence, the raising operator can be used in order to go up the ladder, and the lowering operator to go down the ladder. As the potential $V(x)$ has a minimum value, the ladder is not infinite at its lower end, but is limited from below.

Let the lowest possible eigenfunction be ψ_0 . As it is not possible to go lower than this state, the eigenfunction obtained after using the lowering operator

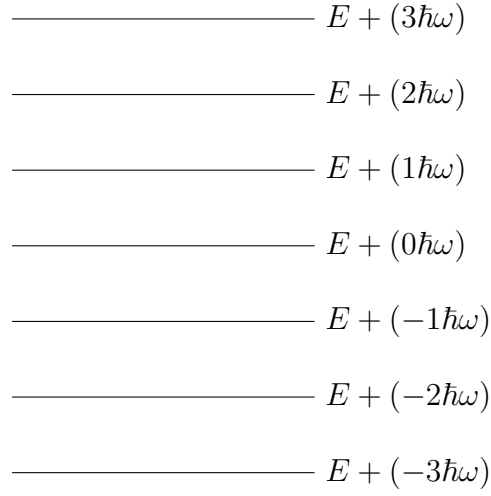


Figure 10: Energy Level Ladder for Quantum Harmonic Oscillator

on ψ_0 must be non-normalizable. Therefore, as $\psi(x) \equiv 0$ is non-normalizable, let

$$\hat{a}_-\psi_0 = 0$$

Therefore,

$$\begin{aligned} 0 &= \frac{1}{\sqrt{2m\hbar\omega}} (m\omega\hat{x} + i\hat{p}) \psi_0 \\ &= \left(m\omega\hat{x} + i(-i\hbar) \frac{\partial}{\partial x} \right) \psi_0 \\ &= \left(m\omega\hat{x} + \hbar \frac{\partial}{\partial x} \right) \psi_0 \end{aligned}$$

Therefore, solving,

$$\psi_0(x) = \left(\frac{m\omega}{\hbar\pi} \right)^{\frac{1}{4}} e^{-\frac{m\omega}{2\hbar}x^2}$$

Therefore, the lowest eigenfunction is a Gaussian function.

All other eigenfunctions can be obtained by applying \hat{a}_+ repeatedly on ψ_0 .

Let the eigenvalue of \hat{H} corresponding to ψ_0 be E_0 . Therefore,

$$\begin{aligned}
 E_0 \psi_0 &= \hat{H} \psi_0 \\
 &= \hbar \omega \left(\hat{a}_+ \hat{a}_- + \frac{1}{2} \right) \psi_0 \\
 &= \hbar \omega \hat{a}_+ \hat{a}_- \psi_0 + \frac{1}{2} \hbar \omega \psi_0 \\
 &= 0 + \frac{1}{2} \hbar \omega \psi_0 \\
 &= \frac{1}{2} \hbar \omega \psi_0
 \end{aligned}$$

Therefore,

$$E_0 = \frac{1}{2} \hbar \omega$$

Therefore,

$$E_n = \hbar \omega \left(\frac{1}{2} + n \right)$$

Let the state obtained after using the raising or lowering operator be

$$\begin{aligned}
 \hat{a}_+ \psi_n &= c_n \psi_{n+1} \\
 \hat{a}_- \psi_n &= d_n \psi_{n-1}
 \end{aligned}$$

Let

$$\begin{aligned}
 f(x) &= \hat{a}_+ \psi_n \\
 g(x) &= \psi_n
 \end{aligned}$$

Therefore, the identity

$$\int_{-\infty}^{\infty} f^*(x) (\hat{a}_{\pm} g(x)) \, dx = \int_{-\infty}^{\infty} (\hat{a}_{\mp} f(x))^* g(x) \, dx$$

implies

$$\begin{aligned}
\int_{-\infty}^{\infty} (\hat{a}_+ \psi_n)^* (\hat{a}_+ \psi_n) dx &= \int_{-\infty}^{\infty} (\hat{a}_- \hat{a}_+ \psi_n)^* \psi_n dx \\
&= \int_{-\infty}^{\infty} \left(\left(\frac{\hat{H}}{\hbar\omega} + \frac{1}{2} \right) \psi_n \right)^* \psi_n dx \\
&= \int_{-\infty}^{\infty} \left(\left(n + \frac{1}{2} + \frac{1}{2} \right) \psi_n \right)^* \psi_n dx \\
&= \int_{-\infty}^{\infty} ((n+1)\psi_n)^* \psi_n dx \\
&= (n+1) \int_{-\infty}^{\infty} \psi_n^* \psi_n dx \\
&= (n+1)(1)
\end{aligned}$$

As ψ_n is normalized,
 $\int_{-\infty}^{\infty} \psi_n^* \psi_n dx = 1$

$$\begin{aligned}
\therefore \int_{-\infty}^{\infty} (c_n \psi_{n+1})^* (c_n \psi_{n+1}) dx &= n+1 \\
\therefore |c_n|^2 \int_{-\infty}^{\infty} \psi_{n+1}^* \psi_{n+1} dx &= n+1
\end{aligned}$$

For ψ_{n+1} to be normalized,

$$\int_{-\infty}^{\infty} \psi_{n+1}^* \psi_{n+1} dx = 1$$

Therefore,

$$|c_n|^2 = n+1$$

Therefore,

$$c_n = \sqrt{n+1}$$

Similarly,

$$d_n = \sqrt{n}$$

Therefore,

$$\begin{aligned}\hat{a}_+\psi_n &= \sqrt{n+1}\psi_{n+1} \\ \hat{a}_-\psi_n &= \sqrt{n}\psi_{n-1}\end{aligned}$$

Therefore,

$$\begin{aligned}\psi_{n+1} &= \frac{1}{\sqrt{n+1}}\hat{a}_+\psi_n \\ &= \frac{1}{\sqrt{n+1}}\hat{a}_+\left(\frac{1}{\sqrt{n}}\hat{a}_+\psi_{n-1}\right) \\ &= \frac{1}{\sqrt{n+1}}\frac{1}{\sqrt{n}}\hat{a}_+\hat{a}_+\psi_{n-1} \\ &= \frac{1}{\sqrt{n}}\frac{1}{\sqrt{n+1}}(\hat{a}_+)^2\psi_{n-1} \\ &\vdots \\ &= \frac{1}{\sqrt{n!}}(\hat{a}_+)^n\psi_0\end{aligned}$$

Therefore, solving,

$$\psi_n(x) = \left(\frac{m\omega}{h\pi}\right)^{\frac{1}{4}} \frac{1}{\sqrt{2^n n!}} e^{-\frac{m\omega x^2}{2\hbar}} H_n\left(\sqrt{\frac{m\omega}{\hbar}}x\right)$$

where H_i are the Hermite polynomials, which are given by

$$H_n(x) = (-1)^n e^{x^2} \frac{d^n}{dx^n} e^{-x^2}$$

Therefore,

$$\begin{aligned}H_0 &= 1 \\ H_1 &= 2x \\ H_2 &= 4x^2 - 2 \\ H_3 &= 8x^3 - 12x \\ H_4 &= 16x^4 - 48x^2 + 12 \\ &\vdots\end{aligned}$$

27.2 Potential Energy

Let the potential energy for the harmonic oscillator be V .

Therefore, the expectation value of V for the n th state of the harmonic

oscillator is

$$\begin{aligned}\langle V \rangle &= \left\langle \frac{1}{2} m \omega^2 \hat{x}^2 \right\rangle \\ &= \frac{1}{2} m \omega^2 \int_{-\infty}^{\infty} \psi_n^* \hat{x}^2 \psi_n \, dx\end{aligned}$$

Also,

$$\hat{x} = \sqrt{\frac{\hbar}{2m\omega}} (\hat{a}_+ + \hat{a}_-)$$

Therefore, substituting and solving,

$$\begin{aligned}\langle V \rangle &= \frac{\hbar\omega}{2} \left(n + \frac{1}{2} \right) \\ &= \frac{E_n}{2}\end{aligned}$$

Exercise 16.

Consider a quantum harmonic oscillator with

$$V(x) = \frac{1}{2} m \omega^2 x^2$$

Assume

$$\psi(x, 0) = \frac{1}{\sqrt{2}} (\psi_0(x) + i\psi_1(x))$$

1. At $t = 0$, what is the probability of measuring the energy to be $\frac{3}{2}\hbar\omega$?
2. At $t = 0$, what is the probability of measuring the energy to be $\frac{11}{2}\hbar\omega$?
3. At $t = 0$, the energy is measured and found to be the maximum possible value. What is the state of the system after measurement?

Solution 16.

1.

$$E_n = \hbar\omega \left(n + \frac{1}{2} \right)$$

Therefore, if

$$E = \frac{3}{2}\hbar\omega$$

then $n = 1$.

Therefore, the probability of measuring the given energy is the probability of n being 1.

Therefore,

$$\begin{aligned} P\left(E = \frac{3}{2}\hbar\omega\right) &= |c_2|^2 \\ &= \left|\frac{i}{\sqrt{2}}\right|^2 \\ &= \frac{1}{2} \end{aligned}$$

2.

$$E_n = \hbar\omega\left(n + \frac{1}{2}\right)$$

Therefore, if

$$E = \frac{11}{2}\hbar\omega$$

then $n = 5$.

As the wave function at $t = 0$ is a combination of ψ_0 and ψ_1 only, the probability of n being 5, and hence the probability of measuring the given energy is 0.

3. As the energy is measured to be the maximum possible, i.e. E_1 , the wave function collapses to $\psi_1(x)$.

Exercise 17.

Consider a quantum particle of mass m in the potential

$$V(x) = \begin{cases} \infty & ; \quad x < 0 \\ \frac{1}{2}m\omega^2x^2 & ; \quad x > 0 \end{cases}$$

Find the allowed energies of the particles.

Solution 17.

If $x < 0$, $V(x) = \infty$.

Therefore,

$$\psi(x) = 0$$

If $x > 0$, the wave function is that for a quantum harmonic oscillator.

Therefore,

$$\psi(x) = \left(\frac{m\omega}{\hbar\pi}\right)^{\frac{1}{4}} \frac{1}{\sqrt{2^n n!}} e^{-\frac{m\omega x^2}{2\hbar}} H_n\left(\sqrt{\frac{m\omega}{\hbar}} x\right)$$

Therefore,

$$\psi(x) = \begin{cases} 0 & : \quad x < 0 \\ \left(\frac{m\omega}{\hbar\pi}\right)^{\frac{1}{4}} \frac{1}{\sqrt{2^n n!}} e^{-\frac{m\omega x^2}{2\hbar}} H_n\left(\sqrt{\frac{m\omega}{\hbar}} x\right) & ; \quad x > 0 \end{cases}$$

For ψ to be continuous,

$$\psi(0) = 0$$

Therefore, solving, the boundary conditions are satisfied by odd n only.

Therefore,

$$E_n = \hbar\omega \left(n + \frac{1}{2}\right)$$

Therefore, the spacing between the energy levels is $2\hbar\omega$

28 3D Harmonic Oscillator

The Hamiltonian operator in 3D is

$$\hat{H} = -\frac{\hbar^2}{2m} \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) + \frac{1}{2} m\omega^2 (x^2 + y^2 + z^2)$$

Let E be the eigenvalue of \hat{H} corresponding to $\psi(x, y, z)$.

Therefore,

$$\hat{H}\psi(x, y, z) = E\psi(x, y, z)$$

1D Schrödinger Equation	3D Schrödinger Equation
$\psi(x, t)$	$\psi(x, y, z, t)$
$\int_{-\infty}^{\infty} \psi ^2 dx = 1$	$\iiint_{-\infty}^{\infty} \psi ^2 dx dy dz = 1$
$\hat{H} = \frac{\hat{p}_x^2}{2m} + V(x)$	$\hat{H} = \frac{\hat{p}_x^2 + \hat{p}_y^2 + \hat{p}_z^2}{2m} + V(x, y, z)$
$\hat{H} = -\frac{i\hbar}{m} \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) + V(x, y, z)$	$\hat{H} = -\frac{i\hbar}{m} \left(\frac{\partial^2}{\partial x^2} \right) + V(x, y, z)$
$\Psi(x, t) = \psi(x)\varphi(t)$	$\Psi(x, y, z, t) = \psi(x, y, z)\varphi(t)$
$\varphi(t) = e^{-\frac{iEt}{\hbar}}$	$\varphi(t) = e^{-\frac{iEt}{\hbar}}$
$\Psi(x, t) = \sum c_n \psi_n(x) e^{-\frac{iEt}{\hbar}}$	$\Psi(x, y, z, t) = \sum c_n \psi_n(x, y, z) e^{-\frac{iEt}{\hbar}}$

Table 1: Comparison of 1D and 3D Schrödinger Equations

Separating the variables, let

$$\psi(x, y, z) = \psi_x(x)\psi_y(y)\psi_z(z)$$

Therefore,

$$\begin{aligned} \left(-\frac{\hbar}{2m} \frac{\partial^2}{\partial x^2} + \frac{1}{2}m\omega^2 x^2 \right) \psi_x(x) &= E_n \psi_x(x) \\ \left(-\frac{\hbar}{2m} \frac{\partial^2}{\partial y^2} + \frac{1}{2}m\omega^2 y^2 \right) \psi_y(y) &= E_m \psi_y(y) \\ \left(-\frac{\hbar}{2m} \frac{\partial^2}{\partial z^2} + \frac{1}{2}m\omega^2 z^2 \right) \psi_z(z) &= E_l \psi_z(z) \end{aligned}$$

Therefore,

$$\begin{aligned} E &= E_n + E_m + E_l \\ &= \frac{3}{2}\hbar\omega + \hbar\omega(n + m + l) \end{aligned}$$

Exercise 18.

Consider a single particle in a 3D infinite well of side a , which is equivalent to a single particle in a box of side a , cornered at $(0, 0, 0)$. The potential is given by

$$V(x, y, z) = \begin{cases} 0 & ; \text{ inside the box} \\ \infty & ; \text{ outside the box} \end{cases}$$

Find the wave function of the particle.

Solution 18.

As the potential outside the box is infinite, for any point outside the box,

$$\psi(x) = 0$$

For a point inside the box,

$$\hat{H} = -\frac{\hbar^2}{2m} \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right)$$

Let E be the eigenvalue of \hat{H} corresponding to $\psi(x, y, z)$. Therefore,

$$\hat{H}\psi(x, y, z) = E\psi(x, y, z)$$

Therefore, separating the variables,

$$\psi(x, y, z) = \psi_x(x)\psi_y(y)\psi_z(z)$$

Therefore,

$$\begin{aligned} \hat{H}\psi(x, y, z) &= E\psi(x, y, z) \\ \therefore \hat{H}(\psi_x(x)\psi_y(y)\psi_z(z)) &= E\psi_x(x)\psi_y(y)\psi_z(z) \\ \therefore -\frac{\hbar^2}{2m}(\psi_x''\psi_y\psi_z + \psi_x\psi_y''\psi_z + \psi_x\psi_y\psi_z'') &= E\psi_x\psi_y\psi_z \\ \therefore \frac{\psi_x''}{\psi_x} + \frac{\psi_y''}{\psi_y} + \frac{\psi_z''}{\psi_z} &= -\frac{2mE}{\hbar^2} \end{aligned}$$

Let

$$\begin{aligned} k_x &= \sqrt{-\frac{\psi_x''}{\psi_x}} \\ k_y &= \sqrt{-\frac{\psi_y''}{\psi_y}} \\ k_z &= \sqrt{-\frac{\psi_z''}{\psi_z}} \\ k &= \sqrt{\frac{2mE}{\hbar^2}} \end{aligned}$$

Therefore,

$$k_x^2 + k_y^2 + k_z^2 = k^2$$

Therefore,

$$\psi_x(x) = A \sin(k_x x) + B \cos(k_x x)$$

$$\psi_y(y) = C \sin(k_y y) + D \cos(k_y y)$$

$$\psi_z(z) = F \sin(k_z z) + G \cos(k_z z)$$

Therefore, as ψ is continuous at $(0, 0, 0)$, solving,

$$B = 0$$

$$D = 0$$

$$G = 0$$

$$k_x a = n_x \pi$$

$$k_y a = n_y \pi$$

$$k_z a = n_z \pi$$

Therefore,

$$\psi_x(x) = A \sin(k_x x)$$

$$\psi_y(y) = C \sin(k_y y)$$

$$\psi_z(z) = F \sin(k_z z)$$

Therefore,

$$\psi(x, y, z) = N \sin\left(\frac{n_x \pi}{a} x\right) \sin\left(\frac{n_y \pi}{a} y\right) \sin\left(\frac{n_z \pi}{a} z\right)$$

Therefore, normalizing,

$$\int_{-\infty}^{\infty} |\psi|^2 dx dy dz = 1$$

Therefore, solving,

$$N = \left(\sqrt{\frac{2}{a}}\right)^3$$

Therefore,

$$\begin{aligned}
E &= \frac{k^2 \hbar^2}{2m} \\
&= \frac{(k_x^2 + k_y^2 + k_z^2) \hbar^2}{2m} \\
&= \frac{\left(\left(\frac{n_x \pi}{a} \right)^2 + \left(\frac{n_y \pi}{a} \right)^2 + \left(\frac{n_z \pi}{a} \right)^2 \right) \hbar^2}{2m}
\end{aligned}$$

Therefore, there is degeneracy in the energies, i.e., for different combinations of n_x , n_y , n_z , there exist different wave functions with same energy.

29 Angular Momentum

Definition 24 (Angular momentum operators). The angular momentum operators in the x , y , z directions are defined as

$$\begin{aligned}
\hat{L}_x &= \hat{y}\hat{p}_z - \hat{z}\hat{p}_y \\
\hat{L}_y &= \hat{z}\hat{p}_x - \hat{x}\hat{p}_z \\
\hat{L}_z &= \hat{x}\hat{p}_y - \hat{y}\hat{p}_x
\end{aligned}$$

where \hat{x} , \hat{y} , \hat{z} are the position operators in the x , y , z directions respectively, and \hat{p}_x , \hat{p}_y , \hat{p}_z are the momentum operators in the x , y , z directions respectively.

30 Commutation Relations for 3D Operators

Theorem 33.

$$\begin{aligned}
[\hat{i}, \hat{j}] &= i\hbar\delta_{ij} \\
[\hat{p}_i, \hat{p}_j] &= 0 \\
[\hat{i}, \hat{p}_i] &= 0
\end{aligned}$$

where i, j can be x, y, z .

30.1 Angular Momentum Operators in Orthogonal Directions

$$\begin{aligned}
[\hat{L}_x, \hat{L}_y] &= (\hat{y}\hat{p}_z - \hat{z}\hat{p}_y)(\hat{z}\hat{p}_x - \hat{x}\hat{p}_x) - (\hat{z}\hat{p}_x - \hat{x}\hat{p}_z)(\hat{y}\hat{p}_z - \hat{z}\hat{p}_y) \\
&= \underbrace{\hat{y}\hat{p}_z\hat{z}\hat{p}_x}_1 - \underbrace{\hat{y}\hat{p}_z\hat{x}\hat{p}_z}_2 - \underbrace{\hat{z}\hat{p}_y\hat{z}\hat{p}_x}_3 + \underbrace{\hat{z}\hat{p}_y\hat{x}\hat{p}_z}_4 \\
&\quad - \underbrace{\hat{z}\hat{p}_x\hat{y}\hat{p}_z}_5 + \underbrace{\hat{z}\hat{p}_x\hat{z}\hat{p}_y}_6 + \underbrace{\hat{x}\hat{p}_z\hat{y}\hat{p}_z}_7 - \underbrace{\hat{x}\hat{p}_z\hat{z}\hat{p}_y}_8
\end{aligned}$$

Therefore, term 2 is

$$\begin{aligned}
[\hat{x}, \hat{p}_z] &= 0 & \hat{y}\hat{p}_z\hat{x}\hat{p}_z &= \hat{y}\hat{x}\hat{p}_z\hat{p}_z \\
[\hat{y}, \hat{x}] &= 0 & &= xyp_zp_x \\
[\hat{y}, \hat{p}_z] &= 0 & &= xp_zyp_z
\end{aligned}$$

Therefore, term 2 is equal to term 7.
similarly, solving,

$$\begin{aligned}
[\hat{L}_x, \hat{L}_y] &= i\hbar\hat{L}_z \\
[\hat{L}_y, \hat{L}_z] &= i\hbar\hat{L}_x \\
[\hat{L}_z, \hat{L}_x] &= i\hbar\hat{L}_y
\end{aligned}$$

30.2 Angular Momentum Operators and Square of Angular Momentum Operator

$$\begin{aligned}
[\hat{L}^2, \hat{L}_x] &= [\hat{L}_x^2 + \hat{L}_y^2 + \hat{L}_z^2, \hat{L}_x] \\
&= \hat{L}_x^3 + \hat{L}_y\hat{L}_y\hat{L}_x + \hat{L}_z\hat{L}_z\hat{L}_x - \hat{L}_x^3 - \hat{L}_x\hat{L}_y\hat{L}_y - \hat{L}_x\hat{L}_z\hat{L}_z
\end{aligned}$$

Therefore,

$$\begin{aligned}
\hat{L}_y\hat{L}_x &= [\hat{L}_y, \hat{L}_x] + \hat{L}_x\hat{L}_y = -i\hbar\hat{L}_z + \hat{L}_x\hat{L}_y \\
\hat{L}_y\hat{L}_y\hat{L}_x &= -i\hbar\hat{L}_y\hat{L}_z + \hat{L}_y\hat{L}_x\hat{L}_y \\
&= -i\hbar\hat{L}_y\hat{L}_z - i\hbar\hat{L}_z\hat{L}_y + \hat{L}_x\hat{L}_y\hat{L}_y
\end{aligned}$$

Similarly,

$$\hat{L}_z\hat{L}_z\hat{L}_x = i\hbar\hat{L}_z\hat{L}_y + i\hbar\hat{L}_y\hat{L}_z + \hat{L}_x\hat{L}_z\hat{L}_z$$

Therefore, substituting and solving,

$$[\hat{L}^2, \hat{L}_x] = 0$$

Similarly,

$$[\hat{L}^2, \hat{L}_y] = 0$$

$$[\hat{L}^2, \hat{L}_z] = 0$$

Therefore, as the commutation relations are zeros, the operators have common eigenfunctions.

Let Y_{lm} be an eigenfunction of the momentum operator \hat{L}_z .

Therefore,

$$\hat{L}_z Y_{lm} = m\hbar Y_{lm}$$

$$\hat{L}^2 Y_{lm} = l(l+1)\hbar^2 Y_{lm}$$

where m and l are real constants.

Let

$$\hat{L}_+ = \hat{L}_x + i\hat{L}_y$$

$$\hat{L}_- = \hat{L}_x - i\hat{L}_y$$

be two operators.

Therefore, the commutation relations on \hat{L}_+ and \hat{L}_- are

$$[\hat{L}_z, \hat{L}_+] = \hbar \hat{L}_+$$

$$[\hat{L}_z, \hat{L}_-] = -\hbar \hat{L}_-$$

$$[\hat{L}^2, \hat{L}_+] = 0$$

$$[\hat{L}^2, \hat{L}_-] = 0$$

30.3 Effect of \hat{L}_\pm on Y_{lm}

$$\begin{aligned} \hat{L}_z \hat{L}_+ Y_{lm} &= (\hbar \hat{L}_+ + \hat{L}_+ \hat{L}_z) Y_{lm} \\ &= \hbar(m+1) \hat{L}_+ Y_{lm} \end{aligned}$$

$$\begin{aligned} \hat{L}_z \hat{L}_+ &= [\hat{L}_z, \hat{L}_+] + \\ \hat{L}_+ \hat{L}_z &= \hbar \hat{L}_+ + \hat{L}_+ \hat{L}_z \end{aligned}$$

Therefore, $\hat{L}_+ Y_{lm}$ is also an eigenfunction of \hat{L}_z , with eigenvalue $\hbar(m+1)$.

Similarly,

$$\begin{aligned} \hat{L}^2 \hat{L}_+ Y_{lm} &= \hat{L}_+ \hat{L}^2 Y_{lm} \\ &= \hbar^2 l(l+1) \hat{L}_+ Y_{lm} \end{aligned}$$

Therefore, $\hat{L}_+ Y_{lm}$ is also an eigenfunction of \hat{L}^2 , but with the same eigenvalue, $\hbar^2 l(l+1)$.

30.4 Ladder of \hat{L}_z

The eigenvalues of \hat{L}_z are separated by \hbar . Therefore, the eigenvalues form a ladder similar to that of the quantum harmonic oscillator, in Figure 10.

As L_z is a component of the angular momentum, its magnitude must be less than or equal to that of the total angular momentum. Hence, this ladder is bounded both at the top and the bottom. Therefore, for the maximum value of m ,

$$\hat{L}_+ Y_{lm} = 0$$

Therefore,

$$\begin{aligned} 0 &= \hat{L}_- \hat{L}_+ Y_{lm} \\ &= (\hat{L}_x - i\hat{L}_y) (\hat{L}_x + i\hat{L}_y) Y_{lm} \\ &= (\hat{L}_x^2 + i\hat{L}_x\hat{L}_y - i\hat{L}_y\hat{L}_x + \hat{L}_y^2) Y_{lm} \\ &= (\hat{L}_x^2 + \hat{L}_y^2 + i[\hat{L}_x, \hat{L}_y]) Y_{lm} \\ &= (\hat{L}_x^2 + \hat{L}_y^2 - \hbar\hat{L}_z) Y_{lm} \\ &= (\hat{L}^2 - \hat{L}_z^2 - \hbar\hat{L}_z) Y_{lm} \\ &= (\hbar^2 l(l+1) - \hbar^2 m^2 - \hbar^2 m) Y_{lm} \\ &= (\hbar^2 l(l+1) - \hbar^2 m(m+1)) Y_{lm} \end{aligned}$$

Therefore,

$$\begin{aligned} \hbar l(l+1) - \hbar m(m+1) &= 0 \\ \therefore m_{\max} &= l \end{aligned}$$

Therefore, each l_i has a ladder, with first step $-l_i\hbar$, last step $l_i\hbar$, and step \hbar corresponding to it.

Let the number of steps in each ladder be N . Therefore,

$$N = 2l + 1$$

Therefore, l can be expressed as

$$l = \frac{N-1}{2}$$

where $N \in \mathbb{N}$.

Therefore, the ladder corresponding to $l = 0$ has one step, with $m = 0$, the ladder corresponding to $l = \frac{1}{2}$ has two steps, with $m = -\frac{1}{2}$ and $m = \frac{1}{2}$, and so on.

31 Angular Momentum in Spherical Coordinates

In spherical coordinates,

$$x = r \sin \theta \cos \varphi$$

$$y = r \sin \theta \sin \varphi$$

$$z = r \cos \theta$$

where $\theta \in [0, \pi]$, $\varphi \in [0, 2\pi)$, $r > 0$.

Therefore,

$$r = \sqrt{x^2 + y^2 + z^2}$$

$$\theta = \cos^{-1} \left(\frac{z}{\sqrt{x^2 + y^2 + z^2}} \right)$$

$$\varphi = \tan^{-1} \left(\frac{y}{x} \right)$$

Therefore,

$$\hat{L}_x = \frac{\hbar}{i} \left(-\sin \varphi \frac{\partial}{\partial \theta} - \frac{\cos \varphi}{\tan \theta} \frac{\partial}{\partial \varphi} \right)$$

$$\hat{L}_y = \frac{\hbar}{i} \left(\cos \varphi \frac{\partial}{\partial \theta} - \frac{\sin \varphi}{\tan \theta} \frac{\partial}{\partial \varphi} \right)$$

$$\hat{L}_z = \frac{\hbar}{i} \frac{\partial}{\partial \varphi}$$

$$\hat{L}^2 = -\hbar^2 \left(\frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \varphi^2} \right)$$

$$\hat{L}_{\pm} = \pm \hbar e^{\pm i\varphi} \left(\frac{\partial}{\partial \theta} \pm i \cot \theta \frac{\partial}{\partial \varphi} \right)$$

Normalizing in spherical coordinates,

$$\iiint |Y_{lm}|^2 r^2 \sin \theta \, dr \, d\theta \, d\varphi = 1$$

As Y_{lm} is independent of r ,

$$\iint |Y_{lm}|^2 \sin \theta \, d\theta \, d\varphi = 1$$

Therefore, as Y_{lm} is the eigenfunction of \hat{L}_z ,

$$\hat{L}_z Y_{lm}(\theta, \varphi) = m\hbar Y_{lm}(\theta, \varphi)$$

Therefore,

$$\begin{aligned} \frac{\hbar}{i} \frac{\partial}{\partial \varphi} Y_{lm}(\theta, \varphi) &= m\hbar Y_{lm}(\theta, \varphi) \\ \therefore \frac{\partial}{\partial \varphi} Y_{lm}(\theta, \varphi) &= im Y_{lm}(\theta, \varphi) \end{aligned}$$

Therefore, solving,

$$\begin{aligned} Y_{lm}(\theta, \varphi) &= N_{lm} e^{im\varphi} f_{lm}(\theta) \\ \hat{L}^2 Y_{lm}(\theta, \varphi) &= \hbar l(l+1) Y_{lm}(\theta, \varphi) \end{aligned}$$

where N_{lm} is the constant of normalization.

For Y_{ll} , the eigenfunction corresponding to the highest step in the ladder,

$$\begin{aligned} 0 &= \hat{L}_+ Y_{ll} \\ &= \hat{L}_+ (N_{ll} e^{il\varphi} f_{ll}(\theta)) \\ &= \hbar e^{i\varphi} \left(N_{ll} e^{il\varphi} \frac{\partial f_{ll}(\theta)}{\partial \theta} + i \cot \theta N_{ll} e^{il\varphi} f_{ll}(\theta) \right) \\ &= \hbar e^{i(l+1)\varphi} N_{ll} \left(\frac{\partial f_{ll}(\theta)}{\partial \theta} - l \cot \theta f_{ll}(\theta) \right) \\ &= \frac{\partial f_{ll}(\theta)}{\partial \theta} - l \cot \theta f_{ll}(\theta) \end{aligned}$$

Therefore, solving,

$$\begin{aligned} \therefore f_{ll}(\theta) &\approx (\sin \theta)^l \\ \therefore Y_{ll}(\theta, \varphi) &= N_{ll} e^{il\varphi} (\sin \theta)^l \end{aligned}$$

Therefore, applying \hat{L}_- on Y_{ll} ,

$$Y_{lm}(\theta, \varphi) = \varepsilon \sqrt{\frac{(2l+1)(l-|m|)!}{4\pi(l+|m|)!}} e^{im\varphi} f_{lm}(\cos \theta)$$

where

$$\varepsilon = \begin{cases} (-1)^m & ; \quad m \geq 0 \\ 1 & ; \quad m \leq 0 \end{cases}$$

and f_{lm} is the associated Legendre function.

32 Hamiltonian Operator in 3D

In 3D,

$$\begin{aligned}\hat{H} &= -\frac{\hbar^2}{2m} \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) + V(x, y, z) \\ &= -\frac{\hbar^2}{2m} \nabla^2 + V(x, y, z)\end{aligned}$$

where m is the mass of the particle, and

$$\begin{aligned}\nabla^2 f &= \frac{1}{r} \frac{\partial^2}{\partial r^2} (rf) + \frac{1}{r^2} \left(\frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial f}{\partial \theta} \right) + \frac{1}{\sin^2 \theta} \frac{\partial^2 f}{\partial \varphi^2} \right) \\ &= \frac{1}{r} \frac{\partial^2}{\partial r^2} (rf) - \frac{\hat{L}^2}{r^2 \hbar^2}\end{aligned}$$

33 Central Potentials

Definition 25. A potential V , such that

$$V(r, \theta, \varphi) = V(r)$$

is called a central potential.

Therefore, for a central potential, the Hamiltonian operator is

$$\hat{H} = -\frac{\hbar^2}{2m} \nabla^2 + V(r)$$

where m is the mass of the particle.

Therefore,

$$\begin{aligned}[\hat{H}, \hat{L}_z] &= 0 \\ [\hat{H}, \hat{L}^2] &= 0\end{aligned}$$

As Y_{lm} is a basis, any function $f(\theta, \varphi)$ can be expressed as

$$f(\theta, \varphi) = \sum_{l=0}^{\infty} \sum_{m=-l}^l c_{lm} Y_{lm}(\theta, \varphi)$$

As for quantum harmonic oscillators,

$$\hat{L}_+ Y_{lm} = A_{lm} Y_{l(m+1)}$$

Similarly,

$$\hat{L}_- Y_{lm} = B_{lm} Y_{l(m+1)}$$

Let

$$\begin{aligned} f(\theta, \varphi) &= \hat{L}_+ Y_{lm} \\ g(\theta, \varphi) &= Y_{lm} \end{aligned}$$

Therefore, the identity

$$\iint f^*(\theta, \varphi) (\hat{L}_\pm g(\theta, \varphi)) \sin \theta \, d\theta \, d\varphi = \iint (L_\mp f(\theta, \varphi))^* g(\theta, \varphi) \sin \theta \, d\theta \, d\varphi$$

implies

$$\begin{aligned} \iint (\hat{L}_+ Y_{lm})^* (\hat{L}_+ Y_{lm}) \sin \theta \, d\theta \, d\varphi &= \iint |A_{lm}|^2 Y_{l(m+1)}^* Y_{l(m+1)} \sin \theta \, d\theta \, d\varphi \\ &= |A_{lm}|^2 \end{aligned}$$

Therefore,

$$\begin{aligned} \iint (\hat{L}_- \hat{L}_+ Y_{lm})^* Y_{lm} \sin \theta \, d\theta \, d\varphi &= \iint \left((\hat{L}_x - i\hat{L}_y) (\hat{L}_x + i\hat{L}_y) Y_{lm} \right)^* Y_{lm} \sin \theta \, d\theta \, d\varphi \\ &= \iint \left((\hat{L}^2 - \hat{L}_z^2 - \hbar \hat{L}_z) Y_{lm} \right)^* Y_{lm} \sin \theta \, d\theta \, d\varphi \\ &= \iint (\hbar^2 l(l+1) - \hbar^2 m^2 - \hbar^2 m) Y_{lm}^* Y_{lm} \sin \theta \, d\theta \, d\varphi \\ &= \hbar^2 (l(l+1) - m(m+1)) \cdot 1 \end{aligned}$$

Therefore,

$$|A_{lm}|^2 = \hbar^2 (l(l+1) - m(m+1))$$

Therefore,

$$\begin{aligned} A_{lm} &= \hbar \sqrt{l(l+1) - m(m+1)} \\ B_{lm} &= \hbar \sqrt{l(l+1) - m(m-1)} \end{aligned}$$

34 General Form of the Uncertainty Principle

Theorem 34 (General Form of the Uncertainty Principle).

$$\sigma_A^2 \sigma_B^2 \geq \left(\frac{1}{2i} \langle [\hat{A}, \hat{B}] \rangle \right)^2$$

34.1

$i++i$

$$\begin{aligned}\hat{L}_+ &= \hat{L}_x + i\hat{L}_y \\ \hat{L}_- &= \hat{L}_x - i\hat{L}_y\end{aligned}$$

Therefore,

$$\begin{aligned}\hat{L}_x &= \frac{\hat{L}_+ + \hat{L}_-}{2} \\ \hat{L}_y &= \frac{\hat{L}_+ - \hat{L}_-}{2}\end{aligned}$$

Let Y_{lm} be an eigenfunction of \hat{L}_z and \hat{L}^2 .
Therefore,

$$\begin{aligned}\langle \hat{L}_x \rangle &= \iint Y_{lm}^* \hat{L}_x Y_{lm} \sin \theta \, d\theta \, d\varphi \\ &= \iint Y_{lm}^* \left(\frac{\hat{L}_+ + \hat{L}_-}{2} \right) Y_{lm} \sin \theta \, d\theta \, d\varphi \\ &= \frac{A_{lm}}{2} \iint Y_{lm}^* Y_{l(m+1)} \sin \theta \, d\theta \, d\varphi + \frac{B_{lm}}{2} \iint Y_{lm}^* Y_{l(m-1)} \sin \theta \, d\theta \, d\varphi\end{aligned}$$

Therefore, as Y_{lm}^* , $Y_{l(m+1)}$, and $Y_{l(m-1)}$ are mutually orthogonal, their products are zero. Therefore,

$$\begin{aligned}\langle \hat{L}_x \rangle &= 0 + 0 \\ &= 0\end{aligned}$$

Similarly,

$$\begin{aligned}\langle \hat{L}_x \rangle &= 0 - 0 \\ &= 0\end{aligned}$$

Also,

$$\begin{aligned}L_- L_+ &= L^2 - L_z^2 - \hbar L_z \\ L_+ L_- &= L^2 - L_z^2 + \hbar L_z\end{aligned}$$

Therefore,

$$\begin{aligned}\hat{L}_x &= \frac{\hat{L}_+ + \hat{L}_-}{2} \\ \therefore \hat{L}_x^2 &= \frac{1}{4} (\hat{L}_+^2 + \hat{L}_-^2 + \hat{L}_+ \hat{L}_- + \hat{L}_- \hat{L}_+)\end{aligned}$$

Therefore,

$$\begin{aligned}\hat{L}_- \hat{L}_+ Y_{lm} &= \hbar^2 (l(l+1) - m(m+1)) Y_{lm} \\ \hat{L}_+ \hat{L}_- Y_{lm} &= \hbar^2 (l(l+1) - m(m-1)) Y_{lm}\end{aligned}$$

Therefore,

$$\begin{aligned}\langle \hat{L}_x^2 \rangle &= \frac{1}{4} \iint Y_{lm}^* (\hat{L}_+ \hat{L}_- + \hat{L}_- \hat{L}_+) Y_{lm} \sin \theta \, d\theta \, d\varphi \\ &= \frac{\hbar^2}{4} (2l(l+1) - 2m^2) \\ &= \frac{\hbar^2}{2} (l(l+1) - m^2)\end{aligned}$$

Similarly,

$$\langle \hat{L}_y^2 \rangle = \frac{\hbar^2}{2} (l(l+1) - m^2)$$

Therefore,

$$\begin{aligned}\sigma_{L_x} &= \frac{\hbar}{\sqrt{2}} \sqrt{l(l+1) - m^2} \\ \sigma_{L_y} &= \frac{\hbar}{\sqrt{2}} \sqrt{l(l+1) - m^2}\end{aligned}$$

Therefore,

$$\sigma_{L_x} \sigma_{L_y} = \frac{\hbar^2}{2} (l(l+1) - m^2)$$

Therefore, by the General Form of the Uncertainty Principle,

$$\begin{aligned}\sigma_{L_x} \sigma_{L_y} &\geq \frac{1}{2i} \langle i\hbar L_z \rangle \\ &\geq \frac{1}{2i} i\hbar \cdot \hbar m\end{aligned}$$

Therefore,

$$\langle L_z \rangle = i\hbar m$$

Therefore,

$$\sigma_{L_x} \sigma_{L_y} \geq \frac{\hbar^2}{2} m$$

Part II

Solid State Physics

1 Lecturer Information

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2 Required Reading

1. Streetman, B. Solid State Electronic Devices

3 Additional Reading

1. Kittel, Introduction to solid state physics, John Wiley & Sons.
2. Pierret. Advanced semiconductor Fundamentals, Prentice Hall.
3. Ashcroft, Solid State Physics, Harcourt college publishers.

4 Electrons

Definition 26 (Particle nature of electrons). An electron behaves as a negatively charged charge carrying particle.

The magnitude of the charge on it is

$$q = 1.602 \times 10^{-19} \text{ C}$$

Its mass is

$$m_0 = 9.11 \times 10^{-31} \text{ kg}$$

Definition 27 (Wave nature of electrons). Electrons exhibit wave-like properties, in addition to particle-like properties.

The energy transmitted by a wave is

$$\begin{aligned} E &= h\nu \\ &= \frac{hc}{\lambda} \end{aligned}$$

where

$$h = \text{Planck's constant } (6.626 \times 10^{-34})$$

ν = frequency

c = speed of light

λ = wavelength

5 Semiconductors

Law 1 (Ohm's Law). *The voltage across two points on a conductor is directly proportional to the current through the conductor. The constant of proportionality is called the resistance of the conductor.*

$$\frac{V}{I} = R$$

Law 2 (Microscopic Ohm's Law).

$$\vec{J} = \sigma \vec{E}$$

where \vec{J} is the current density, σ is the conductivity, \vec{E} is the electric field in the resistor.

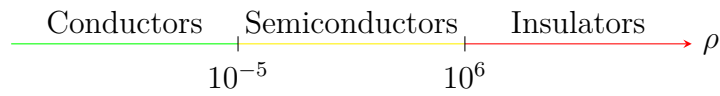
Definition 28 (Resistivity). If

$$R = \rho \frac{L}{A}$$

where R is the resistance of the resistor, L is the length of the resistor, and A is the cross-sectional area of the resistor, then ρ is called the resistivity of the resistor.

$\sigma = \frac{1}{\rho}$ is called the conductivity of the resistor.

They are constant for a particular material.



5.1 Control Factors

The major factors which affect the conductivity of a material are

1. Temperature
2. Chemical composition
 - (a) Atomic bonding
 - (b) Crystal structure
 - (c) Charge carriers in the crystal
3. Optical effects
4. Doping

5.2 Chemical Makeup

II	III	IV	V	VI
	B	C	N	O
	Al	Si	P	S
Zn	Ga	Ge	As	
Cd	In			

Silicon is usually used as a semiconductor, as it

1. is easily available, hence economical
2. performs better at higher temperatures
3. can be converted to silica on heating

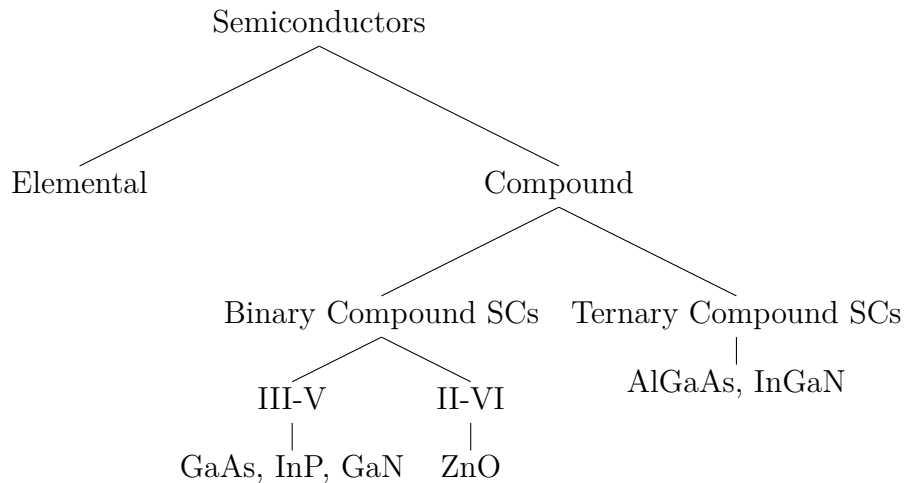


Figure 11: Classification of Semiconductors

Exercise 19.

A sample of Germanium has resistivity $\rho = 0.46\Omega \text{ m}$. The dimensions of the sample are

$$l = 50\mu\text{m}$$

$$h = 0.2\mu\text{m}$$

$$w = 1\mu\text{m}$$

Find the resistance of the sample and the conductivity of the material.

Solution 19.

$$\begin{aligned}\rho &= 0.46\Omega \text{ m} \\ &= 46\Omega \text{ cm}\end{aligned}$$

Therefore,

$$\begin{aligned}\sigma &= \frac{1}{\rho} \\ &= \frac{1}{46\Omega \text{ cm}} \\ &= 0.022 \frac{1}{\Omega \text{ cm}}\end{aligned}$$

Therefore,

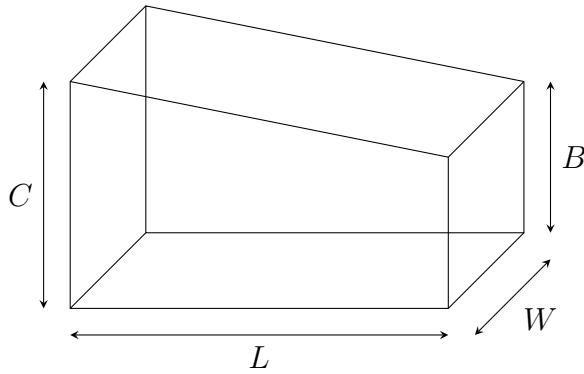
$$\begin{aligned} l &= 50 \mu\text{m} \\ &= 50 \times 10^{-4} \text{cm} \end{aligned}$$

Therefore,

$$\begin{aligned} R &= \rho \frac{l}{A} \\ &= 11500 \times 10^{-4} \Omega \end{aligned}$$

Exercise 20.

A sample of Germanium has resistivity σ . The dimensions of the sample are as shown.



Find the relationship between R and σ .

Solution 20.

Consider a slice with height h , width w , and thickness dx . Therefore, the cross-sectional area of the elemental slice is

$$\begin{aligned} dA &= wh \\ &= w \left(\frac{B-C}{L}x + C \right) \\ &= w \left(\frac{Bx - C(L-x)}{L} \right) \end{aligned}$$

Therefore,

$$\begin{aligned} dR &= \frac{dx}{\sigma wh} \\ &= \frac{L dx}{\sigma w (Bx - C(L-x))} \end{aligned}$$

6 Types of Materials

Atoms tend to arrange themselves in such a way that the resultant energy is minimized.

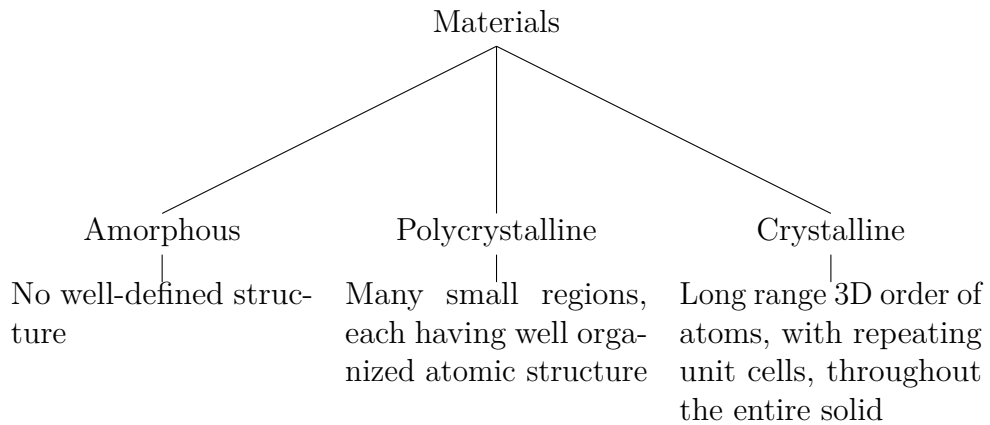


Figure 12: Classification of Materials

Semiconductor devices can use all of these types of materials.

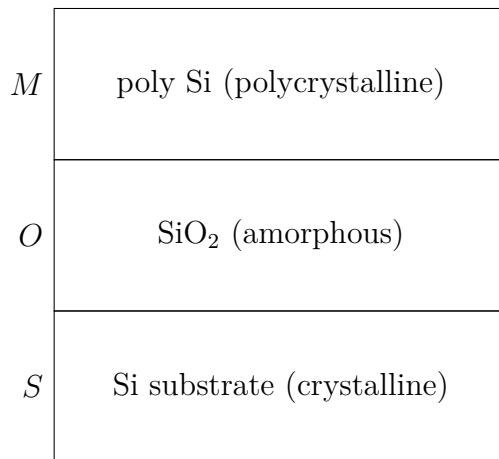


Figure 13: MOS which uses all three types of materials

7 Bohr's Model

According to Bohr's model of the atom, electrons can have discrete energy levels only. The electrons in an atom are arranged in the order of filling

electronic shells, given by the Aufbau Principle.

The energy of a free electron is called E_{vac} . This is used as a reference energy.

8 Atomic Bonding

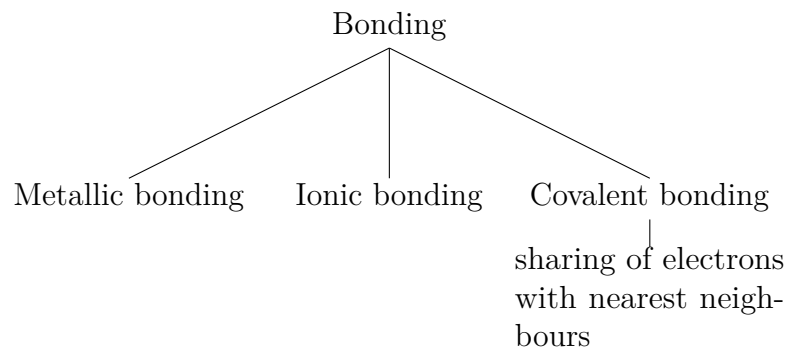


Figure 14: Types of Atomic Bonds

8.1 Covalent Bonds in Silicon

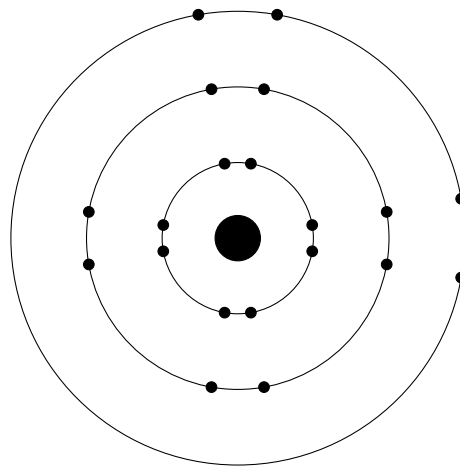


Figure 15: Arrangement of electrons in a silicon atom, in shells according to Bohr's model

In a silicon crystal, the silicon atoms are packed tightly and periodically, i.e. in a repeating pattern. The atomic density for silicon is approximately 10^{23} atoms per cm^3 .

This arrangement of silicon atoms in crystalline form acts as a conductor only

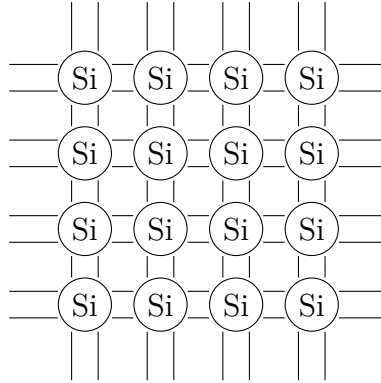


Figure 16: 2D model of bonding in silicon

at certain temperatures. At low temperatures, i.e. around 0 K, the electrons are tightly bound, and hence the crystal behaves as an insulator. However, at higher temperatures, the additional thermal energy may be sufficient to break bonds, and hence produce free electrons.

This property of silicon and other semiconductors makes it possible to control their conductivity.

9 Basics of Crystal Structure

Definition 29 (Crystal lattice). The periodic arrangement of atoms in a crystal is called crystal lattice.

Definition 30 (Lattice constant). The distance between two adjacent atoms in a crystal lattice is called the lattice constant.

It is denoted by a .

It is determined by the attractive and repulsive forces acting on the atoms. The lattice constant corresponds to the lowest energy in the crystal.

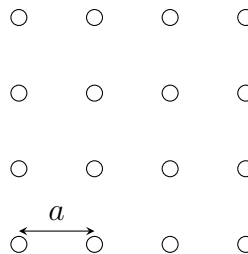
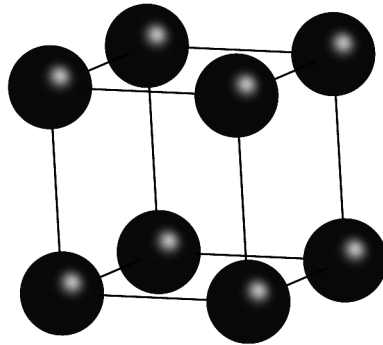


Figure 17: Crystal lattice and lattice constant

Definition 31 (Unit cell). A volume which is repeated throughout the crystal is called the unit cell. It is used to represent the entire crystal lattice.



Definition 32. The atomic packing factor defined to be

$$\text{APF} = \frac{\text{volume filled by atoms}}{\text{total volume of unit cell}}$$

9.1 Simple Cubic Lattice

The number of atoms per unit cell are

$$\begin{aligned} \text{atoms per unit cell} &= \underbrace{\frac{1}{8}}_{\text{atoms per corner}} \times \underbrace{8}_{\text{number of corners}} \\ &= 1 \end{aligned}$$

The atomic packing factor is

$$\begin{aligned} \text{APF} &= \frac{\overbrace{\frac{4}{3}\pi r^3}^{\text{volume of 1 atom}} \times \underbrace{1}_{\text{number of atoms}}}{\underbrace{a^3}_{\text{volume of unit cell}}} = \frac{\overbrace{\frac{4}{3}\pi \left(\frac{a}{2}\right)^3}^{\text{volume of 1 atom}} \times \underbrace{1}_{\text{number of atoms}}}{\underbrace{a^3}_{\text{volume of unit cell}}} \\ &= \frac{\pi}{6} \end{aligned}$$

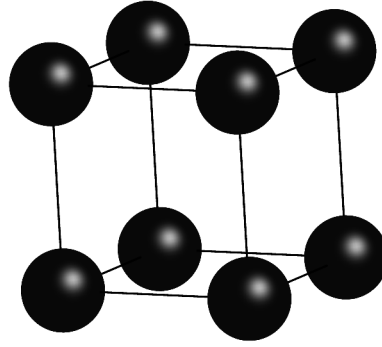


Figure 18: Simple cubic lattice (SC)

9.2 Face Centred Cubic Lattice

The number of atoms per unit cell are

$$\begin{aligned}
 \text{atoms per unit cell} &= \underbrace{\frac{1}{8}}_{\text{atoms per corner}} \times \underbrace{8}_{\text{number of corners}} \\
 &+ \underbrace{\frac{1}{2}}_{\text{atoms per face}} \times \underbrace{6}_{\text{number of faces}} \\
 &= 4
 \end{aligned}$$

The atomic packing factor is

$$\begin{aligned}
 \text{APF} &= \frac{\overbrace{\frac{4}{3}\pi r^3}^{\text{volume of 1 atom}} \times \underbrace{4}_{\text{number of atoms}}}{\underbrace{a^3}_{\text{volume of unit cell}}} = \frac{\overbrace{\frac{4}{3}\pi \left(\frac{a}{2\sqrt{2}}\right)^3}^{\text{volume of 1 atom}} \times \underbrace{4}_{\text{number of atoms}}}{\underbrace{a^3}_{\text{volume of unit cell}}} \\
 &= \frac{\pi}{3\sqrt{2}}
 \end{aligned}$$

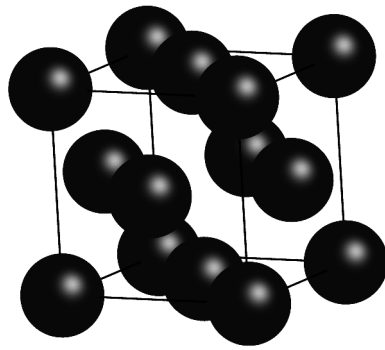


Figure 19: Face centred cubic lattice (FCC)

9.3 Body Centred Cubic Lattice

The number of atoms per unit cell are

$$\begin{aligned}
 \text{atoms per unit cell} &= \underbrace{\frac{1}{8}}_{\text{atoms per corner}} \times \underbrace{8}_{\text{number of corners}} + \underbrace{1}_{\text{atom at centre}} \\
 &= 2
 \end{aligned}$$

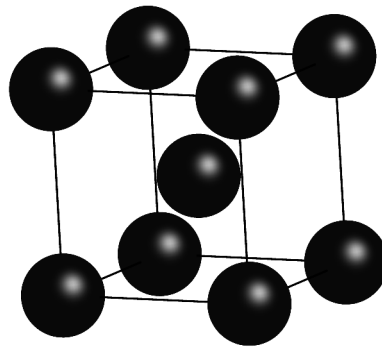


Figure 20: Body centred cubic lattice (BCC)

10 Basics of Crystal Growth

Definition 33. Epitaxy is the growth of layers, called epitaxial layers, on top of a thicker substrate layer.

Homoepitaxy is growing an epitaxial layer of the same material as the substrate. Heteroepitaxy is growing an epitaxial layer of a material different from the substrate.

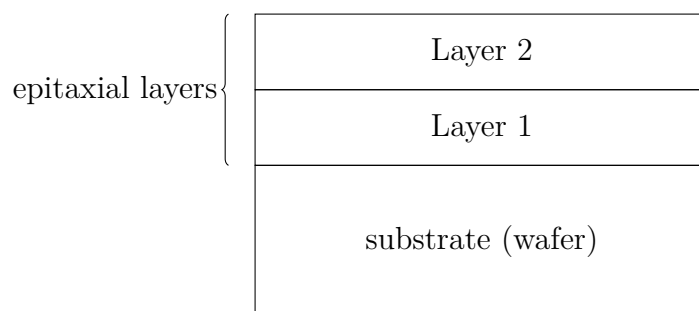


Figure 21: Epitaxial layers

Due to different properties of the materials which make up the substrate

and epitaxial layers, there are certain constraints on the possible combinations. These factors include

1. lattice constants
2. crystal structure

10.1 Lattice Matching

If there is a difference between the lattice constants of the materials of two adjacent layers, there is a stress or strain generated, due to the interaction of the different lattices.

To minimize this stress, a technique called lattice matching is used.

Suppose the lattice constants of AlAs, InAs, and GaAs are as shown in the graph.

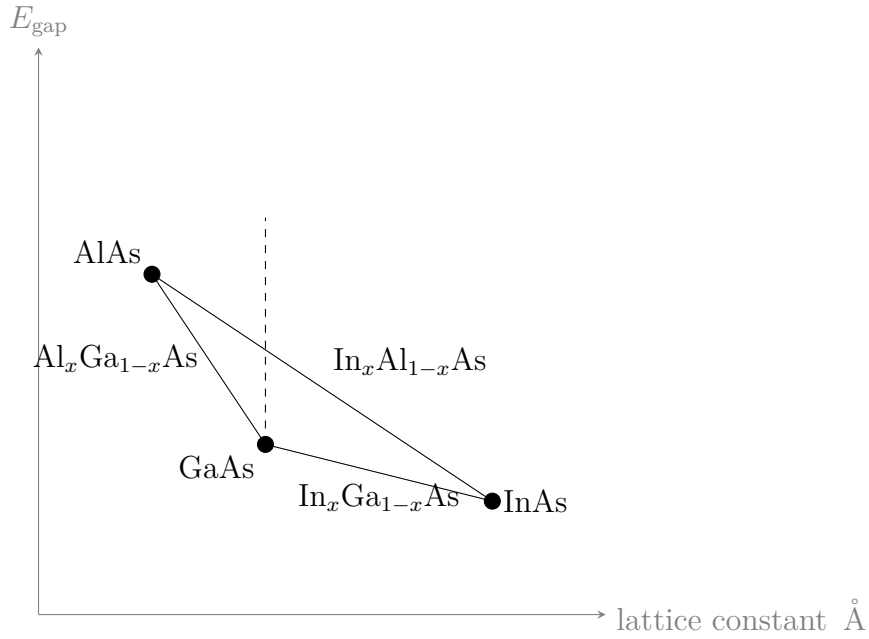


Figure 22: Lattice matching

If the substrate used is GaAs, and the epitaxial layer is made of InAlAs, then the composition of the epitaxial layer of InAlAs can be determined graphically.

The straight line connecting the points representing AlAs and InAs represents all possible values of $\text{In}_x\text{Al}_{1-x}\text{As}$, with x varying from 0 to 1, with 0

corresponding to AlAs, and 1 corresponding to InAs.

Therefore, the ideal material for the epitaxial layer, is the composition of $\text{In}_x\text{Al}_{1-x}\text{As}$ with lattice constant equal to that of GaAs.

11 Thermal Motion

Definition 34 (Ionization energy). The energy needed to break a bond and free an electron is called the binding energy, or the ionization energy.

Definition 35 (Thermal motion). The motion of electrons due to the thermal energy provided, is called thermal motion. This motion is random in nature, with all electrons moving around, rotating, and interacting, i.e. colliding with other electrons.

Theorem 35. *From statistical mechanics, which models the motion of free electrons, the thermal energy of an electron is*

$$E_{\text{thermal}} = \frac{3}{2}kT$$

where k is Boltzmann's constant, and T is the temperature in kelvin. Hence, the average velocity of an electron is

$$v = \sqrt{\frac{3kT}{m}}$$

11.1 Thermal Generation of Carriers

Definition 36 (EHP). A pair of a free electron, and a positive charge called a hole is called an electron-hole pair, or an EHP. Both the electron and the hole are mobile, and can contribute to conduction.

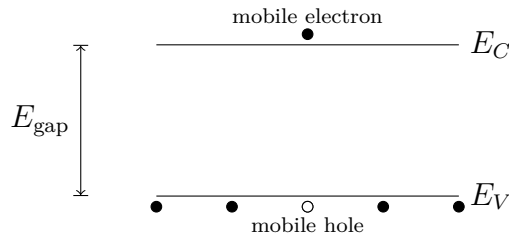


Figure 23: EHP

12 Effective Mass

Definition 37 (Effective mass). To account for influences of the crystal lattice, altered values of particle masses are used.

These values are called effective masses of particles. They indicate how an electron will accelerate in the presence of an external force.

For silicon, the effective masses of electrons and holes are

$$m_n^* = 0.25m_0$$

$$m_p^* = 0.5m_0$$

where m_0 is the real mass of an electron.

13 Intrinsic Semiconductors

Definition 38 (Intrinsic carrier concentration). The concentration of electrons, which is equal to the concentration of holes, in a material is called the intrinsic carrier concentration. It is denoted by n_i .

Therefore,

$$n = p = n_i$$

13.1 Effect of Energy Band Gap on Intrinsic Carrier Concentration

If the energy band gap increases, electrons need more energy to go from the valence band to the conduction band. Therefore, at a specific temperature, the intrinsic carrier concentration decreases as the energy band gap increases.

13.2 Thermal Generation and Recombination

Definition 39 (Generation of an EHP). The process of creation of an EHP, which occurs due to an electron transitioning to the ionization band is called generation of an EHP. The rate of generation of EHPs is denoted by G .

Definition 40 (Recombination). The process of destruction of an EHP, which occurs when an electron transitions back to an empty state, is called recombination. The rate of recombination of EHPs is denoted by R .

For a steady state concentration to be maintained at a given temperature, the recombination of EHPs must take place at the same rate at which they are generated, i.e.,

$$R = G$$

Therefore,

$$\begin{aligned} G &= \alpha_i np \\ &= \alpha_i n_i^2 \\ &= R \end{aligned}$$

where α_i is the recombination coefficient.

At thermal equilibrium,

$$np = n_i^2$$

14 Extrinsic Semiconductors

Definition 41 (Doping). The process of introducing atoms into the crystals of a semiconducting material is called doping.

The newly introduced materials are called dopants or impurities.

The dopant atoms either donate or accept electrons, and hence increase the conductivity, as compared to the intrinsic conductivity.

Definition 42 (Dopant concentration). Dopant concentration is defined as the number of donors or acceptors per unit volume.

Typically the values range from 10^{12} to 10^{19} donors or acceptors per cm^3 .



Figure 24: Typical Energy Diagram for Doped Material

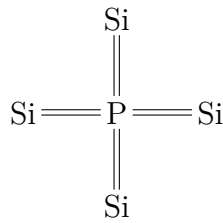


Figure 25: Silicon lattice with phosphorus doping

14.1 N-type Material

Consider a silicon lattice doped with phosphorus atoms. Four of the five valence electrons of the phosphorus atoms form covalent bonds with the adjacent silicon atoms. The extra electron, which is loosely bound, is donated to the lattice, and hence is available for conduction. Such a material is called a N-type material.

14.2 P-type Material

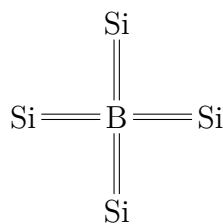


Figure 26: Silicon lattice with boron doping

Consider a silicon lattice doped with boron atoms. The three valence electrons of the boron atoms form covalent bonds with the adjacent silicon atoms. Therefore, there is one missing electron, as compared to the original lattice. Hence, there is a hole per acceptor atom in the lattice, which is available for conduction. Such a material is called a P-type material.

14.3 Thermal Equilibrium

Type of Material	Majority Carriers	Minority Carriers
N-type ($N_d \gg n_i, p$)	Electrons $n = N_d$	Holes $p = \frac{n_i^2}{n}$
P-type ($N_a \gg n_i, n$)	Holes $p = N_a$	Electrons $n = \frac{n_i^2}{p}$

Table 2: Concentration of charge carriers in N-type and P-type materials

14.4 Dependence of Carrier Concentration on Temperature

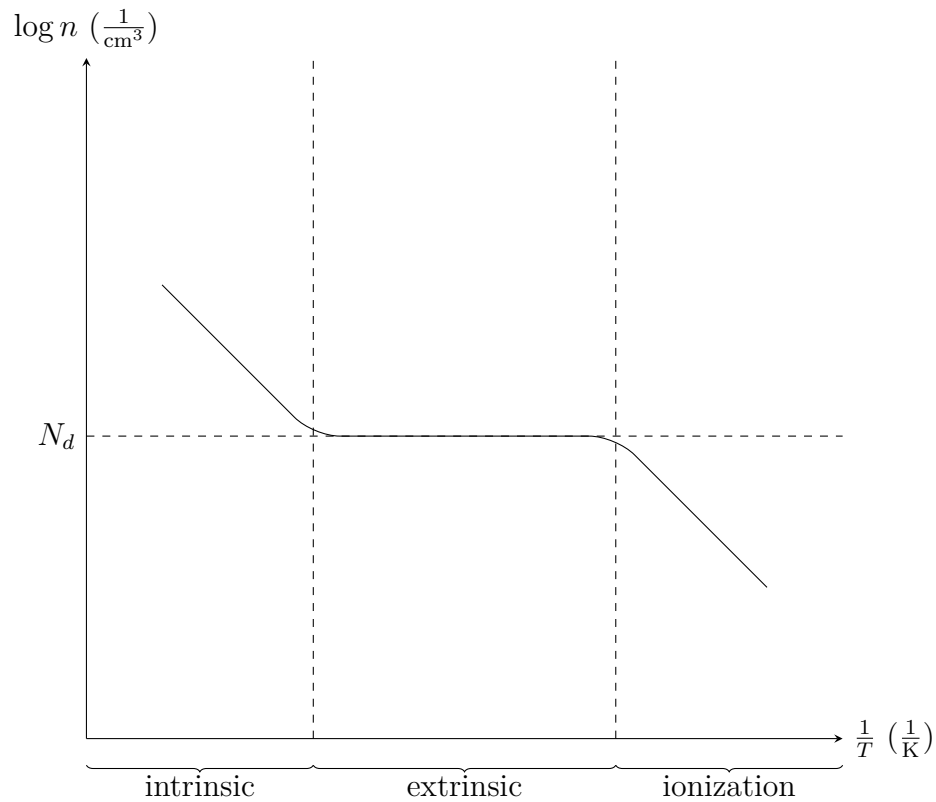


Figure 27: Carrier concentration and temperature regions

At very low temperatures, i.e. close to absolute zero, the number of free electrons are almost zero. As the temperature increases, donor atoms start

ionizing, and hence the number of free electrons increases. This range of temperature, up to the temperature at which all donor atoms are ionized is called the ionization region.

After the temperature at which all donor atoms are ionized, up to the temperature at which the in the host atoms start ionizing and generating EHPs, is called the extrinsic region. In this region, the number of free electrons is equal to the number of electrons donated by the dopant atoms, i.e. $n = N_d$. After the temperature at which the host atoms start ionizing, the number of electrons is equal to the sum of the number of electrons donated by the dopant atoms, and the number of thermally generated EHPs. This region is called the intrinsic region.

15 Energy Band Model

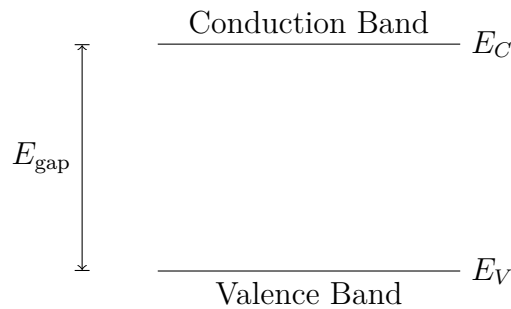


Figure 28: Energy Bands

Definition 43. The energy of a free electron is called E_{vacuum} .

Definition 44. The energy level of the n th shell in an atom is denoted by E_n .

15.1 Splitting of Energy Levels

The energy levels of electrons in an atom can be represented as

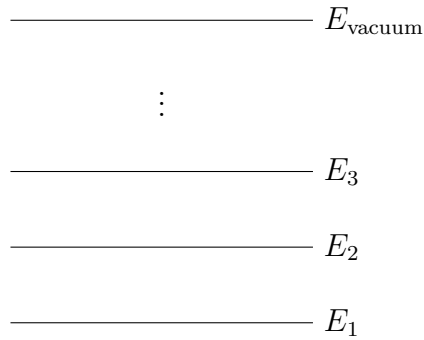


Figure 29: Energy Levels in a Single Atom

According to Pauli's Exclusion Principle, no two electrons can occupy the same energy level. Hence, if two atoms are brought close to each other, the energy levels of the shells split, and form new energy levels, which can be represented as

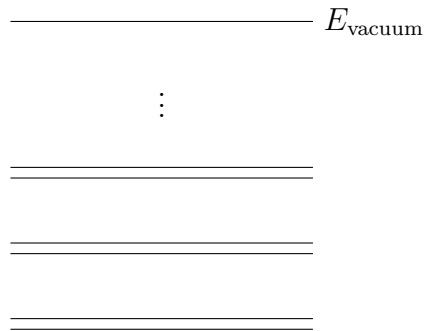


Figure 30: Split Energy Levels in two Atoms

Therefore, in the limiting case, if an infinite number of atoms are brought together to form a crystal, the split energy levels form ranges of energy which can be occupied by electrons. These energy ranges are called energy bands.

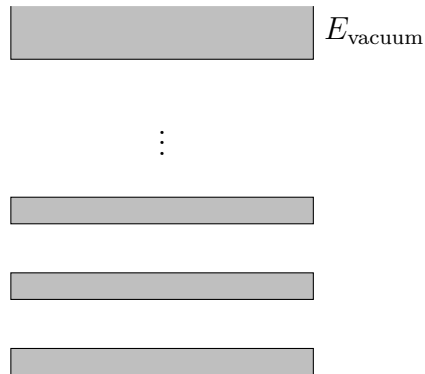


Figure 31: Energy Bands in a Crystal

For a semiconductor, the energy bands can be represented as

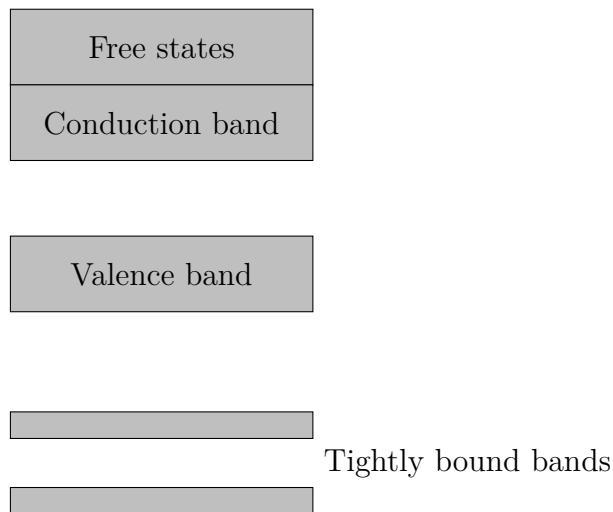


Figure 32: Energy Bands in a Semiconductor Crystal

Definition 45. The energy gap between E_{vacuum} and $E_{\text{conduction}}$ is called the electron affinity for that particular material. It is denoted by $q\chi$.

Definition 46. The intrinsic energy level is denoted by E_i .



Figure 33: Energy Gaps for a Perfect Semiconductor

16 Presence of an Electric Field in the Energy Band Model

Suppose a voltage V is applied across a semiconductor. Therefore,

$$\begin{aligned}\vec{E} &= -\frac{dV}{dx} \\ &= \frac{1}{q} \frac{dE_i}{dx}\end{aligned}$$

where E_i is the intrinsic energy level.

Due to the electric field generated, the electrons move towards the positive terminal and the holes towards the negative terminal.

Let the potential difference be applied such that the electric field generated is directed towards the left.

Hence, the energy bands are modified, and a slope is generated in the bands. Therefore, the energy bands which were originally represented as in Figure 34, will now be represented as in Figure 35. Hence, the electrons in the energy bands tend to move to the right.

This behaviour is analogous to the behaviour of water in tilted pipes, with the water representing the electrons and the air representing the holes.

In general, the direction of the electric field is the same as the slope of the energy bands.

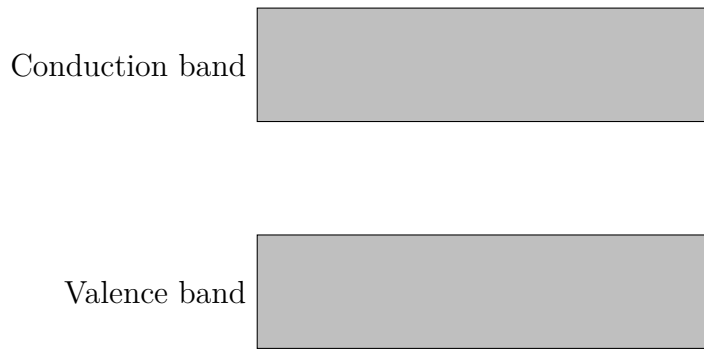


Figure 34: Energy Bands for a Perfect Semiconductor

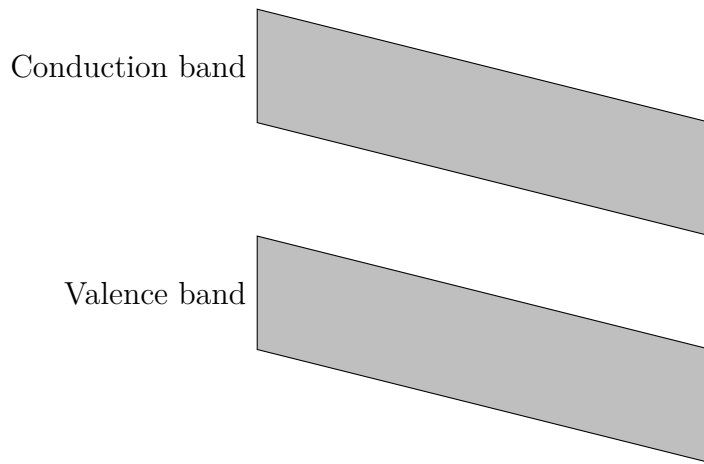


Figure 35: Energy Bands for a Perfect Semiconductor in Presence of an Electric Field

Exercise 21.

Given an electric field

$$\vec{E}(x) = -c \frac{dn}{dx}$$

where

$$c = \frac{kT}{q} = 0.026 \text{ eV (at room temperature)}$$

The doping profile is given to be

$$\begin{aligned} n(x) &= N_d(x) \\ &= N_0 e^{-\alpha x} \end{aligned}$$

1. Find $\vec{E}(x)$.
2. If $\alpha = 1 \frac{1}{\mu\text{m}}$, what is $\vec{E}(x)$?
3. Sketch the energy band diagram.

Solution 21.

1.

$$\begin{aligned}\frac{dn}{dx} &= N_0 \frac{de^{-\alpha x}}{dx} \\ &= -N_0 \alpha e^{-\alpha x}\end{aligned}$$

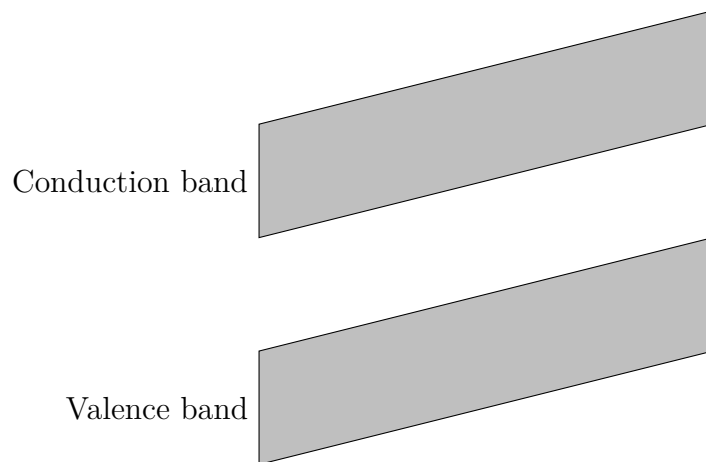
Therefore,

$$\begin{aligned}\vec{E}(x) &= -c \frac{-N_0 \alpha e^{-\alpha x}}{N_0 e^{-\alpha x}} \\ &= \frac{c N_0 \alpha e^{-\alpha x}}{N_0 e^{-\alpha x}} \\ &= c \alpha\end{aligned}$$

2.

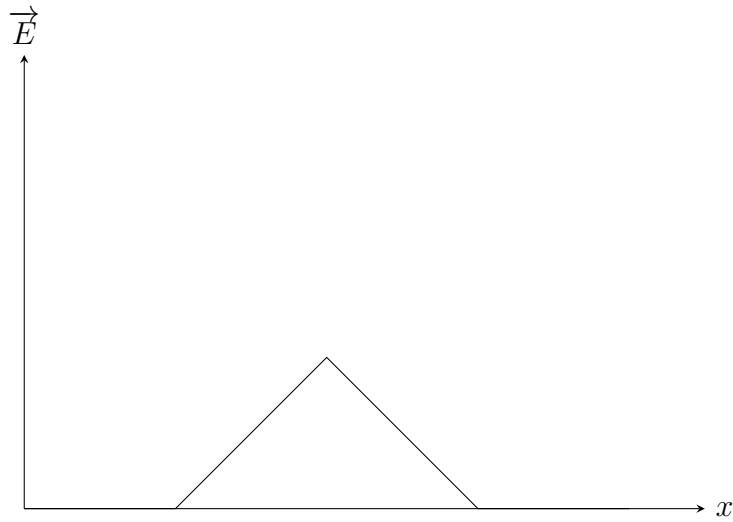
$$\begin{aligned}\vec{E}(x) &= c \alpha \\ &= (0.026\text{V})(1 \frac{1}{\mu\text{m}}) \\ &= 260 \frac{\text{V}}{\text{cm}}\end{aligned}$$

3. As the electric field is positive, the slope of the energy bands also must be positive. As the electric field is constant, the energy bands must be linear. Therefore, the energy band diagram is



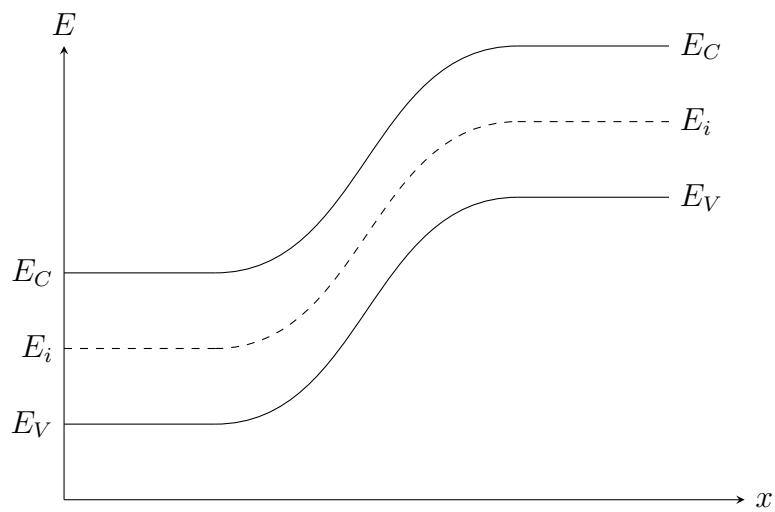
Exercise 22.

The electric field through a sample is given to be as shown.



Sketch the energy levels with respect to x .

Solution 22.



17 Movement of Carriers in Semiconductor Crystals

The motion of charge carriers in a semiconductor crystal is affected by the interactions within the crystal. The collisions between the charge carriers affects their movement.

Definition 47 (Scattering processes). The change in direction of movement of charge carriers, due to collisions is called scattering processes.

Definition 48 (Mean free path). The average distance between collision is called the mean free path. It is denoted by l_m .

Definition 49 (Time between collisions). The average time between collision is called the mean free path. It is denoted by τ_m .

17.1 Drift

Consider an electric field \vec{E} applied on a sample. Therefore the net motion of the electrons is in the direction opposite to that of the electric field, and the net motion of the holes is in the direction of the electric field.

Definition 50 (Drift). The motion of charge carriers which is driven by electric fields is called drift.

The acceleration of the charge carriers due to the external electric field is balanced by the deceleration of the charge carriers due to collisions. Therefore, on a microscopic level, the motion of the charge carriers is random, and on a macroscopic level, the charge carriers move with a constant velocity.

Definition 51 (Average drift velocity). The average velocity of charge carriers, on a macroscopic scale, is called the average drift velocity.

The force on a charge carrier is given by

$$\vec{F} = q\vec{E}$$

Therefore,

$$\begin{aligned}\vec{F}\tau_m &= q\vec{E}\tau_m \\ &= m^*v_d \\ \therefore v_d &= \frac{q\tau_m}{m^*}\vec{E}\end{aligned}$$

Definition 52 (Mobility). The ease with which electrons and holes can flow through the crystal is called mobility. It is denoted by μ . It is defined to be

$$\mu = \frac{q\tau}{m^*}$$

where q is the charge on the particle, τ is the average time between collisions, and m^* is the effective mass.

Mobility is the ratio between the drift velocity of electrons or holes, and the electric field acting on them.

Therefore,

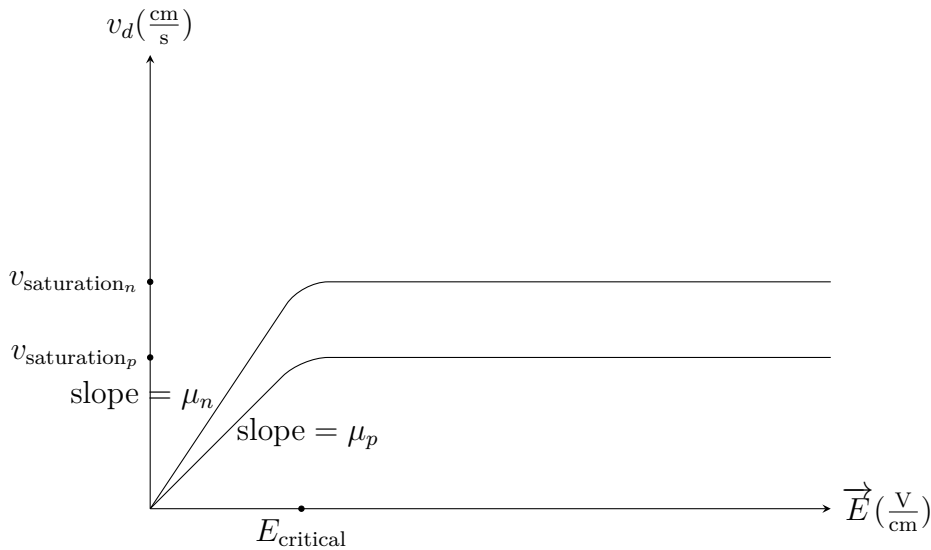
$$\mu = \frac{q\tau_m}{m^*}$$

As the effective mass of an electron is less than that of a hole, the mobility of an electron is more than that of a hole.

For silicon, at room temperature, the mobilities of electrons and holes are

$$\mu_n = 1000 \frac{\text{cm}^2}{\text{V s}}$$

$$\mu_p = 400 \frac{\text{cm}^2}{\text{V s}}$$



17.1.1 Current Density due to Net Drift

Definition 53 (Direction of current due to drift of electrons). The direction of the current generated due to the net drift of electrons is defined to be in the direction opposite to the drift velocity of the electrons.

Definition 54 (Direction of current due to drift of holes). The direction of the current generated due to the net drift of holes is defined to be in the direction of the drift velocity of the holes.

Theorem 36 (Current density due to net drift). *The current density due to the drift of electrons is*

$$\begin{aligned} J_{drift_n} &= -qn v_{d_n} \\ &= qn \mu_n \vec{E} \end{aligned}$$

The current density due to the drift of holes is

$$\begin{aligned} J_{drift_p} &= qp v_{d_p} \\ &= qp \mu_p \vec{E} \end{aligned}$$

Therefore, the net current density is

$$\begin{aligned} J_{drift_{total}} &= J_{drift_n} + J_{drift_p} \\ &= qn \mu_n \vec{E} + qp \mu_p \vec{E} \\ &= q(n \mu_n + p \mu_p) \vec{E} \end{aligned}$$

Theorem 37.

$$\sigma = q(n \mu_n + p \mu_p)$$

Proof.

$$\begin{aligned} J &= \sigma \vec{E} \\ &= q(n \mu_n + p \mu_p) \vec{E} \end{aligned}$$

Therefore,

$$\sigma = q(n \mu_n + p \mu_p)$$

□

Theorem 38. *For a N-type material,*

$$\begin{aligned} \sigma &= \sigma_n \\ &= qn \mu_n \end{aligned}$$

Proof. For a N-type material, $n \gg p$.
Therefore,

$$\begin{aligned}\sigma &= q(n\mu_n + p\mu_p) \\ &= qn\mu_n\end{aligned}$$

□

Theorem 39. For a P-type material,

$$\begin{aligned}\sigma &= \sigma_p \\ &= qp\mu_p\end{aligned}$$

Proof. For a P-type material, $p \gg n$.
Therefore,

$$\begin{aligned}\sigma &= q(n\mu_n + p\mu_p) \\ &= qp\mu_p\end{aligned}$$

□

Exercise 23.

A sample of a P-type material with length $l = 10\mu\text{m}$ is connected to a voltage of $V = 10\text{ V}$. The dopant concentration in the material is $N_a = 10^{17} \frac{1}{\text{cm}^3}$. Find the drift velocity of the holes and the current density.

Solution 23.

$$\begin{aligned}\mu_p &= \frac{q\tau}{m^*} \\ &= \frac{(1.6 \times 10^{-19} \text{ C})(10^{-13} \text{ s})}{0.5 \times 9.11 \times 10^{-31} \text{ kg}} \\ &= 0.035 \frac{\text{m}^2}{\text{V s}}\end{aligned}$$

Therefore,

$$\begin{aligned}v_{\text{drift}_p} &= \mu_p E \\ &= \mu_p \frac{V}{l} \\ &= \left(350 \frac{\text{cm}^2}{\text{V s}}\right) \left(10^4 \frac{\text{V}}{\text{cm}}\right) \\ &= 3.5 \times 10^6 \frac{\text{cm}}{\text{s}}\end{aligned}$$

As the material is P type,

$$\begin{aligned}
J_{\text{drift}} &= J_{\text{drift}_p} \\
&= qp\mu_p E \\
&= qp v_{\text{drift}_p} \\
&= (1.6 \times 10^{-19} \text{ C}) \left(10^{17} \frac{1}{\text{cm}^3} \right) \left(3.5 \times 10^6 \frac{\text{cm}}{\text{s}} \right) \\
&= 5.6 \times 10^4 \frac{\text{A}}{\text{cm}^2}
\end{aligned}$$

17.2 Scattering Mechanisms that Influence Electron and Hole Mobility

Definition 55 (Lattice scattering). The scattering of charge carriers due to interaction with atoms in the crystal lattice is called lattice scattering.

In this scattering, there are collisions of moving carriers due to disturbances in the semiconductor. Therefore, as the velocity of the moving carriers increases with temperature, lattice scattering dominates the mobility at high temperature.

Definition 56 (Impurity scattering). The scattering of charge carriers due to crystal defects is called impurity scattering.

In this scattering, the fixed charges dispersed throughout the crystal exert forces on the charge carriers.

A slow moving carrier is likely to be scattered more strongly due to interaction, than a fast moving carrier. Therefore, impurity scattering dominates the mobility of the charge carriers at low temperature.

The temperature range for the graph in Figure 36 is in the extrinsic region corresponding to the graph in Figure 27. Therefore, the graph of $n\mu$ with respect to $\frac{1}{T}$ is as in Figure 37. Hence, as J and σ are directly proportional to $n\mu$, they exhibit a similar relation to the temperature.

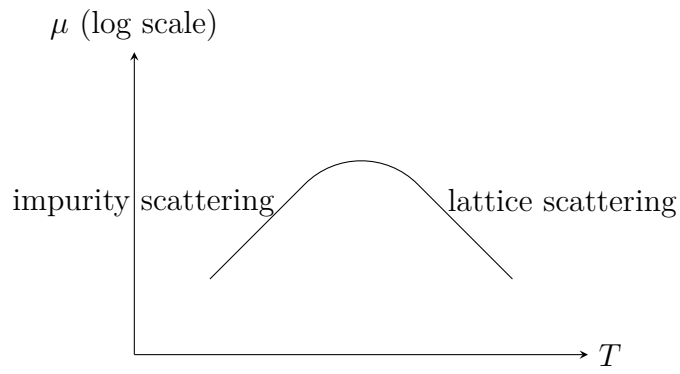


Figure 36: Effect of Temperature on Mobility

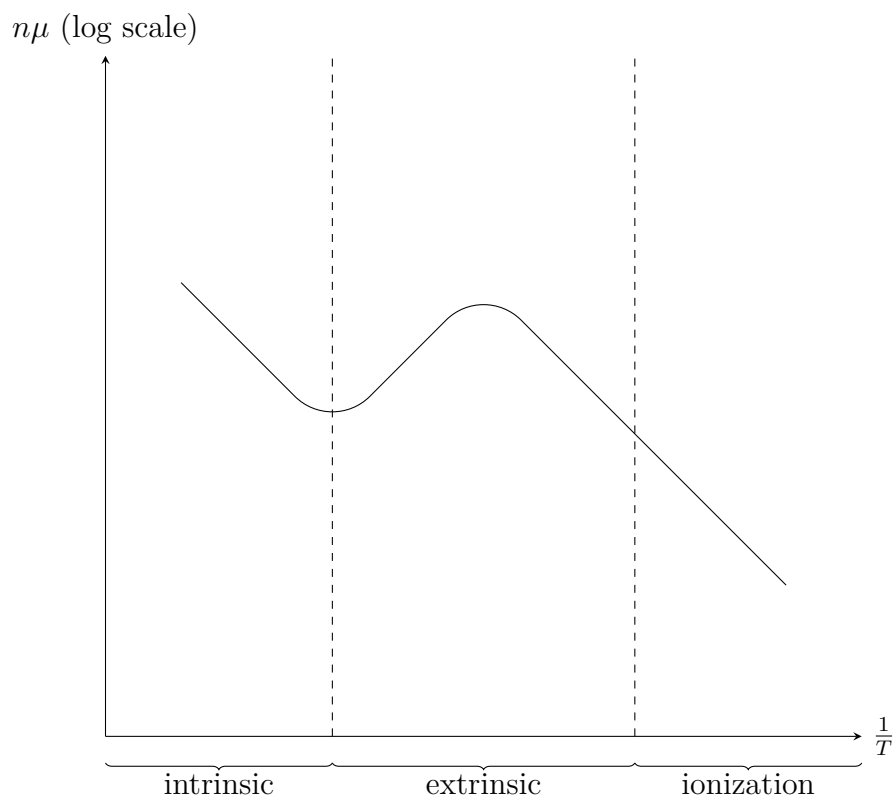


Figure 37: Effect of Doping Concentration and Temperature on Mobility

17.3 Diffusion

Definition 57 (Diffusion). The motion of charge carriers from regions of higher concentration to regions of lower concentration, is called diffusion. It is a gradient driven motion.

Definition 58 (Direction of current due to diffusion of electrons). The direction of the current generated due to the diffusion of electrons is defined to be in the direction opposite to the movement of the electrons.

Definition 59 (Direction of current due to diffusion of holes). The direction of the current generated due to the diffusion of holes is defined to be in the direction of the movement of the holes.

17.3.1 Current Density due to Net Diffusion

Definition 60 (Direction of current due to diffusion of electrons). The direction of the current generated due to the diffusion of electrons is defined to be opposite to the direction of the movement of electrons, and hence in the direction of increase in concentration of electrons with respect to position.

Definition 61 (Direction of current due to diffusion of holes). The direction of the current generated due to the diffusion of holes is defined to be in the direction of the movement of holes, and hence opposite to the direction of increase in concentration of holes with respect to position.

Theorem 40 (Current density due to diffusion). *The current density due to the diffusion of electrons is*

$$J_{diffusion_n} = qD_n \frac{dn}{dx}$$

The current density due to the diffusion of holes is

$$J_{diffusion_p} = -qD_p \frac{dp}{dx}$$

where D_n and D_p are positive constants.

Theorem 41.

$$D_n = \frac{1}{3}v_{thermal_n}l$$

where the motion of electrons is from $x = l$ to $x = -l$.

Proof. Consider a N-type sample with some doping profile. The flux through a cross-section at $x = 0$ is

$$F = \frac{1}{6} v_{\text{thermal}} (n(-l) - n(l))$$

where electrons move from $x = l$ to $x = -l$.
Therefore,

$$\begin{aligned} J_{\text{diffusion}_n} &= -qF \\ &= (-q) \left(-\frac{1}{3} v_{\text{thermal}} l \frac{dn}{dx} \right) \\ &= q \left(\frac{1}{2} v_{\text{thermal}} l \right) \frac{dn}{dx} \end{aligned}$$

The $\frac{1}{6}$ is due to the fact that an electron can move through 3 spatial dimensions, with 2 directions of flow each.

Therefore,

$$D_n = \frac{1}{3} v_{\text{thermal}_n} l$$

□

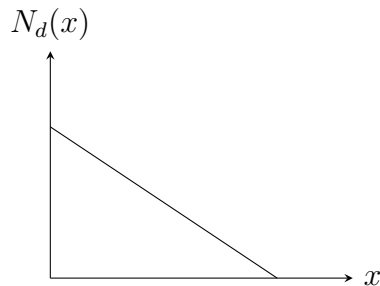
Theorem 42.

$$D_p = \frac{1}{3} v_{\text{thermal}_p} l$$

where the motion of holes is from $x = l$ to $x = -l$.

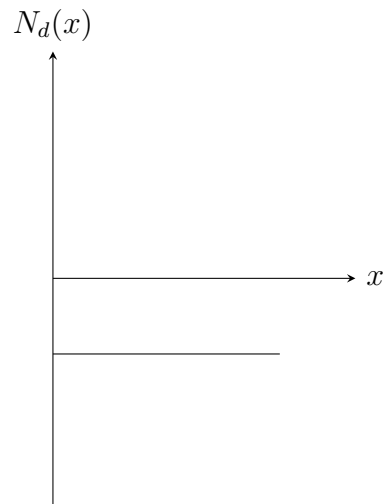
Exercise 24.

A sample of N-type material has doping profile as shown.



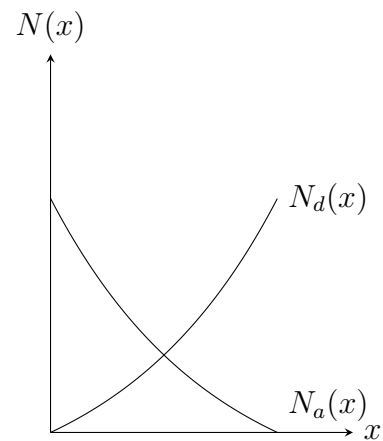
Sketch the graph of the current density with respect to x .

Solution 24.



Exercise 25.

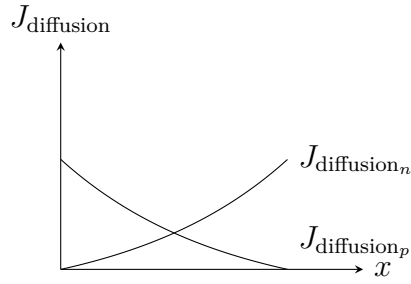
A sample has doping profile as shown.



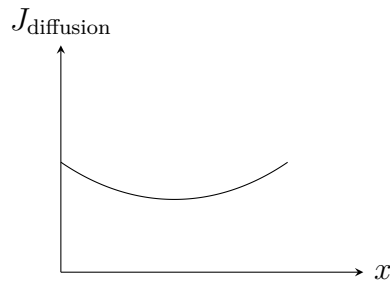
Sketch the graph of the current density with respect to x .

Solution 25.

The current densities due to the diffusion of electrons and holes are



Therefore, the net diffusion current density is



17.4 Transport Equations

The total current densities due to the net motion of electrons and holes is

$$\begin{aligned} J_{\text{total}_n} &= J_{\text{drift}_n} + J_{\text{diffusion}_n} \\ &= qn\mu_n \vec{E} + qD_n \frac{dn}{dx} \end{aligned}$$

$$\begin{aligned} J_{\text{total}_p} &= J_{\text{drift}_p} + J_{\text{diffusion}_p} \\ &= qp\mu_p \vec{E} + qD_p \frac{dp}{dx} \end{aligned}$$

17.5 Einstein's Relation

Theorem 43 (Einstein's Relation).

$$\frac{D}{\mu} = \frac{kT}{q}$$

At room temperature,

$$\frac{D}{\mu} = 0.026 \text{ V}$$

Proof.

$$\begin{aligned}
 v_{\text{thermal}} &= \frac{l}{\tau} \\
 \frac{D}{\mu} &= \frac{\frac{1}{3} v_{\text{thermal}} l}{\mu} \\
 &= \frac{\frac{1}{3} v_{\text{thermal}}^2 \tau}{\mu} \\
 &= \frac{\frac{1}{3} v_{\text{thermal}}^2 \tau}{\frac{q\tau}{m^*}} \\
 \frac{1}{2} m^* v_{\text{thermal}}^2 &= \frac{3}{2} kT \\
 &= \frac{1}{3} \frac{m^* v_{\text{thermal}}^2}{q} \\
 &= \frac{1}{3} \frac{3kT}{q} \\
 &= \frac{kT}{q}
 \end{aligned}$$

□

Exercise 26.

Find D_n for a sample of silicon at room temperature, if the mobility of electrons in the sample is $1000 \frac{\text{cm}^2}{\text{V}\cdot\text{s}}$.

Solution 26.

By Einstein's Relation,

$$\begin{aligned}
 \frac{D_n}{\mu_n} &= 0.026 \text{ V} \\
 \therefore D_n &= 0.026 \mu_n \frac{\text{cm}^2}{\text{s}} \\
 &= 26 \frac{\text{cm}^2}{\text{s}}
 \end{aligned}$$

Exercise 27.

Find D_n for a sample of silicon at 450 K, if the mobility of electrons in the sample is $800 \frac{\text{cm}^2}{\text{V}\cdot\text{s}}$.

Solution 27.

By Einstein's Relation,

$$\begin{aligned}\frac{D_n}{\mu_n} &= \frac{kT}{q} \\ \therefore D_n &= \frac{kT}{q} \mu_n \\ &= 49.7 \frac{\text{cm}^2}{\text{s}}\end{aligned}$$

18 Quantum Wells

Consider a sample made with one semiconductor sandwiched between two samples of another semiconductor. The energy bands in this composite material is as in Figure 38. Hence, there are quantum wells created in the energy bands.

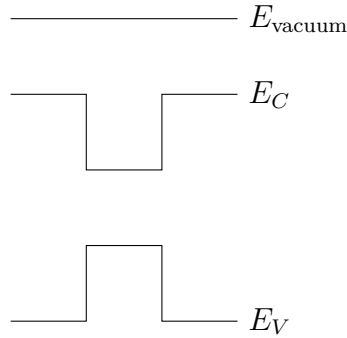


Figure 38: Quantum Wells due to Varying Energy Band Gaps

19 Optical Absorption

Let a sample be illuminated with light of photon energy greater than the energy band gap of the material of the sample.

Therefore, the electrons in the valence band absorb this energy, and move to the conduction band. Such electrons are called hot electrons. Hot electrons then lose energy till they reach E_C . Thereafter, they recombine with holes in the valence band, and in the process give out photons corresponding to the energy band gap of the material.

This emitted light can be observed, and used to find the energy band gap of the material.

20 Optical Generation of Carriers

Consider a sample illuminated by light with photons of energy greater than the energy gap of the material of the sample. This extra energy will lead to generation of EHPs.

Let n_0 and p_0 be the concentrations of electrons and holes, at thermal equilibrium. Let \hat{n} and \hat{p} be the concentrations of electrons and holes, from the optically generated EHPs.

Therefore,

$$n = n_0 + \hat{n}$$

$$p = p_0 + \hat{p}$$

20.1 Carrier Generation Dependent on Illumination

Consider a sample such that

$$N_d = 10^{16} \frac{1}{\text{cm}^3}$$

Therefore the intensity of light affects the concentration of charge carriers. For no light,

$$n = n_0$$

$$= N_d$$

$$= 10^{16} \frac{1}{\text{cm}^3}$$

$$p = p_0$$

$$= \frac{n_i^2}{N_d}$$

$$= 10^4 \frac{1}{\text{cm}^3}$$

For low level injection, i.e. $p_0 < \hat{n} = \hat{p} < n_0$, say $\hat{n} = \hat{p} \approx 10^{14} \frac{1}{\text{cm}^3}$,

$$\begin{aligned} n &= n_0 + \hat{n} \\ &= 10^{16} + 10^{14} \\ &\approx 10^{16} \frac{1}{\text{cm}^3} \\ p &= p_0 + \hat{p} \\ &= 10^4 + 10^{14} \\ &\approx 10^{14} \frac{1}{\text{cm}^3} \end{aligned}$$

For high level injection, i.e. $p_0 < \hat{n} = \hat{p} < n_0$, say $\hat{n} = \hat{p} \approx 10^{20} \frac{1}{\text{cm}^3}$,

$$\begin{aligned} n &= n_0 + \hat{n} \\ &= 10^{16} + 10^{20} \\ &\approx 10^{20} \frac{1}{\text{cm}^3} \\ p &= p_0 + \hat{p} \\ &= 10^4 + 10^{20} \\ &\approx 10^{20} \frac{1}{\text{cm}^3} \end{aligned}$$

Therefore, the concentrations are as shown in Figure 39.

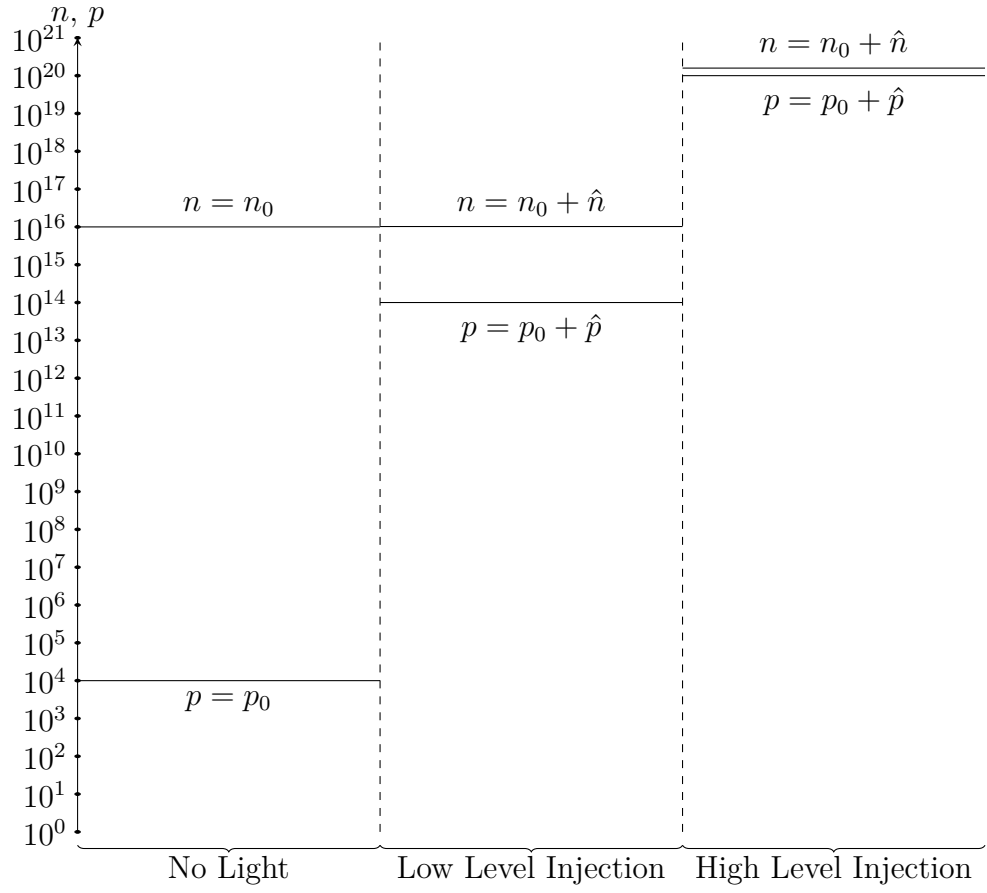


Figure 39: Carrier Concentration in N-type Material under Different Light Conditions

20.2 Minority Carrier Concentration in Illuminated Materials

Consider a N-type material which is illuminated from $t = 0$ onwards. The rate of recombination is defined to be

$$R = \alpha_r np$$

Therefore, in case of low level injection,

$$\begin{aligned}
R &= \alpha_r np \\
&= \alpha_r (n_0 + \hat{n}) (p_0 + \hat{p}) \\
&= \alpha_r (n_0 p_0 + n_0 \hat{p} + p_0 \hat{n} \hat{p} \hat{n}) \\
&= \alpha_r (n_i^2 + n_0 \hat{p}) \\
&= \alpha_r n_i^2 + \alpha_r n_0 \hat{p} \\
&= R_{\text{thermal}} + R_{\text{optical}}
\end{aligned}$$

As $p_0 \hat{n}$ and $\hat{p} \hat{n}$ are small compared to $n_0 p_0$ and $n_0 \hat{p}$.

Therefore,

$$\begin{aligned}
G - R &= G_{\text{thermal}} + G_{\text{optical}} - R_{\text{thermal}} + R_{\text{optical}} \\
&= G_{\text{optical}} - R_{\text{optical}}
\end{aligned}$$

In the steady state, the generation of EHPs is balanced by recombination. Therefore,

$$\begin{aligned}
G &= R \\
\therefore G_{\text{optical}} &= R_{\text{optical}}
\end{aligned}$$

As the rate of recombination cannot increase instantaneously, \hat{p} must be a function of time.

Definition 62 (Lifetime of minority carriers).

$$\tau_p = \frac{1}{\alpha_r n_0}$$

is called the lifetime of the minority carriers

Therefore,

$$\begin{aligned}
R &= \alpha_r n_0 \hat{p} \\
&= \frac{\hat{p}}{\tau_p}
\end{aligned}$$

Therefore, at steady state,

$$\begin{aligned}
G_{\text{optical}} &= R_{\text{optical}} \\
&= \frac{\hat{p}}{\tau_p}
\end{aligned}$$

Therefore,

$$\hat{p} = G_{\text{optical}}\tau_p$$

Therefore, before the steady state,

$$\begin{aligned} G - R &= \frac{d\hat{p}}{dt} \\ \therefore G_{\text{optical}} - \frac{\hat{p}(t)}{\tau_p} &= \frac{d\hat{p}}{dt} \end{aligned}$$

Therefore, solving,

$$\hat{p}(t) = G_{\text{optical}}\tau_p \left(1 - e^{-\frac{t}{\tau_p}}\right)$$

Similarly, for a sample which is stopped being illuminated at $t = 0$,

$$\hat{p}(t) = G_{\text{optical}}\tau_p e^{-\frac{t}{\tau_p}}$$

20.3 Conductivity of Illuminated Materials

When a sample is not illuminated,

$$J = \sigma_0 E$$

When the sample has light shined on it,

$$J = (\sigma_0 + \hat{\sigma}) E$$

Therefore,

$$\begin{aligned} \Delta\sigma &= \hat{\sigma} \\ &= \sigma_{\text{light}} - \sigma_{\text{dark}} \\ &= \left(q\mu_n(n + \hat{n}) + q\mu_p(p + \hat{p})\right) - \left(q\mu_n n + q\mu_p p\right) \\ &= q\mu_p \hat{p} + q\mu_n \hat{n} \\ &= qG_{\text{optical}}\tau_p(\mu_p + \mu_n) \end{aligned}$$

21 Photo-detectors

A device which changes its resistivity when exposed to light is called a photo-detector. The basic structure is as in Figure 40. However, realistically the structure of the device is as in Figure 41.

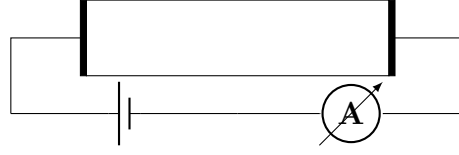


Figure 40: Basic Structure of a Photo-detector

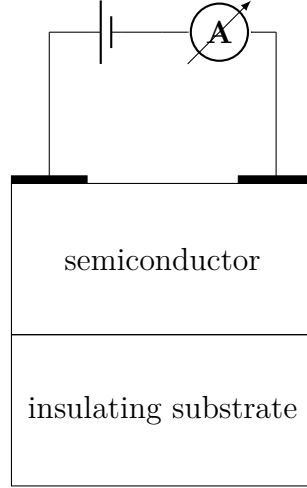


Figure 41: Actual Structure of a Photo-detector

22 Continuity Equations

Consider a slice of thickness Δx in a semiconductor sample. The rate of change of the number of holes in the slice, is the sum of the difference between the number of holes entering this slice and the number of holes exiting this slice, and the net rate of generation of holes.

$$\left. \frac{dp}{dx} \right|_{x \rightarrow x+\Delta x} = \frac{1}{q} \left(\frac{J_p(x) - J_p(x + \Delta x)}{\Delta x} \right) + (G - R)$$

As the temperature is constant, the thermally generated carrier concentrations do not change with time. Therefore, the only change responsible for the change in the net carrier concentration is the change in the optically generated carrier concentration.

Therefore, if the slice is infinitesimally thin,

$$\begin{aligned} \frac{\partial p(x, t)}{\partial t} &= \frac{\partial \hat{p}}{\partial t} \\ &= -\frac{1}{q} \left(\frac{\partial J_p(x, t)}{\partial x} \right) + \left(G_{\text{optical}} - \frac{\hat{p}(x, t)}{\tau_p} \right) \end{aligned}$$

Similarly, for electrons,

$$\begin{aligned}\frac{\partial n(x,t)}{\partial t} &= \frac{\partial \hat{n}}{\partial t} \\ &= \frac{1}{q} \left(\frac{\partial J_n(x,t)}{\partial x} \right) + \left(G_{\text{optical}} - \frac{\hat{n}(x,t)}{\tau_n} \right)\end{aligned}$$

22.1 Uniform Carrier Distribution

Assume

$$\vec{E} = 0$$

If the charge carrier distribution in a semiconductor is zero, i.e.,

$$\begin{aligned}\frac{\partial p}{\partial x} &= 0 \\ \frac{\partial n}{\partial x} &= 0\end{aligned}$$

there is no diffusion in the semiconductor. Therefore,

$$\frac{\partial \hat{p}}{\partial x} = 0$$

If the semiconductor is illuminated,

$$\frac{\partial \hat{p}}{\partial t} = G_{\text{optical}} - \frac{\hat{p}(t)}{\tau_p}$$

Therefore, solving,

$$\hat{p}(t) = G_{\text{optical}}\tau_p \left(1 - e^{-\frac{t}{\tau_p}} \right)$$

If the semiconductor is stopped being illuminated,

$$\frac{\partial \hat{p}}{\partial t} = -\frac{\hat{p}(t)}{\tau_p}$$

Therefore, solving,

$$\hat{p}(t) = \hat{p}_0 e^{-\frac{t}{\tau_p}}$$

where \hat{p}_0 is the initial condition.

22.2 Non-uniform Illumination, Steady State Injection, Illumination only at End

Definition 63 (Diffusion length).

$$L_p = \sqrt{D_p \tau_p}$$

is defined to be the diffusion length for holes.

$$L_n = \sqrt{D_n \tau_n}$$

is defined to be the diffusion length for electrons.

Consider a N-type semiconductor sample, connected to an ohmic contact at one end and illuminated only at $x = 0$.

Assume

$$\vec{E} = 0$$

If the illumination is non-uniform,

$$\begin{aligned} \frac{\partial p}{\partial x} &\neq 0 \\ \frac{\partial n}{\partial x} &\neq 0 \end{aligned}$$

If the illumination is a steady state injection,

$$\begin{aligned} \frac{\partial p}{\partial t} &= 0 \\ \frac{\partial n}{\partial t} &= 0 \end{aligned}$$

Therefore, the continuity equation for holes is

$$\begin{aligned} \frac{\partial \hat{p}(x, t)}{\partial t} &= -\frac{1}{q} \frac{\partial J_p(x, t)}{\partial x} - \frac{\hat{p}(x, t)}{\tau_p} \\ \therefore 0 &= -\frac{1}{q} \frac{\partial J_p(x, t)}{\partial x} - \frac{\hat{p}(x, t)}{\tau_p} \end{aligned}$$

Therefore, as $\vec{E} = 0$, the transport equation for J_p is

$$\begin{aligned} J_p &= J_{\text{diffusion}_p} \\ &= -q D_p \frac{\partial \hat{p}}{\partial x} \end{aligned}$$

Therefore, solving

$$0 = -\frac{1}{q} \frac{\partial}{\partial x} \left(-qD_p \frac{\partial \hat{p}}{\partial x} \right) - \frac{\hat{p}}{\tau_p}$$

$$\therefore \frac{\partial^2 \hat{p}}{\partial x^2} = \frac{\hat{p}}{D_p \tau_p}$$

Therefore,

$$\frac{\partial^2 \hat{p}}{\partial x^2} = \frac{\hat{p}}{L_p^2}$$

Therefore, as the sample is illuminated only at $x = 0$, for $0 < x < L$,

$$G_{\text{optical}} = 0$$

and for $x = 0$,

$$\hat{p}(0) = G_{\text{optical}} \tau_p$$

As there is an ohmic contact at $x = L$, there cannot be any charge accumulated there.

Therefore,

$$\hat{p}(L) = 0$$

Therefore, solving,

$$\hat{p}(x) = A e^{-\frac{x}{L_p}} + B e^{-\frac{x}{L_p}}$$

Therefore, substituting the boundary conditions, and solving,

$$\hat{p}(x) = G_{\text{optical}} \tau_p e^{-\frac{x}{L_p}}$$

If $L \ll L_p$, the recombination rate is negligible. Therefore,

$$\hat{R} = 0$$

Therefore, the differential equation is

$$\frac{\partial^2 \hat{p}}{\partial x^2} = 0$$

Therefore,

$$\hat{p}(x) = A + Bx$$

Therefore, substituting the boundary conditions, and solving,

$$\hat{p}(x) = \hat{p}(0) \left(1 - \frac{x}{L} \right)$$

Exercise 28.

An N-type semiconductor sample of Si at room temperature has the following properties.

$$\begin{aligned}
 N_d &= 10^{16} \frac{1}{\text{cm}^3} \\
 G_{\text{optical}} &= 10^{19} \frac{1}{\text{cm}^3 \text{ s}} \\
 \tau_n &= 1 \mu\text{s} \\
 \tau_p &= 1 \mu\text{s} \\
 D_n &= 40 \frac{\text{cm}^2}{\text{s}} \\
 D_p &= 25 \frac{\text{cm}^2}{\text{s}} \\
 L &= 1 \text{mm}
 \end{aligned}$$

Determine whether the sample is short or long.

Solution 28.

$$\begin{aligned}
 L_p &= \sqrt{D_p \tau_p} \\
 &= \sqrt{\left(25 \frac{\text{cm}^2}{\text{s}}\right) (10^{-6} \text{s})} \\
 &= 5 \times 10^{-3} \text{cm}
 \end{aligned}$$

Therefore,

$$\begin{aligned}
 L &= 1 \text{mm} \\
 &= 0.1 \text{mm} \\
 &= 5 \times 10^2 L_p
 \end{aligned}$$

Therefore, the sample is long.

Exercise 29.

Consider a N-type sample of Si at 300 K. Assuming low level injection, let the part of the sample from $x = -L_1$ to $x = 0$ be exposed to light, and let the part from $x = 0$ to $x = L_2$ be masked to block light. Assuming $L_1 \gg L_p$ and $L_2 \gg L_p$, find $\hat{p}(x)$.

Solution 29.

If the unmasked part was absent, $\hat{p}(x)$ would be

$$\hat{p}(x) = G_{\text{optical}}\tau_p e^{-\frac{x}{L_p}}$$

If the masked part was absent, $\hat{p}(x)$ would be

$$\hat{p}(x) = G_{\text{optical}}\tau_p$$

However, as the concentration must be continuous and differentiable, the graph of $\hat{p}(x)$ must be smooth.

By the steady state diffusion equation, on the illuminated part,

$$\begin{aligned} 0 &= \frac{\partial \hat{p}}{\partial x} \\ &= -\frac{1}{q} \frac{\partial J_p}{\partial x} + \left(G_{\text{optical}} - \frac{\hat{p}}{\tau_p} \right) \end{aligned}$$

As $\vec{E} = 0$, $J_p = J_{\text{diffusion}_p}$. Therefore, by the Transport Equations,

$$J_{\text{diffusion}_p} = -qD_p \frac{\partial \hat{p}}{\partial x}$$

Therefore, substituting,

$$\begin{aligned} 0 &= -\frac{1}{q} \frac{\partial J_{\text{diffusion}_p}}{\partial x} + \left(G_{\text{optical}} - \frac{\hat{p}}{\tau_p} \right) \\ \therefore D_p \frac{d^2 \hat{p}}{dx^2} - \frac{\hat{p}}{\tau_p} &= -G_{\text{optical}} \\ \therefore \frac{d^2 \hat{p}}{dx^2} - \frac{\hat{p}}{D_p \tau_p} &= -\frac{G_{\text{optical}}}{D_p} \\ \therefore \frac{d^2 \hat{p}}{dx^2} - \frac{\hat{p}}{L_p^2} &= -\frac{G_{\text{optical}}}{D_p} \end{aligned}$$

Therefore, for the illuminated part,

$$\hat{p}(x) = C e^{-\frac{x}{L_p}} + D e^{\frac{x}{L_p}} + G_{\text{optical}}\tau_p$$

Similarly, for the masked part, as $G_{\text{optical}} = 0$,

$$\hat{p}(x) = A e^{-\frac{x}{L_p}} + B e^{\frac{x}{L_p}}$$

As the end of the masked part, the concentration must be as if the illuminated part is absent. Therefore, the concentration must be zero.

Therefore,

$$\begin{aligned}\hat{p}(L_2) &= 0 \\ \therefore Ae^{-\frac{L_2}{L_p}} + Be^{\frac{L_2}{L_p}} &= 0\end{aligned}$$

As $L_2 \gg L_p$,

$$Ae^{-\frac{L_2}{L_p}} = 0$$

Therefore,

$$0 + Be^{\frac{L_2}{L_0}} = 0$$

Therefore,

$$B = 0$$

At the end of the illuminated part, the concentration must be as if the masked part is absent. Therefore, the concentration must be $G_{\text{optical}}\tau_p$.

Therefore,

$$\begin{aligned}\hat{p}(-L_1) &= G_{\text{optical}}\tau_p \\ \therefore Ce^{\frac{L_1}{L_p}} + De^{-\frac{L_1}{L_p}} + G_{\text{optical}}\tau_p &= G_{\text{optical}}\tau_p\end{aligned}$$

As $L_1 \gg L_p$,

$$De^{-\frac{L_1}{L_p}} = 0$$

Therefore,

$$Ce^{\frac{L_1}{L_p}} + 0 + G_{\text{optical}}\tau_p = G_{\text{optical}}\tau_p$$

Therefore,

$$C = 0$$

At the junction of the two parts, the concentration must be continuous.

Therefore,

$$\begin{aligned}\hat{p}(0^+) &= \hat{p}(0^-) \\ \therefore Ae^{-\frac{0}{L_p}} &= De^{\frac{0}{L_p}} + G_{\text{optical}}\tau_p \\ \therefore A &= D + G_{\text{optical}}\tau_p\end{aligned}$$

At the junction of the two parts, the diffusion current is continuous. Hence, the derivative of the concentration must be continuous.

Therefore,

$$\begin{aligned}\hat{p}'(0^+) &= \hat{p}'(0^-) \\ \therefore -\frac{Ae^{-\frac{0}{L_p}}}{L_p} &= \frac{De^{\frac{0}{L_p}}}{L_p} \\ \therefore -\frac{A}{L_p} &= \frac{D}{L_p} \\ \therefore A &= -D\end{aligned}$$

Therefore,

$$\begin{aligned}A &= D + G_{\text{optical}}\tau_p \\ \therefore A &= -A + G_{\text{optical}}\tau_p\end{aligned}$$

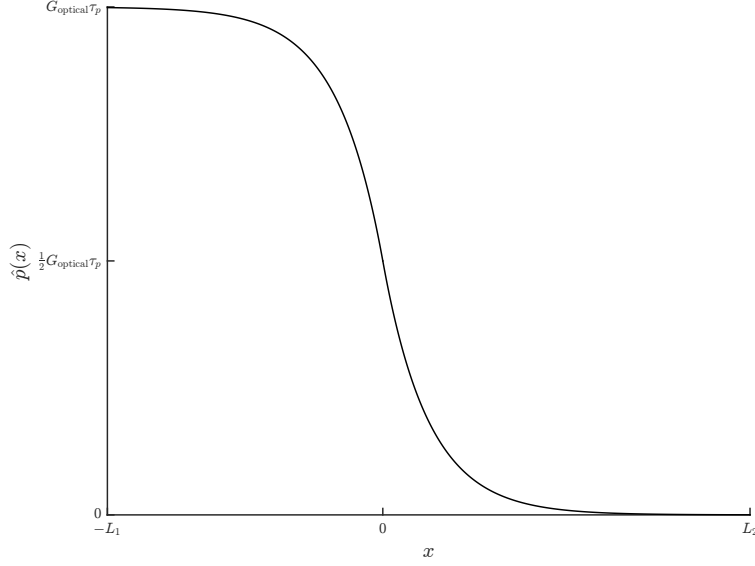
Therefore,

$$\begin{aligned}\therefore A &= \frac{1}{2}G_{\text{optical}}\tau_p \\ \therefore D &= -\frac{1}{2}G_{\text{optical}}\tau_p\end{aligned}$$

Therefore,

$$\hat{p}(x) = \begin{cases} G_{\text{optical}}\tau_p \left(1 - \frac{1}{2}e^{\frac{x}{L_p}}\right) \\ -\frac{1}{2}G_{\text{optical}}\tau_p e^{-\frac{x}{L_p}} + G_{\text{optical}}\tau_p \end{cases} \quad ; \quad x > 0$$

Therefore, the graph is



Consider a N-type semiconductor sample which is masked everywhere, except for a narrow slit at $x = 0$. A pulse of light is shined at $t = 0$.

23 Haynes-Shockley Experiment

A N-type semiconductor sample, of length L , is masked everywhere, except for a narrow slit at $x = 0$. A pulse of light is shined at $t = 0$. If there is no electric field acting on the semiconductor, the concentration of holes is as in Figure 42. If there is an electric field \vec{E} acting on the semiconductor, acting towards the right, the concentration of holes is as in Figure 43.

If recombination occurs, the number of available holes reduces. Therefore, the area under the curve of the excess carrier concentration reduces.

Therefore, the excess carrier concentration is given by

$$\hat{p}(x, t) = \left(\frac{\hat{p}(x, 0)}{2\sqrt{\pi D_p t}} e^{-\frac{x^2}{4D_p t}} \right) e^{-\frac{t}{\tau_p}}$$

Therefore, the effect of the pulse can be felt at the end, such that

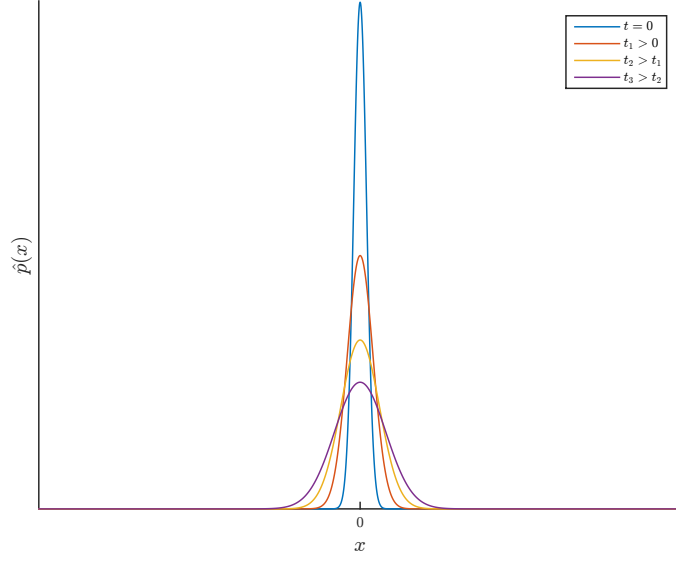


Figure 42: Concentration of Holes Generated due to a Pulse of Light, without Electric Field

$$\begin{aligned}
 V_{\text{drift}} &= \frac{L}{t_d} \\
 \mu_p &= \frac{V_{\text{drift}}}{E} \\
 &= \frac{L}{Et_d}
 \end{aligned}$$

Therefore, the properties of minority carriers can be experimentally measured, using this technique.

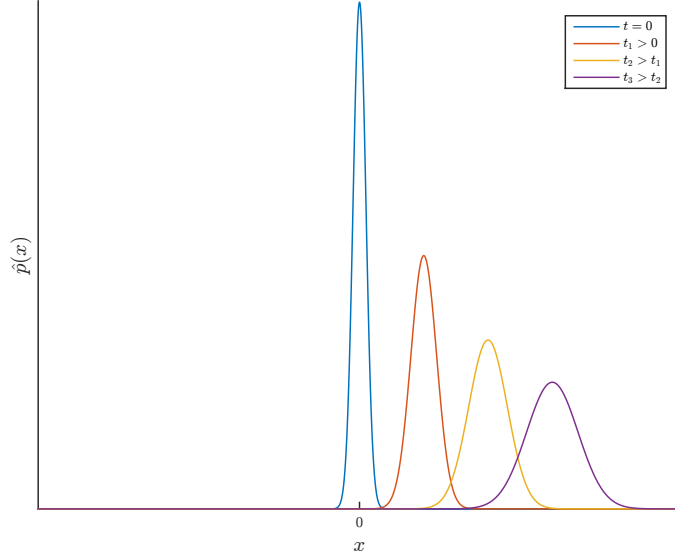


Figure 43: Concentration of Holes Generated due to a Pulse of Light, with Electric Field

24 Quasi-neutrality

24.1 Carrier Concentration Gradient due to Illumination Only

Consider a N-type semiconductor, masked everywhere except for a thin slit in the middle.

By Einstein's Relation,

$$\frac{D_n}{\mu_n} = \frac{kT}{q}$$

$$\frac{D_p}{\mu_p} = \frac{kT}{q}$$

As electrons are lighter than holes, their mobility is higher. Therefore, $\mu_n > \mu_p$, i.e. electrons diffuse faster than holes. Therefore, the excess carrier concentration profiles will be as in Figure 44. However, as electrons and holes have opposite charges, the separation between them creates an internal attractive electric field between them. Therefore, the real excess carrier concentration profiles will be identical, as in Figure 44. This behaviour is called quasi neutrality.

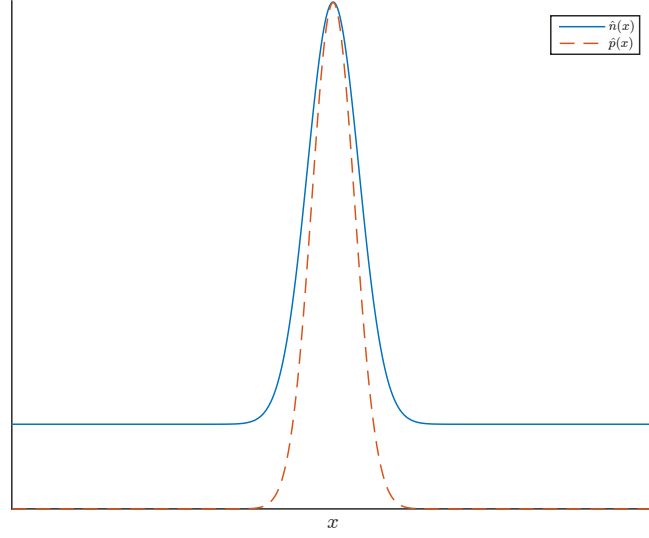


Figure 44: Expected Excess Carrier Concentration in N-type Sample with Constant Illumination Through Narrow Slit Only

In general, if there is a relatively small internal electric field, the excess electrons and holes are forced to drift and diffuse together, i.e., if $\vec{E}_{\text{internal}} \approx 0$,

$$\hat{n}(x) = \hat{p}(x)$$

Such a semiconductor is said to be quasi neutral. In such a case, the transport is dominated by minority carriers, and the excess majority carriers follow the minority.

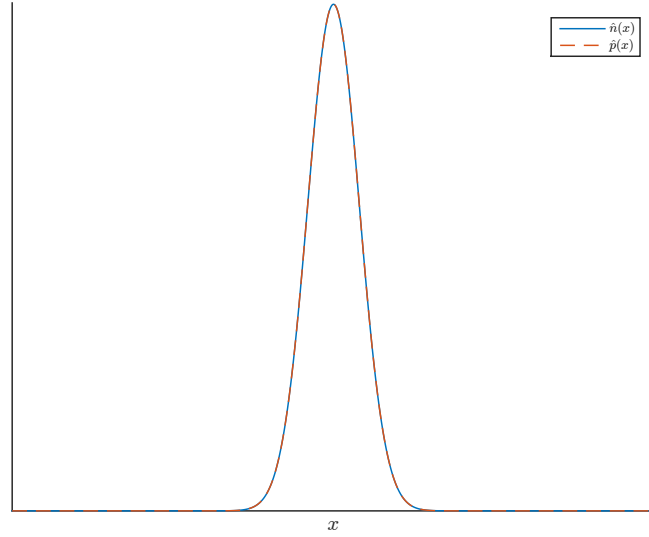


Figure 45: Actual Excess Carrier Concentration in N-type Sample with Constant Illumination Through Narrow Slit Only

Exercise 30.

Consider a P-type semiconductor, masked everywhere except for a thin slit in the middle.

Given $D_p \gg D_n$, what is the direction of the built-in electric fields?

Solution 30.

As $D_p \gg D_n$,

$$\mu_p \gg \mu_n$$

Therefore, as the minority carriers dominate the transport, the holes move according to the electrons. Therefore, the electric field is directed inwards, towards the slit.

24.2 Carrier Concentration Gradient due to Doping Variation Only

Consider a N-type semiconductor sample with concentration gradient

$$n(x) = N_d(x)$$

Let the doping profile be linearly graded, i.e. let $N_d(x)$ grow linearly with x . Let

$$\begin{aligned} N_d(x=0) &= 10^{15} \frac{1}{\text{cm}^3} \\ N_d(x=L) &= 10^{16} \frac{1}{\text{cm}^3} \end{aligned}$$

Therefore, the electrons tend to move towards the left. However, the excess fixed positive charges on the right create an attractive force on the electrons, towards the right. Therefore, the internal electric field, which opposes diffusion is directed towards the left.

24.3 Calculation of Internal Electric Field

Consider a N-type, quasi neutral sample in which the total current density is zero. Therefore,

$$\begin{aligned} 0 &= J_n \\ &= q\mu_n n(x) \vec{E}(x) + qD_n \frac{dn(x)}{dx} \\ \therefore \vec{E}(x) &= -\frac{D_n}{\mu_n} \frac{1}{n(x)} \frac{dn(x)}{dx} \\ &= -\frac{kT}{q} \frac{1}{n(x)} \frac{dn(x)}{dx} \end{aligned}$$

Therefore, as the sample is quasi neutral, $n(x) = N_d(x)$. Therefore,

$$\vec{E}(x) = -\frac{kT}{q} \frac{1}{N_d(x)} \frac{dN_d(x)}{dx}$$

Therefore, the internal voltage drop across the internal field region is

$$\begin{aligned} V &= -\int E dx \\ &= -\frac{kT}{q} \int_{x_1}^{x_2} \frac{1}{N_d(x)} \frac{dN_d(x)}{dx} dx \end{aligned}$$

Therefore, solving,

$$|V| = \frac{kT}{q} \ln \left(\frac{N_d(x_2)}{N_d(x_1)} \right)$$

This built-in voltage is denoted as V_{BI} .

Therefore the intrinsic energy of the sample is given by

$$\begin{aligned}\vec{E} &= - \frac{dV(x)}{dx} \\ &= \frac{1}{q} \frac{dE_i}{dx}\end{aligned}$$

Exercise 31.

The doping profile in a N-type sample is given to be

$$N_d(x) = N_0 e^{-\alpha x}$$

Assuming that the sample is at equilibrium, find the internal electric field.

Solution 31.

As the sample is at equilibrium and is not illuminated, the sample is quasi neutral. Therefore,

$$\begin{aligned}E(x) &= - \frac{kT}{q} \frac{1}{N_d(x)} \frac{dN_d(x)}{dx} \\ &= - \frac{kT}{q} \frac{1}{N_0 e^{-\alpha x}} \frac{d}{dx} N_0 e^{-\alpha x} \\ &= \frac{kT\alpha}{q}\end{aligned}$$