

QUANTUM AND SOLID STATE PHYSICS : ASSIGNMENT

11

AAKASH JOG
ID : 989323563

Exercise 1.

An N-type silicon sample, from $x = -3L$ to $x = 3L$, is illuminated at steady state, from $x = -L$ to $x = L$. The carrier generation rate for $-L < x < L$ is G_{optical} , and is zero outside this window. There are ohmic contacts at the ends of the sample, at $x = -3L$ and $x = 3L$. Assume $L \gg L_p$, low level injection, and no electric field. Set up all the equations and known conditions we would need for solving for $\hat{p}(x)$, across the sample. This should include the differential equations and general solutions in each region, and also continuity and symmetry considerations in order to solve the problem. Hint: Your general solutions, after plugging in boundary conditions, should contain 4 unknowns in total. Draw an approximate plot of $\hat{p}(x)$.

Solution 1.

By the steady state diffusion equation, on the illuminated part,

$$\begin{aligned} 0 &= \frac{\partial \hat{p}}{\partial x} \\ &= -\frac{1}{q} \frac{\partial J_p}{\partial x} + \left(G_{\text{optical}} - \frac{\hat{p}}{\tau_p} \right) \end{aligned}$$

As $\vec{E} = 0$, $J = J_{\text{diffusion}}$. Therefore, by the transport equations,

$$J_{\text{diffusion}_p} = -qD_p \frac{\partial \hat{p}}{\partial x}$$

Therefore,

$$\begin{aligned} 0 &= -\frac{1}{q} \frac{\partial J_{\text{diffusion}_p}}{\partial x} + \left(G_{\text{optical}} - \frac{\hat{p}}{\tau_p} \right) \\ \therefore D_p \frac{d^2 \hat{p}}{dx^2} - \frac{\hat{p}}{\tau_p} &= -G_{\text{optical}} \\ \therefore \frac{d^2 \hat{p}}{dx^2} - \frac{\hat{p}}{D_p \tau_p} &= -\frac{G_{\text{optical}}}{D_p} \\ \therefore \frac{d^2 \hat{p}}{dx^2} - \frac{\hat{p}}{L_p^2} &= -\frac{G_{\text{optical}}}{D_p} \end{aligned}$$

Therefore, for the illuminated part,

$$\hat{p}(x) = Ce^{-\frac{x}{L_p}} + De^{\frac{x}{L_p}} + G_{\text{optical}} \tau_p$$

Date: Thursday 7th January, 2016.

Therefore, for the left non-illuminated part,

$$\hat{p}(x) = Ae^{-\frac{x}{L_p}} + Be^{\frac{x}{L_p}}$$

Therefore, for the right non-illuminated part,

$$\hat{p}(x) = Pe^{-\frac{x}{L_p}} + Qe^{\frac{x}{L_p}}$$

As the sample has ohmic contacts at the ends, the concentration at the ends must be zero. Therefore,

$$\hat{p}(-3L) = 0$$

$$\therefore Ae^{\frac{3L}{L_p}} + Be^{-\frac{3L}{L_p}} = 0$$

$$\hat{p}(3L) = 0$$

$$\therefore Pe^{-\frac{3L}{L_p}} + Qe^{\frac{3L}{L_p}} = 0$$

As $L \gg L_p$,

$$Be^{-\frac{3L}{L_p}} = 0$$

$$Pe^{-\frac{3L}{L_p}} = 0$$

Therefore,

$$Ae^{\frac{3L}{L_p}} = 0$$

$$Qe^{\frac{3L}{L_p}} = 0$$

Therefore,

$$A = 0$$

$$Q = 0$$

Therefore, for the left non-illuminated part,

$$\hat{p}(x) = Be^{\frac{x}{L_p}}$$

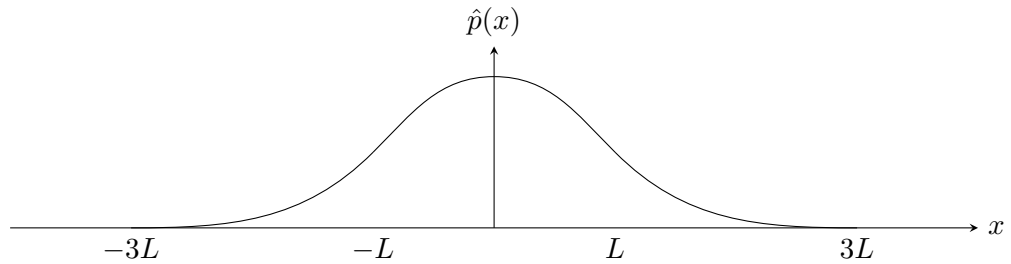
Therefore, for the right non-illuminated part,

$$\hat{p}(x) = Pe^{-\frac{x}{L_p}}$$

As the concentration profile must be continuous at $x = -L$ and $x = L$,

$$Be^{-\frac{L}{L_p}} = Ce^{\frac{L}{L_p}} + De^{-\frac{L}{L_p}} + G_{\text{optical}}\tau_p$$

$$Pe^{-\frac{L}{L_p}} = Ce^{-\frac{L}{L_p}} + De^{\frac{L}{L_p}} + G_{\text{optical}}\tau_p$$



Exercise 2.

Consider a sample of N-type silicon at room temperature, uniformly doped with

$$N_D = 10^{17} \frac{1}{\text{cm}^3}$$

Light is shining uniformly on the sample, at steady state, with an EHP optical generation rate of

$$G_{\text{optical}} = 10^{20} \frac{1}{\text{cm}^3 \text{s}}$$

The minority carrier lifetime is

$$\tau_p = 10^{-5} \text{s}$$

The hole mobility is

$$\mu_p = 350 \frac{\text{cm}^2}{\text{V s}}$$

- (1) What is the steady state hole concentration p across the sample?
- (2) Now suppose the following situation.

The semiconductor extends very far in the positive and negative x directions, and there are no metallic contacts at the ends. At $x = 0$, there is an infinitesimally thin layer with very low lifetime that forces the excess minority carrier concentration to go to 0 at $x = 0$, i.e.,

$$\hat{p}(x = 0) = 0$$

Sketch the form of the expected solution for $\hat{p}(x)$ across the sample. What is the value of \hat{p} as $x \rightarrow \pm\infty$? Label it on your drawing.

- (3) Approximately how far away from $x = 0$ do we need to go so that \hat{p} is approximately its value at $x = \pm\infty$? Explain with words and numerically. Hint: You should know the behaviour of $\hat{p}(x)$. So think about what approximate value represents the distance at which \hat{p} goes to at $\pm\infty$.

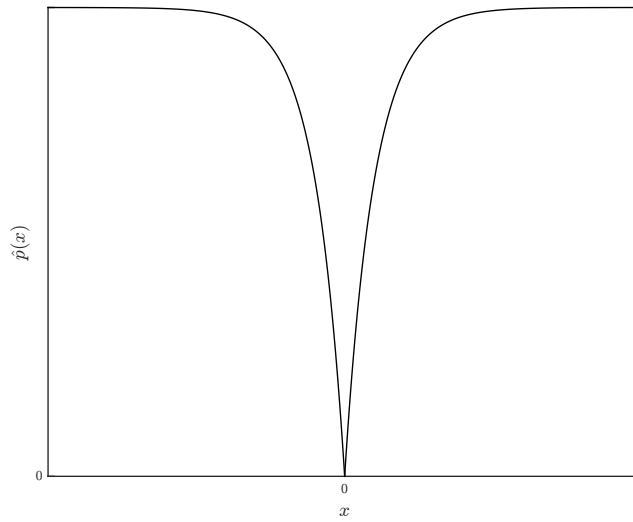
Solution 2.

- (1) As the sample is in steady state,

$$\begin{aligned} \hat{p}(x) &= G_{\text{optical}} \tau_p \\ &= \left(10^{20} \frac{1}{\text{cm}^3 \text{s}} \right) (10^{-5} \text{s}) \\ &= 10^{15} \frac{1}{\text{cm}^3} \end{aligned}$$

Therefore,

$$\begin{aligned}
 p(x) &= p_0 + \hat{p}(x) \\
 &= \frac{n_i^2}{N_d} + \hat{p}(x) \\
 &= \frac{10^{20}}{10^{17}} + 10^{15} \\
 &= 10^3 + 10^{15} \\
 &\approx 10^{15} \frac{1}{\text{cm}^3}
 \end{aligned}$$



- (2) For $x \rightarrow \pm\infty$, the carrier concentration is not affected by the thin slice at $x = 0$. Therefore, the concentration is as if this slice does not exist. Therefore,

$$\begin{aligned}
 \lim_{x \rightarrow \pm\infty} \hat{p}(x) &= G_{\text{optical}} \tau_p \\
 &= \left(10^{20} \frac{1}{\text{cm}^3 \text{s}} \right) (10^{-5} \text{s}) \\
 &= 10^{15} \frac{1}{\text{cm}^3}
 \end{aligned}$$

- (3) The carrier concentration profile $\hat{p}(x)$ for $x > 0$ is of the form

$$\hat{p}(x) = A \left(1 - e^{-\frac{x}{L_p}} \right)$$

and for $x < 0$ is of the form

$$\hat{p}(x) = A \left(1 - e^{\frac{x}{L_p}} \right)$$

Therefore, the concentration at $x \rightarrow \pm\infty$ can be approximated to the concentration at $x = \pm L_p$, where

$$\begin{aligned}
 L_p &= \sqrt{D_p \tau_p} \\
 &= \sqrt{\frac{kT}{q} \mu_p \tau_p} \\
 &= \sqrt{\frac{kT}{q} \left(350 \frac{\text{cm}^2}{\text{V s}} \right) (10^{-5} \text{s})} \\
 &= \sqrt{(0.026 \text{V}) \left(350 \frac{\text{cm}^2}{\text{V s}} \right) (10^{-5} \text{s})} \\
 &= \sqrt{91 \times 10^{-6} \text{cm}^2} \\
 &= 9.54 \times 10^{-3} \text{cm}
 \end{aligned}$$

Exercise 3.

- (1) Show that for a ladder of eigenvalues of L_z , for a given l ,

$$\hat{L}^2 Y_{lm} = \hbar^2 l(l+1) Y_{lm}$$

the lowest step in the ladder corresponds to

$$m_{\text{minimum}} = -l$$

Remember that

$$\hat{L}_z Y_{lm} = \hbar m Y_{lm}$$

- (2) We saw in recitation the following expression.

$$\hat{L}_- Y_{lm} = B_{lm} Y_{l, m-1}$$

Find the coefficient B_{lm} as a function of l and m .

- (3) Explain the following commutation relations, for the case of a central potential. You do not need to calculate anything.

(a) $[\hat{L}^2, \hat{H}] = 0$

(b) $[\hat{L}_z, \hat{H}] = 0$

Solution 3.

- (1) For the lowest step in the ladder,

$$\hat{L}_- \hat{L}_z Y_{lm} = 0$$

Therefore,

$$\begin{aligned}
 \hat{L}_+ 0 &= \hat{L}_+ \hat{L}_- \hat{L}_z Y_{lm} \\
 \therefore 0 &= \hat{L}_+ \hat{L}_- Y_{lm} \\
 &= \left(\hat{L}^2 - \hat{L}_z^2 + \hbar \hat{L}_z \right) Y_{lm} \\
 &= \left(\hbar^2 l(l+1) - \hbar^2 m + \hbar^2 m \right) Y_{lm}
 \end{aligned}$$

Therefore,

$$l(l+1) = m(m-1)$$

$$\therefore m = -l$$

(2)

$$\hat{L}_- Y_{lm} = B_{lm} Y_{l(m-1)}$$

Let

$$f = \hat{L}_i Y_{lm}$$

$$g = Y_{lm}$$

Therefore, the identity

$$\int_{-\infty}^{\infty} f^*(x) (\hat{a}_{\pm} g(x)) dx = \int_{-\infty}^{\infty} (\hat{a}_{\mp} f(x))^* g(x) dx$$

implies

$$\begin{aligned} & \iint \left(\hat{L}_+ \hat{L}_- Y_{lm} \right)^* Y_{lm} \sin \theta d\theta d\varphi \\ &= \iint \left(\hbar^2 l(l+1) - \hbar^2 m^2 + \hbar^2 m^2 \right) Y_{lm}^* Y_{lm} \sin \theta d\theta d\varphi \\ &= \hbar^2 (l(l+1) - m(m-1)) \end{aligned}$$

Also,

$$\iint \left(\hat{L}_- Y_{lm} \right)^* \left(\hat{L}_- Y_{lm} \right) \sin \theta d\theta d\varphi = |B_{lm}|^2$$

Therefore,

$$|B_{lm}|^2 = \hbar^2 (l(l+1) - m(m-1))$$

$$\therefore B_{lm} = \hbar \sqrt{l(l+1) - m(m-1)}$$

(3) (a)

$$[\hat{L}^2, \hat{H}] = 0$$

As this commutation relation is zero, the two operators have common eigenfunctions.

(b)

$$[\hat{L}_z, \hat{H}] = 0$$

As this commutation relation is zero, the two operators have common eigenfunctions.

Exercise 4.

Consider a particle with the following wave function.

$$\Psi(r, \theta, \varphi) = R(r)(2Y_{00} + Y_{11} + 3Y_{10} + 2iY_{1(-1)})$$

where $R(r)$ is a radial function and Y_{lm} are the eigenfunctions of \hat{L}^2 and \hat{L}_z .

- (1) If \hat{L}^2 is measured for this state, what values can it have, and with what probabilities?
- (2) Find the expectation value of \hat{L}^2 for this state.
- (3) If \hat{L}_z is measured for this state, what values can it have, and with what probabilities?
- (4) Find the expectation value of \hat{L}_z for this state.

Solution 4.

(1)

$$\hat{L}^2 Y_{lm} = l(l+1)\hbar^2 Y_{lm}$$

Therefore, the possible values of L^2 are the values of $l(l+1)\hbar^2$ for $l = 0$ and $l = 1$. Therefore, the possible values of L^2 are 0 and $2\hbar^2$.

Let the constant of normalization for Ψ be A .

The probability corresponding to Y_{00} is

$$\begin{aligned} \frac{|c_1|^2}{A} &= \frac{2^2}{A} \\ &= \frac{4}{A} \end{aligned}$$

The probability corresponding to Y_{11} is

$$\begin{aligned} \frac{|c_2|^2}{A} &= \frac{1^2}{A} \\ &= \frac{1}{A} \end{aligned}$$

The probability corresponding to Y_{10} is

$$\begin{aligned} \frac{|c_3|^2}{A} &= \frac{3^2}{A} \\ &= \frac{9}{A} \end{aligned}$$

The probability corresponding to $Y_{1(-1)}$ is

$$\begin{aligned} \frac{|c_4|^2}{A} &= \frac{|2i|^2}{A} \\ &= \frac{4}{A} \end{aligned}$$

Normalizing,

$$\begin{aligned}
 1 &= \sum c_k \\
 &= \frac{4 + 1 + 9 + 4}{A} \\
 &= \frac{18}{A} \\
 \therefore A &= 18
 \end{aligned}$$

Therefore,

$$\begin{aligned}
 P_1 &= \frac{2}{9} \\
 P_2 &= \frac{1}{18} \\
 P_3 &= \frac{1}{2} \\
 P_4 &= \frac{2}{9}
 \end{aligned}$$

(2)

$$\begin{aligned}
 \langle L^2 \rangle &= \sum P_k L_k^2 \\
 &= \frac{2}{9} 0 + \frac{1}{18} 2\hbar + \frac{1}{2} 2\hbar + \frac{2}{9} 2\hbar \\
 &= \frac{7}{9} 2\hbar \\
 &= \frac{14\hbar}{9}
 \end{aligned}$$

(3)

$$\hat{L}_z Y_{lm} = m\hbar Y_{lm}$$

Therefore, the possible values of L_z are the values of $m\hbar$ for $m = 0$, $m = 1$, and $m = -1$. Therefore, the possible values of L_z are 0 , \hbar , and $-\hbar$.

Let the constant of normalization for Ψ be A .

The probability corresponding to Y_{00} is

$$\begin{aligned}
 \frac{|c_1|^2}{A} &= \frac{2^2}{A} \\
 &= \frac{4}{A}
 \end{aligned}$$

The probability corresponding to Y_{11} is

$$\begin{aligned}
 \frac{|c_2|^2}{A} &= \frac{1^2}{A} \\
 &= \frac{1}{A}
 \end{aligned}$$

The probability corresponding to Y_{10} is

$$\begin{aligned}\frac{|c_3|^2}{A} &= \frac{3^2}{A} \\ &= \frac{9}{A}\end{aligned}$$

The probability corresponding to $Y_{1(-1)}$ is

$$\begin{aligned}\frac{|c_4|^2}{A} &= \frac{|2i|^2}{A} \\ &= \frac{4}{A}\end{aligned}$$

Normalizing,

$$\begin{aligned}1 &= \sum c_k \\ &= \frac{4 + 1 + 9 + 4}{A} \\ &= \frac{18}{A} \\ \therefore A &= 18\end{aligned}$$

Therefore,

$$\begin{aligned}P_1 &= \frac{2}{9} \\ P_2 &= \frac{1}{18} \\ P_3 &= \frac{1}{2} \\ P_4 &= \frac{2}{9}\end{aligned}$$

(4)

$$\begin{aligned}\langle L_z \rangle &= \sum P_k L_{zk} \\ &= \frac{2}{9}0 + \frac{1}{18}\hbar + \frac{1}{2}0 + \frac{2}{9}(-\hbar) \\ &= \frac{\hbar}{18} - \frac{4\hbar}{18} \\ &= -\frac{3\hbar}{18}\end{aligned}$$