

Mécanica Cuántica Ejercicios

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Question 1: Schwabl 2.3

Using the Bohr–Sommerfeld quantization rules, determine the energy eigen- states of a particle of mass m moving in an infinitely high potential well:

$$V(x) = \begin{cases} 0 & 0 \leq x \leq a \\ \infty & \text{otherwise} \end{cases}$$

Respuesta. Fuera de la barrera de potencial la función de onda es cero, dentro el potencial vale 0, cumpliendo la ecuación diferencial

$$-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} = E\psi$$

Escribiéndola de manera alterna tenemos

$$\frac{d^2\psi}{dx^2} = -k^2\psi, \text{ donde } k \equiv \frac{\sqrt{2mE}}{\hbar}$$

Así identificamos la solución para un oscilador armónica

$$\psi(x) = A \sin kx + B \cos kx$$

y aplicando las siguientes condiciones iniciales

$$\psi(0) = \psi(a) = 0.$$

Encontramos:

$$\psi(0) = A \sin 0 + B \cos 0 = B = 0.$$

$$\psi(x) = A \sin kx$$

$$\psi(a) = A \sin ka = 0 \rightarrow \sin ka = 0$$

$$ka = n\pi$$

$$k = \frac{n\pi}{a}$$

Sabiendo que el exponente se comporta indicialmente, la energía también tiene este comportamiento y está asociada de la siguiente manera:

$$E_n = \frac{\hbar^2 k_n^2}{2m} = \frac{n^2 \pi^2 \hbar^2}{2ma^2}$$

Question 2: Schwabl 2.4

1. $p^\dagger = p$, como p es un operador, debe cumplir: $\langle \psi | p | \phi \rangle = \langle \phi | p | \psi \rangle^*$

$$\begin{aligned}\langle \psi | p | \phi \rangle &= \int \phi^* (x\psi) dx \\ \langle \phi | p | \psi \rangle^* &= \int (\phi p \psi^*)^* dx \\ &= \int \left(\phi - i\hbar \frac{\partial \psi^*}{\partial x} \right)^* dx \\ &= i\hbar \int \phi^* \frac{\partial \psi}{\partial x} dx\end{aligned}$$

Integrando por partes :

$$\begin{aligned}&= [\phi^* \psi]_{-\infty}^{\infty} - i\hbar \int \psi^* \frac{\partial \phi}{\partial x} dx \\ &= \int \phi^* \left(-i\hbar \frac{\partial}{\partial x} \right) \psi dx \\ &= \int \phi^* p \psi dx = \langle \psi | p | \phi \rangle \quad \ominus\end{aligned}$$

2. $(AB)^\dagger = B^\dagger A^\dagger$. Desarrollando el producto:

$$A = \begin{bmatrix} a_{11} & a_{12} & \dots & a_{1n} \\ a_{21} & a_{22} & \dots & a_{2n} \\ \dots & \dots & \dots & \dots \\ a_{n1} & a_{n2} & \dots & a_{nn} \end{bmatrix} \quad B = \begin{bmatrix} b_{11} & b_{12} & \dots & b_{1n} \\ b_{21} & b_{22} & \dots & b_{2n} \\ \dots & \dots & \dots & \dots \\ b_{n1} & b_{n2} & \dots & b_{nn} \end{bmatrix}$$

$$\begin{aligned}A \cdot B &= \begin{bmatrix} a_{11}b_{11} + \dots + a_{1n}b_{n1} & a_{11}b_{12} + \dots + a_{1n}b_{n2} & \dots & a_{11}b_{1n} + \dots + a_{1n}b_{nn} \\ a_{21}b_{11} + \dots + a_{2n}b_{n1} & a_{21}b_{12} + \dots + a_{2n}b_{n2} & \dots & a_{21}b_{1n} + \dots + a_{2n}b_{nn} \\ \dots & \dots & \dots & \dots \\ a_{n1}b_{11} + \dots + a_{nn}b_{n1} & a_{n1}b_{12} + \dots + a_{nn}b_{n2} & \dots & a_{n1}b_{1n} + \dots + a_{nn}b_{nn} \end{bmatrix} \\ (A \cdot B)^\dagger &= \begin{bmatrix} a_{11}b_{11} + \dots + a_{1n}b_{n1} & a_{21}b_{11} + \dots + a_{2n}b_{n1} & \dots & a_{n1}b_{11} + \dots + a_{nn}b_{n1} \\ a_{11}b_{12} + \dots + a_{1n}b_{n2} & a_{21}b_{12} + \dots + a_{2n}b_{n2} & \dots & a_{n1}b_{12} + \dots + a_{nn}b_{n2} \\ \dots & \dots & \dots & \dots \\ a_{11}b_{1n} + \dots + a_{1n}b_{nn} & a_{21}b_{1n} + \dots + a_{2n}b_{nn} & \dots & a_{n1}b_{1n} + \dots + a_{nn}b_{nn} \end{bmatrix}^*\end{aligned}$$