

Patrón de Interferencia

$$\Psi(x,t) = A e^{i(\vec{k} \cdot \vec{x} - \omega t)} \rightarrow E(x,t)$$

$$\vec{E}(x,t) = \vec{E}_1(x,t) + \vec{E}_2(x,t). \text{ Intensidad } I \propto |\vec{E}|^2 = |\vec{E}_1|^2 + |\vec{E}_2|^2 + 2\vec{E}_1 \cdot \vec{E}_2$$

Para Complejos $(E_1 + E_2)(E_1^* + E_2^*) = |E_1|^2 + |E_2|^2 + \underbrace{E_1 E_2^* + E_1^* E_2}_{2 \operatorname{Re}(E_1 E_2^*)}$

Función de onda

Ecuación de Schrödinger hipótesis $P(\vec{x},t) d^3x = |\Psi(\vec{x},t)|^2 d^3x$
 \hookrightarrow Densidad de probabilidad

• Ecuación de movimiento

$$1. \frac{\partial^2 \Psi}{\partial t^2} = \frac{E}{m} \rightarrow x \xrightarrow{\frac{d}{dt}} \dot{x}(t)$$

$$\Psi(x,t_0) \rightarrow \Psi(x,t) \text{ Ecuación Dif. de primer orden para } \Psi$$

$$2. \text{ Lineal } L\Psi = 0 \quad 3. \int |\Psi(x,t)|^2 dx = 1$$

$$4. \Psi(x,t) = A e^{i(kx - \omega t)} = A e^{\frac{i}{\hbar}(\hbar kx - \hbar \omega t)} \quad \frac{\partial \Psi}{\partial t} = L\Psi + c$$

$$\Psi(x,t) = A e^{\frac{i}{\hbar}(\hbar kx - \hbar \omega t)}$$

$$\frac{\partial \Psi}{\partial t} = -\frac{i\omega}{\hbar} \Psi, \quad \nabla \Psi = \frac{i}{\hbar} \hbar k \Psi \rightarrow \nabla^2 \Psi = -\frac{\hbar^2 k^2}{\hbar^2} \Psi$$

$$\frac{\hbar^2}{2m} \nabla^2 \Psi = -\frac{\hbar^2 k^2}{2m} \Psi \rightarrow \frac{\partial \Psi}{\partial t} = -\frac{i\hbar^2 k^2}{2m\hbar} \Psi = -\frac{i\hbar k^2}{2m} \Psi = -\frac{iE}{\hbar} \Psi$$

$$\hbar i \frac{\partial \Psi}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \Psi \quad \text{Ec. de Schrödinger partícula libre.}$$

Para una solución general $\rightarrow \Psi(x,t) = \sum_{k=1}^n A(k) e^{i(\frac{\hbar k}{m}x - \frac{\hbar^2 k^2}{2m}t)}$
 \hookrightarrow Función de peso

Paso al continuo

$$\Psi(x,t) = \int \frac{d^3p}{(2\pi\hbar)^3} \varphi(p) e^{\frac{i}{\hbar}(p \cdot x - \frac{p^2}{2m}t)}$$

Función de peso ~ Dist. Gaussiana

$$\varphi(p) = A e^{-\frac{(\vec{p} - \vec{p}_0)^2}{2\hbar^2}}$$

$$\Psi(x,t) = \int \frac{d^3p}{(2\pi\hbar)^3} A e^{\frac{i}{\hbar}(p \cdot x - \frac{p^2}{2m}t) - \frac{i}{\hbar}(p_0^2 - 2\vec{p} \cdot \vec{p}_0) \frac{t}{2m}} \Rightarrow 1 \text{ dimensión } \int \frac{dp}{2\pi\hbar} A e^{(-a p^2 + b p + c)}$$

$$-a p^2 + b p + c = -\left(\frac{p^2}{\hbar^2} - \frac{p_0^2}{2m\hbar}\right) \frac{t}{2m} + \left(\frac{p_0}{\hbar} + \frac{x}{2\hbar}\right) p - \left(\frac{p_0^2}{\hbar^2}\right) \frac{t}{2m}$$

$$\tilde{F}(w) = \int \frac{dt}{\sqrt{2\pi}} e^{-iwt} f(t)$$

$$\Psi(x,t) = \int \frac{dp}{2\pi\hbar} A e^{\frac{i}{\hbar}(p \cdot x - \frac{p^2}{2m}t) - \frac{i}{\hbar}(p_0^2 - 2\vec{p} \cdot \vec{p}_0) \frac{t}{2m}} \Rightarrow \int \frac{dp}{2\pi\hbar} A e^{(-a(p - \frac{b}{2a})^2 + \frac{b^2}{4a} - c)}$$

$$|\Psi(x,t)|^2 = \frac{A^2 \pi}{(2\pi\hbar)^2} e^{\frac{b^2 - 4ac}{4a}} = \frac{A^2 \pi}{(2\pi\hbar)^2} e^{\frac{-(x-vt)^2}{4(\sigma^2 + \Delta^2)}} = e^{\frac{-(x-vt)^2}{2(\sigma^2 + \Delta^2)}}$$

$$\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \frac{d^3p}{(2\pi\hbar)^3} \frac{d^3p'}{(2\pi\hbar)^3} \varphi(p) \varphi(p') e^{\frac{i}{\hbar}(p \cdot x - \frac{p^2}{2m}t) - \frac{i}{\hbar}(p_0^2 - 2\vec{p} \cdot \vec{p}_0) \frac{t}{2m}} e^{\frac{i}{\hbar}(p' \cdot x - \frac{p'^2}{2m}t) - \frac{i}{\hbar}(p_0^2 - 2\vec{p}' \cdot \vec{p}_0) \frac{t}{2m}}$$

$$|\Psi(x,t)|^2 = \frac{A^2 \pi}{(2\pi\hbar)^2} e^{\frac{-(x-vt)^2}{2(\sigma^2 + \Delta^2)}}$$

$$2Re\left(\frac{b^2 - 4ac}{4a}\right) = 2Re\left(\frac{b^2 - 4ac}{4a}\right)$$

$$\int dx |\Psi(x,t)|^2 = 1 = \frac{A^2}{4\pi\hbar^2} \int_{-\infty}^{\infty} e^{\frac{-(x-vt)^2}{2(\sigma^2 + \Delta^2)}} dx$$

$$= \frac{A^2}{4\pi\hbar^2} \int_{-\infty}^{\infty} e^{-\frac{(x-vt)^2}{2(\sigma^2 + \Delta^2)}} dx = \frac{A^2}{4\pi\hbar^2} \sqrt{2\pi(\sigma^2 + \Delta^2)} = \frac{A^2}{4\pi\hbar^2} \sqrt{2\pi} \sqrt{\sigma^2 + \Delta^2}$$

$$|a| = \frac{1}{4\pi\hbar^2} \sqrt{\frac{\pi}{2(\sigma^2 + \Delta^2)}} = \frac{1}{4\pi\hbar^2} \sqrt{\frac{\pi}{2}} \frac{1}{\sqrt{\sigma^2 + \Delta^2}}$$

$$\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \frac{d^3p}{(2\pi\hbar)^3} \frac{d^3p'}{(2\pi\hbar)^3} \varphi(p) \varphi(p') e^{\frac{i}{\hbar}(p \cdot x - \frac{p^2}{2m}t) - \frac{i}{\hbar}(p_0^2 - 2\vec{p} \cdot \vec{p}_0) \frac{t}{2m}} e^{\frac{i}{\hbar}(p' \cdot x - \frac{p'^2}{2m}t) - \frac{i}{\hbar}(p_0^2 - 2\vec{p}' \cdot \vec{p}_0) \frac{t}{2m}}$$

$$= \frac{A^2}{4\pi\hbar^2} \sqrt{\frac{\pi}{2(\sigma^2 + \Delta^2)}} = \frac{A^2}{4\pi\hbar^2} \sqrt{\frac{\pi}{2}} \frac{1}{\sqrt{\sigma^2 + \Delta^2}}$$

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$$\frac{A^2}{4\pi\hbar^2} \sqrt{\frac{\pi}{2(\sigma^2 + \Delta^2)}} = \frac{A^2}{4\pi\hbar^2} \sqrt{\frac{\pi}{2}} \frac{1}{\sqrt{\sigma^2 + \Delta^2}}$$

Condición de normalización

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2:25 p. m.

$$\begin{aligned} \int |\Psi(x,t)|^2 dx &= \int \left| \frac{A}{2\pi\hbar^2} e^{\frac{i(x-x_0)^2}{2d^2(1+\Delta^2)}} \right|^2 dx = \frac{A^2}{4\pi^2\hbar^2} \cdot \sqrt{\frac{2\pi}{d^2(1+\Delta^2)}} = A^2 \sqrt{\frac{2\pi d^2 (1+\Delta^2)}{16\pi^2 d^2 (1+\Delta^2)}} = A^2 \sqrt{\frac{1}{8\pi d^2}} = 1 \\ 4\pi\hbar^2 \sqrt{\frac{d^2}{\hbar^4} + \frac{1}{4\hbar^2 m^2}} &= 4\pi\hbar^2 \sqrt{\frac{4d^2\hbar^2 m^2 + 1}{4\hbar^6 m^2}} \\ &= \frac{4\pi^2 \hbar^2}{2\hbar^2 m} = \frac{2\pi}{\hbar m} \sqrt{4d^2\hbar^2 m^2 + 1} \\ &= \frac{4\pi^2 d^2 \cancel{\hbar^2} \sqrt{1+\Delta^2}}{\hbar m} \\ &= 4\pi^2 d^2 \sqrt{1+\Delta^2} \end{aligned}$$

$$A^2 = \sqrt{8\pi} d$$