

Novel Process for Screen-Printed Selective Area Front Polysilicon Contacts for TOPCon Cells Using Laser Oxidation

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Abstract—The efficiency potential of double-side tunnel oxide passivated contact (DS-TOPCon) solar cells is limited by parasitic absorption in the front poly-Si layer, despite excellent passivation and high V_{OC} . The use of patterned poly-Si only under the front metal grid lines can significantly reduce the parasitic absorption loss without sacrificing voltage. In this work, we demonstrate a simple, manufacturing-friendly method of patterning the front poly-Si using a nanosecond UV (355 nm) laser. We found that with laser powers ≥ 3 W at a 400 mm/s scan speed, an estimated 1–4 nm thick stoichiometric SiO_2 layer was grown on TOPCon. This served as a mask for KOH-etching of 200 nm poly-Si, allowing for patterning of poly-Si fingers required for selective TOPCon. While laser powers above 3 W caused substantial deterioration in passivation quality, the resulting damage in J_0 was largely recovered by subsequent PECVD SiN_x deposition. At 3 W, the full area J_0 was found to be $36.8 \text{ fA}\cdot\text{cm}^{-2}$. This translates to $1.68 \text{ fA}\cdot\text{cm}^{-2}$ for 4.48% coverage from the wing area of the polyfinger lines (100 lines–100 μm wide and 30 μm metal) contributing to a total front J_0 of $\sim 10 \text{ fA}\cdot\text{cm}^{-2}$, well suited for 25% efficient solar cells.

Index Terms—Double-side TOPCon, laser oxidation, passivating contacts, screen-printed contacts, selective area contacts.

I. INTRODUCTION

THE vast majority of commercial crystalline silicon solar cells use phosphorus diffused emitters/junctions on the front side that limit the efficiency by the heavily doped regions and metal-induced recombination at the contacts. Carrier-selective passivating contacts can mitigate this loss by displacing the diffused and metalized regions outside the Si wafer

and providing appropriate band bending for collecting carriers while preventing minority carrier recombination at the metal–semiconductor and oxide–semiconductor interfaces. Tunnel oxide passivated contacts (TOPCon) utilize an ultrathin tunneling SiO_2 layer capped with doped poly-Si to provide passivation and carrier selectivity, respectively [1]. TOPCon-based cells have the added benefit of being compatible with traditional fire-through screen-printed metallization, making them compatible with preexisting solar cell production lines. While having full-area poly-Si contacts on both sides of a solar cell can yield V_{OC} as high as 728 mV [2], using thick poly-Si on the front of a solar cell, results in significant current loss due to parasitic light absorption. On the other hand, thinning the front poly-Si increases damage during fire-through metallization and leads to higher J_0 [3]. This has restricted the use of TOPCon only on the rear side of the screen-printed solar cell utilizing a traditionally diffused front junction with record efficiencies around 25% [4].

Using a patterned TOPCon only under the metal grid on the front can significantly reduce parasitic absorption and boost cell efficiency [5]. Moreover, for rear-junction solar cells, efficiency is nearly independent of front sheet resistance, allowing for the use of polyfingers or selective area front TOPCon with a dielectric passivated undiffused Si wafer in the field region. This concept works because for this cell design with appropriate bulk resistivity and lifetime carrier transport on the front side occurs through bulk Si without appreciable resistive loss [6]. This can be accomplished by forming a patterned mask on the full area poly-Si, etching the poly-Si in the field region around the mask to define polyfingers, and stripping the mask before metallization. A schematic of a solar cell based on this design is shown in Fig. 1. This process can be cumbersome and expensive and may reduce yield if done by conventional methods [7]. Alternatively, Chen et al. [8], used a self-aligned metal mask and TMAH etching to create such a device. However, their process involves thermally evaporated metal, which is not suitable for industrial use. In this article, we reveal a simple, manufacturable, clean, and rapid process involving laser oxidation to fabricate selective TOPCon on the front side of a solar cell.

In this work, we demonstrate the use of a pulsed ultraviolet (UV) nanosecond laser to form a patterned oxide to serve as a poly-Si etch mask on the full area TOPCon. We show that at certain power densities of our UV laser it is possible to

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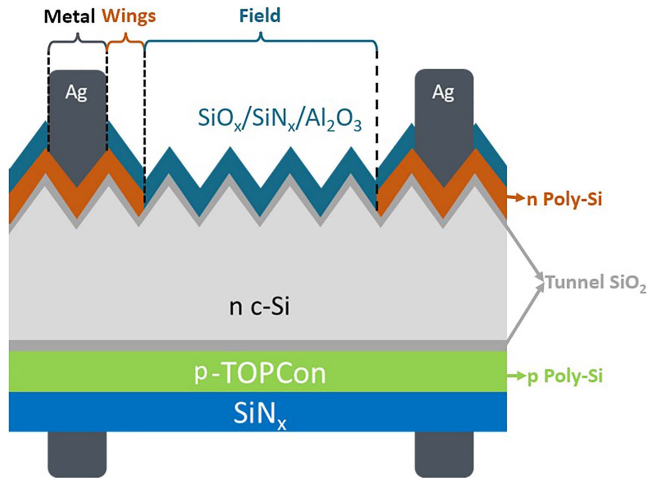


Fig. 1. Schematic of a rear junction selective area double side TOPCon solar cell. The metallization is only performed on selectively etched fingers of polysilicon, enabling excellent light transparency in the field region.

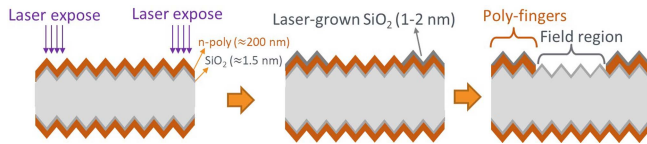


Fig. 2. General process flow for laser-grown oxide mask patterning and subsequent poly-Si etching.

grow a very thin (few nm thick) silicon oxide layer that is sufficient to prevent etching of the underlying poly-Si during the poly-Si etch between the finger pattern. Schäfer et al. [9] used a picosecond laser to achieve similar etch-masking behavior. However, due to significantly lower power densities involved, their etch resistance, while assisted by oxygen, was attributed to amorphization of the surface of the poly-Si [10].

II. FABRICATION AND CHARACTERIZATION OF TOPCON SYMMETRIC STRUCTURES

Symmetric n-poly-Si/SiO₂/n-Si/SiO₂/n-poly-Si test structures were grown on $\sim 3 \Omega\text{-cm}$ n-type Cz-Si wafers textured on both sides with random pyramids. Tunnel oxide was grown using the NAOS process [11] on both sides of the wafers in HNO₃ at 100 °C for 15 min, resulting in a thickness of about 15 Å. Subsequently, ~ 200 nm n-type poly-Si was grown in a Tystar LPCVD tube at 588 °C. Test samples for sheet resistance measurements were made with the same process on p-type wafers. Layers for optical measurements were grown on polished CZ wafers with 150 nm of PECVD SiO_x for better ellipsometry.

A. UV Laser Patterning

Fig. 2 shows a schematic of the laser patterning process. First, the poly-Si coated wafer was patterned to define grid lines with laser irradiation from a Coherent Avia 355-10 nanosecond pulsed UV laser (355 nm) as a function of laser power controlled

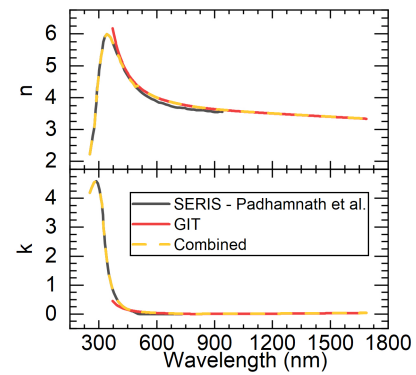


Fig. 3. Combined refractive index (n) and extinction coefficient (k) obtained from measurements at Georgia Tech (this work) and by Padhamnath et al. [14].

by the pulse repetition frequency. All laser scans were performed at 400 mm/s with a spot size of 70 μm . After patterning, the wafers were immersed in 9%_{vol} KOH at 40 °C to remove the poly-Si from the field region, resulting in patterned poly-Si fingers. For lifetime test structures, laser patterning was only done on one side of the wafer. Subsequently, ~ 70 nm of SiN_x was deposited at a temperature of 450 °C and at a pressure of 1700 mtorr for 700 s on both sides of the samples in a Centrotherm PECVD reactor to mimic the SiN_x required for the fire-through screen-printed contact process. The objective is to define 100–150 μm wide poly-Si fingers/TOPCon on top of which 30–40 μm screen-printed metal fingers will be formed and aligned by an industrial-compatible screen-printing process.

B. Measurements

Optical constants of the grown poly-Si layers were obtained through spectroscopic ellipsometry measurements on a J.A. Woollam M-2000 ellipsometer. Carrier lifetimes J_0 and iV_{oc} of the n-type test structures were measured by photoconductance decay measurement using a Sinton WCT-120. Sheet resistances on the n-poly-Si/SiO₂/p-Si/SiO₂/n-poly-Si structures were measured using a four-point probe. Further, the laser-grown oxide was characterized using a Thermo K-alpha XPS. Surface morphologies of the laser-processed samples were observed using a laser scanning confocal microscope (Keyence VK-X3000).

III. RESULTS AND DISCUSSION

A. Optical Modeling of Selective Area Front Polysilicon Solar Cells

For accurate optical modeling, the refractive index and extinction coefficient of the grown n-poly-Si were obtained through spectroscopic ellipsometry. The raw data were fit to two Tauc-Lorentz oscillators [12] to account for visible range absorption and a Drude oscillator [13] for free carrier absorption in the near IR range. Due to limitations of the ellipsometer, measurements are made between 370 and 1687 nm. However, for a poly-Si layer at the front of a solar cell, quantifying UV range absorption is crucial. As shown in Fig. 3, in the range between 400 and

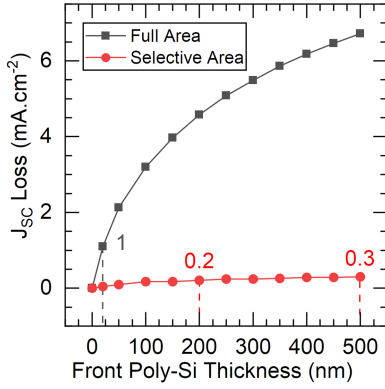


Fig. 4. Simulated J_{SC} loss due to parasitic absorption of light in the front poly-Si for full area and selective TOPCon structures.

700 nm, our optical constants are very similar to those measured by Padhamnath et al. [14]. Therefore, for optical modeling, Padhamnath's data were used between 250 and 500 nm, and our measurements were utilized in the remaining wavelength range.

The refractive index and extinction coefficient data discussed previously were used in PVLighthouse's OPAL 2 optical simulation program [15], [16] to model the parasitic absorption of light in the front poly-Si similar to a process described by Huang et al. [5]. For full-area front TOPCon, the parasitic absorption in poly-Si ($J_{abs, Poly}$) can be calculated as the difference of generated current in the absorber with ($J_{gen, Poly}$) and without ($J_{gen, NoPoly}$) poly-Si, given as follows:

$$J_{abs, Poly} = J_{gen, NoPoly} - J_{gen, Poly}. \quad (1)$$

For 100 μm wide poly-Si fingers at a pitch of 1.56 mm shaded with 30 μm metal lines at the front of a solar cell, the total area of poly-Si exposed to light is 4.48%. The component of the generated current absorbed in this "wing area" is given as follows:

$$J_{abs, Wing} = J_{abs, Poly} \times 4.48\% \quad (2)$$

Therefore, the generation current in such a selective TOPCon structure can be quantified as follows:

$$J_{gen, Selective} = J_{gen, NoPoly} - J_{abs, wing} \quad (3)$$

From these equations, it is inferred that the loss in J_{SC} in selective and full area font TOPCon cells are approximately equal to $J_{abs, Wing}$ and $J_{abs, Poly}$, respectively. Fig. 4 shows the simulated J_{SC} loss as a function of front poly-Si thickness for both full area and selective area TOPCon. We find that exceeding 20 nm of full area poly-Si on the front of the solar cell will incur a current loss over 1 $\text{mA}\cdot\text{cm}^{-2}$. On the other hand, using a selective TOPCon structure renders the front parasitic absorption far less sensitive to the thickness of the poly-Si. For 200-nm thick selective poly-Si fingers, only 0.2 $\text{mA}\cdot\text{cm}^{-2}$ current density is lost in poly-Si, allowing for easy and high-quality fire-through screen-printed metallization. In fact, increasing the thickness by a factor of 2.5–500 nm results in a loss of only an additional

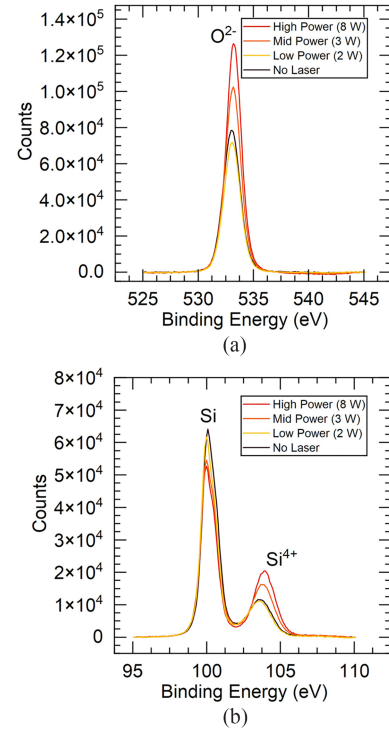


Fig. 5. (a) O^{2-} and (b) Si^{4+} peaks on XPS for various laser powers. Higher laser powers grow thicker and more stoichiometric oxides at a given scan speed.

0.1 $\text{mA}\cdot\text{cm}^{-2}$, indicating that the current loss in this device structure is very forgiving to the increased poly-Si thickness while potentially improving passivation quality (lower J_0) with fire-through screen-printed contacts [3].

B. Laser-Induced Oxidation of Polysilicon

XPS measurements in Fig. 5(a) show an increase in the intensity of the divalent oxygen peak (O^{2-}) at 532 eV, which, in our case, is related to the thickness of surface SiO_2 present [17]. This demonstrates growth of silicon oxide with the UV laser and shows that increased laser power increases the oxide thickness beyond that of native oxide. Laser-grown oxide thickness is estimated to be in the range of 1–4 nm. Moreover, the shift in the Si^{4+} peak (at 103.5 nm) in Fig. 5(b) indicates that laser oxide is more stoichiometric than the native oxide [18]. This suggests that the oxide growth mechanism is similar to thermal oxidation as governed by the Deal–Grove model. Similar observations on oxidation with a nanosecond UV laser were made by Orlowski et al. [19].

C. Demonstration of Laser Oxide As Etch Mask

We used a 9% $_{\text{vol}}$ KOH solution (DI water:45% KOH \equiv 5:1) at 40 $^{\circ}\text{C}$ to etch poly-Si between the laser patterned oxide to form poly-Si fingers. This solution is highly selective in etching silicon over SiO_x . It was observed that 200 nm of poly-Si can be etched at 40 $^{\circ}\text{C}$ in \sim 90 s. In addition, the laser-grown oxide

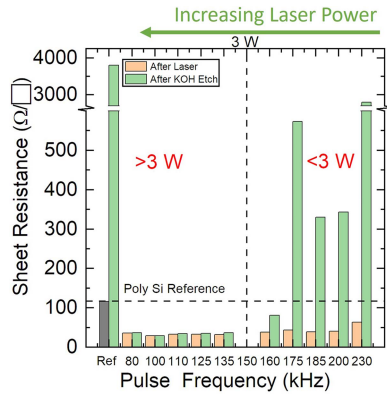


Fig. 6. Measured sheet resistance of Poly-Si films before and after KOH etch. For laser powers greater than 3 W, the sheet resistance does not change with KOH etching, indicating protection of the poly-Si by the laser grown oxide.

was successful in preventing etching of the underlying poly-Si, serving as an effective etching mask to form polyfingers.

The presence of poly-Si after 2 min of KOH immersion was verified through sheet resistance measurement of 25 mm square lasered areas on the n-type poly-Si films grown on p-type wafers, as shown in Fig. 6. Without any laser treatment, the poly-Si showed a sheet resistance of 120 Ω/\square . After two min of KOH etching, this increased to $\sim 3800 \Omega/\square$ indicating complete removal of poly-Si. For all tested laser powers, the sheet resistance decreased significantly to $\sim 30 \Omega/\square$ after laser treatment. This suggests a laser-induced change in surface dopant concentration of the deposited poly-Si. More importantly, for powers over 3 W (pulse frequencies below 150 kHz), the sheet resistance does not change before and after 2 min of KOH poly-Si etch, demonstrating successful protection of the phosphorus doped poly-Si layer under the laser-grown oxide. Conversely, at powers below 3 W, we observed that the post-etch sheet resistance increased, indicating thinning of the poly-Si layer. Therefore, below 3 W of laser power, the oxide grown by the laser, if any, was inadequate to protect the doped layer from KOH etching.

D. Laser-Induced Surface Damage, Degradation of Passivation, and Its Recovery By Subsequent SiN_x Deposition

Fig. 7 shows laser confocal micrographs of the textured surface of a silicon wafer at different laser powers. Compared with the reference textured wafer, the sample at 8 W shows significant rounding of pyramids. At 3 W of laser power, while the damage is less, there is still observable rounding of the tips of the pyramids. Change in the morphology of the 3–5 μm sized pyramids implies that the damage depth of the laser exceeds the thickness of the poly-Si. This suggests that the recombination at the $\text{Si}/\text{SiO}_x/\text{Poly-Si}$ interface would be affected by the laser.

The passivation of the absorber is quantified using J_0 obtained from photoconductance measurements (Table 1). For the reference (not lasered) sample, the measured total J_0 of the as deposited symmetric TOPCon test structure was $64.5 \text{ fA}\cdot\text{cm}^{-2}$, which was reduced to $12.8 \text{ fA}\cdot\text{cm}^{-2}$ after deposition of 70 nm PECVD SiN_x on both sides. Since this is a symmetric structure, $J_{0,\text{poly}} = 6.4 \text{ fA}\cdot\text{cm}^{-2}$ for each surface.

TABLE I
STRUCTURE J_0 WITH DIFFERENT LASER POWERS

Laser Power	J_0 (Both Sides)		
	No Laser	After Laser	After SiN_x Deposition
Reference	$64.5 \text{ fA}\cdot\text{cm}^{-2}$	-	$12.8 \text{ fA}\cdot\text{cm}^{-2}$
1.8 W	-	$518 \text{ fA}\cdot\text{cm}^{-2}$	$15.8 \text{ fA}\cdot\text{cm}^{-2}$
3 W	-	$1860 \text{ fA}\cdot\text{cm}^{-2}$	$43.2 \text{ fA}\cdot\text{cm}^{-2}$
8 W	-	$2309 \text{ fA}\cdot\text{cm}^{-2}$	$194 \text{ fA}\cdot\text{cm}^{-2}$

Fig. 8 shows that laser oxidation caused significant deterioration in passivation increasing the measured total J_0 values in the range of 500–2300 $\text{fA}\cdot\text{cm}^{-2}$. However, after the subsequent SiN_x deposition, which is required for screen-printed metallization, the J_0 value decreased dramatically from $15 \text{ fA}\cdot\text{cm}^{-2}$ at 1.8 W to $194 \text{ fA}\cdot\text{cm}^{-2}$ at 8 W. No further change was observed after belt furnace firing without screen-printed metal at a peak temperature of 760 $^\circ\text{C}$.

This improvement is likely due to hydrogenation of the surface resulting from the nitride growth [20]. At 3 W of laser power, above which the laser-grown oxide acts as an etch mask, the J_0 of the test structure was $43.2 \text{ fA}\cdot\text{cm}^{-2}$ after SiN_x deposition. Since the laser treatment was performed only on one side of the wafer, the J_0 of the lasered surface can be obtained by subtracting the J_0 of the nonlasered surface from the total J_0 , giving us $36.8 \text{ fA}\cdot\text{cm}^{-2}$ for the laser-treated surface. However, this translates into only $1.65 \text{ fA}\cdot\text{cm}^{-2}$ for J_0 contribution from the unmetallized wings of polyfingers because of 4.48% area coverage on the front of the solar cell.

E. Device Simulations of Laser Processed Selective TOPCon

The efficiency potential of the cell structure shown in Fig. 1 was studied through 2-D device simulations in Quokka 2 [21]. The parameters used in the model are detailed in Table II. The J_0 of the wing areas of the polyfingers was approximated to $40 \text{ fA}\cdot\text{cm}^{-2}$ based on the discussion in Section III-D. Furthermore, with the increasing popularity of busbarless solar cells and measurement methods, we have adopted that design in our simulations [22]. The device characteristics for this simulated solar cell are shown in Fig. 9 with efficiency approaching 25%. Using this model as a baseline, the dependence of performance on various inputs is studied in the next sections.

1) *Dependence of Solar Cell Performance on Polyfinger Thickness:* Fig. 10(a) shows the efficiency of the solar cell as a function of thickness of the front selective poly-Si under the metal fingers. As suggested by our OPAL 2 model discussed in Section III-A, the J_{sc} and hence the efficiency of a selective area front TOPCon structure, is relatively insensitive to the thickness of the poly-Si. In fact, we observe that increasing the thickness from 200–500 nm incurs only an additional 0.05% loss in absolute efficiency, potentially allowing less damaging metallization without incurring a significant drop in current.

2) *Dependence of Solar Cell Performance on Polyfinger Width:* Polyfinger width affects device performance in two ways: (1) wider polyfingers incur higher current loss due to parasitic light absorption in the wing area and (2) increased contribution of the $J_{0,\text{FrontPoly}}$ to the total J_0 . Because the

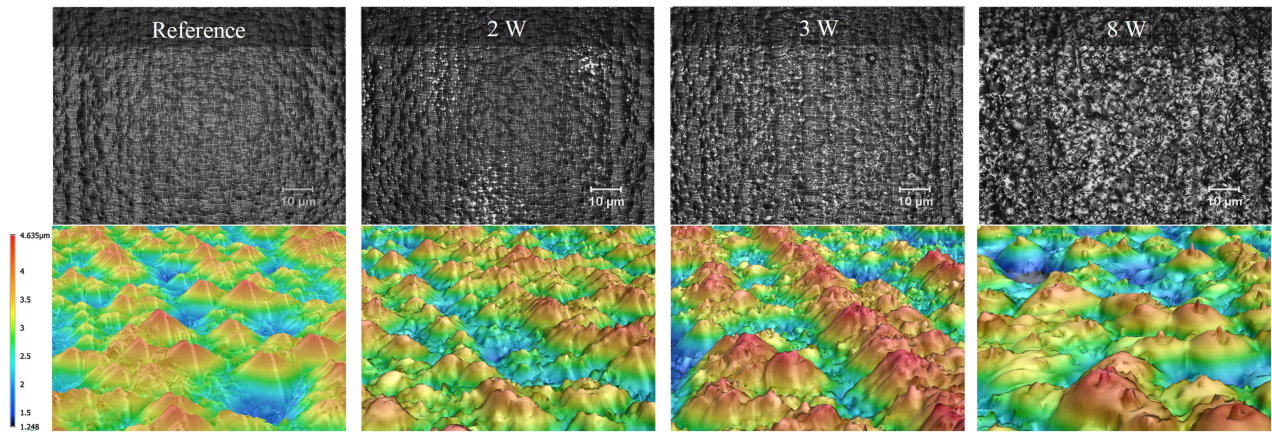


Fig. 7. Laser confocal microscopy images in plan view and isometric view showing increasing damage to the surface morphology at higher laser powers.

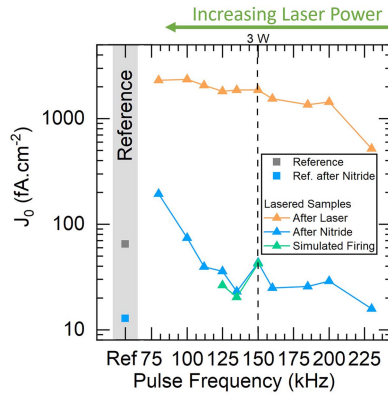


Fig. 8. Deterioration of J_0 after laser incidence and subsequent recovery of passivation after SiN_x deposition.

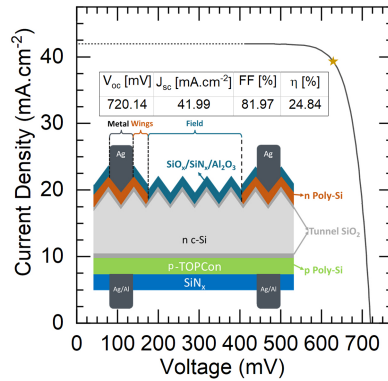


Fig. 9. Current–voltage characteristics of the simulated selective-front double-side TOPCon solar cell.

$J_{0,\text{FrontPoly}}$ is significantly higher than $J_{0,\text{field}}$ ($= 5 \text{ fA}\cdot\text{cm}^{-2}$ calculated from SRV of $3.8 \text{ cm}\cdot\text{s}^{-1}$) due to laser damage, increasing the width causes deterioration in V_{OC} . However, due to the low overall area fraction, the effect of this deterioration is diminished. We observe that doubling the width from 100 to 200 μm decreases the V_{OC} by 2.1 mV and the J_{SC} by $0.4 \text{ mA}\cdot\text{cm}^{-2}$, together amounting to an efficiency loss of only 0.2%, as shown

TABLE II
INPUT PARAMETERS FOR QUOKKA 2 DEVICE SIMULATION

	Parameter	Value	Reference
Cell Data	V_{OC} [mV]	720.15	
	J_{SC} [$\text{mA}\cdot\text{cm}^{-2}$]	41.99	
	FF [%]	81.98	
	Efficiency [%]	24.84	
	Total J_0 [$\text{fA}\cdot\text{cm}^{-2}$]	34	
Front Selective TOPCon	Front Poly Type	N (textured)	
	Selective area width [μm]	100 (variable)	
	Front Poly thickness [nm]	200 (variable)	
	Front field SRV [$\text{cm}\cdot\text{s}^{-1}$]	3.8	[23]
	$J_{0,\text{frontpoly,pass}}$ [$\text{fA}\cdot\text{cm}^{-2}$]	40	Section III-D
	$J_{0,\text{frontpoly,metal}}$ [$\text{fA}\cdot\text{cm}^{-2}$]	300 (variable)	In house
Bulk	$J_{0,\text{front,total}}$ [$\text{fA}\cdot\text{cm}^{-2}$]	12	
	Bulk Type	N	
	Bulk thickness [μm]	170	
	Bulk resistivity [$\Omega\cdot\text{cm}$]	1	[5]
	SRH τ_{p0} [ms]	5	[5]
	SRH τ_{p0} [ms]	5	[5]
Rear TOPCon	$J_{0,\text{bulk,total}}$ [$\text{fA}\cdot\text{cm}^{-2}$]	11	
	Rear Poly Type	P (planar)	
	Rear Poly thickness [nm]	200	
	Rear Poly sheet resistance [Ω/\square]	200	In house
	$J_{0,\text{rearpoly,pass}}$ [$\text{fA}\cdot\text{cm}^{-2}$]	4.6	[24]
	$J_{0,\text{rearpoly,metal}}$ [$\text{fA}\cdot\text{cm}^{-2}$]	112	[24]
Front grid	$J_{0,\text{rear,total}}$ [$\text{fA}\cdot\text{cm}^{-2}$]	11	
	Number of lines	100	
	Pitch [mm]	1.56	
	Line thickness [μm]	30	
	Metal coverage	1.92%	
	Contact resistivity [$\text{m}\Omega\cdot\text{cm}^2$]	2 (variable)	[25]
Rear grid	Number of lines	200	
	Pitch [mm]	0.78	
	Line thickness [μm]	30	
	Metal coverage	3.84%	
	Contact resistivity [$\text{m}\Omega\cdot\text{cm}^2$]	3	[24]
	External R_s [$\Omega\cdot\text{cm}^2$]	0.009	
	J_{gen} [mA/cm^2]	43.26	Section III.A

in Fig. 10(b). This can help relax the metal grid alignment requirements if needed.

3) *Dependence of Performance on $J_{0,\text{metal}}$ on the Front of the Solar Cell:* The reference simulation above uses $J_{0,\text{metal}-100\%} = 300 \text{ fA}\cdot\text{cm}^{-2}$. This was obtained through initial experiments on laser-processed symmetric n-TOPCon samples based on the methodology described by Huang et al. [26]. For

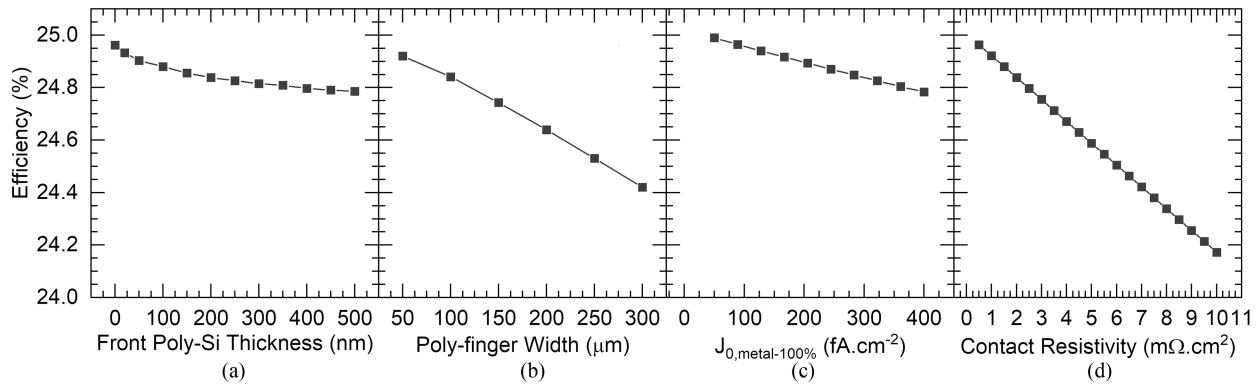


Fig. 10. Variation of simulated solar cell efficiency as a function of (a) polyfinger thickness, (b) polyfinger width, (c) $J_{0,\text{metal}}$, and (d) contact resistivity.

30 μm wide metal lines, amounting to a coverage fraction of 1.92%, this amounts to an effective contribution of 5.76 $\text{fA}\cdot\text{cm}^{-2}$. Fig. 10(c) shows that decreasing the full area $J_{0,\text{metal}}$ from 300 – 50 $\text{fA}\cdot\text{cm}^{-2}$ (0.96 $\text{fA}\cdot\text{cm}^{-2}$ for 1.92% metal coverage) can increase the efficiency of the cell from 24.84% to 25%. This shows that when the metal coverage is low, the efficiency does not vary significantly with $J_{0,\text{metal}}$.

4) *Dependence of Performance on Contact Resistivity of Metal/TOPCon Stack:* For optimally metalized high-quality n-TOPCon layers, achieving a contact resistivity of 2 $\text{m}\Omega\cdot\text{cm}^2$ is not uncommon [25], [27]. However, low temperature or underfiring of the metallization can increase contact resistivities. The simulations show that for this cell structure every 2 $\text{m}\Omega\cdot\text{cm}^2$ increase in contact resistivity decreases the fill factor by $\sim 0.5\%$ and the efficiency by $\sim 0.2\%$. This dependence is shown in Fig. 10(d).

IV. CONCLUSION

In this work, we have demonstrated the growth of 1–4 nm silicon oxide on TOPCon using a 355-nm nanosecond UV laser. We have shown that for laser powers above 3 W at 400 mm/s, the grown oxide layer can protect the underlying selective TOPCon from KOH etching. Despite the appreciable degradation in J_0 from the laser treatment, passivation quality is largely recovered after subsequent PECVD SiN_x deposition, which is required for screen printed contacts. At 3 W of laser power, a full area J_0 of 36.8 $\text{fA}\cdot\text{cm}^{-2}$ was achieved, which translates to only 1.65 $\text{fA}\cdot\text{cm}^{-2}$ for 4.48% coverage of polyfingers (100 polyfingers, 100 μm wide, and with 30 μm of metal). Based on our Quokka 2 modeling of double-side TOPCon cells with selective area poly-Si on the front with practically achievable material parameters, the performance of a back junction bifacial fully screen-printed solar cell can approach 25%. Further, we have found that a front selective area poly-Si-based solar cell is very insensitive to the passivation quality of the front contact. These advantages combined with the ease of fabricating this selective area front TOPCon structure using the novel laser-oxide patterning technology make it a strong contender for the next generation of high-efficiency silicon solar cells.

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