

Non-vanishing Yukawa interactions in asymptotically safe quantum gravity

Quantum Spacetime and the Renormalization Group

Heidelberg, April 1, 2025

Marc Schiffer, Radboud University Nijmegen

In collaboration with

G. de Brito, M. Reichert: **arXiv: 2504.XXXXX**

Radboud Universiteit



Towards the Yukawa sector of the Standard Model

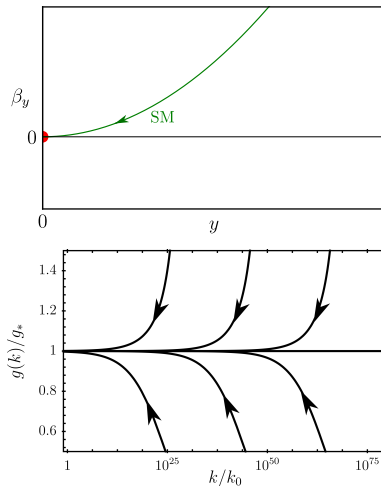
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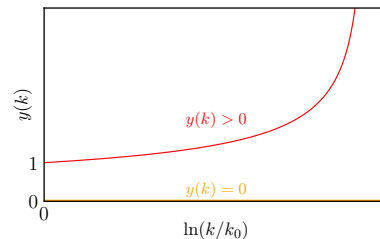
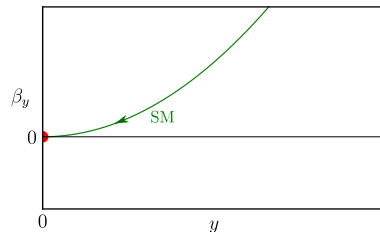


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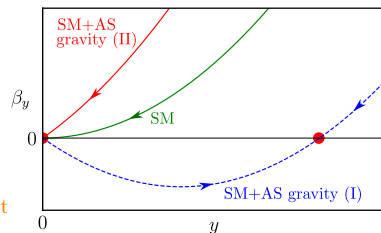
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$$\beta_y = -f_y g_N y + \#_{\text{SM}} y^3 + \mathcal{O}(y^4)$$

$$g_N f_y \begin{cases} = \text{const.} > 0, & \text{for } k > M_{\text{Pl}} \wedge \Lambda < \Lambda_{\text{crit}} \\ = \text{const.} < 0, & \text{for } k > M_{\text{Pl}} \wedge \text{for } \Lambda > \Lambda_{\text{crit}} \\ \approx 0, & \text{for } k < M_{\text{Pl}} \end{cases}$$



[Oda, Yamada; 2015], [Eichhorn, Held, Pawłowski; 2016],
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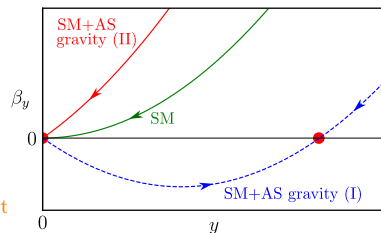
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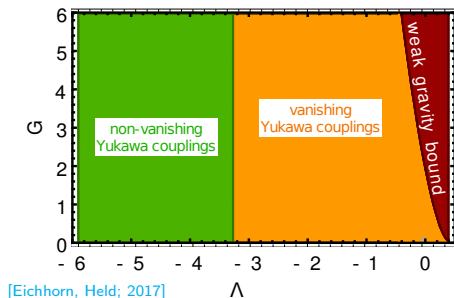
UV completion of the simple Yukawa system: constraints on gravitational dynamics

Additionally: top mass might be retro-dicted [Eichhorn, Held, 2017]

Simple Yukawa system: state of the art (LO)

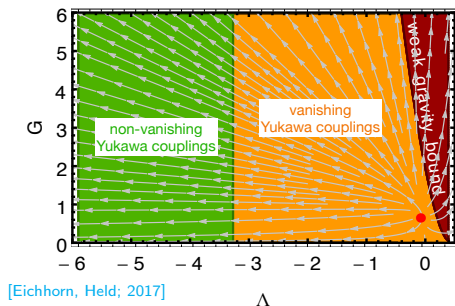
- Simplest approximation:

$$\Lambda_{crit} \approx -3.3$$



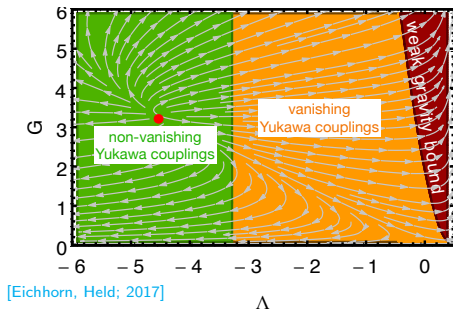
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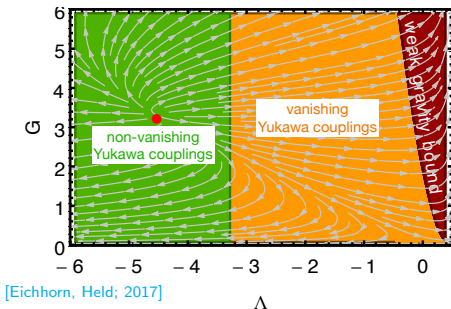


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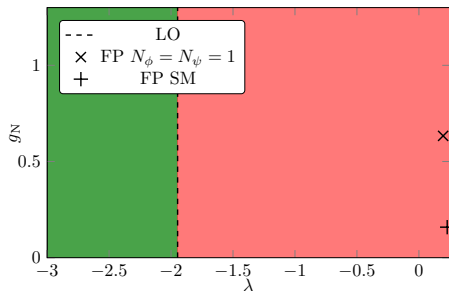


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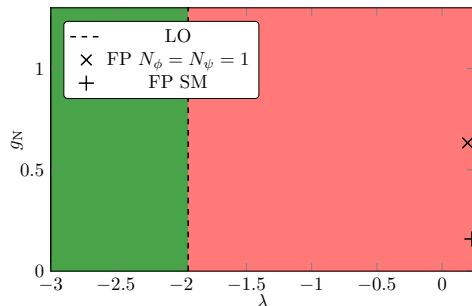
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- But: SM fixed point:
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- Different scheme:
(fluctuation computations)
 $f_{y,*} < 0$
[see, e.g., \[Christiansen et al.; 2019\]](#)

Possible ways out?

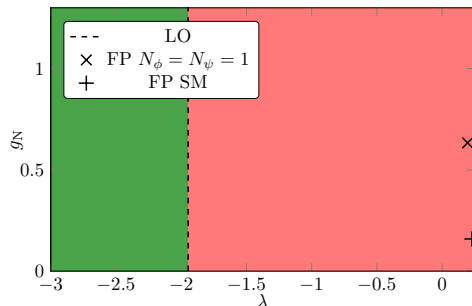


1. FP shifts when including higher-order operators

see e.g. [Eichhorn, MS; 2022], [Pawlowski, Reichert; 2023] for reviews

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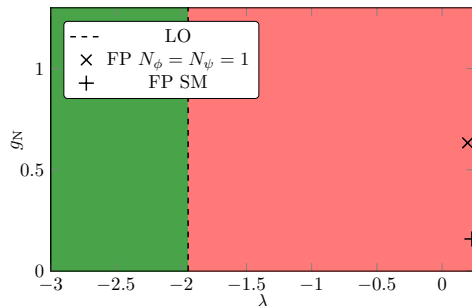
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2. Region of relevance shifts when including higher-order operators

- more precisely, need $\Theta_y > 0$, for $y_{\text{IR}} > 0$

$$\Theta_I = -\text{eig}(M_{ij}) , \quad \text{with} \quad M_{ij} = \left(\frac{\partial \beta_{g_j}}{\partial g_i} \right) \Big|_{\mathbf{g}=\mathbf{g}^*}$$

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- At LO: $\Theta_y = f_y g_N$;
- NLO: g_N^2 contributions to Θ_y ;
- Possible sources:
 - ▶ direct contributions to M_{11}
 - contribution $\sim y g_{\text{ind}}$ to β_y with $g_{\text{ind},*} \neq 0$
 - ▶ off-diagonal contributions in M_{ij} - contribution to β_y can be independent of y itself

Direct contributions: induced operators

$$\beta_y = \left(-f_y g_N \right) y + \#_{\text{SM}} y^3 + \mathcal{O}(y^4)$$

- Example: Chiral fermion ψ

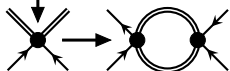
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$$\Gamma_k^{\text{kin}} = \int d^4x \sqrt{g} \ (\bar{\psi} \not{D} \psi)$$



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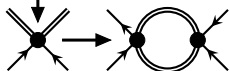
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- Schematically:

$$\beta_{\lambda_1} = C_0(g_N) + C_1(g_N) \lambda_1 + C_2 \lambda_1^2$$

- Interacting fixed point: $\lambda_{1,*} \sim g_N^2$

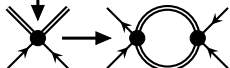
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- In general:

two kinds of such contributions

$$g_{\text{self},*} \sim g_N^2, \quad g_{n-m,*} \sim g_N$$

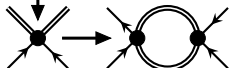
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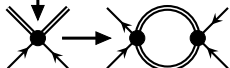
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$$g_{\text{self},*} \sim g_N^2, \quad g_{n-m,*} \sim g_N$$

- g_{ind} does not induce y
 $\Rightarrow \beta_{g_{\text{ind}}}$ do not contribute
 to Θ_y via M_{ij}

- Consider new coupling g_{ho} such that

$$M = \begin{bmatrix} -f_y g_{\text{N}} + \#_{\text{ind}} g_{\text{N}}^2 & \#_1 g_{\text{N}} \\ \#_2 g_{\text{N}} & \dim_{\text{ho}} + \#_3 g_{\text{N}} \end{bmatrix}.$$

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- Includes operators with $g_{\text{ho},*} = 0$!

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- **contributions to M_{11}**
- contributions to β_y with $g_* \neq 0$

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$(\bar{\psi}\psi)^2$	$(\psi\gamma_\mu\psi)^2$	λ_{V}	6
	$(\psi\gamma_\mu\gamma_5\psi)^2$	λ_{A}	6
$\phi^2\psi\psi$	$\partial_\mu\phi\partial^\mu\phi\bar{\psi}\overleftrightarrow{D}\psi$	$\chi_{1/2}$	8
ϕ^4	$(\partial_\mu\phi\partial^\mu\phi)^2$	K_2	8

see [Eichhorn, MS; 2022] for references

contribute via η_i only \Rightarrow neglect

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Induced operators shift FP, but admit lower-triangular stability sub-matrix

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Modulo momentum-dependences: No further g_N^2 contribution to Θ_y

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Non-minimal gravity-matter system: fixed points

- Solve EH + induced couplings
(bilocal anomalous dimensions and non-minimal couplings)

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$$(g_N, \lambda, \rho_{\text{Ric}}, \rho_{\text{R}}, \sigma_{\text{Ric}}, \sigma_{\text{R}}, \lambda_+, \lambda_-)^* = (0.65, 0.19, -0.13, 0.50, -0.058, 0.42, -0.63, 0.63) .$$

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$$(\eta_h(k^2), \eta_h(0), \eta_\phi(k^2), \eta_\phi(0), \eta_\psi(k^2), \eta_\psi(0))^* = (0.229, 1.42, 0.588, -0.226, 0.988, 0.0819) .$$

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- both cases: well-behaved extensions of the Reuter FP!

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- Plug FP-values into non-induced operators, and compute critical exponents at

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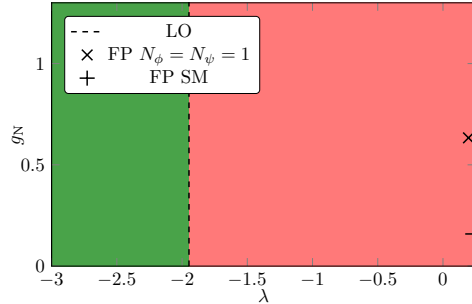
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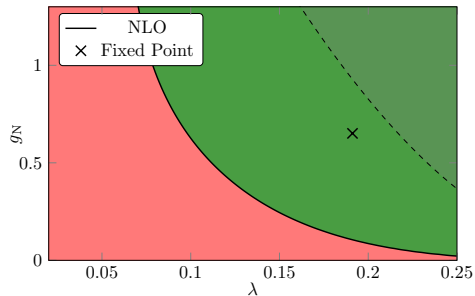
UV completion of the simple Yukawa system via NLO operators!

Yukawa coupling at NLO



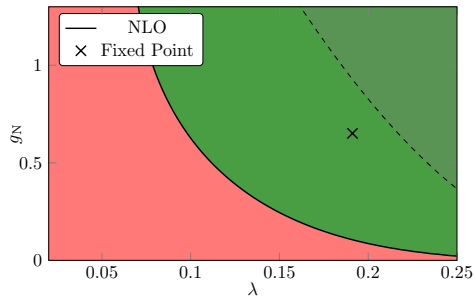
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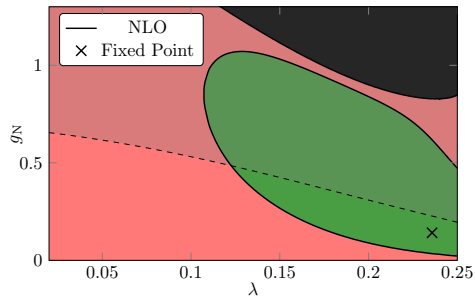


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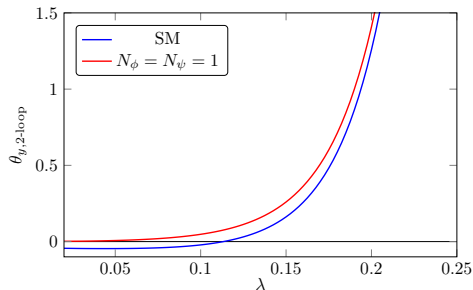


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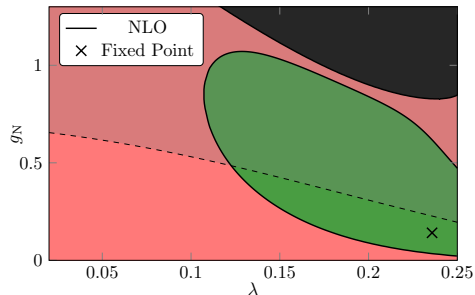


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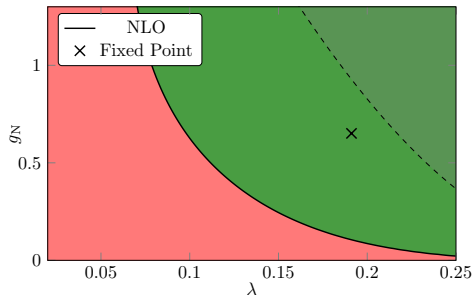
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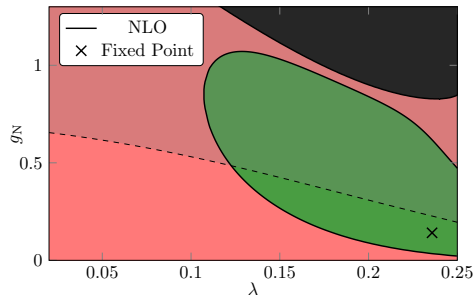
- NLO contributions to stability matrix: generate new regime where $\Theta_y > 0$
- Fixed point lies inside that new regime!

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- NLO contributions to stability matrix: generate new regime where $\Theta_y > 0$
- Fixed point lies inside that new regime!
- LO: $\Theta_y < 0$, NLO: $\Theta_y > 0$; What about NNLO (i.e., g_N^3 -contributions)?

Estimating NNLO effects ('error bars')

- previously observed: $R \phi \bar{\psi} \psi$, $R^2 \phi \bar{\psi} \psi$ and $C^2 \phi \bar{\psi} \psi$ dominate
- estimate effect of $R^3 \phi \bar{\psi} \psi$ -type operator

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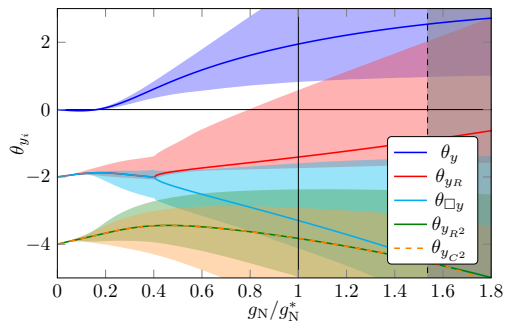
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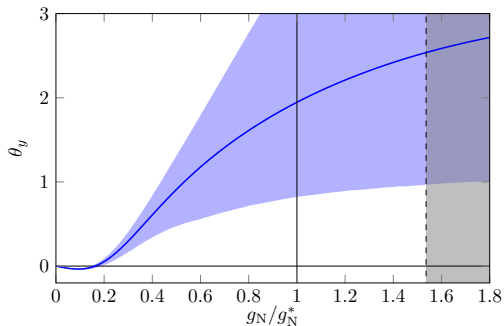
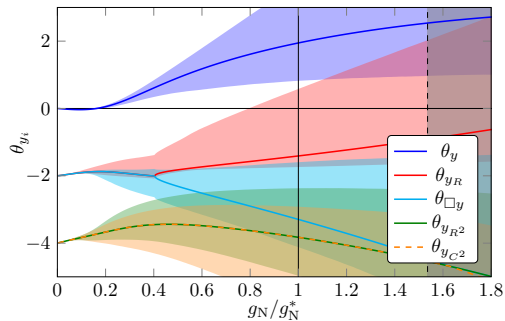
Results of (N)NLO simulations

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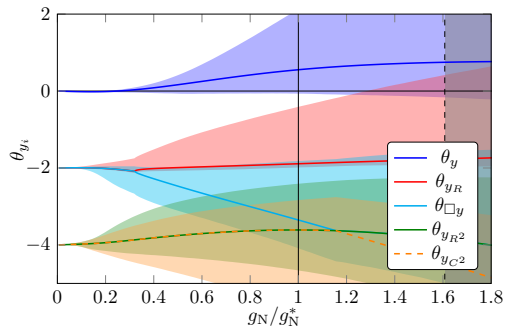
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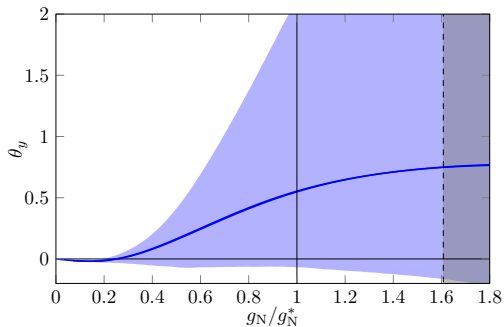
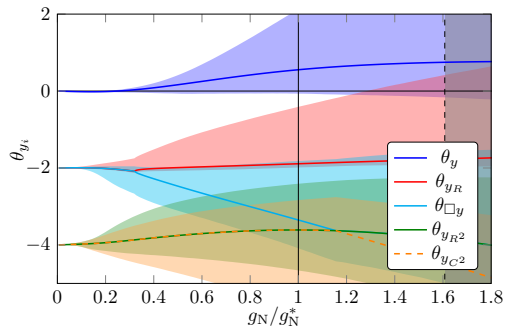
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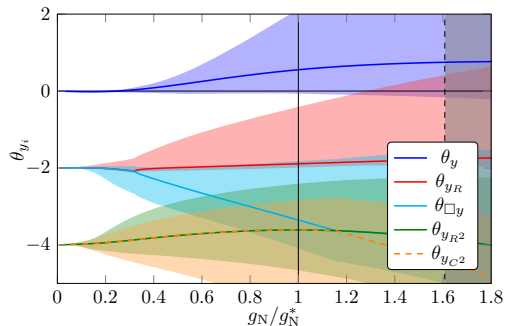
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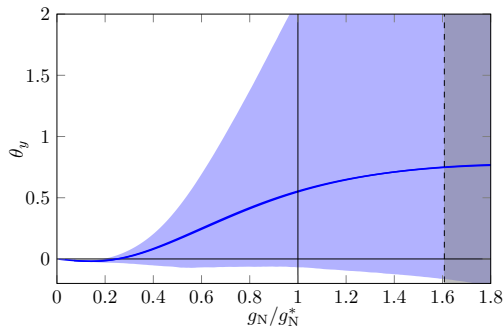
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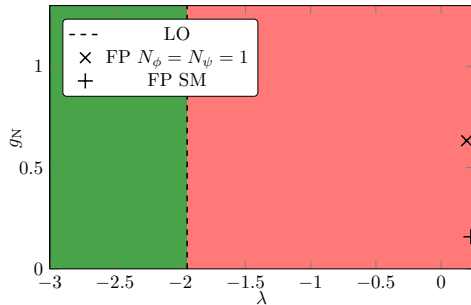
- (N)NLO simulations: insights into robustness
- but: large(ish) error-bars

Summary

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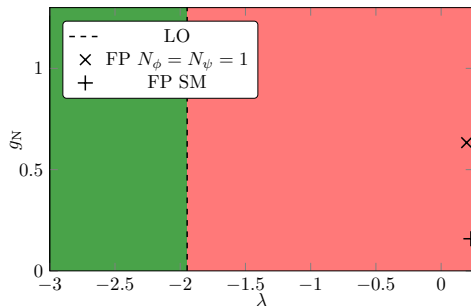
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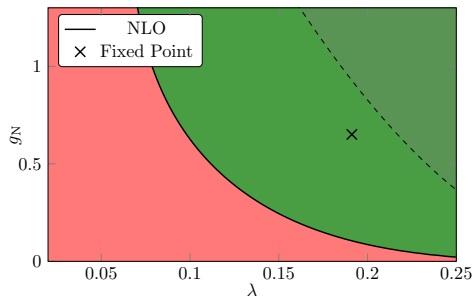
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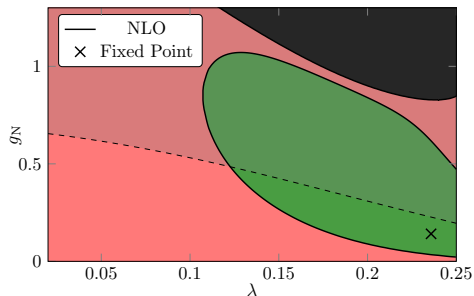
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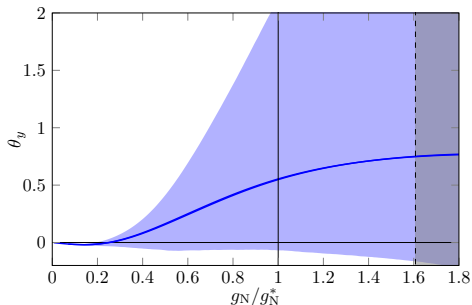
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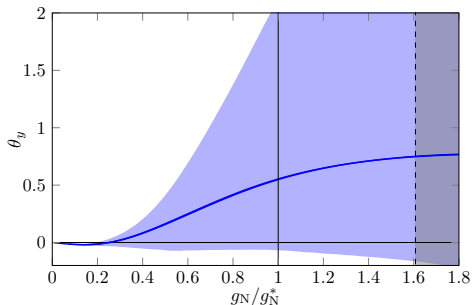
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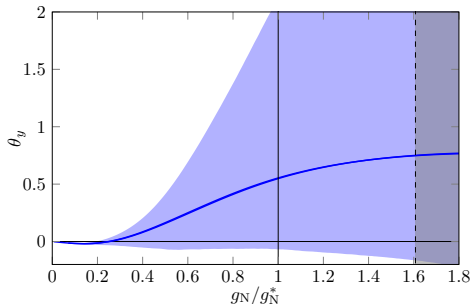
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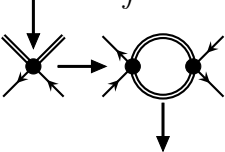


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Thank you for your attention!

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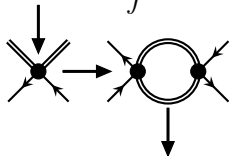
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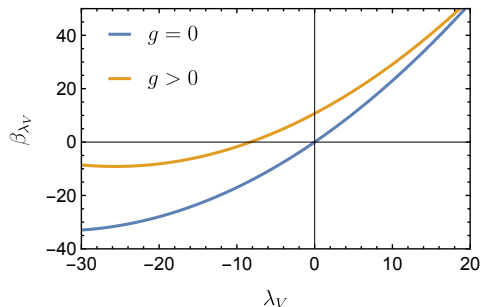
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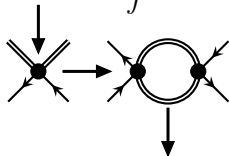
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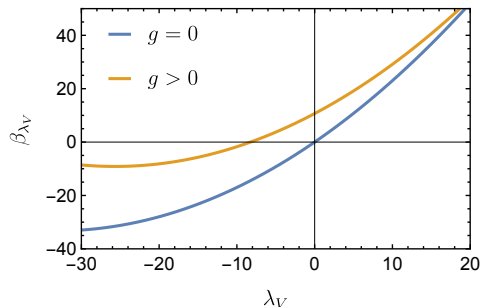
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- Schematically:

$$\beta_{\lambda_V} = B_0(g) + \lambda_V B_1(g) + \lambda_V^2 B_2$$

- For $g > 0$: $\lambda_{V,*} \neq 0$