

## A Note on Spherical Co-ordinates

In cartesian co-ordinates the volume element  $d^3r = dx dy dz$ . In spherical coordinates  $(r, \theta, \phi)$ ,  $d^3r = r^2 dr \sin \theta d\theta d\phi$ . The volume element is the dot product of an area element and the displacement in a perpendicular direction. The area element in  $y - z$  plane is defined by  $d\vec{A}_{yz} = dy\hat{y} \times dz\hat{z}$ . The volume element is  $dx\hat{x} \cdot d\vec{A}_{yz} = dx dy dz$ .

The spherical coordinates are defined by the radial distance  $r$  from the center, polar angle  $\theta$  (the angle  $\vec{r}$  makes with  $z$  axis) and the azimuthal angle  $\phi$  (the angle the projection of  $\vec{r}$  on  $x - y$  plane makes with  $x$ -axis). With these definitions it is straight forward to work out the cartesian coordinates to be

$$\begin{aligned}x &= r \sin \theta \cos \phi \\y &= r \sin \theta \sin \phi \\z &= r \cos \theta.\end{aligned}$$

We consider three mutually perpendicular displacements. The radial displacement  $dr$  occurs when we hold  $\theta$  and  $\phi$  fixed. This is along the radial unit vector  $\hat{r}$ , whose direction changes from point to point (unlike the cartesian unit vectors). The displacement along  $\theta$  occurs when we hold  $r$  and  $\phi$  fixed. This is given by  $r d\theta \hat{\theta}$ . On the surface of the earth, this is equivalent to moving along a longitude. The corresponding unit vector  $\hat{\theta}$  is tangential to the longitude. The displacement along  $\phi$  occurs when we hold  $r$  and  $\theta$  fixed. This is given by  $r \sin \theta d\phi \hat{\phi}$ . On the surface of the earth, this is equivalent to moving along a latitude, whose radius is  $r \sin \theta$ . The unit vector  $\hat{\phi}$  is tangential to the latitude. In spherical co-ordinates, the position vector is  $\vec{r} = r\hat{r}$  but the most general displacement is

$$d\vec{r} = dr\hat{r} + r d\theta \hat{\theta} + r \sin \theta d\phi \hat{\phi}.$$

When the position vector changes, in general, both its magnitude  $r$  and its direction  $\hat{r}$  change. The change in  $\hat{r}$  is what leads to the additional terms in  $d\vec{r}$ .

The three unit vectors  $\hat{r}$ ,  $\hat{\theta}$  and  $\hat{\phi}$  form a right-handed triad of unit vectors, just as  $\hat{x}$ ,  $\hat{y}$  and  $\hat{z}$  do. Based on the above definitions, you should be able to derive

$$\hat{r} = \sin \theta \cos \phi \hat{x} + \sin \theta \sin \phi \hat{y} + \cos \theta \hat{z}$$

$$\begin{aligned}\hat{\theta} &= \cos \theta \cos \phi \hat{x} + \cos \theta \sin \phi \hat{y} - \sin \theta \hat{z} \\ \hat{\phi} &= -\sin \phi \hat{x} + \cos \phi \hat{y}.\end{aligned}$$

The area element on the surface of a sphere is given by the cross product of the two independent displacements on it,

$$d\vec{A}_{\theta-\phi} = r d\theta \hat{\theta} \times r \sin \theta d\phi \hat{\phi} = r^2 \sin \theta d\theta d\phi \hat{r}.$$

This quantity is sometimes written as  $r^2 d\Omega$ , where  $d\Omega$  is called the **solid angle** subtended by the area element  $dA_{\theta-\phi}$  at the origin. The volume element is given by

$$d^3r = dr \hat{r} \cdot d\vec{A}_{\theta-\phi} = r^2 dr \sin \theta d\theta d\phi.$$

If we want to cover the whole space, the range of  $r$  is from 0 (origin) to  $\infty$ . The range of  $\theta$  is from 0 (along positive  $z$ -axis) to  $\pi$  (along negative  $z$ -axis). The range of  $\phi$  is from 0 to  $2\pi$ . Quite often, we use the variable  $\cos \theta$  rather than  $\theta$ . Then the volume element is written as

$$d^3r = r^2 dr d(\cos \theta) d\phi,$$

with the range of  $\cos \theta$  going from  $-1$  to  $+1$ . It is, of course, true that  $\sin \theta d\theta$  is  $-d(\cos \theta)$  but the negative sign is taken care of by flipping the limits of integration. The rule of thumb is that if we use  $\theta$  as the variable of integration, then we integrate from lower value of  $\theta$  to higher value of  $\theta$ . If we use  $\cos \theta$  as the variable of integration then we integrate from lower value of  $\cos \theta$  to higher value of  $\cos \theta$ . This automatically takes care of the negative sign.

From the notion of displacement vector

$$d\vec{r} = dx \hat{x} + dy \hat{y} + dz \hat{z},$$

we can define the differential operator

$$\nabla = \hat{x} \frac{\partial}{\partial x} + \hat{y} \frac{\partial}{\partial y} + \hat{z} \frac{\partial}{\partial z}.$$

In spherical co-ordinates we use exactly the same idea. From

$$d\vec{r} = dr \hat{r} + r d\theta \hat{\theta} + r \sin \theta d\phi \hat{\phi},$$

we define

$$\nabla = \hat{r} \frac{\partial}{\partial r} + \hat{\theta} \frac{1}{r} \frac{\partial}{\partial \theta} + \hat{\phi} \frac{1}{r \sin \theta} \frac{\partial}{\partial \phi}.$$

In general relativity (which uses the language of differential geometry)  $d\vec{r}$  is called *contra-variant vector* and  $\nabla\phi$  is called a *co-variant vector* where  $\phi$  is a scalar function. From the above definition of  $\nabla$ , we can compute

$$\nabla^2 = \nabla \cdot \nabla = \left( \hat{r} \frac{\partial}{\partial r} + \hat{\theta} \frac{1}{r} \frac{\partial}{\partial \theta} + \hat{\phi} \frac{1}{r \sin \theta} \frac{\partial}{\partial \phi} \right) \cdot \left( \hat{r} \frac{\partial}{\partial r} + \hat{\theta} \frac{1}{r} \frac{\partial}{\partial \theta} + \hat{\phi} \frac{1}{r \sin \theta} \frac{\partial}{\partial \phi} \right).$$

A similar definition in the cartesian case gives us only three terms because the unit vectors  $\hat{x}$ ,  $\hat{y}$  and  $\hat{z}$  do not change in direction as we move from point to point. That is not true in the case of the unit vectors in spherical polar coordinates.  $\hat{r}$  and  $\hat{\theta}$  are functions of both  $\theta$  and  $\phi$  whereas  $\hat{\phi}$  is a function of  $\phi$ . This dependence has to be kept in mind while taking the derivatives. If we do the calculation carefully, we will get

$$\nabla^2 = \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2}{\partial \phi^2}.$$