



Kepler's Equation Validity



Orbital Limits of Kepler's Equation

Conceptually, Kepler's equation is **valid** for all values of 'e' including **1.0**.

However, we note from **E - θ** relation that as value of 'e' increases, value of '**E**' decreases, for same ' **θ** '.



Orbital Limits of Kepler's Equation

We also **see** that for **e = 1**, Kepler's equation **reduces** to the following **form**.

$$E - \sin E = 0 \quad (a = \infty; \quad n = 0; \quad M = n \cdot \Delta t = 0)$$

Thus, the only **possible** solution is **E = 0**, for all ' Δt '.

In view of the **above**, we need an **alternate** approach.



TOF for Parabolic Trajectories



Parabolic Trajectory Time Solution

We know that ' **Δt** ' for any **conic** section can be **obtained** from the following **relation**.

$$TOF = t_B - t_A = \frac{h^3}{\mu^2} \int_{\theta_A}^{\theta_B} \frac{1}{(1 + e \cos \theta)^2} d\theta$$



Parabolic Trajectory Time Solution

Thus, for **e = 1**, it can be **re-written** as,

$$t_B - t_A = \frac{h^3}{4\mu^2} \int_{\theta_A}^{\theta_B} \frac{d\theta}{\cos^4(\theta/2)}$$

‘ **Δt** ’ can be **obtained** by carrying out the **integration**.



Time Solution

Integration can be **performed** by using the trigonometric **identity** for $\sec^2(\theta/2)$, as shown below.

$$\sec^4\left(\frac{\theta}{2}\right) = \sec^2\left(\frac{\theta}{2}\right) \left[1 + \tan^2\left(\frac{\theta}{2}\right)\right]$$
$$t_B - t_A = \Delta t = \frac{h^3}{2\mu^2} \left(\tan \frac{\theta}{2} + \frac{1}{3} \tan^3 \frac{\theta}{2} \right)_{\theta_A}^{\theta_B}$$

‘**t**’ is ∞ at apogee ($\theta = 180^\circ$), where ‘**v**’ = ‘0’ and ‘**r**’ = ∞ .



Nature of Parabolic Δt Solution

An important feature of Δt **solution** for a parabolic path is that **unlike** ellipse, angle θ is always less than $\pm 180^\circ$, due to apogee being **undefined** (or infinite).

Thus, **practical** solutions for ' Δt ' (or ' $\Delta \theta$ ') on a **parabolic** path are for a **finite** distance, ' r '.



Summary

As was **noted** earlier, TOF **solution** for parabolic trajectories can be **obtained** by direct integration of the **time** equation.

We also **note** that Kepler's equation, as **formulated** for an elliptic orbit, **breaks** down for $e = 1$.



TOF for Hyperbolic Trajectories



Hyperbolic Trajectories

If **spacecraft** have $\varepsilon > 0$, these move on a **hyperbolic** path.

Similar to **parabolic** case, in these **cases** also, we **need** to know how **long** it would take to **reach** a specific point and thus, **need** a process to arrive at **hyperbolic** Δt .



Hyperbolic Trajectory Relation

Hyperbola is **characterized** by $\varepsilon > 0$, resulting in '**a**' < 0 while, '**p**', which is a physical **distance**, is positive.

Thus, we need to **rewrite** the basic conic section **relations**, in the context of **hyperbola**.



Hyperbolic Trajectory Relation

The **applicable** conic relations are as **follows**.

$$p = (-a)(e^2 - 1), \quad r_p = (-a)(e - 1); \quad r_a = (a)(e + 1)$$

We see that while, both '**p**' and '**r_p**' are positive, '**r_a**' is **negative**, because '**r**' **crosses** '**∞**' at $\cos\theta = (-1/e)$.



Hyperbolic Δt Solution Strategy

One way to **arrive** at the **solution** for hyperbola is to perform **explicit** integration of **time** equation for $e > 1$.

In this regard, we can **examine** the elliptic orbit **integral** for its **possible** re-interpretation, when $e > 1$.



Hyperbolic Δt Solution Strategy

Another option is to **examine** the possibility of using **Kepler's** equation in the context of a **hyperbolic** path.

This can be done by **identifying** an applicable auxiliary **geometry** that can be **used** to set up ' Δt ' **relations**.

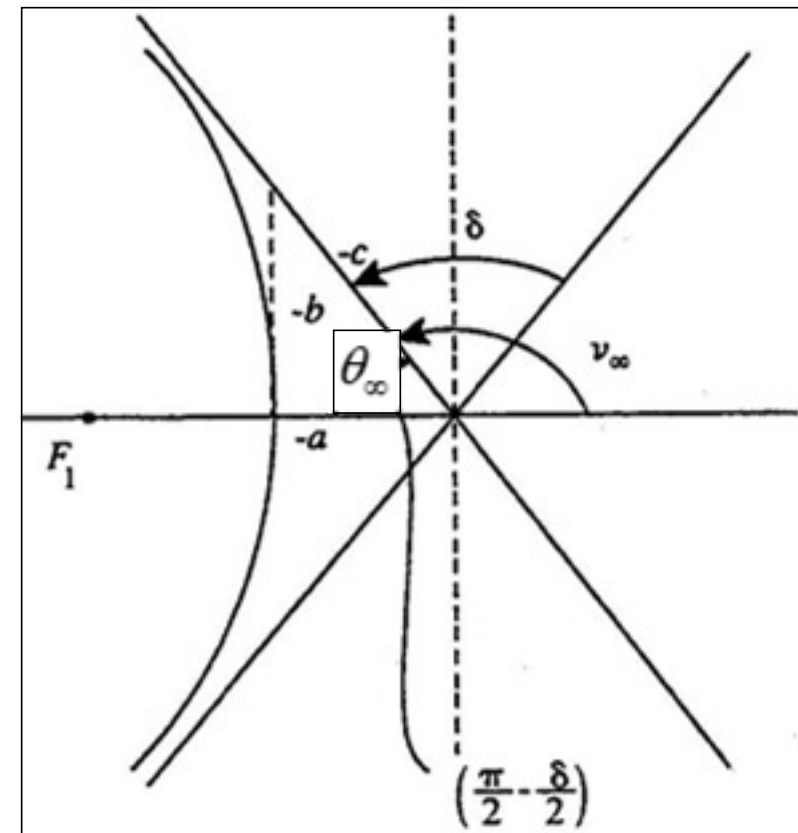


Hyperbolic Trajectory Features

Applicable **auxiliary geometry** is arrived at by considering hyperbolic **features** as shown alongside.

Here, ' δ ' is angle **between** asymptotes at $\mathbf{r} = \pm\infty$, and is related to ' e ', as follows.

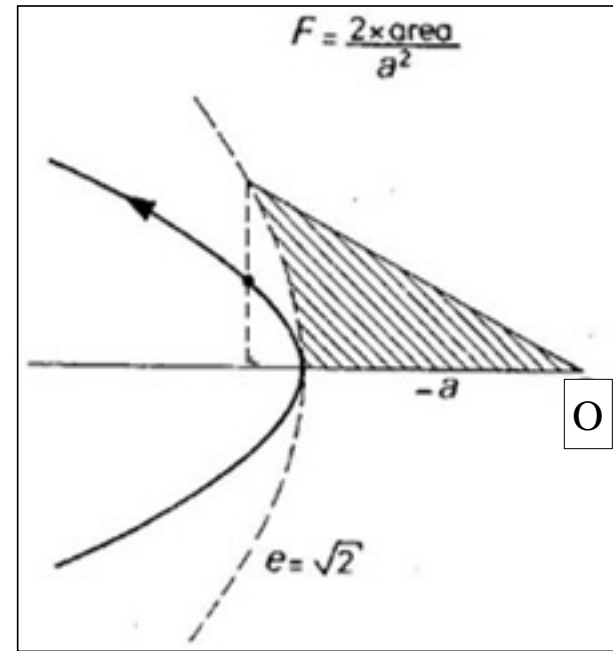
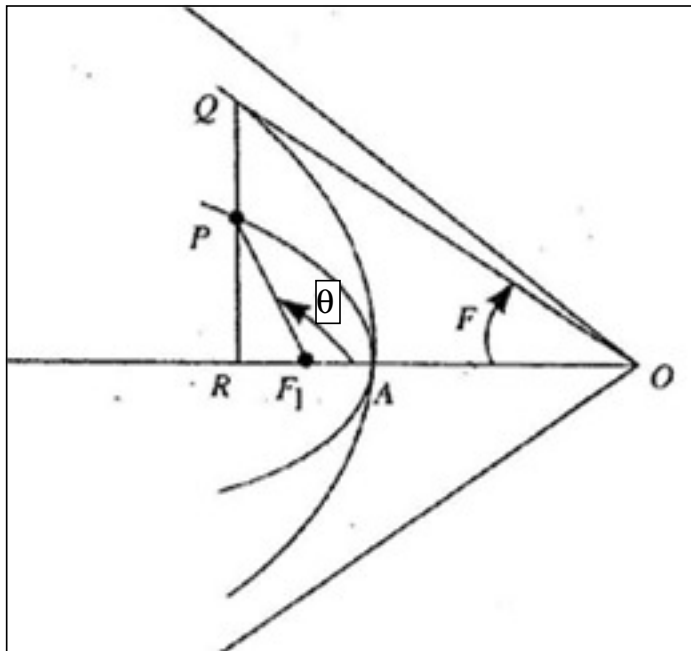
$$\delta = 2 \sin^{-1} \frac{1}{e}$$





Auxiliary Hyperbola Concept

In this case, we employ an **equilateral** (rectangular) hyperbola ($e = \sqrt{2}, \delta = 90^\circ$), as shown below.





Hyperbola Features

We see that the **direction** of eccentricity vector **reverses**, with origin being **outside** of the hyperbola.

This is same as **$a < 0$** and $a(1 - e) < r < \infty$, which permits us to define '**F**' such that $r = a(1 - e \cosh F)$, **similar** to what is done for **elliptic** orbits ($r = a\{1 - e \cos \theta\}$).



Hyperbola Features

It is seen that **minimum** value of 'coshF' is **1**, so that **minimum** value of r is ' $(-a)(e - 1)$ ', which is r_p .

It is to be noted that here, ' F ' can take **any** value between $-\infty$ and $+\infty$.



Hyperbolic Δt Formulation

We can **employ** ellipse – hyperbola (or **conic** section) **similarity** to arrive at **E – F** transformation, as follows.

$$E = \pm iF, i = \sqrt{-1}; \quad 0 < E < 2\pi, \quad -\infty < F < +\infty$$



Hyperbolic Kepler's Equation

This **permits** us to write **hyperbolic** form of Kepler's **equation**, as given below.

$$\cosh F = \frac{e + \cos \theta}{1 + e \cos \theta}; \quad M = e \sinh F - F$$

The sign of '**F**' is resolved on the **basis** of sign of '**θ**'. i.e. **F** < **0** for **θ** < 0, under the **constraint** that get **M** > **F** > **θ**.



Hyperbolic Δt Solution Features

We can now **find** the explicit **expression** ' Δt ' based on the **hyperbolic** relations, by using **applicable** hyperbolic and trigonometric **identities**, which is as given **below**.

$$\Delta t = \left\{ \sqrt{\frac{(-a)^3}{\mu}} \right\} M = \left\{ \sqrt{\frac{(-a)^3}{\mu}} \right\} \left\{ \frac{e(\sqrt{e^2-1}) \sin \theta}{1 + e \cos \theta} - \ln \frac{\sqrt{e+1} + \sqrt{e-1} \tan(\theta/2)}{\sqrt{e+1} - \sqrt{e-1} \tan(\theta/2)} \right\}$$

It is **interesting** to note that above **expression** is also the exact time **integral** for $e > 1$.



Summary

In **conclusion**, we see that TOF for **hyperbolic** trajectories is obtained by a **suitable** extension of the Kepler's equation **derived** for elliptic orbits.

This is **due** to the fact that both these are the **solutions** of the same conic **section** relation.