



# Flight Mechanics/Dynamics

(Course Code: AE 305/305M/717)

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Time: 180 Minutes

**End-Semester Examination**

Total Points: 100

## Instructions

- All questions are mandatory.
- In case a question is missing some data/information, assume the same suitably and clearly mention it in your answer sheet.
- You are only allowed to open lecture slides of the course, any other form of help/reference is not permitted.
- In cases where the answers of two students are found to be copied, both of them will be awarded zero marks for that particular question.
- Answer sheets need to be submitted in a single “Roll\_Number.pdf” format on Moodle.
- You will get 15 minutes duration for submission of your answer sheet on Moodle after the exam time.

1. The coefficients of the characteristic polynomial corresponding to lateral-directional stability of an aircraft are

$$A = 1, \quad B = 9.42, \quad C = 9.48 + N_v, \quad D = 10.29 + 8.4N_v, \quad E = 2.24 - 0.39N_v.$$

Find the range of values of  $N_v$  for which the aircraft will be laterally dynamically stable.

[10]

**Solution:** For dynamic stability of an aircraft, the coefficients of the characteristic equation are required to satisfy the Routh's criteria of stability. As per the Routh's criteria, we get  $A, B, D, E > 0$  and  $R > 0$  where  $R$  is the Routh's discriminant. It can be noted from the given data that  $A > 0$  and  $B > 0$ . Enforcing the conditions on  $D$  and  $E$ , we get

$$N_v > -\frac{10.29}{8.4}, \text{ and } N_v < \frac{2.24}{0.39} \implies N_v \in (-1.225, 5.7435).$$

Furthermore, enforcing the necessary and sufficiency conditions on

$$R = D(BC - AD) - B^2E$$

by substituting for the given data, we get

$$(10.29 + 8.4N_v)(9.42 \times (9.48 + N_v) - 10.29 - 8.4N_v) - 9.42^2(2.24 - 0.39N_v) > 0$$

$$(10.29 + 8.4N_v)(79.01 + 1.02N_v) - 198.77 + 34.61N_v > 0$$

$$813.01 + 674.18N_v + 8.568N_v^2 - 198.77 + 34.61N_v > 0$$

$$8.568N_v^2 + 708.79N_v + 614.24 > 0$$

$$N_v^2 + 82.72N_v + 71.69 > 0$$

Analyzing the roots of the above polynomial, one can conclude that  $R > 0 \forall N_v \in (-\infty, -81.85) \cup (-0.8759, \infty)$ . Therefore, the feasible set of  $N_v$  for which the aircraft has lateral dynamical stability is

$$N_v \in (-1.225, 5.7435) \cap (-0.8759, \infty) \implies N_v \in (-0.8759, 5.7435).$$

2. Consider an aircraft equipped with accelerometers and gyroscopes to measure accelerations and body rates.

- (a) If the aircraft is undergoing a steady rotation with angular velocity components in body axes system  $p = 10$  deg/s,  $q = 2$  deg/s, and  $r = 5$  deg/s, then determine the corresponding Euler angle rates at the time instant where the Euler angles are  $\psi = -30$  deg,  $\theta = 10$  deg, and  $\phi = 15$  deg.
- (b) If the aircraft is flying at an angle of attack of  $10^\circ$ , sideslip of  $5^\circ$ , and a bank angle of  $10^\circ$  and the onboard accelerometers record  $a_{xb} = 10$  ft/s<sup>2</sup>,  $a_{yb} = 5$  ft/s<sup>2</sup>, and  $a_{zb} = -5$  ft/s<sup>2</sup>, then determine the acceleration components in the wind axes system.

[7.5+7.5]

### Solution:

- (a) The given data is as follows:

$$p = 10^\circ/s, \quad q = 2^\circ/s, \quad r = 5^\circ/s, \quad \psi = -30^\circ, \quad \theta = 10^\circ, \quad \phi = 15^\circ.$$

The Euler angle rates can be computed using

$$\begin{bmatrix} \dot{\psi} \\ \dot{\theta} \\ \dot{\phi} \end{bmatrix} = \begin{bmatrix} \frac{q \sin \phi + r \cos \phi}{\cos \theta} \\ q \cos \phi - r \sin \phi \\ p + \tan \theta (q \sin \phi + r \cos \phi) \end{bmatrix}.$$

On substitution of the given data, we get

$$\begin{bmatrix} \dot{\psi} \\ \dot{\theta} \\ \dot{\phi} \end{bmatrix} = \begin{bmatrix} 5.429 \\ 0.637 \\ 10.942 \end{bmatrix} \text{ }^\circ/\text{s}.$$

(b) The rotation matrix from the wind axes to body axes is given by

$$C_w^b = \begin{bmatrix} \cos \alpha \cos \beta & -\cos \alpha \sin \beta & -\sin \alpha \\ \cos \beta \sin \alpha \sin \phi + \sin \beta \cos \phi & -\sin \beta \sin \alpha \sin \phi + \cos \beta \cos \phi & \cos \alpha \sin \phi \\ \cos \beta \sin \alpha \cos \phi - \sin \beta \sin \phi & -\sin \beta \sin \alpha \cos \phi - \cos \beta \sin \phi & \cos \alpha \cos \phi \end{bmatrix}.$$

The side-slip angle  $\beta$ , angle of attack  $\alpha$ , and bank angle  $\phi$  are given as  $5^\circ$ ,  $10^\circ$ , and  $10^\circ$ , respectively. Substituting the same in the above transformation matrix, we get

$$C_w^b = \begin{bmatrix} 0.9811 & -0.0858 & -0.1736 \\ 0.1159 & 0.9784 & 0.1710 \\ 0.1552 & -0.1879 & 0.9698 \end{bmatrix}.$$

The acceleration in the wind axes can be written as  $a_w = C_b^w a_b$ , where  $a_w$  and  $a_b$  are the acceleration in the wind and body axes. Therefore,

$$a_w = C_b^w a_b = [C_w^b]^{-1} a_b = [9.6138 \quad 4.9733 \quad -5.7307]^T \text{ ft/s}^2.$$

3. Answer the following:

(a) Determine the missing elements of the following direction cosine matrix:

$$C = \begin{bmatrix} 0.1587 & c_{12} & 0.4858 \\ 0.8595 & -0.1218 & c_{23} \\ c_{31} & 0.4963 & 0.7195 \end{bmatrix}$$

(b) Consider a rotation of a vector, using quaternion, about an axis defined by the vector  $(1, 0, 0)$  through an angle of  $2\pi/3$ .

(i) Obtain the quaternion  $Q$  to perform this rotation.

(ii) Compute the effect of rotation on the basis vector  $\mathbf{k} = (0, 0, 1)$ .

**Solution:**

- (a) Using the properties that the norm of any column or row of a DCM matrix is unity. Moreover, these column vectors of a DCM matrix are linearly independent. Therefore,

$$c_{12} = \sqrt{1 - 0.1587^2 - 0.4858^2} \approx \pm 0.8595$$

$$c_{23} = \sqrt{1 - 0.8595^2 - 0.1218^2} \approx \pm 0.4963$$

$$c_{31} = \sqrt{1 - 0.4963^2 - 0.7195^2} \approx \pm 0.4858$$

However, utilizing the linear independence property, we can eliminate one of the roots. Hence, after some analysis, we get

$$c_{12} \approx -0.8595, \quad c_{23} \approx -0.4963, \quad c_{31} \approx 0.4858.$$

- (b) Given axis of rotation :  $\hat{q} = (1, 0, 0)$  and angle of rotation:  $\theta = \frac{2\pi}{3}$ . Therefore,

- (i) The resultant Quaternion  $Q$  is given by

$$\begin{aligned} Q &= \cos \frac{\theta}{2} + \hat{q} \sin \frac{\theta}{2} \\ &= \frac{1}{2} + \frac{\sqrt{3}}{2} \hat{i} \end{aligned}$$

- (ii) Effect of rotation on  $\mathbf{k} = (0, 0, 1)$  is given by  $\mathbf{k}' = [Q] \mathbf{k} [Q^*]$ , which can be simplified to

$$\begin{aligned} \mathbf{k}' &= \cos \theta \mathbf{v} + (1 - \cos \theta) (\hat{q} \cdot \mathbf{v}) \hat{q} + \sin \theta (\hat{q} \times \mathbf{v}) \\ &= -\frac{1}{2} \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} + \left(1 + \frac{1}{2}\right) \cdot (0) + \frac{\sqrt{3}}{2} \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix} \times \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} \\ &= \begin{pmatrix} 0 \\ 0 \\ -1/2 \end{pmatrix} - \frac{\sqrt{3}}{2} \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} \\ &= -\frac{\sqrt{3}}{2} \hat{j} - \frac{1}{2} \hat{k} \end{aligned}$$

4. Consider the speed controller, as shown in Fig. 1, with the system output,  $\mathbf{y} = [u \ \gamma]^T$  and the control vector  $\mathbf{c} = [\delta_e \ \delta_p]^T$ . Note that only the elevator input is in feedback

loop. The desired and actual speeds are denoted by  $u_c$  and  $u$ , respectively. Transfer function matrix  $\mathbf{G}(s)$  which relates the control vector,  $\mathbf{c}$ , to the output,  $\mathbf{y}$ , is represented as

$$\mathbf{G}(s) = \begin{bmatrix} G_{u\delta_e} & G_{u\delta_p} \\ G_{\gamma\delta_e} & G_{\gamma\delta_p} \end{bmatrix}.$$

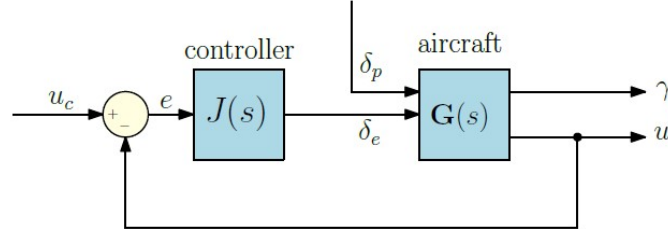


Figure 1: Speed controller

Explicitly derive the *closed-loop transfer functions*, denoted by  $G_{u\delta_p}^*$  and  $G_{\gamma\delta_p}^*$ , corresponding to the throttle input and the outputs, i.e.,  $\delta_p \rightarrow u$  and  $\delta_p \rightarrow \gamma$ .

[10+10]

**Solution:** Using the transfer function between the control vector and output vector, we can write

$$u(s) = G_{u\delta_e}\delta_e + G_{u\delta_p}\delta_p, \quad (1)$$

$$\gamma(s) = G_{\gamma\delta_e}\delta_e + G_{\gamma\delta_p}\delta_p. \quad (2)$$

From Fig. 1 and the above result, we can write

$$\delta_e = J(s)e = J(s)[u_c(s) - u(s)]. \quad (3)$$

Substituting (3) in (1), we get

$$u(s) = G_{u\delta_e}J(s)[u_c(s) - u(s)] + G_{u\delta_p}\delta_p, \quad (4)$$

$$[G_{u\delta_e}J(s) + 1]u(s) = G_{u\delta_e}J(s)u_c(s) + G_{u\delta_p}\delta_p, \quad (5)$$

$$\Rightarrow u(s) = \left[ \frac{G_{u\delta_e}J(s)}{G_{u\delta_e}J(s) + 1} \right] u_c(s) + \left[ \frac{G_{u\delta_p}}{G_{u\delta_e}J(s) + 1} \right] \delta_p. \quad (6)$$

Therefore, the transfer function between the  $u$  and  $\delta_p$  is given by

$$G_{u\delta_p}^* = \frac{G_{u\delta_p}}{G_{u\delta_e}J(s) + 1}.$$

From (2) and (3), we can write

$$\gamma(s) = G_{\gamma\delta_e} J(s) [u_c(s) - u(s)] + G_{\gamma\delta_p} \delta_p.$$

Substituting for (6), we get

$$\begin{aligned} \gamma(s) &= G_{\gamma\delta_e} J(s) \left[ \left( \frac{1}{G_{u\delta_e} J(s) + 1} \right) u_c(s) - G_{u\delta_p}^* \delta_p \right] + G_{\gamma\delta_p} \delta_p, \\ &= \left[ \frac{G_{\gamma\delta_e} J(s)}{G_{u\delta_e} J(s) + 1} \right] u_c(s) - G_{\gamma\delta_e} J(s) G_{u\delta_p}^* \delta_p + G_{\gamma\delta_p} \delta_p, \\ &= \left[ \frac{G_{\gamma\delta_e} J(s)}{G_{u\delta_e} J(s) + 1} \right] u_c(s) + \left[ G_{\gamma\delta_p} - G_{\gamma\delta_e} J(s) G_{u\delta_p}^* \right] \delta_p. \end{aligned}$$

Therefore, the transfer function between the  $\gamma$  and  $\delta_p$  is given by

$$G_{\gamma\delta_p}^* = G_{\gamma\delta_p} - G_{\gamma\delta_e} J(s) G_{u\delta_p}^*.$$

5. Show that the small-disturbance equations, with  $\theta_0 = 0$  and neglecting all  $Y$  force derivatives, yields the following approximation for the lateral displacement:

$$\Delta y_E(t) = g \int_0^t \int_0^t \phi(\tau) d\tau dt.$$

[10]

**Solution:** After ignoring all the  $Y$  derivative terms in the lateral small-disturbance equation and substituting for  $\theta_0 = 0$ , we get

$$\begin{aligned} \dot{v} &= -u_0 r + g\phi, \\ \dot{\psi} &= r, \\ \Delta \dot{y}_E &= u_0 \psi + v. \end{aligned}$$

Differentiating  $\Delta \dot{y}_E$  with respect to time, we get

$$\begin{aligned} \Delta \ddot{y}_E &= u_0 \dot{\psi} + \dot{v} \\ &= u_0 r - u_0 r + g\phi \\ &= g\phi. \end{aligned}$$

Therefore, we can write the displacement in the  $y$  direction as

$$\Delta y_E = g \int_0^t \int_0^t \phi(\tau) d\tau dt.$$

6. Consider a stable system whose transfer function is

$$\mathbf{G}(s) = \frac{Y(s)}{U(s)} = \frac{N(s)}{f(s)} = \frac{1}{(s^2 + 3s + 2)(s^2 + 7s + 12)}.$$

The roots of the characteristic equation are denoted as  $\lambda_1, \lambda_2, \lambda_3$  and  $\lambda_4$ . Find these roots and derive the expression for output  $y(t)$  of the system in time domain, using the partial fraction expansion, for an input given by  $u(t) = e^{i2t}$ . Also, show that the *steady state* output becomes a scaled version of the input along with some phase change depending on the transfer function.

[10]

**Solution:** The roots of the characteristic equation can be found as  $\lambda_1 = -1$ ,  $\lambda_2 = -2$ ,  $\lambda_3 = -3$ , and  $\lambda_4 = -4$ . The Laplace transformation of the input is given by

$$u(s) = \frac{1}{s - 2i}.$$

The output can be written as

$$Y(s) = G(s)u(s) = \frac{1}{(s + 1)(s + 2)(s + 3)(s + 4)(s - 2i)}.$$

Therefore, the control input adds a root  $\lambda_5 = 2i$ . Using partial fraction expansion, the output in time domain can be obtained as

$$\begin{aligned} y(t) &= \sum_{r=1}^5 \left[ \frac{(s - \lambda_r)}{(s + 1)(s + 2)(s + 3)(s + 4)(s - 2i)} \right]_{s=\lambda_r} e^{\lambda_r t} \\ &= \frac{1}{(2i + 1)(2i + 2)(2i + 3)(2i + 4)} e^{i2t} + \frac{1}{6(-1 - 2i)} e^{-t} \\ &\quad - \frac{1}{2(-2 - 2i)} e^{-2t} + \frac{1}{2(-3 - 2i)} e^{-3t} - \frac{1}{6(-4 - 2i)} e^{-4t}, . \end{aligned}$$

As  $t \rightarrow \infty$ , the terms  $e^{-t}$ ,  $e^{-2t}$ ,  $e^{-3t}$  and  $e^{-4t}$  tend to zero. Therefore the steady state response of the system is given by

$$\begin{aligned} y(t) &= \frac{1}{(2i + 1)(2i + 2)(2i + 3)(2i + 4)} e^{i2t}, \\ &\approx 0.0098 \times e^{-i(\phi_1 + \phi_2 + \phi_3 + \phi_4)} e^{i2t}, \\ &\approx 0.0098 \times e^{-i(2.944)} \times u(t), \end{aligned}$$

where  $\phi_1, \phi_2, \phi_3$ , and  $\phi_4$  are the phase angles corresponding to the polar representation of  $2i + 1$ ,  $2i + 2$ ,  $2i + 3$ , and  $2i + 4$ , respectively. Therefore, it can be seen that the steady state response is a scaled version of the input with the output magnitude 0.0098 times the input and phase lagging by  $\phi_1 + \phi_2 + \phi_3 + \phi_4 \approx 168.69$  degrees.

7. Consider the following approximate system of equations, corresponding to the short period mode of an aircraft, given by

$$\begin{aligned}\dot{w} &= \frac{Z_w}{m}w + u_0q, \\ \dot{q} &= \frac{1}{I_y} \left[ M_w + \frac{M_{\dot{w}}Z_w}{m} \right] w + \frac{1}{I_y} [M_q + M_{\dot{w}}u_0] q,\end{aligned}$$

where  $u_0$  is the nominal speed and  $m$  is the mass of the aircraft.

- Find the natural frequency and damping ratio corresponding to the short period mode as a function of  $m, u_0, I_y, M_w, M_q, M_{\dot{w}}$  and  $Z_w$ .
- If an aircraft weighing  $2 \times 10^6$  N is moving at a nominal speed of 230 m/s has the following structural and aerodynamic parameters:  $I_y = 0.5 \times 10^8$  kg m<sup>2</sup>,  $M_w = -1.563 \times 10^4$  Nm,  $M_q = -1.521 \times 10^7$  Nm,  $M_{\dot{w}} = -1.702 \times 10^4$  Nm and  $Z_w = -9.030 \times 10^4$  N, then use the relationships derived in the previous part to compute the natural frequency and damping ratio corresponding to short-period mode of the aircraft. Assume the gravitational acceleration to be 9.81 m/s<sup>2</sup>.

[20]

### Solution:

- The short period dynamics given in the question can be written in the state space form  $\dot{\mathbf{x}} = \mathbf{A}\mathbf{x}$ , where the matrix  $\mathbf{A}$  is given by

$$\mathbf{A} = \begin{bmatrix} Z_w/m & u_0 \\ \frac{1}{I_y} \left[ M_w + \frac{M_{\dot{w}}Z_w}{m} \right] & \frac{1}{I_y} [M_q + M_{\dot{w}}u_0] \end{bmatrix}.$$

The corresponding characteristic equation can be obtained by simplifying

$$\begin{aligned}& \begin{vmatrix} \lambda - Z_w/m & -u_0 \\ -\frac{1}{I_y} \left[ M_w + \frac{M_{\dot{w}}Z_w}{m} \right] & \lambda - \frac{1}{I_y} [M_q + M_{\dot{w}}u_0] \end{vmatrix} = 0 \\ \Rightarrow & \lambda^2 - \left( \frac{Z_w}{m} + \frac{1}{I_y} [M_q + M_{\dot{w}}u_0] \right) \lambda + \frac{Z_w}{mI_y} [M_q + M_{\dot{w}}u_0] - \frac{u_0}{I_y} \left[ M_w + \frac{M_{\dot{w}}Z_w}{m} \right] = 0 \\ \Rightarrow & \lambda^2 - \left( \frac{Z_w}{m} + \frac{1}{I_y} [M_q + M_{\dot{w}}u_0] \right) \lambda + \frac{1}{I_y} \left[ \frac{M_qZ_w}{m} - u_0M_w \right] = 0\end{aligned}$$

On comparison with  $\lambda^2 + 2\zeta\omega\lambda + \omega^2 = 0$ , we get the natural frequency as

$$\omega = \sqrt{\frac{1}{I_y} \left[ \frac{M_qZ_w}{m} - u_0M_w \right]}.$$



Moreover, the damping coefficient is given by

$$\begin{aligned}\zeta &= -\frac{1}{2\omega} \left( \frac{Z_w}{m} + \frac{1}{I_y} [M_q + M_{\dot{w}}u_0] \right) = -\frac{1}{2\omega} \left( \frac{Z_w I_y + M_q m + M_{\dot{w}} u_0 m}{m I_y} \right) \\ &= -\frac{1}{2\sqrt{m I_y}} \left( \frac{Z_w I_y + M_q m + M_{\dot{w}} u_0 m}{\sqrt{M_q Z_w - u_0 M_w m}} \right)\end{aligned}$$

- (b) The mass of the aircraft is given by  $m = (2/9.81) \times 10^6$ . The natural frequency and damping coefficient of the short period are

$$\begin{aligned}\omega &= \sqrt{\frac{1}{I_y} \left[ \frac{M_q Z_w}{m} - u_0 M_w \right]} = 0.4546, \\ \zeta &= -\frac{1}{2\sqrt{m I_y}} \left( \frac{Z_w I_y + M_q m + M_{\dot{w}} u_0 m}{\sqrt{M_q Z_w - u_0 M_w m}} \right) = 0.9079.\end{aligned}$$