



Central Force Motion Solution



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In **order** to solve the two-body **equations**, we first define the vector **product** concept employed for **N-body** case, as shown below.

$$\begin{aligned}\vec{H} &= \vec{r} \times \dot{\vec{r}}; \quad \left(\ddot{\vec{r}} + \frac{\mu}{r^3} \vec{r} \right) \times \vec{H} = 0 \rightarrow \ddot{\vec{r}} \times \vec{H} = -\frac{\mu}{r^3} \vec{r} \times \vec{H} \\ \frac{d}{dt}(\dot{\vec{r}} \times \vec{H}) &= -\frac{\mu}{r^3} [\vec{r} \times (\vec{r} \times \dot{\vec{r}})] = -\frac{\mu}{r^3} [(\vec{r} \cdot \dot{\vec{r}}) \vec{r} - (\vec{r} \cdot \vec{r}) \dot{\vec{r}}] \\ \frac{d}{dt}(\dot{\vec{r}} \times \vec{H}) &= \frac{\mu}{r} \dot{\vec{r}} - \frac{\mu \dot{r}}{r^2} \vec{r} = \frac{d}{dt} \left(\frac{\mu \vec{r}}{r} \right) \rightarrow \dot{\vec{r}} \times \vec{H} - \frac{\mu \vec{r}}{r} = \mu \vec{e}\end{aligned}$$



Central Force Motion Trajectory

Next, we take **scalar** product of the result with **vector** ‘**r**’ and **simplify** the expression, as follows.

$$\vec{r} \cdot \left(\dot{\vec{r}} \times \vec{H} - \frac{\mu \vec{r}}{r} = \mu \vec{e} \right) \rightarrow \vec{r} \cdot (\dot{\vec{r}} \times \vec{H}) - \frac{\mu}{r} = \mu (\vec{r} \cdot \vec{e})$$
$$\vec{r} \cdot (\dot{\vec{r}} \times \vec{H}) = (\vec{r} \times \dot{\vec{r}}) \cdot \vec{H} = h^2; \quad \mu (\vec{r} \cdot \vec{e}) = \mu r e \cos \theta$$
$$r = \frac{\left(\frac{h^2}{\mu} \right)}{1 + e \cos \theta}; \quad \vec{e} : \text{Constant Vector}; \quad \theta: \angle \text{ between } \vec{r} \text{ \& } \vec{e}$$

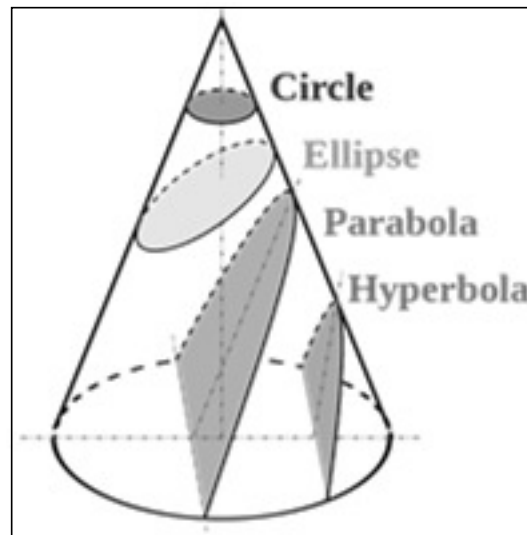
We now have a solution for **radius** magnitude in terms of ‘**θ**’, with ‘**e**’ and ‘**μ**’ as **constants**.



Central Force Motion Features

Solution obtained for ' r ' as a **function** of ' θ ', represents the equation of a '**conic**' section in polar **coordinates**.

'**Conic**' sections are **geometries** that are created from **intersection** of a **plane** with a **cone**, as shown below.





Energy Conservation Solution

We can also obtain the **solution** for energy **conservation**, as outlined below.

$$\begin{aligned}\dot{\vec{r}} \cdot \left(\ddot{\vec{r}} + \frac{\mu}{r^3} \vec{r} \right) &= 0 \rightarrow \dot{\vec{r}} \cdot \ddot{\vec{r}} + \frac{\mu}{r^3} \dot{\vec{r}} \cdot \vec{r} = \frac{d}{dt} \left(\frac{1}{2} \dot{\vec{r}} \cdot \dot{\vec{r}} \right) + \frac{\mu}{r^3} (r\dot{r}) \\ \frac{d}{dt} \left(\frac{1}{2} V^2 \right) + \frac{d}{dt} \left(-\frac{\mu}{r} \right) &= 0 \rightarrow \frac{1}{2} V^2 - \frac{\mu}{r} = \varepsilon \rightarrow \text{A constant}\end{aligned}$$

Conic section solution, along with **energy** conservation solution, are sufficient to **examine** spacecraft motion.



Conic Section Features

$r(t)$ is a **vector** drawn from **focus**, with ' $\theta(t)$ ' (Positive anti-clockwise), **measured** with respect to vector '**e**', which is taken as **one axis** of plane.

Ellipse is the most generic '**conic section**' that is used in orbit **solutions**, as it is able to **capture** the features of the other **conic** sections e.g. circle, parabola and **hyperbola**.

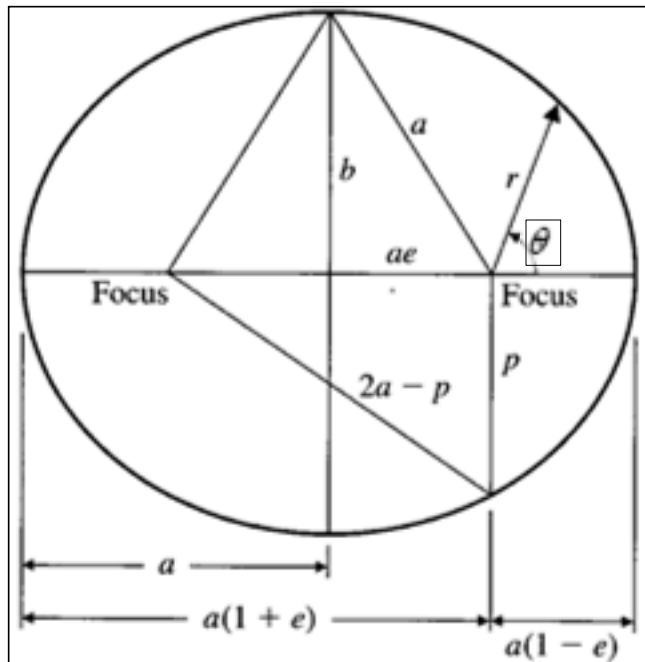


Basic Orbital Solution



Elliptic Orbit Parameters

Consider **general** elliptical **geometry**, as shown below



We can **relate** the ellipse **parameters** with the motion **variables**, as given below.

$$r = \frac{p \left(= \frac{h^2}{\mu} \right)}{1 + e \cos \theta}; \quad p = a(1 - e^2)$$



Ellipse as Generic Conic Section

Ellipse is the basic **conic** section that is applicable to **orbits**. Following are the orbit related **parameters**.

$$\begin{aligned} h &= \sqrt{\mu a(1-e^2)} = r_p v_p = r_a v_a; \quad r_p = a(1-e) \\ r_a &= a(1+e); \quad a = \frac{r_a + r_p}{2}; \quad e = \frac{r_a - r_p}{r_a + r_p} \\ \varepsilon &= \frac{1}{2}v_p^2 - \frac{\mu}{r_p} = \frac{1}{2}v_a^2 - \frac{\mu}{r_a} = -\frac{\mu}{2a}; \quad e = \sqrt{1 + \frac{2\varepsilon h^2}{\mu^2}} \end{aligned}$$

Here, ' ε ' & ' h ' are related to the burnout **parameters** of the corresponding **ascent mission** performance.



Orbit Nature & Parameters

We see that '**a**', which denotes orbit **size**, depends only on total mechanical **energy** imparted by the **ascent mission**.

However, we note that '**e**', which denotes the **shape** of the orbit, depends on both **energy** & angular momentum.



Orbit Nature & Parameters

We also **note** that all **missions** that have 'e' between **0** and **1**, will form the **orbits**.

Therefore, we can **arrive** at the **conditions** for either \bar{r} and \bar{v} , or for **h** and ϵ , for forming an **orbit**, for designing the **ascent** mission.



Bound on Orbits

It should be **noted** here that **$e = 0$** represents lower **limit** that degenerates **into** a circle, as shown below.

$$r_{\text{circular}} = \frac{\left(\frac{h^2}{\mu}\right)}{1 + e \cos \theta} \Big|_{e=0} = \left(\frac{h^2}{\mu}\right) = \frac{r_{\text{circular}}^2 v_{\text{circular}}^2}{\mu} \rightarrow v_{\text{circular}} = \sqrt{\frac{\mu}{r_{\text{circular}}}}$$

Similarly, **$e = 1$** represents the upper **limit** that degenerates into a **parabola**, as shown below.

$$e = 1 \rightarrow \varepsilon = 0 \rightarrow \frac{1}{2} v_{\text{parabolic}}^2 = \frac{\mu}{r_{\text{parabolic}}} \rightarrow v_{\text{parabolic}} = \sqrt{\frac{2\mu}{r_{\text{parabolic}}}}$$



Non-orbital Trajectories

In case of $e > 1$, equation of **conic** corresponds to a **hyperbola**, as follows.

$$e > 1 \rightarrow r = \infty \text{ for } 1 + e \cos \theta = 0 \rightarrow \theta < 180^\circ \rightarrow \text{hyperbola}$$

In **all** such cases, the **spacecraft** escapes from earth's **gravitational** field and attains **inter-planetary** path.



Non-orbital Trajectories

However, in case of $e < 0$, definitions of ' r_a ' and ' r_p ' **inter-change**, resulting in a **reformulation** of problem.

This can be **understood** by assuming that $(2\epsilon h^2/\mu^2) < -1$ so that $e^2 < 0$, making ' e ' a complex number, and **leading** to situations where $a < R_E$.

Thus, in such **cases**, the object will **fall** back to earth, as per the various **forces** and entry initial **conditions**.



Summary

We see that two-body **equations** can be solved in the **closed** form implicitly, through **simple** vector operations.

We **further** note that ellipse is the **basic** geometry that is applicable in the **context** of space objects forming **orbits**.