

# Gompertz Processes: A Theory of Ageing

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## Abstract

Ageing is a accelerating failure process that is universal to all biological systems. Motivated by considering infinitesimal stochastic accelerations of time, we hypothesize a general explanation for the emergent determinism of ageing processes in the theory of Gompertz processes: Poisson processes subordinated by integrated Geometric Brownian motion.

## 1 Preliminaries

Basic science experiments in biology have ubiquitously observed that organisms respond to environmental stresses, including communicable diseases and exposures to toxins, with an acceleration  $a$  in their failure time  $ah(at)$ , where  $h$  is the bare hazard rate of the failure event. Outside of the controlled setting of a laboratory the environmental stresses occur stochastically, resulting in a sequence of accelerations  $a_n$  of the organism's metabolic time. The accelerations can either increase the rate of failure  $a_n > 1$ , an exacerbation of the stresses, or decrease the rate of failure  $a_n < 1$ , an alleviation from the stresses. However, even in the controlled setting of a laboratory it is experimentally challenging to directly measure the organism's accelerated metabolic time of the organism, instead we only have access to the failure times in bare units. Thus a theory of ageing must be one that studies the subordination of the failure time by an accelerated metabolic time that is stochastically incremented.

Over the course of an organism's lifetime it will encounter exacerbations and alleviations that accelerate  $a_n, \dots, a_0 = 1$  metabolic time at times  $t_n, \dots, t_0 = 0$ . Furthermore, because the accelerations are all positive  $a_i > 0$  for each acceleration we can find a finite real valued generator  $w_n, \dots, w_0 = 0$  such that  $a_i = e^{w_i}$ . It follows that the accelerated metabolic time  $y_{t_n}$  at time  $t_n$  is the sum of the products of the accelerations up to time  $t_{n-1}$  and the elapsed bare time steps  $t_i - t_{i-1}$ :

$$y_{t_n} = \sum_{i=1}^n e^{\sum_{j=0}^{i-1} w_j} (t_i - t_{i-1}) \quad (1)$$

Provided no great explosions of acceleration occur in any small time scale, like say when an actual explosion occurs, the generators  $w_i$  will become infinitesimal

on the same order as  $t_i - t_{i-1} \rightarrow 0$ :

$$\mathcal{O}(w_i) = \mathcal{O}(t_i - t_{i-1}) \quad (2)$$

The sum of products then becomes a stochastic process  $Y_t$  that is an integral of a random geometric infinitesimal generator process  $e^{W_u}$ :

$$Y_t = \int_0^t e^{W_u} du \quad (3)$$

This is essentially the continuous part of the Kolmogorov's characterization of stochastic processes as being composed of either a finite number of discrete jumps or an infinite number of continuous changes in a span of time. If we assume as a first approximation that the exacerbations and alleviations, and their respective accelerations, are independent and stationary over time then by the Lévy-Khintchine characterization the only infinitesimal generator of accelerated metabolic time that is Lévy and continuous “*jump free*” is Brownian motion  $W_u$ . The stochastic process of accelerated metabolic time is better known as integrated geometric Brownian motion.

## 2 Gompertz

Motivated by the preceding heuristic derivation we formally define the Gompertz process on time  $t \geq 0$ .

**Definition 1** (Gompertz Process). A Gompertz process  $G_t$  is a subordinated Poisson process  $N_t$ , with rate  $\lambda$ , where the subordinating process is integrated geometric Brownian motion  $Y_t$ , with drift  $\mu$  and diffusion  $\sigma$ :

$$G_t = N_{Y_t} \quad (4)$$

given:

$$Y_t = \int_0^t X_s ds \quad (5)$$

$$= \int_0^t e^{\mu s + \sigma W_s} ds \quad (6)$$

Phenomenologically the finite real stochastic process  $\mu t + \sigma W_s$  is the infinitesimal acceleration at time  $t$  that generates a non-negative geometric stochastic process  $X_t$  of accumulated accelerations up to time  $t$  and whose integral  $Y_t$  is a strictly increasing stochastic process of accelerated metabolic time up to time  $t$ .

From the definition of the Gompertz process  $G_t$  conditioning on the history of a sample path of integrated geometric Brownian motion  $Y_t$  yields a conditional Poisson process:

$$\mathbb{P}[G_t = n \mid Y_t] = \frac{(\lambda Y_t)^n}{n!} e^{-\lambda Y_t} \quad (7)$$

$$\mathbb{E}[G_t \mid Y_t] = \lambda Y_t \quad (8)$$

And the dual relationship that conditioning on the Gompertz process estimates the elapsed stochastically accelerated metabolic time:

$$\mathbb{P}[Y_t \leq y \mid G_t] = \int_0^{\lambda y} \frac{u^{G_t-1} e^{-u}}{(G_t-1)!} du \quad (9)$$

$$\mathbb{E}[Y_t \mid G_t] = \frac{G_t}{\lambda} \quad (10)$$

Reflecting that integrated geometric Brownian motion  $Y_t$  is truly measuring the elapsed stochastically accelerated metabolic time.

To start our exploration of the rich subtleties of Gompertz processes we will briefly review the properties of integrated geometric Brownian motion which are salient to developing our theory. This is by no means a comprehensive compendium. Much of the material I will cover has been deeply and thoroughly explored in the quantitative finance literature in the theory of pricing Asian options, and the stochastic processes syllabus on subordinated Poisson processes, usually in the context of deriving the characteristic function of the subordinating process.

Our first observation is that the increments of  $Y_t$  can be factored by its carrier process  $X_t$ , for times  $t > s$ :

$$\mathbb{P}[Y_t - Y_s] = \mathbb{P}[X_s] \mathbb{P}[Y_{t-s}] \quad (11)$$

where the process  $X_s$  is independent of the process  $Y_{t-s}$ . Phenomenologically the increments of  $Y_t - Y_s$  are equivalent to a process that starts with acceleration  $X_s$ . Analogous to the Fundamental Theorem of Arithmetic this factorization allows us to reach many useful inferences; for example we can immediately observe that for times  $t > s$ :

$$\mathbb{E}[(Y_t - Y_s)^n \mid X_s] = X_s^n \mathbb{E}[Y_{t-s}^n] \quad (12)$$

We will liberally exploit this technique of arbitraging the accelerated metabolic time  $Y_t$  against the accumulated acceleration  $X_t$  to reduce expectations down to the well known standard terms for  $X_t$  and  $Y_t$ :

$$\mathbb{E}[X_t^n] = e^{(n\mu + n^2\sigma^2/2)t} \quad (13)$$

$$\mathbb{E}[Y_t] = \frac{e^{(\mu + \sigma^2/2)t} - 1}{\mu + \sigma^2/2} \quad (14)$$

Note that in the last equation we have implicitly invoked the Fubini-Tonelli Theorem to switch the order of integration, and will broadly continue to use this theorem throughout this work.

**Lemma 1** (Acceleration Lemma). *The expectation of non-negative integer powers  $n \geq 0$  of the carrier process of geometric Brownian motion  $X_t$  conditioned on integrated geometric Brownian motion  $Y_t$  is given by:*

$$\mathbb{E}[X_t^n \mid Y_t] = \frac{\mathbb{E}[X_t^n]}{\mathbb{E}[Y_t^n]} Y_t^n \quad (15)$$

*Proof.* We proceed with induction on  $n$ .

1. For  $n = 0$  clearly  $\mathbb{E}[X_t^0 \mid Y_t] = 1$
2. Now assume that up to  $n$  the relation holds.
3. For  $u < t$  consider the integrable function

$$f(u, t, y) = \mathbb{E}[X_u X_t^n \mid Y_t = y] \quad (16)$$

4. It follows from the Fubini-Tonelli Theorem and the induction assumption that for  $n$  the relation holds

$$\int_0^t f(u, t, y) du = y^{n+1} \frac{\mathbb{E}[X_t^n]}{\mathbb{E}[Y_t^n]} \quad (17)$$

5. Thus by the Fundamental Theorem of Calculus the integrand has the decomposition  $f(u, t, y) = y^{n+1} f(u, t)$
6. By the Law of Conditional Expectations we have

$$\mathbb{E}[X_u X_t^n] = \mathbb{E}[\mathbb{E}[X_u X_t^n \mid Y_t]] \quad (18)$$

$$= f(u, t) \mathbb{E}[Y_t^{n+1}] \quad (19)$$

7. Which immediately yields the decomposed integrand as

$$f(u, t) = \frac{\mathbb{E}[X_u X_t^n]}{\mathbb{E}[Y_t^{n+1}]} \quad (20)$$

8. Finally taking  $u = t$  yields the desired result

$$\mathbb{E}[X_t^{n+1} \mid Y_t] = f(t, t, t) Y_t^{n+1} \quad (21)$$

$$= \frac{\mathbb{E}[X_t^{n+1}]}{\mathbb{E}[Y_t^{n+1}]} Y_t^{n+1} \quad (22)$$

□

A simple corollary follows from a nearly trivial derivation.

**Corollary 1** (Counterpoint Corollary). *For a times  $t > s$  we have the following conditional expectation:*

$$\mathbb{E}[Y_t^n \mid Y_s] = \frac{\mathbb{E}[Y_t^n]}{\mathbb{E}[Y_s^n]} Y_s^n \quad (23)$$

*Proof.* Factoring and applying the previous lemma:

$$\mathbb{E}[Y_t^n \mid Y_s] = \mathbb{E}[X_s^n \mid Y_s] \mathbb{E}[Y_{t-s}^n] \quad (24)$$

$$= \frac{\mathbb{E}[Y_{t-s}^n] \mathbb{E}[X_s^n]}{\mathbb{E}[Y_s^n]} Y_s^n \quad (25)$$

$$= \frac{\mathbb{E}[Y_t^n]}{\mathbb{E}[Y_s^n]} Y_s^n \quad (26)$$

□

Even with the factorization observation we are still in need of a means of reducing the expectation of the powers  $Y_t^n$ . A small lemma suffices to provide the means of finding powers:

**Lemma 2** (Recursion-Convolution Lemma). *The expectation of non-negative integer powers  $n \geq 0$  of integrated geometric Brownian motion  $Y_t$  is given by:*

$$\mathbb{E}[Y_t^n] = n \int_0^t \mathbb{E}[X_u^n] \mathbb{E}[Y_{t-u}^{n-1}] du \quad (27)$$

*Proof.* Expanding the expectant as a multi-variable integral we have, applying the Fubini-Tonelli Theorem we have, and then factoring we have:

$$\mathbb{E}[Y_t^n] = \int_0^t \cdots \int_0^t \mathbb{E}[X_t^n] du_1 \dots du_n \quad (28)$$

$$= \binom{n}{1} \int_0^t \mathbb{E}[X_u (Y_t - Y_u)^{n-1}] du \quad (29)$$

$$= n \int_0^t \mathbb{E}[X_u^n] \mathbb{E}[Y_{t-u}^{n-1}] du \quad (30)$$

□

With the preceding lemma in hand we have the sufficient tools required to estimate all the usual statistics involving powers of  $Y_t$ , including the expectation, variance, and covariances.

### 3 Martingale

By our counterpoint corollary the amortized integrated geometric Brownian motion:

$$\hat{Y}_t = \frac{Y_t}{\mathbb{E}[Y_t]} \quad (31)$$

forms a martingale, so that for times  $t > s$ :

$$\mathbb{E}[\hat{Y}_t \mid \hat{Y}_s] = \hat{Y}_s \quad (32)$$

We can deduce that the amortized Gompertz process:

$$\hat{G}_t = \frac{G_t}{\mathbb{E}[Y_t]} \quad (33)$$

is a martingale, so that for times  $t > s$ :

$$\mathbb{E}[\hat{G}_t \mid \hat{G}_s] = \hat{G}_s \quad (34)$$

This allows us to leverage optional stopping time theorems to evaluate expectations of stopped versions of Gompertz processes.

## 4 Markov

From the Martingale  $Y_{s \uparrow t}$  we can infer that  $Y_t$  is a 2 step Markov process, and that the increments  $Y_t - Y_s$ , with  $t > s$  are Markov. To verify this consider the sequence of bare times  $0 = t_0 < \dots < t_{n+1} = t$ , and accelerated times  $0 = y_0 < \dots < y_{n+1} = y$ , working through the conditional probability we have:

$$\begin{aligned} \mathbb{P}[Y_{t_{n+1}} - Y_{t_n} = y_{n+1} - y_n \mid Y_{t_n} - Y_{t_{n-1}} = y_n - y_{n-1}, \dots, Y_{t_1} - Y_{t_0} = y_1 - y_0] \\ = \int_0^\infty \int_0^\infty \mathbb{P}\left[X_{t_n - t_{n-1}} = \frac{u}{v} \frac{y_{n+1} - y_n}{y_n - y_{n-1}}\right] \\ \cdot \mathbb{P}[Y_{t_n - t_{n-1}} = u] \\ \cdot \mathbb{P}[Y_{t_{n+1} - t_n} = v] dudv \end{aligned} \quad (35)$$

$$= \mathbb{P}[Y_{t_{n+1}} - Y_{t_n} = y_{n+1} - y_n \mid Y_{t_n} - Y_{t_{n-1}} = y_n - y_{n-1}] \quad (36)$$

In practice we cannot directly measure the accelerated metabolic time  $Y_t$  instead we have access to the stopping times  $T_n$  of the passages  $G_{T_n} = n$  of the Gompertz process.

To study the Markov properties of the stopping times of the passages of the Gompertz process we revisit the recovery of the Poisson process from the Gompertz process by conditioning on the history of the sample path of integrated geometric Brownian motion. By differentiating the subordinated cumulative distribution of a single stopping time we have the conditional probability distribution:

$$\mathbb{P}[T_n = t \mid Y_{T_n}] = \frac{\mathbb{E}[X_t]}{\mathbb{E}[Y_t]} \frac{(\lambda Y_t)^n}{(n-1)!} e^{-\lambda Y_t} \quad (37)$$

Finding a closed form for the distribution of the stopping times of the passages of the Gompertz process will require a deeper understanding of the characteristic function of integrated geometric Brownian motion, which we will develop later. In the meantime there are many fruits to be plucked from a study of stopping times of the passages of the Gompertz process.

The central statistic of study in the longitudinal analysis of biological systems is the latency  $T_{1+G_t} - T_{G_t} = T_{1+G_t} - t$  between consecutive stopping times of the passages of the Gompertz process. Due to immortal time bias “*we cannot see indefinitely into the past*” we have subordinated the stopping times  $T_{n+G_t}$  of the passages of the Gompertz process by the increments of the Gompertz process  $n+G_t$  from a sentinel event  $G_t$ . In practice observational studies, particularly in clinical research and epidemiology, are only able to observe consecutive passages of the Gompertz process from a fixed sentinel event without the knowledge of how many events have occurred before the sentinel event.

From the preceding probability density we can immediately deduce the tail probability and hence the expectation of the latency  $s > 0$  conditioned on a

sentinel event at time  $t > 0$ :

$$\mathbb{P}[T_{1+G_t} - T_{G_t} \geq s \mid T_{G_t} = t] = \mathbb{E}[e^{-\lambda X_t Y_s}] \quad (38)$$

$$\mathbb{E}[T_{1+G_t} - T_{G_t} \mid T_{G_t} = t] = \int_0^\infty \mathbb{E}[e^{-\lambda X_t Y_s}] ds \quad (39)$$

where  $Y_s$  is the increment process independent of the acceleration  $X_t$  at the start of the increment.

Remarkably the stopping times of the passages for the Gompertz process are Markov in their absolute values, even though the accelerated metabolic time of integrated geometric Brownian motion is only Markov in its increments. Specifically knowledge of the age  $t_0 > 0$  at a sentinel event  $G_{t_0}$  in the history of the Gompertz process is sufficient to determine the distribution of the latency  $t > 0$  to next event  $T_{1+n+G_{t_0}}$ . To see this consider the subordinated stopping times  $T_{n+G_{t_0}} = t_n, \dots, T_{G_{t_0}} = t_0$ , the cumulative probability of the latency conditioned of the previous events is:

$$\begin{aligned} \mathbb{P}[T_{1+n+G_t} - T_{n+G_t} \geq t \mid T_{n+G_{t_0}} = t_n, \dots, T_{G_{t_0}} = t_0] \\ = \mathbb{E}\left[e^{-\lambda(Y_{t+t_n} - Y_{t_n})} \mid T_{n+G_{t_0}} = t_n, \dots, T_{G_{t_0}} = t_0\right] \end{aligned} \quad (40)$$

$$= \mathbb{E}[e^{-\lambda X_{t_n} Y_t}] \quad (41)$$

$$= \mathbb{P}[T_{1+G_{t_n}} - T_{G_{t_n}} \geq t \mid T_{G_{t_n}} = t_n] \quad (42)$$

where  $Y_t$  is the increment process independent of the acceleration  $X_{t_n}$  at the start of the increment.

TODO conditional correlation

The immediate prediction from these results is that age of an organism introduces a correlation between the latencies of consecutive events, and that this correlation can be fully removed by conditioning on the age of the organism at the sentinel event. However to date not a single longitudinal study of a biological system has properly accounted for the Gompertz process of ageing.

Analogous to the Hubble constant which describes the intrinsic expansion of space-time the Gompertz processes introduces an irreducible exponential dependence on age *“older organisms are red-shifted with respect to younger organisms”* that cannot be removed or linearized by a coordinate transform, just as no co-ordinate transform can remove the intrinsic red-shift of space-time. The serious consequence of this observation is that every single longitudinal study that has followed outcomes whose latencies are on the same timescale as the lifespan of the investigated organism has introduced spurious correlations into their longitudinal analysis of event latency which are purely an artefact of the exponential age dependence of the Gompertz process. Fortunately the majority of the studies of biological systems have either been cross sectional or of short enough duration that the hazard of the Gompertz process is approximately constant over the timescale of the followup on the organism.

## 5 Lévy

Integrated geometric Brownian motion  $Y_t$  and its carrier process of geometric Brownian motion  $X_t$  can be embedded into the Lie algebra of  $2 \times 2$  upper triangular matrices  $\mathfrak{h}_2$  by means of the factorization observed earlier, so that the increments are independent under matrix multiplication, for times  $t > s$ :

$$\begin{bmatrix} 1 & Y_t \\ 0 & X_t \end{bmatrix} = \begin{bmatrix} 1 & Y_{t-s} \\ 0 & X_{t-s} \end{bmatrix} \begin{bmatrix} 1 & Y_s \\ 0 & X_s \end{bmatrix} \quad (43)$$

where the increment processes of  $Y_{t-s}$  and  $X_{t-s}$  are independent of the processes  $Y_s$  and  $X_s$ .

## 6 Fokker-Planck

The upper triangular Lévy process of integrated geometric Brownian motion satisfies the stochastic differential equation:

$$d \begin{bmatrix} 1 & Y_t \\ 0 & X_t \end{bmatrix} = \begin{bmatrix} 0 & X_t \\ 0 & \left(\mu + \frac{\sigma^2}{2}\right) X_t \end{bmatrix} dt + \begin{bmatrix} 0 & 0 \\ 0 & \sigma X_t \end{bmatrix} dW_t \quad (44)$$

It follows from Fokker-Planck that the probability density of the joint process  $p = \mathbb{P}[Y_t = y, X_t = x]$  satisfies the partial differential equation:

$$\frac{\partial p}{\partial t} = -\frac{\partial}{\partial x} \left( \mu + \frac{\sigma^2}{2} \right) xp - \frac{\partial}{\partial y} xp + \frac{\partial^2}{\partial x^2} \frac{\sigma^2}{2} x^2 p \quad (45)$$

We can restate this as an Eigen evolution equation:

$$\frac{\partial p}{\partial t} + \left( \mu + \frac{3}{2}\sigma^2 \right) x \frac{\partial p}{\partial x} + x \frac{\partial p}{\partial y} - \frac{\sigma^2}{2} x^2 \frac{\partial^2 p}{\partial x^2} = \left( -\mu + \frac{\sigma^2}{2} \right) p \quad (46)$$

Marginalizing over the probability of  $X_t$  and applying the acceleration lemma yields the first order partial differential equation for the distribution of  $Y_t$ :

$$\frac{\mathbb{E}[Y_t]}{\mathbb{E}[X_t]} \frac{\partial p}{\partial t} - y \frac{\partial p}{\partial y} = p \quad (47)$$

Which by trial solution of separation of variables has the general solution:

$$p = \mathbb{E}[Y_t] f_{\mu,\sigma}(y \mathbb{E}[Y_t]) \quad (48)$$

for any analytic  $f_{\mu,\sigma}$  dependent on the drift and diffusion of the carrier process  $X_t$ . Note that we are offloading the dimensional analysis into the analytic function.



## 7 Hazard Rate

Consider the first passage stopping time  $T_1$  of the Gompertz process  $G_t$ , its cumulative distribution is the characteristic function of  $Y_t$ :

$$\mathbb{P}[T_1 \geq t] = \mathbb{E}[e^{-\lambda Y_t}] \quad (49)$$

Thus the hazard rate  $h$  of  $T_1$ :

$$h = \mathbb{P}[T_1 = t | T_1 \geq t] \quad (50)$$

satisfies the partial differential Eigen equation:

$$\frac{\mathbb{E}[Y_t]}{\mathbb{E}[X_t]} \frac{\partial h}{\partial t} - \lambda \frac{\partial h}{\partial \lambda} = \left( \frac{\mathbb{E}[Y_t]}{\mathbb{E}[X_t]} \right)^2 \left( \frac{\partial}{\partial t} \frac{\mathbb{E}[X_t]}{\mathbb{E}[Y_t]} \right) h \quad (51)$$

Which by trial solution of separation of variables has the general solution:

$$h = \lambda^2 \mathbb{E}[X_t] \mathbb{E}[Y_t] g_{\mu, \sigma}(\lambda \mathbb{E}[Y_t]) \quad (52)$$

for any analytic  $g_{\mu, \sigma}$  dependent on the drift and diffusion of the carrier process  $X_t$ . Taking the limit to deterministic subordination yields the constraints on  $g_{\mu, \sigma}$ :

Table 1: Sequential Boundary Conditions

Boundary	Condition	Removes
$\sigma = 0$	$h = \lambda e^{\mu t}$	diffusion
$\mu = 0$	$h = \lambda$	then drift
$\lambda = 0$	$h = 0$	finally jumps

Sequential boundary conditions on the hazard rate  $h$  derived from the limits to deterministic subordination.

From the boundary conditions we can immediately deduce that in the deterministic limit of  $\sigma \rightarrow 0$  we have:

$$g_{\mu, \sigma}(x) \xrightarrow{\sigma=0} \frac{1}{x} \quad (53)$$

However this alone cannot be the solution as it results in the characteristic function in  $\lambda$  of a purely deterministic  $Y_t$ . Equating the general solution for the hazard rate to the Laplace of the Fokker-Planck solution yields the implicit equation in  $f_{\mu, \sigma}$  and  $g_{\mu, \sigma}$ :

$$-\frac{\partial}{\partial t} \ln \int_0^\infty e^{\frac{\lambda u}{\mathbb{E}[Y_t]}} f_{\mu, \sigma}(u) du = \lambda^2 \mathbb{E}[X_t] \mathbb{E}[Y_t] g_{\mu, \sigma}(\lambda \mathbb{E}[Y_t]) \quad (54)$$

Dimensional analysis provides an inference for a solution, which remains an open problem:

**Proposition 1** (Gompertz Anomaly). *The analytic function  $g_{\mu,\sigma}$  is simply the exponential divided by its argument, so that the hazard  $h$  is given by:*

$$h = \lambda \mathbb{E}[X_t] e^{\frac{\sigma^2/2}{\mu+\sigma^2/2} \lambda \mathbb{E}[Y_t]} \quad (55)$$

*Proof.* The parity multiplied derivatives of the exponential of the integral of the hazard rate satisfies the recursion-convolution lemma, generating the moments of  $Y_t$ , and hence is the characteristic function of  $Y_t$ .  $\square$

## 8 Discussion

In most circumstances  $\lambda$  is the new born infant mortality due to ageing alone, and is less than 1 in 32000 person-years. As such the hazard is very close to the original hazard observed by Gompertz. Intuitively, when  $\mu \gg \sigma^2/2$  the process becomes approximately deterministic due to the large impact of the drift.

Conversely, the central conjecture in the preliminary material is that ageing is driven by stochastic accelerations, hence requiring that the drift vanish  $\mu = 0$ . In this case the anomalous Gompertz hazard rate simplifies to:

$$h = \lambda e^{t\sigma^2/2} e^{\lambda \frac{e^{t\sigma^2/2}-1}{\sigma^2/2}} \quad (56)$$

We have observed the anomalous Gompertz hazard rate in mortality rates in Alberta in the Hazard Rate Zoo experiment. Specifically when we remove the dominate exponential process of a doubling of mortality every 7 years from an infant mortality of 1 in 32000 person-years:

$$h = \frac{2^{t/7}}{32000} e^{\frac{7(2^{t/7}-1)}{32000 \ln 2}} \quad (57)$$

there remains a residual anomalous growth in mortality with ageing, reflecting the higher order affect of the stochastic accelerations.