

→ CAAM 336 FINAL REVIEW ←

→ THINGS FROM THE 1st HALF TO REMEMBER:

- Eigenvalues / Eigenfunctions

↳ be able to compute them etc.

- Best Approximation from a subspace:

especially notice that $\psi_1, \psi_2, \psi_3, \dots$ forms a basis for V of eigenfunctions, then $f \in V$ can be written as:

$$f = \sum_{k=1}^{\infty} \frac{(f, \psi_k)}{(\psi_k, \psi_k)} \psi_k$$

- Weak form of diff. eq. & Finite Element Method (steady state)

↳ Final may cover this topic begins just before midterm.

→ 2nd HALF OF COURSE: TIME DEPENDENT PROBLEMS

GENERIC MATRIX PROBLEM:

$$u'(t) = Au(t)$$

Solution: $u(t) = e^{tA} u(0)$

↳ Matrix exponential e^{tA}

- If $A = V\Lambda V^{-1}$, then

$$e^{tA} = Ve^{t\Lambda}V^{-1}, \quad \Lambda = \begin{bmatrix} \lambda_1 & & 0 \\ & \ddots & \\ 0 & & \lambda_n \end{bmatrix}$$

$$e^{t\Lambda} = \begin{bmatrix} e^{t\lambda_1} & & 0 \\ & \ddots & \\ 0 & & e^{t\lambda_n} \end{bmatrix}$$

OR

$$e^{tA} = I + (tA) + \frac{1}{2}(tA)^2 + \frac{1}{3!}(tA)^3 + \dots$$

How does $u(t)$ behave as $t \rightarrow \infty$?

- How does e^{tA} behave as $t \rightarrow \infty$?

- If ALL eigenvalues of A have neg. real part,

THEN... $e^{tA} \rightarrow 0$ (decay)

↳ $(\text{Re}(\lambda_j) < 0 \forall j)$

(If A is symmetric, λ_j is real,
So $\text{Re}(\lambda_j) = \lambda_j$)

- If ANY eigenvalue of A has positive real part,

THEN... e^{tA} blows up as $t \rightarrow \infty$

- If $\text{Re}(\lambda_j) \leq 0 \forall \lambda_j$, and $\text{Re}(\lambda_k) = 0$ for some k ,

THEN... e^{tA} neither decays nor blows up!

Think about...

$A = -M^{-1}K$ from heat equation

$\frac{du}{dt} A = \begin{bmatrix} 0 & I \\ M^{-1}K & 0 \end{bmatrix}$ from wave equation

PROBLEMS WITH FORCING FUNCTIONS:

$$u'(t) = Au(t) + f(t)$$

General solution:

$$u(t) = e^{tA} u(0) + \int_0^t e^{(t-s)A} f(s) ds$$

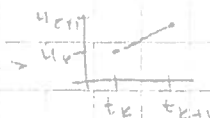
APPROXIMATE SOLUTION IN TIME:

$$u'(t) = Au(t); u_k \approx u(\Delta t k)$$

• Forward Euler:

$$u_{k+1} = u_k + \Delta t (A u_k)$$

$$\Rightarrow u_k = (I + \Delta t A)^k u_0$$



• Backward Euler:

$$u_{k+1} = u_k + \Delta t A u_{k+1}$$

$$(I - \Delta t A) u_{k+1} = u_k$$

$$\Rightarrow u_{k+1} = (I - \Delta t A)^{-1} u_k$$

$$\Rightarrow u_k = ((I - \Delta t A)^{-1})^k u_0$$

"implicit system"

$$\rightarrow \text{MATLAB: } u_{k+1} = (I - \Delta t A) \backslash u_k$$

Both methods give problems like:

$$u_{k+1} = B^k u_0, \text{ for some matrix } B.$$

- When does $B^k \rightarrow 0$ as $k \rightarrow \infty$?

When all eigenvalues λ of B satisfy $|\lambda| < 1$

- When does B^k blow up as $k \rightarrow \infty$?

If $|\lambda| > 1$ for any eigenvalue λ of B .

OK.

Let $\lambda_1, \dots, \lambda_n$ be eigenvalues of A

FW Euler

$$B = I + \Delta t A \rightarrow \text{eigs of } B: \beta_j = 1 + \Delta t \lambda_j$$

BW Euler

$$B = (I - \Delta t A)^{-1} \rightarrow \text{eigs of } B: \beta_j = \frac{1}{1 - \Delta t \lambda_j}$$

Suppose A is symmetric w/ negative eigenvalues, $\lambda_j < 0 \quad \forall j$

THEN $e^{tA} \rightarrow 0$ as $t \rightarrow \infty$

Thus, $u(t) = e^{tA} u(0)$ = exact solution decays.

We want the numerical approximation u_k to decay as well.

• FW Euler decays provided: $|1 + \Delta t \lambda_j| < 1 \quad \forall j$

$$-1 < 1 + \Delta t \lambda_j < 1$$

$$-2 < 1 + \Delta t \lambda_j < 0$$

$$\Rightarrow \frac{-2}{\lambda_j} > \Delta t > 0 \quad \} \text{ We must pick } \Delta t \text{ s.t. } 0 < \Delta t < \frac{-2}{\lambda_j} \quad \forall j$$

↳ if some λ_j are large & negative this imposes a severe restriction on Δt .

• BW Euler decays provided: $|\frac{1}{1 - \Delta t \lambda_j}| < 1$

which holds for any Δt !

SOLUTION OF TIME DEPENDENT PROBLEMS BY EIGENFUNCTIONS:

L = linear operator, symmetric

$\lambda_1, \lambda_2, \dots$ eigenvalues, real

ψ_1, ψ_2, \dots eigenfunctions, orthogonal.

We wish to solve $u_t = -Lu$

① Write the solution in the form $u(x, t) = \sum_{j=1}^{\infty} a_j(t) \psi_j(x)$

② Determine ODE's that the coefficients a_j must satisfy.

Substitute $u(x, t)$ into $u_t = -Lu$

$$\frac{\partial}{\partial t} \left(\sum_{j=1}^{\infty} a_j(t) \psi_j(x) \right) = -L \left(\sum_{j=1}^{\infty} a_j(t) \psi_j(x) \right)$$

$$\begin{aligned} \sum_{j=1}^{\infty} a_j'(t) \psi_j(x) &= - \sum_{j=1}^{\infty} a_j(t) L \psi_j(x) \\ &= - \sum_{j=1}^{\infty} a_j(t) \lambda_j \psi_j(x) \end{aligned}$$

TAKE I.P. w/ ψ_k :

$$(\sum a_j'(t) \psi_j(x), \psi_k(x)) = - (\sum a_j(t) \lambda_j \psi_j(x), \psi_k(x))$$

$$\sum a_j'(t) (\psi_j, \psi_k) = - \sum a_j(t) \lambda_j (\psi_j, \psi_k)$$

AND... $(\psi_j, \psi_k) = 0$ w/ $j \neq k$ so!

$$\boxed{a_k'(t) = -\lambda_k a_k(t)}$$

③ Solve the ODE's w/ initial condition from $u(x, 0) = u_0$

$$a_j(0) = \frac{(u_0, \psi_j)}{(\psi_j, \psi_j)} \quad \} \text{ Best approx!}$$

$$\text{SOLUTION: } \boxed{a_k(t) = e^{-\lambda_k t} a_k(0)} \Rightarrow \boxed{u_{x,t}(x, t) = \sum_{j=1}^{\infty} \left(e^{-\lambda_j t} \frac{(u_0, \psi_j)}{(\psi_j, \psi_j)} \right) \psi_j}$$

Example: Heat equation w/ Neumann B.C.'s

$$L: C^2_N[0,1] \rightarrow C[0,1]$$

$$Lu = -u'', \quad u'(0) = u'(1) = 0, \quad \text{Eigenvalues: } \lambda_0 = 0$$

$$-\psi''(x) = \lambda \psi(x)$$

$$\Rightarrow \psi(x) = A \sin(\sqrt{\lambda} x) + B \cos(\sqrt{\lambda} x)$$

$$\psi'(0) = 0, \quad \psi'(1) = 0$$

$$\psi'(x) = \sqrt{\lambda} A \cos(\sqrt{\lambda} x) - \sqrt{\lambda} B \sin(\sqrt{\lambda} x)$$

$$0 = \psi'(0) = \sqrt{\lambda} A \Rightarrow A = 0$$

$$0 = \psi'(1) = -\sqrt{\lambda} B \sin(\sqrt{\lambda}) \Rightarrow \sin(\sqrt{\lambda}) = 0$$

$$\Rightarrow \sqrt{\lambda} = n\pi \rightarrow \lambda_n = n^2 \pi^2$$

λ_0 is special:

$$-\psi''(x) = 0$$

$$\Rightarrow \psi(x) = A + Bx \rightarrow \psi'(x) = B$$

$$0 = \psi'(0) = B \Rightarrow B = 0$$

So! $\lambda_1 = \pi^2, \psi_1(x) = \sqrt{2} \cos(\pi x)$

$$0 = \psi'(1) = B \Rightarrow B = 0$$

$$\lambda_0 = 0, \psi_0(x) = 1$$

Take $\psi(x) = A = 1$ for $\lambda = 0$

Soln: $u(x,t) = \sum e^{-\lambda_j t} \frac{(u_0, \psi_j)}{(\psi_j, \psi_j)} \psi_j(x)$

$$= e^{-\lambda_0 t} \frac{(u_0, \psi_0)}{(\psi_0, \psi_0)} \psi_0 + \sum_{j=1}^{\infty} e^{-\lambda_j t} \frac{(u_0, \psi_j)}{(\psi_j, \psi_j)} \psi_j$$

$$= \boxed{\frac{(u_0, \psi_0)}{(\psi_0, \psi_0)} \cdot 1} + \boxed{\sum_{j=1}^{\infty} e^{-\lambda_j t} \frac{(u_0, \psi_j)}{(\psi_j, \psi_j)} \psi_j}$$

Constant!
 $\int_0^1 u_0(x) dx$

$\rightarrow 0$ as $t \rightarrow \infty$
since $\lambda_j = j^2 \pi^2 > 0$

FINITE ELEMENT METHOD FOR $u_t = -Lu$ (HEAT EQUATION)

(Turn PDE into system of ODE's)

① Weak Form

Ex. $u_t = u_{xx}, \quad u(0,t) = u(1,t) = 0$

Pick $V = C^2_0[0,1]$, For $v \in V$

$$\frac{\partial}{\partial t} (u, v) = (u_t, v) = (u_{xx}, v)$$

$$\text{IBP} = (u_{xx}, v) = - \int_0^1 u_x(x,t) v_x(x) dx = -a(u, v)$$

energy inner product

$$\boxed{\frac{\partial}{\partial t} (u, v) = -a(u, v) \quad \forall v \in V}$$

WEAK FORM

② Galerkin Approximation

$$V_N = \text{span} \{ \phi_1, \dots, \phi_N \} \subseteq V$$

$$\text{Find } u_N \in V_N \times C[0, \infty] \text{ s.t.}$$

$$\frac{\partial}{\partial t} (u_N, v) = -a(u_N, v) \quad \forall v \in V_N$$

③ Turn into a linear algebra problem:

$$\text{Write } u_N(x, t) = \sum_{j=1}^N c_j(t) \phi_j(x)$$

$$\text{Test at } v = \phi_k, k=1, \dots, N$$

$$\rightarrow \sum_{j=1}^N c_j'(t) (\phi_j, \phi_k) = - \sum_{j=1}^N c_j(t) a(\phi_j, \phi_k)$$

$$\begin{bmatrix} (\phi_1, \phi_1) & \dots & (\phi_N, \phi_1) \\ \vdots & & \vdots \\ (\phi_1, \phi_N) & \dots & (\phi_N, \phi_N) \end{bmatrix} \begin{bmatrix} c_1' \\ c_2' \\ \vdots \\ c_N' \end{bmatrix} = \begin{bmatrix} a(\phi_1, \phi_1) & \dots & a(\phi_N, \phi_1) \\ \vdots & & \vdots \\ a(\phi_1, \phi_N) & \dots & a(\phi_N, \phi_N) \end{bmatrix} \begin{bmatrix} c_1 \\ \vdots \\ c_N \end{bmatrix}$$

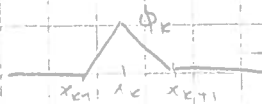
M (mass) c' $= -K$ (stiffness) c

④ Solve for $c(t)$: $\begin{bmatrix} c_1(t) \\ \vdots \\ c_N(t) \end{bmatrix}$

$$\boxed{c'(t) = -M^{-1}Kc(t)} \rightarrow \boxed{c(t) = e^{-M^{-1}Kt} c(0)} \quad \begin{matrix} \text{Approx in space,} \\ \text{exact in time.} \end{matrix}$$

If ϕ_1, \dots, ϕ_N = hat functions, $c_j(0) = u_0(x_j)$

$$\phi_j(x_k) = \begin{cases} 1 & j=k \\ 0 & j \neq k \end{cases}$$



$$u_N(x_k, t) = \sum_{j=1}^N c_j(t) \phi_j(x_k) = c_k(t)$$

$$c_k(0) = u_N(x_k, 0) = u_0(x_k)$$

Alternative: FW Euler: eigs of $M^{-1}K \approx j^2 \pi^2$, $j=1, \dots, N$

Roughly $\Delta t \in (0, \frac{2}{N^2 \pi^2})$

BW Euler: ANY $\Delta t > 0$ is stable.

Rule of thumb: If you double N ,
cut Δt to $\frac{\Delta t}{4}$

WAVE EQUATION ($u_{tt} = -Lu$)

L is symmetric linear operator

$\lambda_1, \lambda_2, \dots$ eigenvalues (real)

$\varphi_1, \varphi_2, \dots$ eigenfunctions orthogonal

① Write $u(x,t) = \sum a_j(t) \varphi_j(x)$

② Substitute into PDE to get ODE for coefficients

$$\begin{aligned}\sum a_j''(t) \varphi_j(x) &= -\sum a_j(t) L\varphi_j(x) \\ &= -\sum a_j(t) \lambda_j \varphi_j(x)\end{aligned}$$

inner product \Rightarrow

$$a_k''(t) = -\lambda_k a_k(t)$$

③ Solve ODE's for a_k

$$a_k(t) = A_k \sin(\sqrt{\lambda_k} t) + B_k \cos(\sqrt{\lambda_k} t)$$

with initial conditions:

$$u_0(x) = u(x,0) = \sum \frac{(u_0, \varphi_j)}{(\varphi_j, \varphi_j)}$$

$$v_0(x) = u_t(x,0) = \sum \frac{(v_0, \varphi_j)}{(\varphi_j, \varphi_j)}$$

$$\begin{aligned}\rightarrow u_0(x) &= \sum_{j=1}^{\infty} A_j \sin(\sqrt{\lambda_j} 0) + B_j \cos(\sqrt{\lambda_j} 0) \varphi_j \\ &= \sum_{j=1}^{\infty} B_j \varphi_j(x) \Rightarrow \boxed{B_j = \frac{(u_0, \varphi_j)}{(\varphi_j, \varphi_j)}}\end{aligned}$$

$$\begin{aligned}\rightarrow v_0(x) &= \sum_{j=1}^{\infty} \sqrt{\lambda_j} (A_j \cos(\sqrt{\lambda_j} 0) - B_j \sin(\sqrt{\lambda_j} 0)) \varphi_j(x) \\ &= \sum_{j=1}^{\infty} \sqrt{\lambda_j} A_j \varphi_j(x) \Rightarrow \boxed{A_j = \frac{1}{\sqrt{\lambda_j}} \frac{(v_0, \varphi_j)}{(\varphi_j, \varphi_j)}} \quad \lambda_j \neq 0\end{aligned}$$

$$u(x,t) =$$

$$\sum_{j=1}^{\infty} \left(\frac{1}{\sqrt{\lambda_j}} \frac{(v_0, \varphi_j)}{(\varphi_j, \varphi_j)} \sin(\sqrt{\lambda_j} t) + \frac{(u_0, \varphi_j)}{(\varphi_j, \varphi_j)} \cos(\sqrt{\lambda_j} t) \right) \varphi_j(x)$$

oscillator!

FINITE ELEMENT METHOD w/ WAVE EQ

Find $u \in V \times C(0, t)$

Weak form: $\frac{\partial^2}{\partial t^2} (u, v) = -a(u, v) \quad \forall v \in V$

Galerkin: $u_N(x, t) = \sum c_j(t) \psi_j(x)$

Find u_N s.t. $\frac{\partial^2}{\partial t^2} (u_N, v) = -a(u_N, v)$

$\forall v \in V_N = \text{span} \{ \phi_1, \phi_2, \dots, \phi_N \}$

$\rightarrow \boxed{M \ddot{c} = -Kc}$

$\frac{\partial}{\partial t} \begin{bmatrix} \dot{c} \\ c \end{bmatrix} = \begin{bmatrix} 0 & I \\ -M^{-1}K & 0 \end{bmatrix} \begin{bmatrix} \dot{c} \\ c \end{bmatrix}$

Solve w/ Mix exponential

FW Euler

BW Euler

Trapezoid

If eigs of $M^{-1}K$ are real, positive then...

eigs of $\begin{bmatrix} 0 & I \\ -M^{-1}K & 0 \end{bmatrix}$ are purely imaginary

d'Alembert's Solution on Unbounded Domains

$u_{tt}(x, t) = c^2 u_{xx}(x, t) \quad x \in (-\infty, \infty)$

$u(x, 0) = \psi(x) \quad u_t(x, 0) = \gamma(x)$

$u(x, t) = \frac{1}{2} (\psi(x+ct) + \psi(x-ct)) + \frac{1}{2c} \int_{x-ct}^{x+ct} \gamma(s) ds$

Also review domain of dependence!

