

LEC - 1

Pre-Modern era of Physics - 4000 BC - 1900 AD

Celestial Mechanics \longrightarrow Universal Law of gravitation

↓
Observation
of celestial
bodies

* After the development
of mechanics

④ General relativity

↓
Differential
Geometry

Concept of
a field (E & M)

Early calendars \rightarrow developed based on motion

of heavenly bodies \rightarrow But sunrise to sunrise time not always constant. Later refined as
Day \Rightarrow Sunrise to Sunrise $\underline{\text{midday to midday}} \Rightarrow$ More accurate

Year \Rightarrow Time in which the stars regain the same position
in the night sky

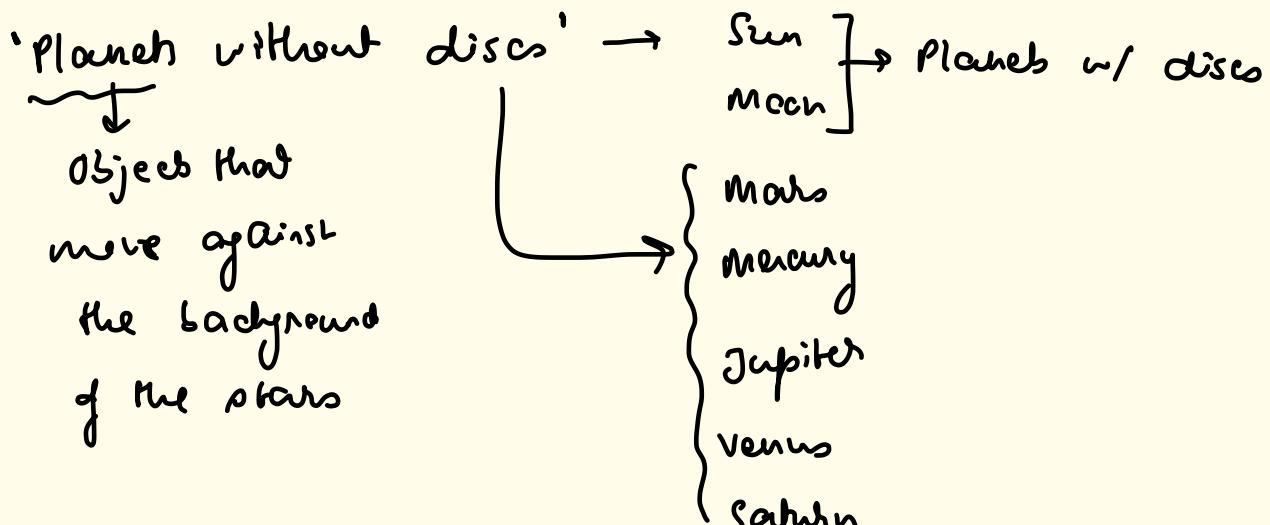
Sub-units of a year



Month
- Lunar
cycle
(New moon
to New moon)

Week
- Why 7 days?

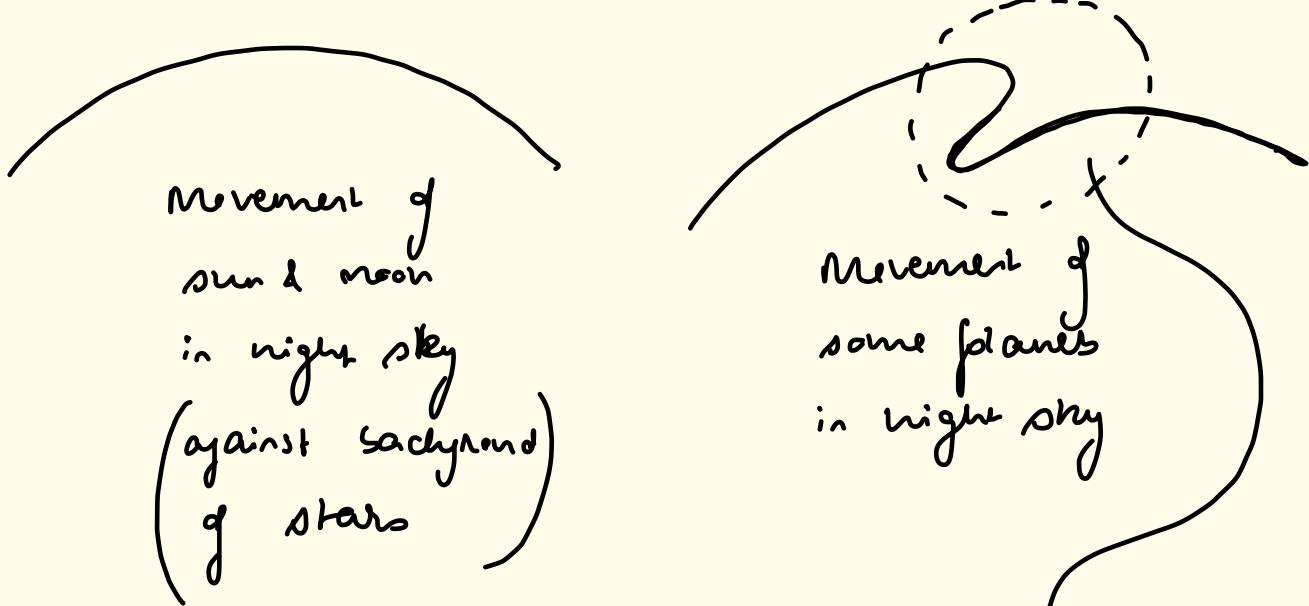
LEC - 5



Current celestial objects

- Sun & moon that move on their own spheres which also rotate wrt Earth at their own speed
- Stars fixed on celestial sphere that rotate wrt Earth

However, none of these explain retrograde motion in planes



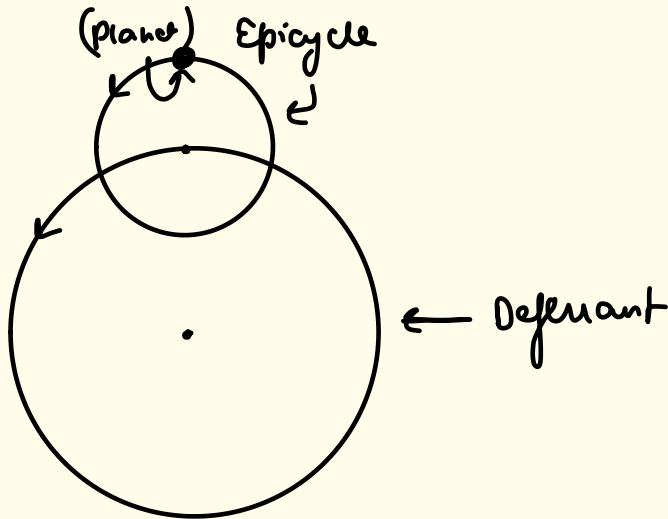
Aristotle - 'Motion in the heaven's is perfect'

How would you explain retrograde motion being confined to Aristotelian philosophy?

Appollonius & Hipparchus
came up w/ an
explanation



Epicycles



It was believed that planetary rotation in epicycle
as well as the motion of the cycle itself were
constant

However, data over the next few centuries contradicted
this theory → Ptolemy came up with

(Still Aristotelian)
theory

↓
Concept of
Equant

2nd century AD

There was an elephant in the room



CATHOLIC CHURCH → for some reason
they liked Aristotle's
philosophy

Still kinda
greens of
Aristotelian
philosophy

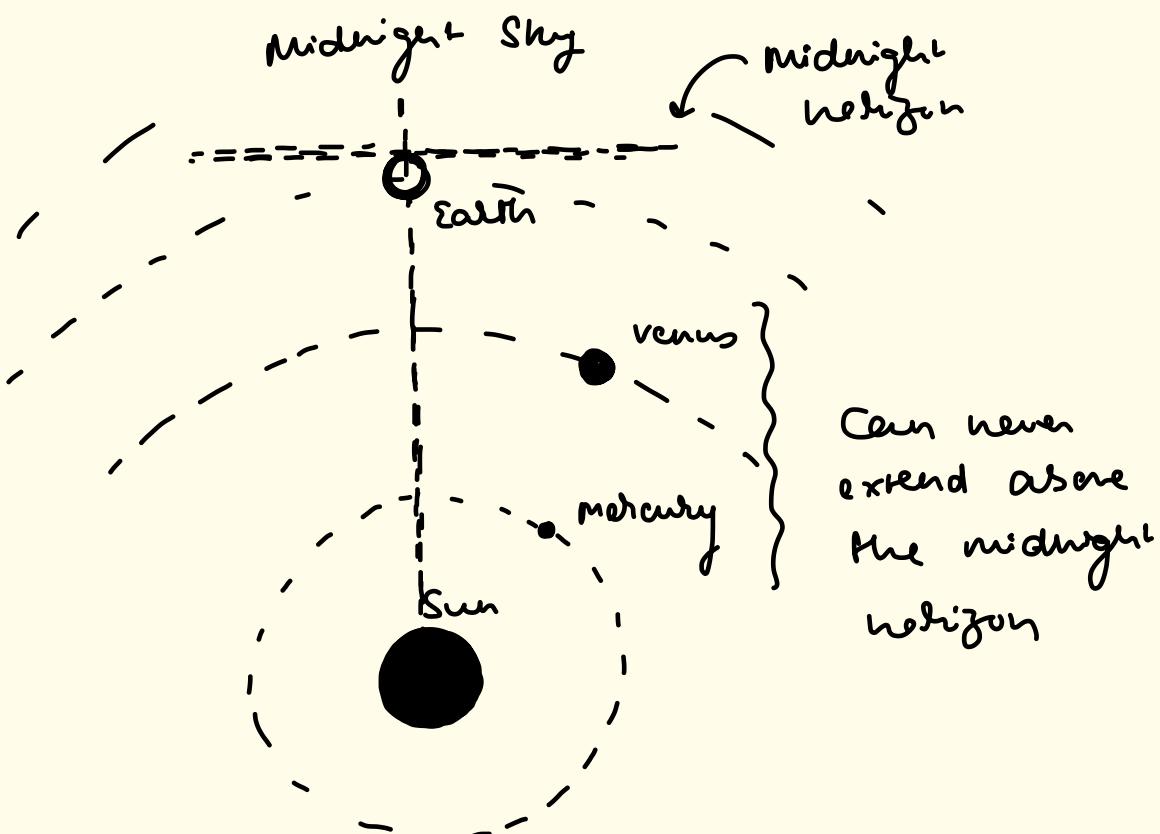
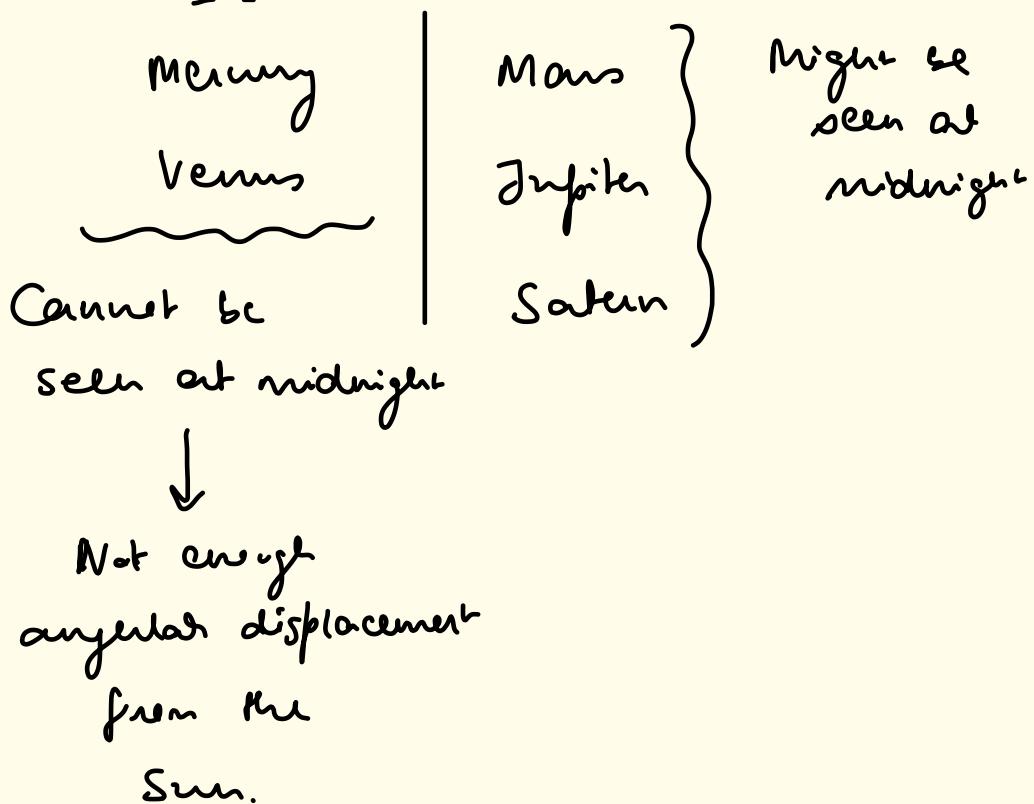
And so
was the story
until

15th century AD

Nicholas Copernicus → Heliocentric
theory

Plants (incl. Earth excl. Moon) revolve
in a circular path around the Sun
w/ constant angular philosophy

* Even observationally, we can make a distinction b/w planets



Tycho Brahe → created one of the best naked eye
observatories w/ help from King of Denmark

He then pissed himself to death in aroyal banque
and Johannes Kepler got hold of his observations

Kepler was the first one to bin Aristotle and any
pre-conceived notions, which led him to]



Kepler's Laws

of Planetary
Motion

→ Popularised
by Galileo

→ Discovery of Law
of Gravitation by Newton

LEC-G

Themes of this course so far

1. Motion in Heavens is perfect
- * 2. Only those ideas supported by data are given the distinction of being called theories
3. There are universal laws which hold true in heaven and on Earth

Galileo (16th century)

- Italian, hence devout catholic
- Started to question Aristotle

↓
 "Heavier objects fall faster" (Allegedly)

Realised through experimentation and observation that time period of a pendulum is independent of its mass

Galileo also came up w/ eqns of motion by letting objects fall down inclined plane

Galileo thought of as first modern scientist ↓

"Only ideas supported by data are valid theories"

Galileo's astronomical discoveries using a telescope

- ↳ Milky way
- ↳ 4 moons of Jupiter
- ↳ Features of surface of the moon } Against Aristotle

Then dude wrote this → 'Dialogue on Two World Systems'

↳ wanted to bitchslap Aristotle once and for all.

- No equations
- Conversation b/w three people
- First pop-sci book
- Argued for heliocentrism in a manner the general public could understand

* The people arguing for the respective theories

- Heliocentrism - Sagredo (Sage)
- Geocentrism - Simplicio (Dumbass)

Church kinda didn't like that... there was a trial
and he had to apologise so the higher-ups at
Church didn't get their nips in a twist (and so
Galileo didn't get tortured)



then kept on
house arrest for
the rest of his life

Next dude → Huygens (17th century)

→ Analysed uniform circular
motion and identified centripetal
force → (before Newton)

→ Built an Tautochrone

④ Pendulum clocks
depending on simple
harmonic oscillation
were unreliable

- If amplitude large,
time period differs from
expected time period
- If amplitude small,
difficult to practically
deal with problems caused
by friction

kind of
clock

↓
Used oscillations
in a cycloidal
shape to produce
constant time period
irrespective of
amplitude

Newton

→ Childhood trauma

→ fairly brilliant student, got into Cambridge

→ College closed due to plague outbreak, Newton sent home

→ Bored Newton looked at apple falling and moon in the sky ↓

Apple falling → Apple accelerating
Moon revolving around Earth → Moon accelerating

Newton decided to work on relation b/w the two

He found

$$\frac{a_{\text{apple}}}{a_{\text{moon}}} = \frac{g}{w^2 R} = \left(\frac{R}{R_{\text{Earth}}}\right)^2$$

$$a_{\text{apple}} \propto \frac{1}{(\text{dist. of apple from center of Earth})^2}$$

Same case for moon

During this time, he also casually invented calculus

Used it to find a relation b/w inverse square law and Kepler's laws (They both, in fact, imply each other)

" There are universal laws which hold true in heaven and on Earth.

LEC-7

Up till now, Galileo → Huygens → Newton

→ Revolutionised

the very way people
thought of Physics

Law of inertia (Newton's 1st law)

An object which does not have
a force acting on it is either
at rest or in a state of
uniform velocity

But really we
came up w/ it
first

Law of kinematics

Emphasised the
role of time \Rightarrow

The concept of time
as an independent
variable that helps
describing motion of
objects

Galileo

Uniform Circular
motion

→ Huygens

↓

Requires centripetal force

{ These were
concepts known
before Newton }

Now comes the man himself \rightarrow Newton's 2nd Law

↓

Newton's 3rd
Law

$$\bar{F} = m \bar{a}$$

Describes the
motion of every single
moving body in the
universe.

Every action has
an equal and opposite
reaction.

$$\bar{F}(\bar{r}) = m \cdot \frac{d^2 \bar{r}}{dt^2}$$

↓
 "Ordinary differential
 equation
 in modern language."

Ever since Newton, physicists
 solve stuff by making differential \Rightarrow Newton is one
 equations Aristotle is

Theorems regarding differential equations

\rightarrow Ordinary Diff Eqn of nth order has 'n' independent
 solutions

\rightarrow To obtain a unique solution, you need 'n' initial
 conditions.

* Independent soln \rightarrow No linear combination of the
 two solns will be zero for every value in their domain

"Now after praising Newton so much let us look
 at the creativity he did" - Prof Uma Sanhar

Looking at the eqn for 2nd law,

\bar{r} is only defined at a single point

CM of the body

But forces acting on surface of the body, not on CM

So what? \bar{F} = Sum of forces acting on the body.

Even in Law of gravity, when Newton described motion of Earth and Sun, he considered their entire mass at their COM



Later Newton
proved that for
all gravitational purpose,
solid sphere \equiv point at GM

Moreover, another complication w/ IRL objects



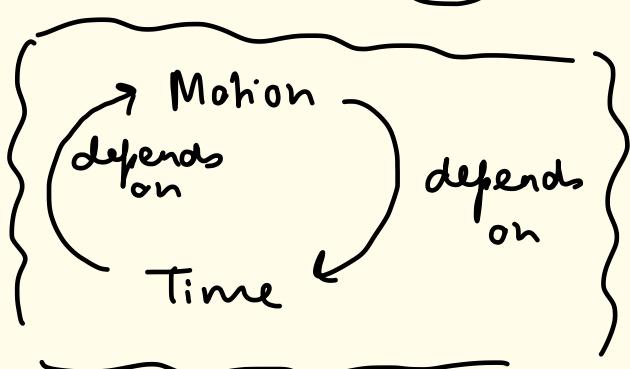
Solid objects
rotate \Rightarrow Later solved
by Euler.

Now, Newton's actual cheating

$$\bar{F}(\vec{r}) = \frac{d^2\vec{r}}{dt^2}$$

\rightarrow Time \Rightarrow measured by looking
at motion of celestial
objects

or motion of objects in
an inertial frame



Circular argument

Newton just hacked the problem to pieces



Said there was 'Absolute time' which is independent of any motion

Now, we knew $\bar{F} = m\bar{a}$ true only in "inertial frames".
 Newton simply declared that there exists 'Absolute space', and that any frame moving with constant velocity wrt that frame is hence inertial.

↓
 what is that?

$$\bar{F}_{NI} = \underbrace{\bar{F}_{\text{pseudo}}}_{\text{Restoration}} + \bar{F}_{\text{Real}} = m\bar{a}_{NI}$$

Restoration

of Newton's law
for Non-inertial
frames

* Newton's law of gravity

$$\bar{F}_{gr} = -\frac{G Mm}{r^2} \hat{r}$$

Important discovery because of Newton's Law



Discovery of Uranus

on observation over a period of time, people realise Uranus was not moving in a Kepler orbit. So either



Newton
was wrong
(GOD PLEASE NO)



there is another massive body very close to Uranus that feeds it orbit up

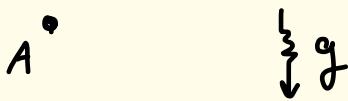


found to be Nephrite

LEC 8

Bernoulli \longrightarrow Brachistochrone problem

D'Alembert



Euler }
Lagrange } $\rightarrow *$

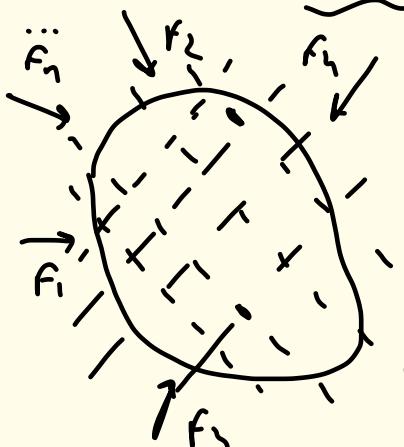
Hamilton

find curve such that
you could travel from A
to B in the least time
of time.

Euler derived that Newton's 2nd law can
be reduced to the following form



$$\bar{F}_{\text{res}} = M \left(\frac{d^2 \vec{r}_{\text{cm}}}{dt^2} \right)$$



$N (\gg 1)$ pieces

for i th piece = point mass,

$$m_i \cdot \frac{d^2 \vec{r}_i}{dt^2} = \bar{f}_{i,\text{ext}} + \bar{f}_{i,\text{int}}$$

$$\therefore \sum m_i \cdot \frac{d^2 \vec{r}_i}{dt^2} = \underbrace{\sum \bar{f}_{i,\text{ext}}}_{\sigma} + \underbrace{\sum \bar{f}_{i,\text{int}}}_{\sigma}$$

$$\Rightarrow \frac{d^2}{dt^2} \left(\sum m_i \vec{r}_i \right) = \bar{F}_{\text{res}} \quad \Rightarrow \frac{M \cdot d^2 \vec{r}_{\text{cm}}}{dt^2} = \bar{F}_{\text{res}}$$

* Euler's Law for motion of a rigid body



Arbitrary motion
involves translation of
CM and rotation
about CM

$$m_i \cdot \frac{d^2 \vec{r}_i}{dt^2} = \vec{f}_{i\text{ext}} + \vec{f}_{i\text{int}}$$

$$\therefore \vec{r}_i \times m_i \cdot \frac{d^2 \vec{r}_i}{dt^2} = \vec{r}_i \times (\vec{f}_{i\text{ext}} + \vec{f}_{i\text{int}})$$

$$\vec{r}_i \times \frac{d}{dt} (m_i \vec{v}_i) = \vec{r}_i \times (\vec{f}_{i\text{ext}} + \vec{f}_{i\text{int}})$$

$$\vec{v}_i = \vec{\omega} \times \vec{r}_i = \vec{\omega} \times \vec{r}_i$$

→ All particles have same $\vec{\omega}$

$$\therefore \vec{r}_i \times \frac{d}{dt} (m_i (\vec{\omega} \times \vec{r}_i)) = \vec{r}_i \times (\vec{f}_{i\text{ext}} + \vec{f}_{i\text{int}})$$

$$\vec{l}_i = \vec{r}_i \times \vec{p}_i$$

$$= \vec{r}_i \times (m_i \vec{v}_i)$$

$$\begin{aligned} \therefore \frac{d \vec{l}_i}{dt} &= \frac{d \vec{r}_i}{dt} \times (m_i \vec{v}_i) + \vec{r}_i \times \frac{d}{dt} (m_i \vec{v}_i) \\ &= \vec{r}_i \times m_i \left(\frac{d^2 \vec{r}_i}{dt^2} \right) \end{aligned}$$

$$\begin{aligned}
 \frac{d\bar{\omega}}{dt} &= \frac{d}{dt} \sum_i \vec{f}_i \times (m_i \cdot \vec{\omega} \times \vec{r}_i) \\
 &= \frac{d}{dt} \left[I(\vec{\omega}) \right] = \sum_i \vec{f}_i \times (\vec{f}_{i,\text{ext}} + \vec{f}_{i,\text{int}}) \\
 &= \underbrace{\sum_i \vec{f}_i \times \vec{f}_{i,\text{ext}}}_{\bar{\tau}_{\text{ext}}} + \underbrace{\sum_i \vec{f}_i \times \vec{f}_{i,\text{int}}}_{0}
 \end{aligned}$$

$$\frac{d\bar{\omega}}{dt} = \bar{\tau}_{\text{ext}}$$

All internal forces appear as "act-react" pairs
net moment of such pairs about any axis is zero

* $\bar{p} = m\bar{v}$ Are m & \bar{v} similar quantities?

$\bar{L} = I \bar{\omega}$
They appear to be so,
but not really

\bar{p} & \bar{v} must be in same direction

\bar{L} & $\bar{\omega}$ may not be in same direction

Scalar \rightarrow Number

Vector \rightarrow Set of 3 numbers
 Components
 (which we then treat as scalars)

$$\begin{bmatrix} p_1 \\ p_2 \\ p_3 \end{bmatrix} = m \cdot \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} \xrightarrow{\text{Tensor}} \underbrace{\text{Tensor}}_{\text{3x1}}$$

$$\begin{aligned}
 (3 \times 1) &= \text{scalar. } (3 \times 1) \\
 (3 \times 1) &= (3 \times 3) \cdot (3 \times 1)
 \end{aligned}$$

$$\begin{bmatrix} L_1 \\ L_2 \\ L_3 \end{bmatrix} = \begin{bmatrix} \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot \end{bmatrix} \begin{bmatrix} \omega_1 \\ \omega_2 \\ \omega_3 \end{bmatrix}$$

$$\bar{A} = A_1 \hat{i} + A_2 \hat{j} + A_3 \hat{k}$$

$$\bar{B} = B_1 \hat{i} + B_2 \hat{j} + B_3 \hat{k}$$

$$\begin{aligned}
 (\bar{A} \cdot \bar{B}) &= (A_1 \hat{i} + A_2 \hat{j} + A_3 \hat{k}) \cdot (B_1 \hat{i} + B_2 \hat{j} + B_3 \hat{k}) \\
 &= \underbrace{A_1 B_1 \hat{i} \hat{i} + A_1 B_2 \hat{i} \hat{j} \dots}_{\text{9 terms}} \Rightarrow \text{9 terms} \\
 &\quad \downarrow \\
 &\text{Tensor (of order 2)}
 \end{aligned}$$

$$*\text{Young's Modulus} = \frac{\text{Stress}}{\text{Strain}}$$

Strain = $\frac{\text{force}}{\text{Area}} \rightarrow$ Impossible to take ratio of vectors

$$\Rightarrow \text{force} = (\text{strain}) \text{ Area}$$

\downarrow

Tensor-2

Tensor-2

$$\text{Strain} = \frac{\text{Change in length}}{\text{length}} \Rightarrow \text{Change in length} = \text{Strain} \cdot \text{length}$$

$$\text{Similarly, Stress} = (Y) \cdot \text{Strain}$$

\downarrow

T-2

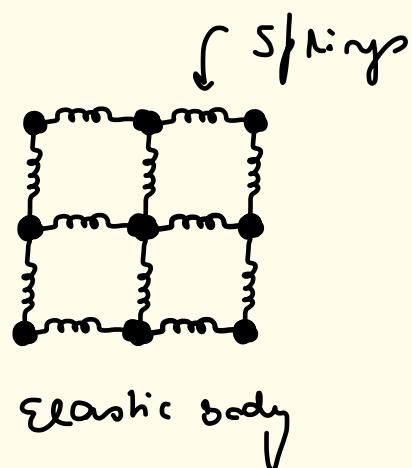
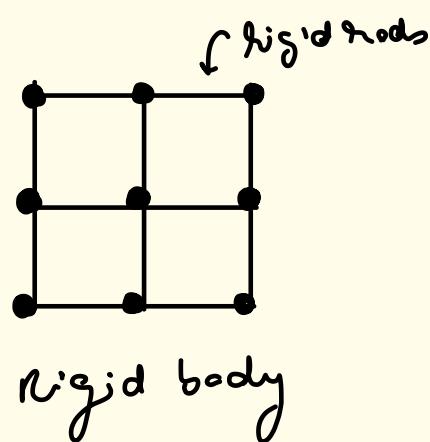
\downarrow

T-2

Tensor
of order 4

Rigid body \rightarrow No matter how it moves, the distance b/w any two mass points on a rigid body remains constant.

But all solids do deform sooner or later, so what happens when these points translate w.r.t each other the distance b/w them changes



* index notation

$$\bar{A} = A_i \quad (i=1,2,3)$$

↓
implicit

Strain tensor \rightarrow fn of position of mass point

$$g_{ij}(u)$$

↓

Related to non-Euclidean geometry

$$\text{Eucl. geo} \rightarrow ds^2 = [dx \ dy \ dz] \begin{bmatrix} dx \\ dy \\ dz \end{bmatrix}$$

$$\text{Non Eucl. geo} \rightarrow [dx \ dy \ dz] \cdot f(u) \begin{bmatrix} dx \\ dy \\ dz \end{bmatrix}$$

In Euclidean geometry, displacement is irrespective of where the point initially was



In non Euclidean geometry, over the displacement is a function of the position of the point

Fluid Dynamics

Technical defⁿ of a fluid

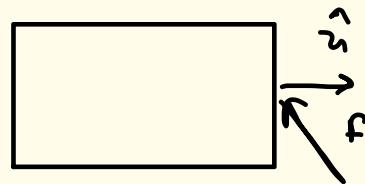
↓

Fluids can't
support shear
stress

$$\bar{f} = [\sigma] \bar{A}$$

$$\begin{bmatrix} f_1 \\ f_2 \\ f_3 \end{bmatrix} = \begin{bmatrix} \sigma_{11} & \cdot & \cdot \\ \cdot & \sigma_{22} & \cdot \\ \cdot & \cdot & \sigma_{33} \end{bmatrix} \begin{bmatrix} A \\ 0 \\ 0 \end{bmatrix}$$

Area having
only
one component



Shear stress

⇒ Component of
stress due to force
applied perpendicular
to area

$$\sigma_{11} \rightarrow f_x \Leftrightarrow A_x$$

$$\sigma_{21} \rightarrow f_y \Leftrightarrow A_y$$

$$\sigma_{31} \rightarrow f_z \Leftrightarrow A_z$$

... and so on

Fluids



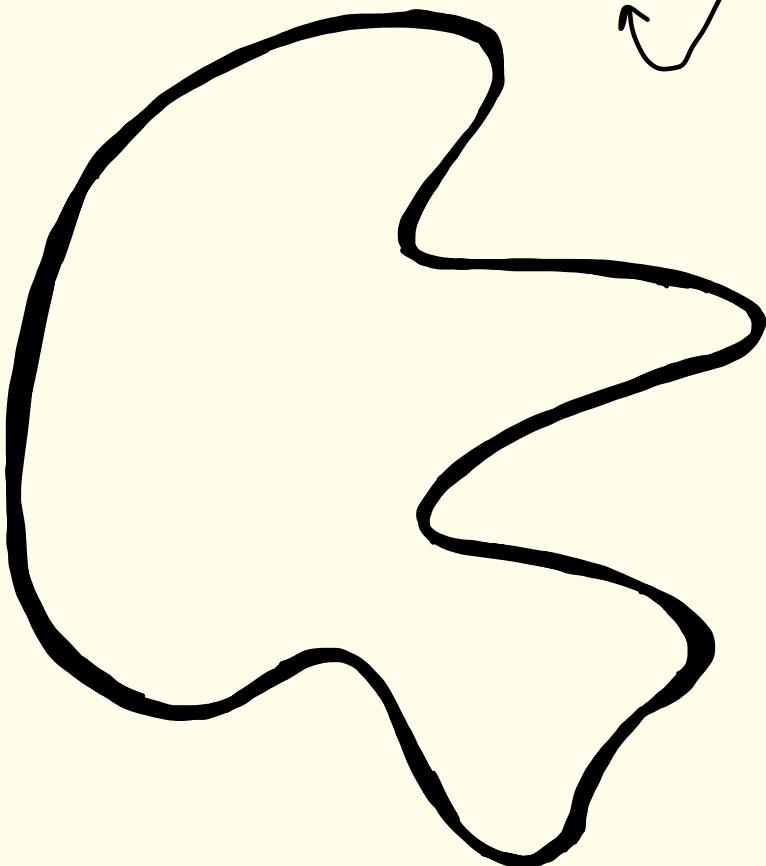
Liquid

Do not compress
under normal
stress



Compressible
fluids

Formalism for fluid mechanics \rightarrow Defines a macroscopic body of fluid called a continuum



No gaps in between

* Transcendental no.s



No.s which cannot be obtained as the root of an algebraic eqn

No. of integers

No. of fractions

II

(No. of algebraic no.s)

Obtained as root of alg.-exp.

However,

No. of transcendental no.s

All of these

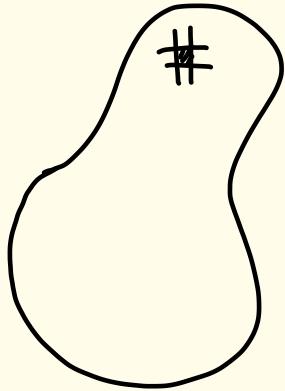


Real no. line is filled with transcendental numbers.

Integers, fractions, surds

=

Even : rationals



Taking a volume element

$$d^3 u_i$$

with mass

$$dm_i = \rho(u_i) d^3 u_i$$

velocity pattern
of fluids



$$\bar{v}(x, y, z)$$

for simplicity's
sake, let us

assume there is →
no explicit dependence
of velocity of time

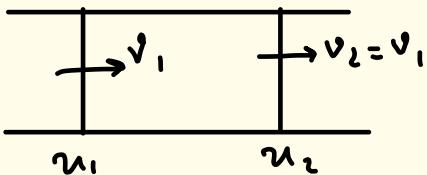
$$V = \int d^3 u$$

$$M = \int \rho(u) d^3 u$$

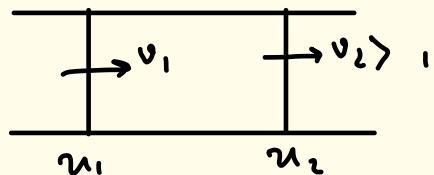
No change in
flow pattern
with time



Steady State
flow



Considering no
leaks/loss/gain
of water



Perhaps there is somehow
a source of water in the pipe

↓
rel b/w surface
integral of velocity field
and its source

Any function
depending on
spatial coordinates

$$\underset{S}{\text{closed surface}} \oint \bar{V} \cdot d\bar{A} = (V_2 - V_1) A > 0 \quad (\text{so } \delta \text{ is the claim})$$

Physics - Matter cannot be created or destroyed (c. 1930ish)

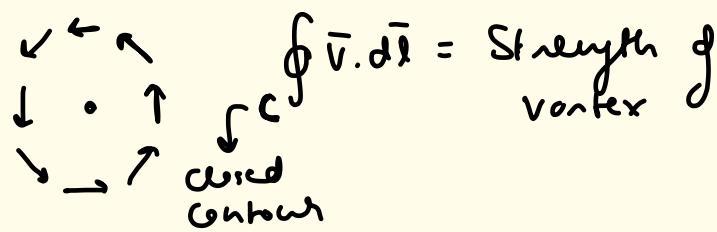
mass flowing out of the volume is greater than mass flowing into the volume

\therefore our velocity flux is related somehow to strength of source pumping water into the canal

we see this in E&M

Physical action causing such motion

* Vortex



Here, we abruptly end our conversation on fluid mechanics and jump into...

\therefore Integrals of a vector field over a closed loop or surface contain crucial physical info about what produces it.

Electricity and Magnetism

$$\bar{F}_{12} = \left(\frac{1}{4\pi\epsilon_0} \right) \frac{q_1 q_2}{|\bar{r}_1 - \bar{r}_2|^3} (\bar{r}_1 - \bar{r}_2) \quad \bar{F}_{21} = -\bar{F}_{12}$$

↓
Force on q_1 due to q_2

Similarly, Ampere's law for force b/w closed loop,

Newton said - figure out what the force is, then figure out the movement caused by such forces



Both Coulomb

& Ampere work
their eqns in → under Newton's
terms of force influence

But Faraday's POV



If there is q_1 at \vec{r}_1 ,
then electric field $\vec{\epsilon}_1$ at
 \vec{r}_1 is given as follows

$$\vec{\epsilon}_1 = \left(\frac{1}{4\pi\epsilon_0} \right) \frac{q_1 (\vec{r}_1 - \vec{r}_2)}{|\vec{r}_1 - \vec{r}_2|^3}$$

If there happens to be a charge q_1 at \vec{r}_1 ,
then force on $q_1 = q_1 \cdot \vec{\epsilon}_1(\vec{r}_1)$

LEC - 10

Concept of field introduced by Faraday
usefulness → Allowed us to think of electrostatic forces
as a charge somehow interacting with an 'aura'
produced by a 'source' charge

thanks to fluid mech,
we had vast amount of mathematical
machinery to deal with these
kinds of problems

Similarly, a magnetic field was also defined

$$d\bar{B}(\vec{r}_1) = \frac{\mu_0}{4\pi} \cdot I_2 \cdot d\vec{l}_2 \times (\vec{r}_1 - \vec{r}_2)^{-3}$$

↓
'Auro' produced at \vec{r}_1 due to tiny length of $d\vec{l}_2$ carrying current I_2 located at \vec{r}_2 .

$$d\bar{F}_{12} = I_1 \cdot d\vec{l}_1 \times d\bar{B}(\vec{r}_1)$$

↳ This formulation of magnetic force on current elements doesn't satisfy Newton's third law \Rightarrow small current elements do not exist in physical reality, current always flows in loops.

Integrating this expression over the loop results in equal and opposite forces on both bodies

we know source \Rightarrow field

But also, field \Rightarrow source

"Given a field on a surface, how can we figure out the sources?"

Necessary mathematical armory \rightarrow Calculus of fields

Field = some physical quantity whose value is a function of its location.

for a fn in two variables $f(u, y)$,

$$\begin{aligned}\Delta f &= f(u + \Delta u, y + \Delta y) - f(u, y) \\ &= \underbrace{f(u + \Delta u, y + \Delta y) - f(u + \Delta u, y)}_{\frac{\partial f}{\partial y} \cdot \Delta y} \\ &\quad + \underbrace{f(u + \Delta u, y) - f(u, y)}_{\frac{\partial f}{\partial u} \cdot \Delta u}\end{aligned}$$

now, for single variable fn $g(u)$,

$$\Delta g = \frac{dg}{du} \Delta u$$

similarly, we can argue there must be a similar reln b/w Δf & Δg

$$\text{i.e. } \Delta f = \bar{\nabla} f \cdot \bar{\Delta} u$$

del operator

$$\bar{\nabla} = \hat{u} \cdot \frac{d}{du} + \hat{y} \cdot \frac{d}{dy}$$



Differential operator

$$\bar{\nabla} f = \frac{df}{du} \hat{u} + \frac{df}{dy} \hat{y}$$

$$\bar{\Delta} u = \Delta u \hat{u} + \Delta y \hat{y}$$

Takes one f^n as input
and gives another
function

as output

$$\therefore \Delta f = \frac{df}{du} \cdot \Delta u + \frac{df}{dy} \cdot \Delta y$$

How does the del operator work on vector fields?

Considering a vector field $\vec{V} = V_1 \hat{i} + V_2 \hat{j} + V_3 \hat{k}$

$$\nabla \bar{v} = \left(\hat{x} \frac{\partial}{\partial x} + \hat{y} \frac{\partial}{\partial y} + \hat{z} \frac{\partial}{\partial z} \right) (v_1 \hat{x} + v_2 \hat{y} + v_3 \hat{z})$$

→ Tensor - 2

$$\nabla \cdot \vec{V} = \frac{\partial V_1}{\partial x} \hat{i} + \frac{\partial V_2}{\partial y} \hat{j} + \frac{\partial V_3}{\partial z} \hat{k} \rightarrow \text{Divergenz}$$

$$\nabla \times \vec{v} = \text{curl } \vec{v}$$

our jump \rightarrow Divergence and curl are
"sources" of a vector field

Helmholtz theorem - If divergence & curl of a vector field are known, then the vector field can be reconstructed provided the field vanishes as $\frac{1}{r^2}$ as $r \rightarrow \infty$ (or any higher exponent)

Electric Static field

Σ remains constant

$$\oint_S \vec{E} \cdot d\vec{A} = \frac{\rho_{in}}{\epsilon_0} \xrightarrow{\text{w/ time.}} \frac{1}{\epsilon_0} \int_V \rho(\vec{r}) dV \xrightarrow{\downarrow} \int \left(\vec{J} \cdot \vec{\varepsilon} - \frac{\rho(\vec{r})}{\epsilon_0} \right) dV = 0$$

$\xrightarrow{\text{Gauss Thm}}$

$$\oint_S \bar{\Sigma} \cdot d\bar{A} = \nabla \cdot \bar{\Sigma}$$

Gauss law

True for any arbitrary
vector field

Vanishes for
arbitrarily small
closed loop

only if integrand
vanishes

$$\therefore \nabla \cdot \bar{\epsilon} = \frac{\rho(\bar{r})}{\epsilon_0}$$

\Rightarrow Divergence of electric field at a point \propto charge density at that point

By similar argument,

LEC-11

$$(\nabla \times \bar{B}) = \mu_0 \cdot \bar{J}(\bar{r})$$

(By Stokes' Thm)

Two forms of vector calculus

$$\rightarrow \text{Gauss' Thm} - \oint_S \bar{v} \cdot dA = \int_V \nabla \cdot \bar{v} dV$$

$$\rightarrow \text{Stokes' Thm} - \oint_C \bar{v} \cdot d\bar{l} = \int_S (\nabla \times \bar{v}) \cdot dA$$

True for any arbitrary vector field

$$\nabla \cdot \bar{\epsilon} = \frac{\rho(\bar{r})}{\epsilon_0} \quad \nabla \times \bar{\epsilon} = 0 \quad (\text{for electrostatic field})$$

$$\nabla \cdot \bar{B} = 0 \quad \nabla \times \bar{B} = \mu_0 \bar{J}(\bar{r})$$

Two more forms of vector calculus

$$\rightarrow \underbrace{\nabla \times \nabla \phi(\bar{r})}_{\text{curl of gradient}} = 0$$

curl of gradient of any arbitrary scalar field is zero

$$\rightarrow \underbrace{\nabla \cdot (\nabla \times \bar{v}(\bar{r}))}_{\text{Divergence of curl}} = 0$$

Divergence of curl of any arbitrary vector field is zero.

We know

$$\nabla \times \vec{\xi} = 0 \quad \text{Defn of potential}$$

$$\therefore \vec{\xi} = -\nabla \cdot \phi(\vec{r})$$

$$\nabla \cdot (-\nabla \cdot \phi(\vec{r})) = \frac{\rho(\vec{r})}{\epsilon_0}$$

$$\Rightarrow \boxed{\frac{\rho(\vec{r})}{\epsilon_0} = -\nabla^2 \phi(\vec{r})}$$

our beloved 2nd order
differential eqn just
like Newton

Also,

$$\vec{\nabla} \cdot \vec{B} = 0$$

$\therefore \vec{B} = \vec{\nabla} \times \vec{A} \rightarrow$ some vector field

$$\nabla \times (\nabla \times \vec{A}) = \mu_0 \vec{J}(\vec{r})$$

$$= \nabla(\nabla \cdot \vec{A}) - \nabla^2 \vec{A}$$

["Truth me bho"
- Prof]

Poisson's
Equation

$$\mu_0 \vec{J}(\vec{r}) = -\nabla^2 \vec{A}(\vec{r})$$

How to solve Poisson's eqn \rightarrow Analogue w/ Newton's 2nd law

$$\frac{d^2 u}{dt^2} = \frac{f(u)}{m} \rightarrow \text{we need}$$

$u(0), \dot{u}(0) \rightarrow$ Condition at initial pt

To get a unique eqn \rightarrow

or

$u(0), u(T)$

$\dot{u}(0), \dot{u}(T)$

Conditions at
boundary pt

$$\nabla^2 \phi = \frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} + \frac{\partial^2 \phi}{\partial z^2}$$

$$= -\frac{\rho(\vec{r})}{\epsilon_0}$$

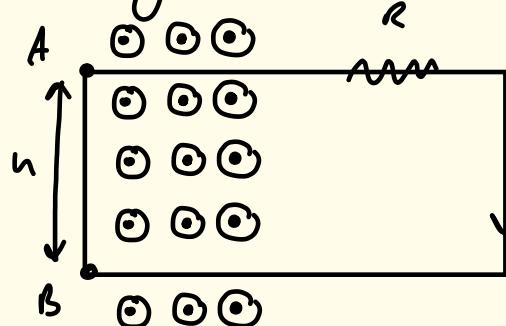
Boundary
conds

Dirichlet generalised
this concept over
higher dimensions

In two dimensions, we need to specify
conditions over a closed contour to
get a unique soln (can extend to 3D)

LEC - 12

Talking this experimental setup



when wire is moved with velocity \vec{v} , a current starts flowing through the wire

Lorentz force

$$\begin{aligned}\bar{F}_{\text{Lor}} &= q(\vec{v} \times \vec{B}) \\ &= q(v\hat{i} \times B\hat{j}) \\ &= -qvB\hat{j}\end{aligned}$$



what causes this charge flow?

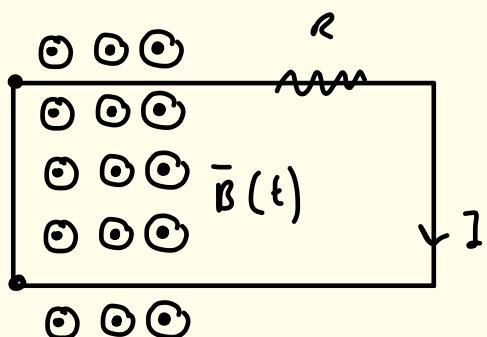
$$\begin{aligned}\text{EMF (potential diff)} &= \int_A^B \frac{\bar{F}}{q} \cdot d\bar{l} = vBh \\ &= \frac{dI}{dt} \cdot B \cdot h \\ &= B \cdot \frac{dA}{dt}\end{aligned}$$

} Ignoring signs

$\text{EMF} = -\frac{d\phi_B}{dt}$

Lenz's law

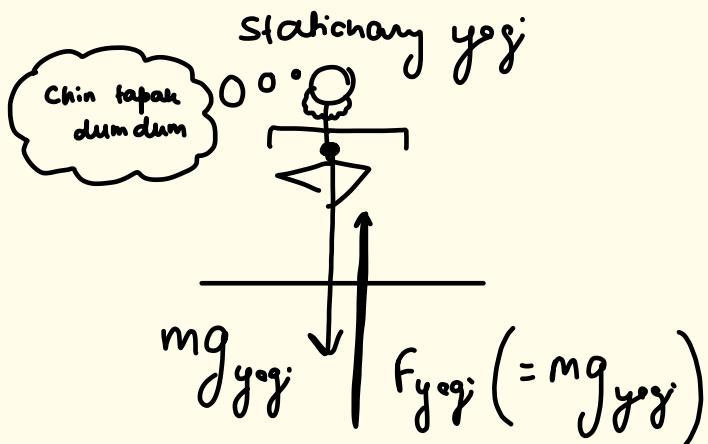
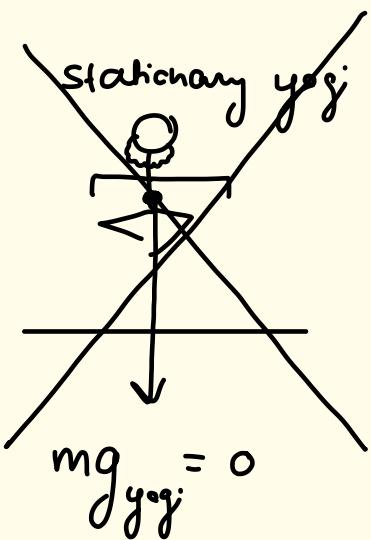
However, now taking a similar but different experimental setup



Magnetic field is taken as a function of time.

Current flows through the wire.
What is the force that causes the current to flow?

faraday's proposal



faraday merely concluded that the force causing charges to move inside the wire is produced by the changing magnetic field.

↓

this changing magnetic field produced an induced electric field that applied force on charges inside the wire

$$\oint_C \bar{\epsilon} \cdot d\bar{l} = \int_A (\nabla \times \bar{\epsilon}) \cdot d\bar{a} = - \frac{d}{dt} [\bar{B} \cdot \bar{A}]$$

↖ Area vector

$$= \int_A - \frac{d\bar{B}}{dt} \cdot d\bar{a}$$

$$\therefore \int_A (\nabla \times \bar{\epsilon} + \frac{d\bar{B}}{dt}) \cdot d\bar{a} = 0 \Rightarrow \boxed{\nabla \times \bar{\epsilon}_{ind} = - \frac{d\bar{B}}{dt}}$$

Slight cheating
but okay
faraday's law

$$\text{Meanwhile, } \nabla \cdot \bar{\Sigma}_{\text{ind}} = 0$$

Static electric field \rightarrow Non-zero divergence but zero curl
 Induced electric field \rightarrow Non-zero curl but zero divergence
 field

↓
 Resembles
 magnetic
 field

$$\left. \begin{aligned} \bar{\Sigma} &= \bar{\Sigma}_{\text{ind}} + \bar{\Sigma}_{\text{stat}} \\ \nabla \cdot \bar{\Sigma} &= \nabla \cdot \bar{\Sigma}_{\text{stat}} = \frac{\rho}{\epsilon_0} \\ \nabla \times \bar{\Sigma}_{\text{ind}} &= -\frac{\partial \mathbf{B}}{\partial t} \end{aligned} \right\}$$

Now,

$$\nabla \times \bar{\mathbf{B}} = \mu_0 \bar{\mathbf{J}}$$

$$\nabla \cdot (\nabla \times \bar{\mathbf{B}}) = \mu_0 \nabla \cdot \bar{\mathbf{J}} = 0$$

Must necessarily
 be true \Rightarrow
divergence of a
 curl is zero

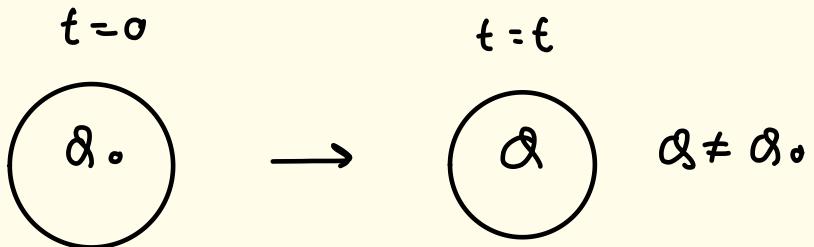
However, $\nabla \cdot \bar{\mathbf{J}}$ is only zero when
 $\bar{\mathbf{J}}$ is independent of time \Rightarrow Ampere's
 law doesn't hold for time
 varying $\bar{\mathbf{J}}$.

Maxwell's modification

$$\nabla \times \bar{\mathbf{B}} = \mu_0 (\bar{\mathbf{J}} + \bar{\mathbf{J}}_0)$$

↓
 "unfortunate
 name"
 Displacement
 current

$$\nabla \cdot \bar{\mathbf{J}} + \nabla \cdot \bar{\mathbf{J}}_0 = 0$$



$$\therefore \frac{d\delta}{dt} + \int \bar{J} \cdot d\bar{A} = 0$$

$$\frac{d}{dt} \int P \cdot dv + \int (\nabla \cdot \bar{J}) dv = 0$$

$$\Rightarrow \int_V \left(\nabla \cdot \bar{J} + \frac{\partial P}{\partial t} \right) dv = 0 \rightarrow \text{As it holds for arbitrarily small volumes,}$$

$$\nabla \cdot \bar{J}_0 + \frac{d}{dt} \left(\epsilon_0 \cdot \nabla \cdot \bar{\epsilon} \right) = 0$$

$$\nabla \cdot \bar{J}_0 + \frac{dP}{dt} = 0$$

$$\Rightarrow \bar{\nabla} \cdot \left(\bar{J}_0 + \epsilon_0 \frac{d\bar{\epsilon}}{dt} \right) = 0$$

$$\therefore \bar{J}_0 = \epsilon_0 \frac{d\bar{\epsilon}}{dt}$$

\Rightarrow

$$\therefore \nabla \times \bar{B} = \mu_0 \bar{J} + \mu_0 \epsilon_0 \frac{d\bar{\epsilon}}{dt}$$

Meanwhile Faraday's manipulations were far more imaginative

\hookrightarrow Gave rise to a new kind of $\bar{\epsilon}$ which resembles \bar{B} rather than regular static $\bar{\epsilon}$

Maxwell's manipulation

\downarrow

We get 'same kind' of $\bar{B} \Rightarrow \nabla \cdot \bar{B} = 0$

$\nabla \times \bar{B} = \text{something}$
 \downarrow
Just the value changes

Imagination starts here



In a space, there exists

$$\bar{\Sigma} (\vec{r}, t)$$

$$\bar{B} (\vec{r}, t)$$

$$\left. \begin{array}{l} \rho = 0 \\ j = 0 \end{array} \right\} \text{in immediate environment}$$

$$\nabla \times \bar{\Sigma} = - \frac{d\bar{B}}{dt}$$

$$\nabla \times \bar{B} = \mu_0 \epsilon_0 \frac{d\bar{\Sigma}}{dt}$$

$$\nabla \times (\nabla \times \bar{\Sigma}) = - \nabla \times \left(\frac{d\bar{B}}{dt} \right)$$

$$= - \frac{d}{dt} (\nabla \times \bar{B})$$

$$= - \mu_0 \epsilon_0 \frac{d^2 \bar{\Sigma}}{dt^2}$$



$$\boxed{\nabla^2 \bar{\Sigma} = \mu_0 \epsilon_0 \frac{d^2 \bar{\Sigma}}{dt^2}}$$

$$\nabla (\nabla \cdot \bar{\Sigma}) - \nabla^2 \bar{\Sigma}$$

$$\stackrel{0}{(\rho = 0)}$$

Now, it's one dimensional
wave

$$\frac{d^2 f}{dz^2} = \frac{1}{v^2} \frac{d^2 f}{dt^2} \rightarrow \text{wave eqn in one dimension}$$

Imagination



Light is an
electromagnetic
wave consisting
of $\bar{\Sigma}$ & \bar{B} moving
at speed c

Describes a wave w/
wave function $\bar{\Sigma}$ moving
with speed $\frac{1}{\sqrt{\mu_0 \epsilon_0}} = c$

Maxwell proved his leap in imagination by deriving laws of reflection and refraction from electromagnetic equations ↓

Maxwell's next leap in imagination ↓

There exist other kinds of EM waves other than (visible) light having their own wavelength & frequency →

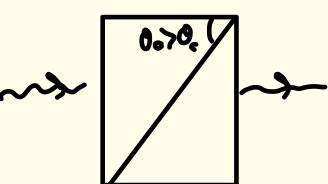
Radio waves

Produced by Hertz

Earlier accepted to be empirical laws ↓
Eqns derived purely from a large number of observations

Soon enough, microwaves were discovered

↓
Experimented → Crucial in proving Maxwell's split on by JC Bose



wave eqn in the form of

↓

fields must be continuous at all points



reflection

Transmission?

derivatives

Consequence of req'n of the fields to be continuous

TIR \Rightarrow 100% energy reflected

field exists, but dissipates quickly

No energy propagation

LEC-13

Maxwell's EM Theory

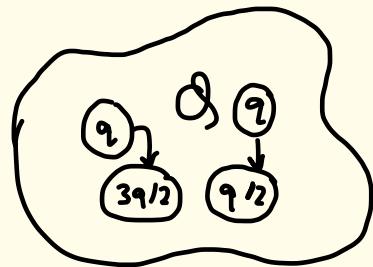
→ Ampere's law is not consistent with local charge conservation

→ formulation of displacement current

However, charge conservation is also applicable in infinitesimally small volumes

↓
locality

↓
we can trace the path of movement of a charge.



charge conservation in large volume

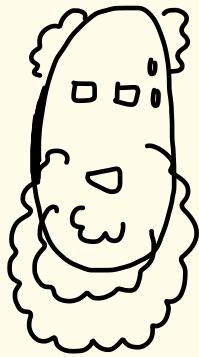
In regions of $\rho = 0, \bar{J} = 0,$
 $\bar{\epsilon}, \bar{B}$ satisfy

$$\frac{1}{c^2} \frac{\partial \bar{\epsilon}^2}{\partial t} - \nabla^2 \bar{\epsilon} = 0 \Rightarrow \text{light is an EM wave}$$

↓
wave eqn

Lots of people weren't
comfy with
this idea

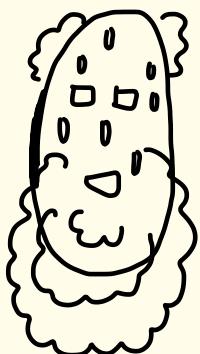
Up till that point, every wave had a medium it propagated in, so what medium did light propagate through?



Umm... there is a medium, it's called...
Aether... yup... *hodor*

So nervous
that it passes +
through all
living matter

Such strong
mechanical properties
that it could transmit
a wave of speed
of the order 10^8 ms^{-1}



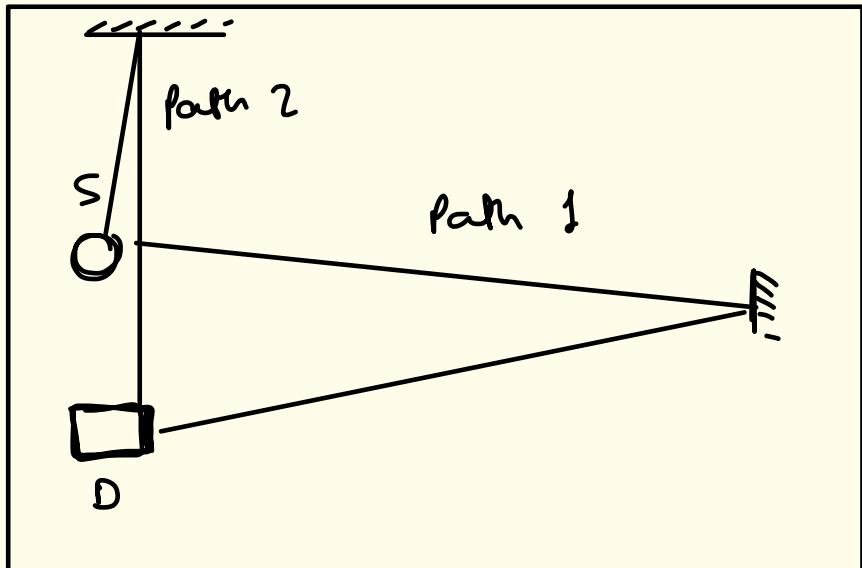
Yeah... of course...
what? Why doesn't
that make sense?
It makes sense, trust
me bro...

Scientists trusted the bro

Now, how to detect this Aether thing?



Michelson
Experiment



Path 1 \perp Path 2

$$l_1 = l_2$$

↓

Constructive interference should → occur at detector

Some people argued
↳ Orient path 1 along motion of the Earth

Speed of travelling light in path 1 would be distorted by movement of Earth, we wouldn't get perfectly constructive interference, there would be some fringe shift

resolution of Earth around sun ($v_{\text{orb}} \gg v_{\text{rel}}$)

Expected → 0.4 Error → ± 0.1 Observed → 0

over the years, Michelson improved his experiment

Expected $\rightarrow 40$ Error $\rightarrow \pm 0.1$ Observed $\rightarrow 0$

So what is happening with Aether?



- Uh well actually-

Shut up you're dead now

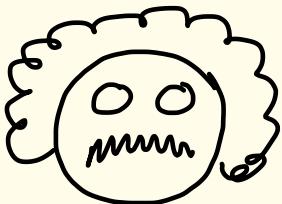
Proposed explanation - Lorentz - Fitzgerald Contraction



Could explain
Michelson's results,
but not further
experiments

'Lengths along direction
of motion are
contracted'

But then came the man, the myth, the
one and only...



- No Aether lol

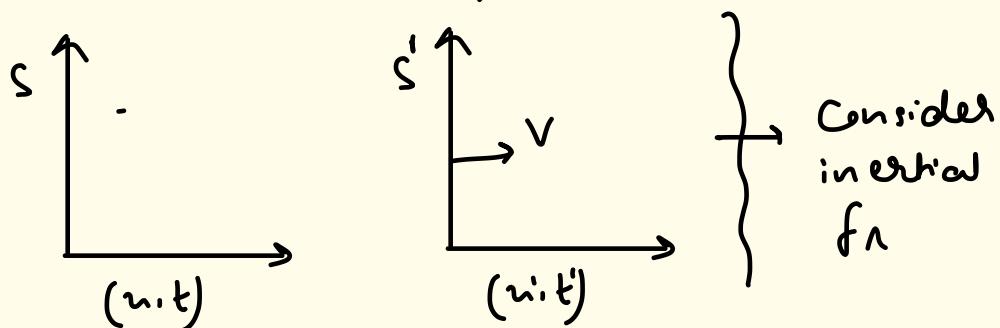
why try to forcibly make sense of something
that doesn't make sense?

Concept of inertial frame \rightarrow laws of Physics should have same form in two different frames namely
 constant relative velocity

↙
 Another way
 to state Newton's
 1st Law

\Downarrow
 Uniqueness
 version of what
 Einstein stated

Galilean Relativity



$$u' = u - vt$$

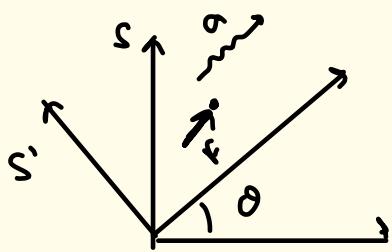
$f = m\bar{a}$
 $F = m\bar{a}'$

force
 wouldn't
 change

Acceleration
 is seen as
 the same in
 the two frames

\bar{a} remains constant
 on transformation
 from S to S' , just
 like \bar{F}

* \bar{F} & \bar{a} are invariant
 under Galilean transformations



$$u' = u \cos\theta + y \sin\theta$$

$$y' = u \sin\theta + y \cos\theta$$

$$\bar{F} = m a_n \hat{i} + m a_y \hat{j}$$

$$\bar{F}' = m a_n \hat{i} + m a_y \hat{j}$$

* \bar{F} & \bar{a} are covariant under rotation

$$\bar{F} \rightarrow \bar{F}'$$

$$\uparrow \downarrow \text{Same transformation}$$

$$a_n \leftrightarrow a'_n$$

$$a_y \leftrightarrow a'_y$$

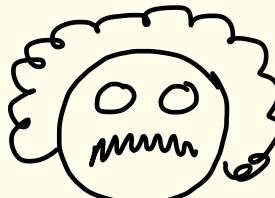
Now, looking at $\Sigma \Delta M$ eqns

$$\nabla \cdot \bar{\mathbf{E}} = \frac{f}{\epsilon_0} \quad \nabla \times \bar{\mathbf{E}} = - \frac{\partial \bar{\mathbf{B}}}{\partial t}$$

$$\nabla \cdot \bar{\mathbf{B}} = 0 \quad \nabla \times \bar{\mathbf{B}} = \mu_0 \left[\bar{\mathbf{J}} + \epsilon_0 \frac{\partial \bar{\mathbf{E}}}{\partial t} \right]$$

Under rotation,
they are covariant

However, under Galilean transformations,
they are neither
invariant nor covariant



- Why y'all
be worried
'bout some
stupid ass
Action?

Einstein pointed
out this very
crisis

How does Einstein solve this crisis? Stay tuned,
we shall be back after a short (2 weeks) break

LEC-14

Theory of Relativity

Arose from conflict b/w Newton's first law and

EM Theory

Neither invariant
nor covariant under
Galilean transformation

laws of Physics
should have the
same form in all
inertial frames
of reference

Newton defined inertial frames in terms of an
"absolute space" → critiqued by Mach



④ Foucault's pendulum

People solved problems by
considering a frame attached
to Earth as a good enough
approximation for absolute
space

However, Foucault argued ...

Objects on surface of Earth affected by acceleration
(due to Earth's rotation)

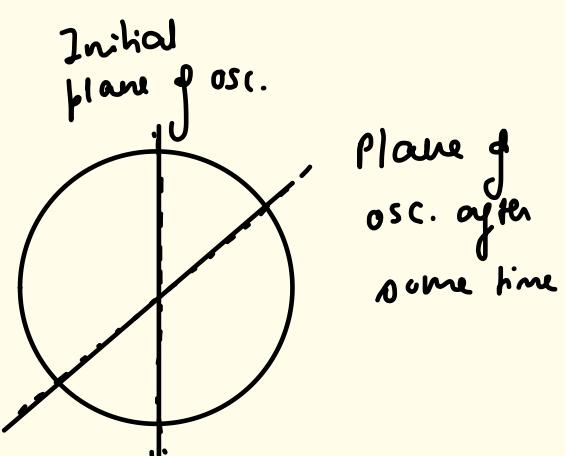
→ Centrifugal force

→ Coriolis force

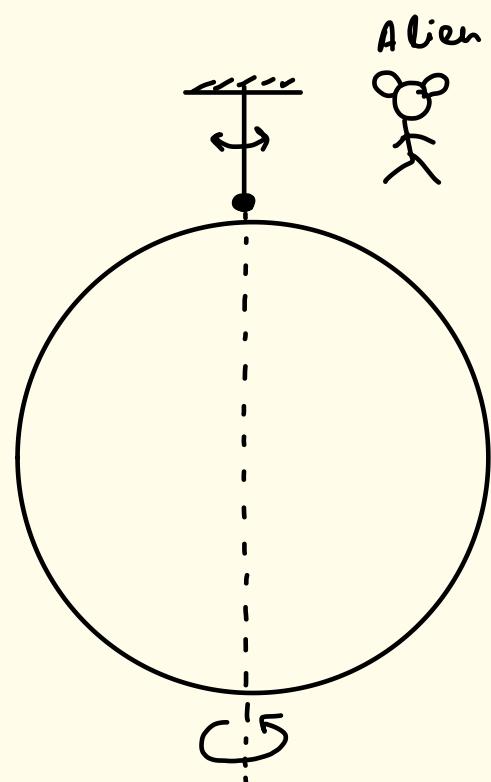
Foucault argued we should start seeing
effects of Coriolis force with a large enough pendulum
over a large enough time period

Earth is
a rotating
frame

↓
Turned
out to be
true



Time period of revolution
of plane = $24 \text{ hrs} \times \sin \theta$



Time period of revolution of
plane of osc. = 24 hrs

But an alien present in
a coordinate system such
that it sees the Earth's rotation
observes the plane of oscillation
as constant.

↪ How does the pendulum
'know' to keep its pl. of osc. constant?

Galilean transform

$$\left. \begin{array}{l} \rightarrow u' = u - vt \\ \rightarrow y' = y \\ \rightarrow z' = z \\ * \rightarrow t' = t \end{array} \right\} \approx$$

Can be shown that EM forces change on Galilean transformation.

Woldemar Voigt found a transformation such that EM eqns were invariant

$$\rightarrow u' = u - vt$$

$$\rightarrow y' = \frac{y}{\gamma}$$

$$\boxed{\gamma = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}}}$$

$$\rightarrow z' = \frac{z}{\gamma}$$

$$\rightarrow \boxed{t' = t - \frac{vu}{c^2}}$$

why?

Others who worked on this problem



Fitzgerald

Lorentz

Poincaré

Recalling the electromagnetic wave eqn,

$$\mu_0 \epsilon_0 \frac{\partial^2 \bar{\epsilon}}{\partial t^2} = \nabla^2 \bar{\epsilon}$$

Now, if there was an observer moving at speed c , it is logical to assume that they see a static solution \rightarrow however, no static soln exists for $\bar{\epsilon}$.

Einstein did similar math to Voigt and came up w/ same soln $\rightarrow \underline{\underline{t' \neq t}}$

examining the consequences of such a relation, Einstein published two papers on special relativity.

1st: $\underline{\underline{u' = \gamma(u - vt)}}$

$\underline{\underline{y' = y}}$

$\underline{\underline{z' = z}}$

$\underline{\underline{t' = \gamma \left(t - \frac{vu}{c^2} \right)}}$

$\xrightarrow{\text{Lorentz transformations}}$

↓

name given
first by Poincaré

2nd: $\underline{\underline{E = mc^2}}$

Einstein did here two postulates, the acceptance of which allows us to derive any relativistic phenomena

- The laws of physics must have the same form in all inertial frames (i.e. covariant)
- Speed of light is constant in all inertial frames.

LEC-16

for $\underline{c \rightarrow \infty}$, Lorentz Transfⁿ \rightarrow Galilean Transfⁿ
 !!!

$v \ll c \rightarrow$ True for most day-to-day scenarios

But in Einstein's 2nd postulates, how does the speed of light suddenly pop up?



People had other points of view that tried to do away with Einstein's 2nd postulate \rightarrow Proposed replacement
 ↓
 { Transformation equations should form a group
 ↓
 Eqs linking $n' \rightarrow n$ & $t' \rightarrow t$



This postulate implies the existence of a universal velocity

\hookrightarrow Universal velocity unmeasurably large
 ↓

Galilean transfⁿ

\hookrightarrow Universal velocity = c
 ↓

Lorentz transfⁿ

* However, there must always be a time delay while making a measurement
 \hookrightarrow we get different values in the same frame
 ↓
 we assume there to be an infinite number of observers located at every point in spacetime, all with synchronised clocks

- Confusing shit about special relativity
- Time dilation → Pole - Galilei's paradox
 - Length contraction → Loss of Simultaneity
 - Twin Paradox

* Lorentz invariant interval

$$\rightarrow (\Delta s)^2 = c^2 (\Delta t_{21})^2 - (\Delta u_{21})^2$$

True if $c|\Delta t_{21}| > |\Delta u_{21}|$

Symbol
to denote
RHS

↳ Not always
true

Looking at Lorentz transf.,

$$u' = \gamma(u \cdot v t)$$

$v \rightarrow$ speed of signal

$$t' = \gamma\left(t - \frac{v u}{c^2}\right)$$

b/w events 1 & 2

→ Any events in universe

$$= \left| \frac{\Delta u_{21}}{\Delta t_{21}} \right| < c \text{ if } LII > 0$$

→ Light pulse transmitted from one event to another

$$= c \text{ if } LII = 0$$

$$> c \text{ if } LII < 0$$

↳ Events 1 and 2 are 'disconnected'

Causal Connection



One event occurred due to another event occurring

✳ If events are causally connected,
 LII must be positive.
 Inverse not necessarily true

i.e. no causal connection

LFC-17

Length Contraction

$L = \text{Proper length}$

\downarrow
length of a
rod wrt a

frame where it
appears at
rest

$$L_{\text{moving}} = \frac{L}{\gamma}$$

$$\gamma = \frac{1}{\sqrt{1 - v^2/c^2}}$$

Time Dilation \rightarrow Can be applied only in very specific circumstances



Event 1: first Hd
 $(u_1, t_1), (u'_1, t'_1)$

Event 2: 2nd Hd
 $(u_2, t_2), (u'_2, t'_2)$

$$\Delta t'_{21} = \gamma (\Delta t_{21} - \frac{v}{c^2} \Delta u_{21})$$

$$\Delta u_{21} = 0 \quad \Delta u'_{21} = u'_2 - u'_1$$

$$\therefore T' = \gamma T$$

$$\Delta t_{21} = T \quad \Delta t'_{21} = T'$$

* Works only when $\Delta u_{21} = 0$
i.e. the time interval b/w
two events is dilated in any
other frame only when
the two events occur at the
same spatial position in
some frame

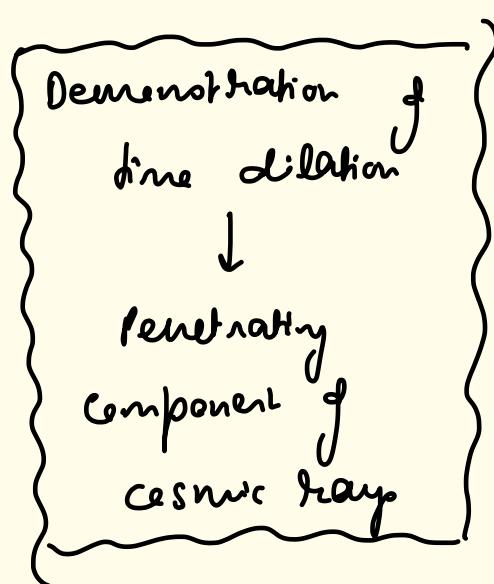
* whenever $|L| > 0$ b/w two events, then there exists a frame where the two events occur at the same spatial position i.e. $\Delta x_{21} = 0$. In such a frame, the period b/w the two events would be minimum

\Rightarrow In any other frame,

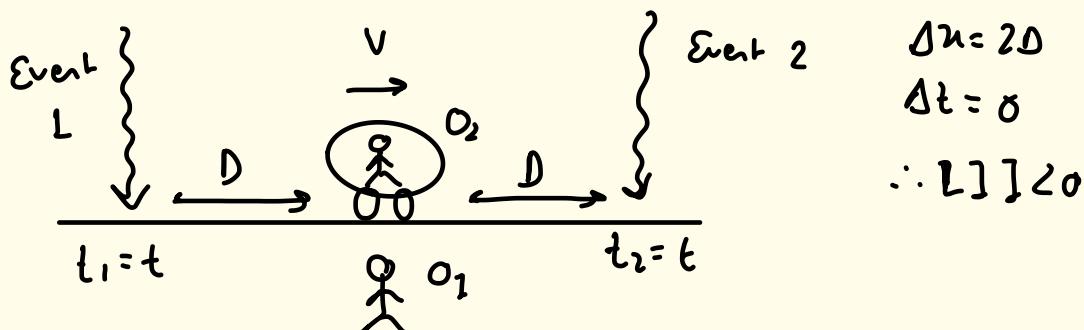
$$\Delta t'_{21} = \gamma \Delta t_{21} \rightarrow \text{Always } > 1$$

$$\downarrow \Delta t'_{21} \text{ & } \Delta t_{21}$$

always have the same sign, the occurrence of the two events is always in the same order in every frame of reference



* It is impossible to truly say that two spatially separated events are 'simultaneous'



$$\begin{aligned}\Delta x &= 2D \\ \Delta t &= 0 \\ \therefore L &\neq 0\end{aligned}$$

O_1 in stationary frame sees events occurring simultaneously
 O_2 in moving frame sees Event 2 occurring first
 Other observers see Event 1 first

When $L \neq 0$, the order of events depends on the reference frame

* Pole vaulter's paradox & Twin paradox \rightarrow HC Verma

LEC-18

Some further topics to be discussed

→ Non-invariance / Covariance of Newton's 2nd law
on Lorentz transformations



Discussion
of energy and
momentum

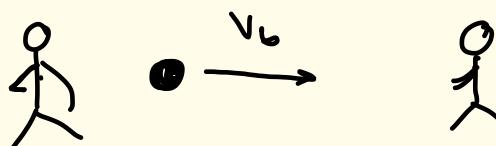
→ Velocity addition

Lorentz transformations

$$\Delta u' = \gamma (\Delta u - v \Delta t)$$

$$\Delta t' = \gamma \left(\Delta t - \frac{v \Delta u}{c^2} \right)$$

Imagine



wrt frame at rest,

$$v_b = \frac{\Delta u}{\Delta t}$$

WRT frame moving with speed v ,

$$v_b' = \frac{\Delta u'}{\Delta t'}$$

$$v_b' = f(v_b, v) \text{ such that if } v_b = c, v_b' = c$$

This condition
gives us velocity
addition function

However, looking at Lorentz transformations,

' Δu ' signifies direction of velocity of the object

$$\bar{v} = v_x \hat{i} + v_y \hat{j}$$

↳ We cannot take

two Lorentz transformation,
one along \hat{i} & one
along \hat{j}

$$\bar{v} = v_x \hat{i} + v_y \hat{j} = \underbrace{\sqrt{v_x^2 + v_y^2}}_{\hat{v}} \quad \begin{array}{l} \text{Take Lorentz transf.} \\ \text{along } \hat{v} \end{array}$$

Moreover,

$$v_y \text{ (rather } v_1 \text{ direction of Lorentz transf)} = \frac{\Delta y'}{\Delta t'} = \underbrace{\frac{\Delta y}{\Delta t}}_{\hat{v}} + \underbrace{\frac{\Delta y}{\Delta t}}$$

All the more reason
to take Lorentz transf.
in direction of velocity

Energy & Momentum in Relativity

$$\bar{p}_{\text{non-rel.}} = m \bar{v} \rightarrow \boxed{\bar{p}_{\text{rel.}} = \gamma m \bar{v}} \quad \gamma = \frac{1}{\sqrt{1 - v^2/c^2}}$$

$$E_{\text{rel.}} = \gamma m c^2$$

* Def" of mass in rel.
Rest mass $\rightarrow m_0$

'Relativistic mass' $\rightarrow m = \gamma m_0$

We do not talk about this

$$\text{For } v \rightarrow 0, \underline{\underline{p_{\text{rel}}} \rightarrow 0}$$

$\underline{\underline{E_{\text{rel}} \rightarrow mc^2}}$ → A particle with mass has an intrinsic energy

* E_{rel} & 3 components of $\underline{\underline{p_{\text{rel}}}}$ form a four vector

$$\begin{aligned} E^2 - \underline{\underline{p^2 c^2}} &= \gamma^2 m^2 c^4 - \gamma^2 m^2 v^2 c^2 \\ &= \frac{c^2}{c^2 - v^2} \cdot m^2 c^2 (c^2 - v^2) \\ &= m^2 c^4 \end{aligned}$$

↓
has a Lorentz invariant interval

$$\therefore E = \sqrt{m^2 c^4 + \underline{\underline{p^2 c^2}}}$$

$$= mc^2 \left(1 + \frac{\underline{\underline{p^2}}}{(mc)^2} \right)^{1/2}$$

$$= mc^2 \left(1 + \frac{\underline{\underline{p^2}}}{2(mc)^2} \right) \quad \left[\text{in limit of } \underline{\underline{p}} \ll mc \right]$$

$$E = \underline{\underline{mc^2}} + \frac{\underline{\underline{p^2}}}{2m} \rightarrow KE$$

$$\underbrace{\qquad\qquad\qquad}_{\downarrow}$$

Energy content of
a body in non-relativistic
limit

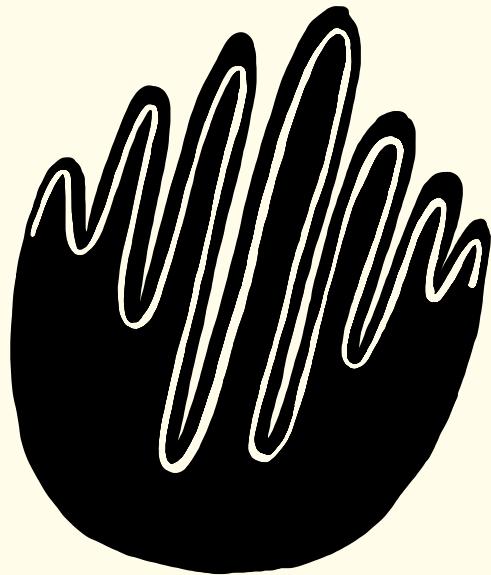
Newton's 2nd law

$$\underline{\underline{F}} = \underline{\underline{\frac{dp}{dt}}}$$

* $\frac{|\underline{\underline{p}}|c}{E} = \frac{|\underline{\underline{v}}|}{c}$

General Relativity - A glimpse

Not much, just some non-Euclidean geometry type shit.



LEC-19

Ancient mathematics — natural no. operations — -ve nos bad — solution of quadratic eqn — we don't talk about roots of -ves

Shift to ~1500 years later \Rightarrow Now the drama starts

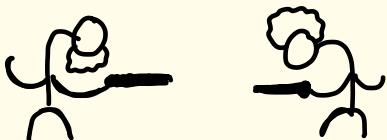
Scipione del Ferro \rightarrow figured out how to solve a reduced cubic



Didn't publicise his solution



Monetary reasons

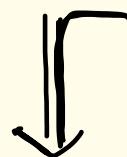


Mathematical duels

Being able to solve a problem no one could solve was a bag ooh up your sleeve

$$x^3 + px + q = 0$$
$$p, q \in \mathbb{Z}$$

Only told one guy — his son-in-law Antonio Fiore



Made a bunch of money winning duels over this

Eventually beat by a guy Nicolo Tartaglia \rightarrow figured it out over the course of the duel

Eventually Tartaglia told this solution to Girolamo Cardano on the condition he keep it secret



Cardano used
the result & further
generalised the
problem to
solve $\rightarrow u^3 + au^2 + bu + c = 0$

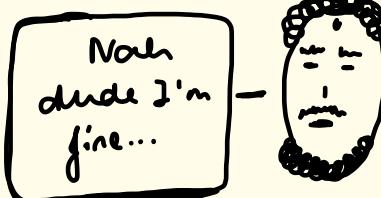
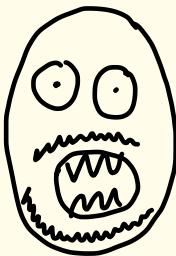
However, Cardano figured some important mathematical results in the meanwhile which he wanted to publish, but was unable to do so as it used Tartaglia's result.

Then he realised \rightarrow If Ferre posed the problem he must know the solution



Tartaglia was ←
pissed...

Talked to
him, and figured it
was okay to
publish Ferre's result
instead of Tartaglia's
result



okay new drama over back to math

Cardano's advancement → Proved that equations involving square roots of -ve ws were solvable

↓
His algorithm
for solving cubic
equations

1. Integers \Rightarrow Start

:

m. $\sqrt{-1} \Rightarrow$ Some intermediate

:

n. Rational \Rightarrow end

Earlier people would just give up whenever they came across the square root of a -ve w.

↓
Cardano proved these ws could be worked with

Eventually quartics were solved

~1800 it was proved by Abel that it was impossible to find the solution to a general quintic eqn

Galois also proved the same for quintic using diff argument
that could also be extended for arbitrarily higher order equations

↓
Ended up dead at 23 because he fucked around and fund out (+ government conspiracy) (+ duels [non-mathematical])

↓
Introduced the concept of a Group

Group

→ A set

→ May have finite/infinite no. of elements

→ Must satisfy a set of properties

- $a, b \in \text{group } G \Rightarrow \underline{axb} \in G$

"Group multiplication"

but not really multiplication,
it could be any operation where
we get a third element from
combination of two elements

- $\exists \underline{e} \in G$ s.t. $\underline{axe} = \underline{e}xa \forall a \in G$

identity

- $\exists \underline{a^{-1}} \in G$ s.t. $\underline{axa^{-1}} = \underline{a^{-1}xa} = e \forall a \in G$

inverse

- Associativity $\rightarrow ax(bxc) = (axb)xc$

* axb need not be bxa

If $axb = bxa \forall a, b \in G$,

then $G \rightarrow$ Abelian Group

$C_1 \rightarrow$ group w/ 1 element

$C_2 \rightarrow$

$C_3 \rightarrow$

cyclic
groups of order
 n

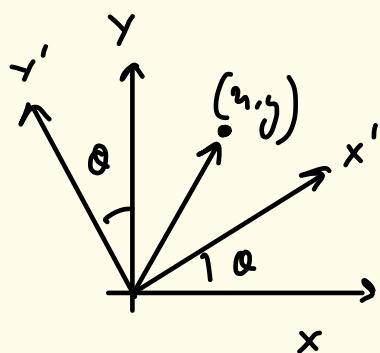
A group w/ n elements
is not possible. But groups
w/ 1, 2, 3 elements possible -

$$\left\{ \begin{array}{l} C_1 (1) \rightarrow n-1=0 \\ C_2 (1, -1) \rightarrow n^2-1=0 \\ C_3 (1, \omega, \omega^2) \rightarrow n^3-1=0 \end{array} \right.$$

Continuous parameter groups



Set of $n \times n$ matrices with $\det \neq 0$ → called $GL(n, \mathbb{R})$
 General linear matrix
 of n dimensions on
 the set of real no.s



$$\begin{bmatrix} u' \\ y' \end{bmatrix} = \begin{bmatrix} \cos\theta & \sin\theta \\ -\sin\theta & \cos\theta \end{bmatrix} \begin{bmatrix} u \\ y \end{bmatrix}$$

$$= \underbrace{R(\theta)}_{\downarrow} \begin{bmatrix} u \\ y \end{bmatrix}$$

Similarly,

$$\begin{bmatrix} u' \\ y' \end{bmatrix} = \begin{bmatrix} u \\ y \end{bmatrix} R^t(\theta)$$

Defines a relational matrix to transform a vector from one coordinate system to another.

$$\begin{bmatrix} u' \\ y' \end{bmatrix} \begin{bmatrix} u' \\ y' \end{bmatrix} = \begin{bmatrix} u \\ y \end{bmatrix} R^t(\theta) R(\theta) \begin{bmatrix} u \\ y \end{bmatrix}$$



'length' of
vector

$$= \begin{bmatrix} u \\ y \end{bmatrix} \begin{bmatrix} u \\ y \end{bmatrix} \rightarrow R^t(\theta) \cdot R(\theta) = I$$

$$= R(\theta) R^t(\theta)$$

Consider set of matrices R s.t.

$$\underbrace{R^t R = R R^t = I}_{\text{Orthogonality condition}} \rightarrow \text{Set of } R \text{ matrices forms a group}$$

* $\det(R) = +1$ or -1

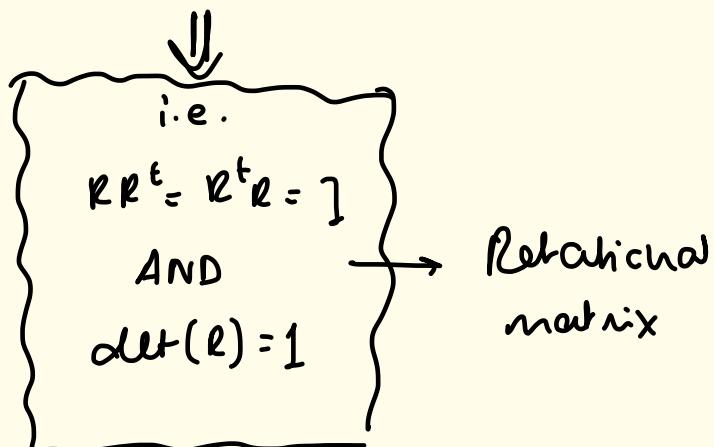
↪ R s.t. $\det(R) = +1 \Rightarrow$ is also a group \rightarrow group of rotations

↪ R s.t. $\det(R) = -1 \Rightarrow$ NOT A GROUP

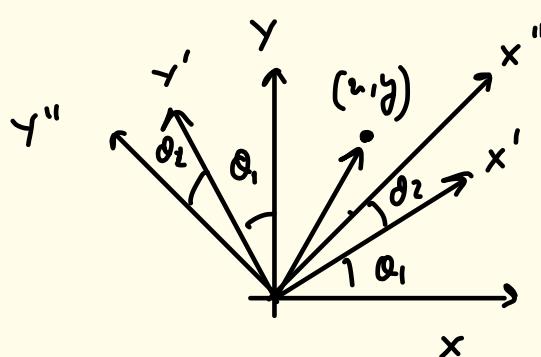


group of matrices
having $\det = -1$ cannot
have identity

\therefore Rotation matrices are those which satisfy orthogonality equation and have determinant = 1



Applying a second rotation to the given axis



$$\begin{bmatrix} x'' \\ y'' \end{bmatrix} = \underbrace{R(\theta_2) R(\theta_1)}_{R(\theta_1 + \theta_2)} \begin{bmatrix} x \\ y \end{bmatrix}$$

$$= R(\theta_1) R(\theta_2)$$

↓
Rotations in 2D commute

* $\text{Cryp} \rightarrow \underline{\text{SO}(2)}$

Special orthogonal of 2 dimensions

$$\det = 1 \quad RR^t = R^t R \\ = I$$

* $R = \begin{bmatrix} a & b \\ c & d \end{bmatrix} \rightarrow 4 \text{ nos.}$

However, rotational matrices need only one parameter
i.e. the angle of rotation

∴ How does one reduce 4 points of variation to 1

↓
Apply three
constraints

$$RR^t = \begin{bmatrix} 1 & c \\ 0 & 1 \end{bmatrix}$$

for arbitrary vector product, we would
need 6 constraints to solve this eqn

However, we must remember that

RR^t is a symmetric
↳ 3 independent
constraints

Counting the constraints,
we have 2 diagonal constraints
and 1 off-diagonal constraint

Remember, so far, purely mathematical
arguments, no physics has actually
started

Now let's move a dimension up

Taking set of 3×3 matrices satisfying constraints of rotational matrices

$$R R^t = I \rightarrow R = 3 \times 3 \text{ matrix}$$

$$= \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \quad \begin{array}{l} \text{Using argument of} \\ \text{symmetry, we} \\ \text{now have } \underline{6 \text{ independent}} \\ \underline{\text{constraints}} \end{array}$$

9 numbers $\xrightarrow{6 \text{ ind. const.}}$ 3 points of variation

Similarly, in
4 dimensions
 \downarrow
 $4c_2 = 6$ points
of variation

x, y, z
Rotation in
different planes
 \downarrow
 $x_1, y_2, z_3 \rightarrow \underline{3c_2}$ points
of variation

\therefore In three dimensions, we have 3 independent rotations $\rightarrow \theta_x, \theta_y, \theta_z$ (Rotation abt each axis)

$$R(\theta_x, \theta_y, \theta_z) = R(0, \theta_x, 0) R(0, \theta_y, 0) R(0, 0, \theta_z)$$

↓

$$SO(3) = \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos\theta_x & \sin\theta_x \\ 0 & -\sin\theta_x & \cos\theta_x \end{bmatrix} \begin{bmatrix} \cos\theta_y & 0 & -\sin\theta_y \\ 0 & 1 & 0 \\ \sin\theta_y & 0 & \cos\theta_y \end{bmatrix}$$

$$\begin{bmatrix} \cos\theta_z & \sin\theta_z & 0 \\ -\sin\theta_z & \cos\theta_z & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

* Rotations in 3d DO NOT COMMUTE

↓

As well as
higher dimensions

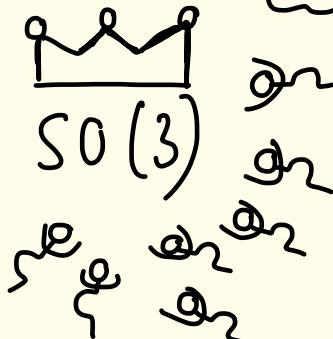
Non-Abelian
Groups

$$R(0, \theta_x, 0) R(0, \theta_y, 0) - R(0, \theta_y, 0) R(0, \theta_x, 0) \neq 0$$

The value of this difference follows a certain structure that rules the entire landscape of Physics

All hail SO(3)

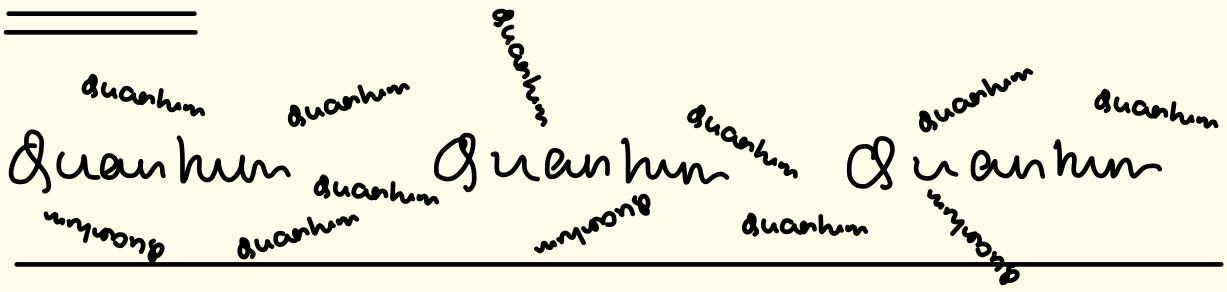
Worship SO(3)
with your
heart and soul



SO(3) is the
truth of life and
the universe

Probability

Gaussian distribution stuff



One mole of ideal gas at temperature T

$$\text{One mole of He at room temp.} \Rightarrow U = \frac{3}{2} RT \\ (300\text{K}) \quad \sim 4000\text{J}$$

Taken to a height of 100m

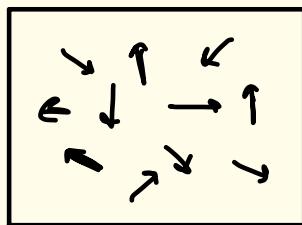
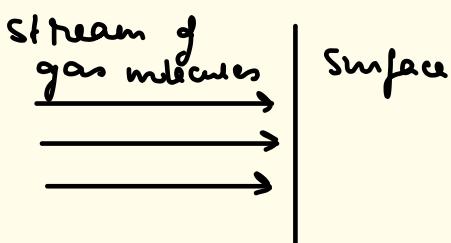
Mass of 1 mol He = 4g

$$\text{Potential energy} = mgh \\ = 4\text{J}$$

Total energy = 4004J, out of which we can utilise only 4J as work

Internal energy \rightarrow due to KE of gas molecules

We can't get any energy out of it as motion of gas molecules is completely random



the system as a whole does work on the surface } All gas molecules have a certain momentum, which changes uniformly on collision with a surface

Net momentum of system = 0
large no. of randomly moving particles \rightarrow their cumulative impact cancels out

Entropy

$$\Delta S = T \Delta S$$

Helmholtz Free Energy

$$H = U - TS$$

= 0 for a stationary gaseous system

* Entropy is crucial in our conception of temperature



we need internal energy as well as entropy



we can only define temperature of a system when there is thermodynamic equilibrium



Maximum possible value of $\rightarrow H=0$ entropy

$$v \in (0, \infty)$$

Probability distribution function

$$P(v) dv = (1) e^{-\frac{\frac{1}{2}mv^2}{kT}} dv$$

Probability of a particle to have speed in the range v to $v+dv$

Boltzmann factor $\equiv e^{-\frac{E}{kT}}$ Generalised entropy in any form

Blackbody Radiation

↓
Perfectly
absorbs & emits
light of
all frequencies

* Stefan's Law
Power emitted = σT^4

↓
 Σ Energy emitted
over all Frequencies

Emissivity = $A f(\nu, T)$

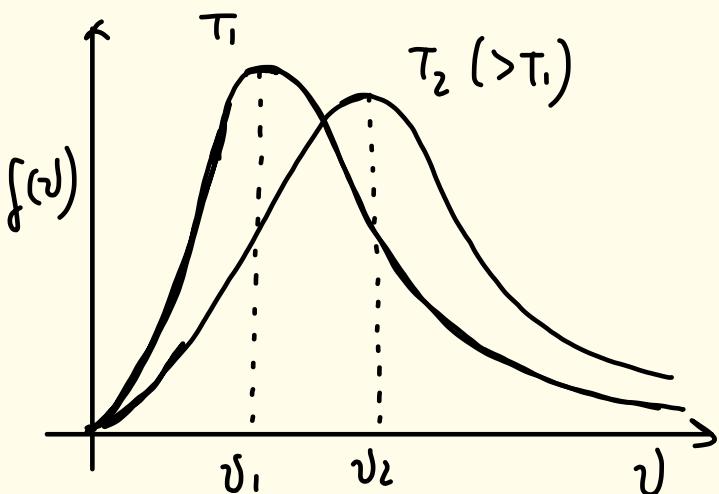
* Wien's Law



Black object
immersed in an
environment of radiation,
in thermal equilibrium
radiation at temp. T

$f(\nu, T) \sim (\nu)^3 e^{-\frac{E}{kT}}$

$e^{-\frac{E}{kT}} \leftrightarrow$ maybe tried to
imply a relation
between energy & freq.



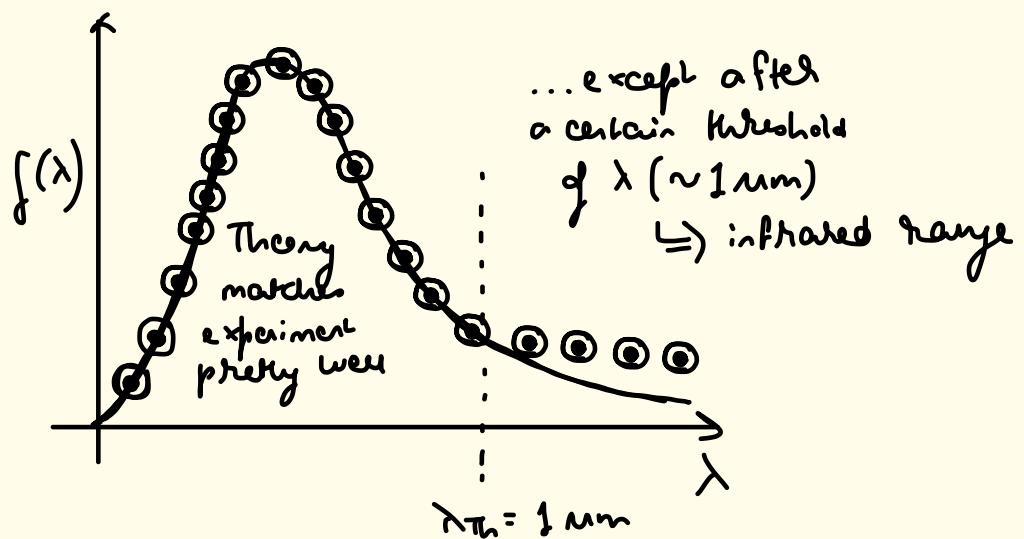
went against
ideas of classical
electromagnetism
↓

Energy related to
amplitude of electromagnetic
wave

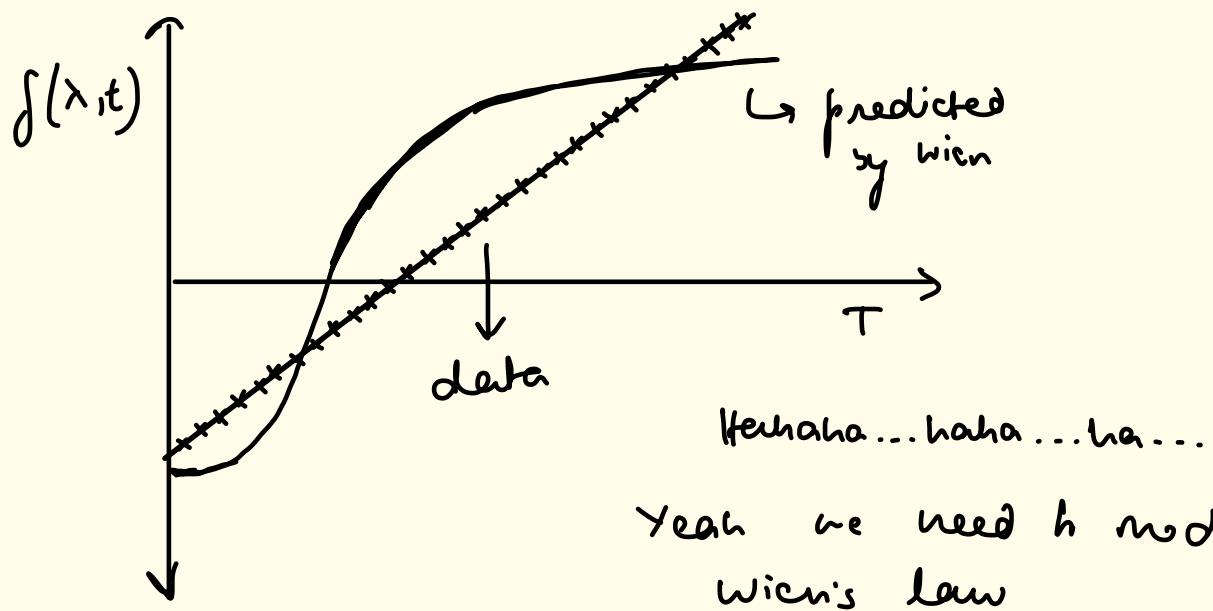
Wien's displacement law →

$\lambda_{max} T = \text{constant}$

Comparing theory with experiment,



Moreover, at far infrared frequencies, fixing λ and merely looking at variation w/ temperature



Max Planck came up with a solution

$$f(v, T) = \left(\right) v^3 \frac{1}{e^{\frac{hv}{kT}} - 1} : \begin{array}{l} v \uparrow \Rightarrow \text{Wien's original law} \\ v \downarrow \Rightarrow \text{Linear eqn} \end{array}$$

How did Planck justify this change without
disobeying laws of statistics



Replace \int with \sum

[Declared that energy exists
in quantised bundles
of min. unit $\rightarrow h\nu$]

Einstein went further in his explanation of the
photoelectric effect



[Energy only
exists in 'packets' of
 $h\nu$, unit size is
never $2h\nu$, $3h\nu$ or
so on.]

→ In order to
maximise
entropy

Energy packets = photons



LEC-22

Why do we need Quantum Theory?

- (a) Explain blackbody radiation
- (b) Atoms radiate at discrete frequencies
- (c) Photoelectric effect

Mathematical backbone arose from (a) & (c),
but (b) is perhaps the most significant.

↓

Hydrogen spectra

Dark lines in absorption spectrum
Bright lines in emm. sp.

Lydsey formula →

$$E_{nm} = (-13.6 \text{ eV}) \left(\frac{1}{n^2} - \frac{1}{m^2} \right)$$

$m > n$

↓

As energy levels are quantised, so is frequency

Niels Bohr attempted to explain the absorption & emission spectra



First stated that classical rules merely don't apply at atomic scales

- * Bohr's postulates

 1. Angular momentum of electrons is quantised

$$L = nh$$

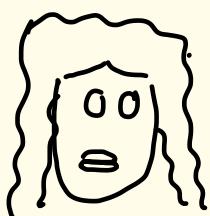
 2. Only orbits satisfying above rule are allowed

Considering electrostatic force to provide centrifugal force, the eqn for Energy we get

$$E_n = \frac{-k}{n^2}, \quad k = 13.6 \text{ eV}$$

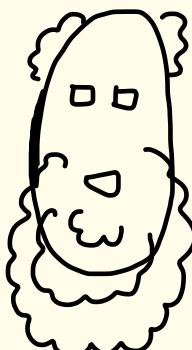
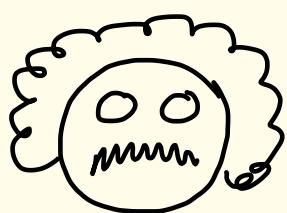
Matched eqns known at the time

Bohr ended up being totally wrey, but he introduced the idea of quantisation into atomic models



↓
Light doh
be corpuscular

Watch me double slit
experiment all over your
theory



Yup, plus
look at
Electromagnetism
wooooow

He's a quantum
she's a quantum
Light's a quantum
Your mom's a quantum

Light do be pretty wavy

Einstein said → Light is a particle in that it appears as a bundle of energy

→ Light is a wave in that it has frequency & wavelength

now a guy chap called de Broglie
looked at electrons and thought "Hmmm..."

de Broglie → ~~light~~^{electron} is a particle in that it
said appears as a bunch of energy
→ ~~light~~^{electron} is a wave in that it has
frequency? Hmmm...



Eventually other
wave properties
of electrons were
discovered

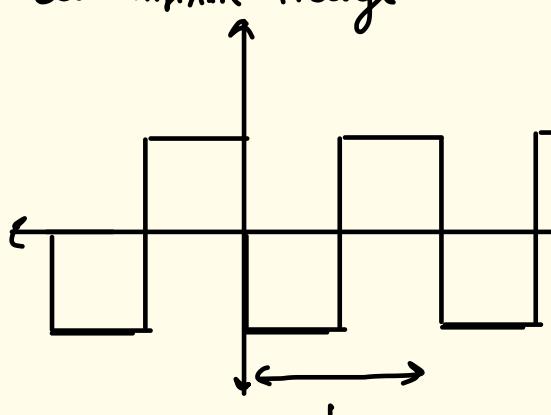


found a relation
for energy of an
electron in terms
of its momentum,
which could be related
to its v & λ

A mathematical solution to the problem was
already available \Rightarrow Fourier Analysis

1. Fourier Series

Infinite series over
an infinite range

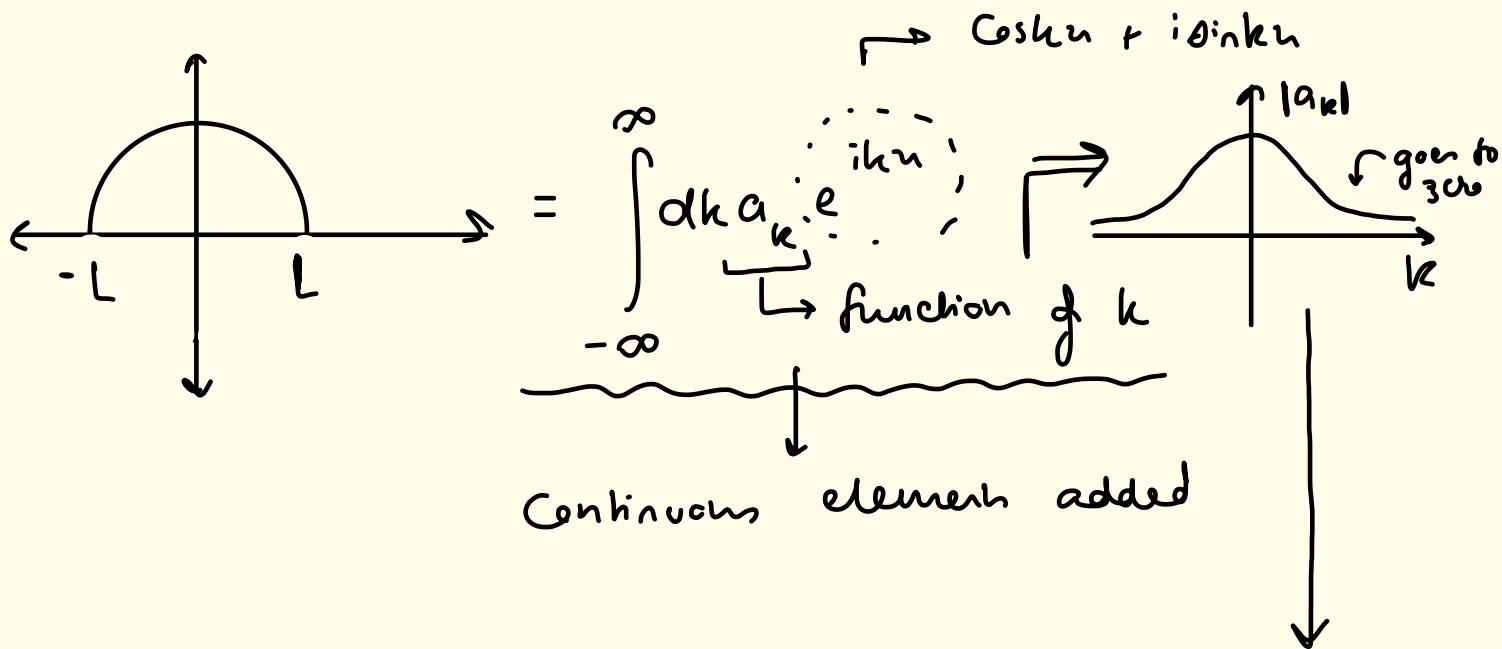


$$= \sum_{n=0}^{\infty} \left[A_n \cos\left(\frac{2\pi n u}{L}\right) + B_n \sin\left(\frac{2\pi n u}{L}\right) \right]$$

Adding discrete harmonics

$$f(u+L) = f(u)$$

2. Fourier Integral



Taking $f(u) \rightarrow$ has nonzero values in a limited range

Most important thing of Fourier analysis

Corresponding to it we have a function a_k

↓
nonzero values in limited range → falls to zero faster

Say $f(u): (-L, L)$

$a_k: (-k_0, k_0)$

Range of f in x

Range of a in k

!!

Number of \Rightarrow single digit
order 1

Representing an electron by a wave function
↓
↓
uncertainty Δn

Superposition of
infinite w. of waves
with wave numbers

in range (k_{\min}, k_{\max})

↓
Uncertainty
 $k_{\max} - k_{\min} = \Delta k$

$\Delta n, \Delta k \rightarrow \text{Order 1}$

LEC-23

Major developments in quantum theory

1. Energy is quantised (Planck & Einstein)
↓
Radiation

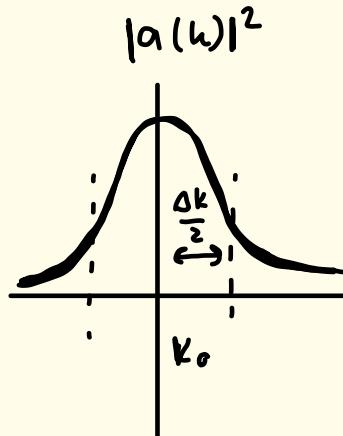
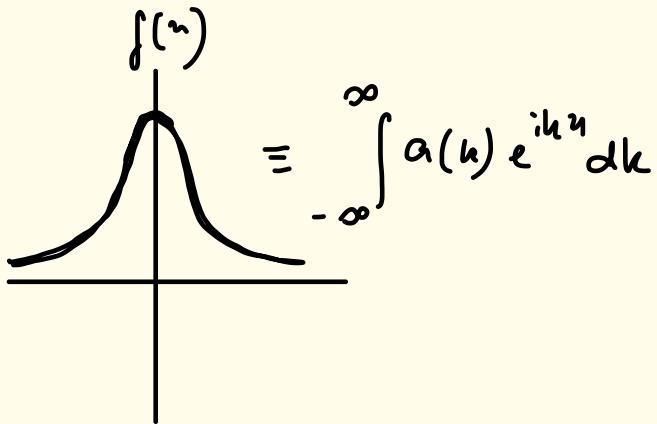
2. Angular momentum is quantised (Bohr)

↓
Energy of
electrons is
quantised

* Stern Bohr orbits \Rightarrow fit an integer no. of
de Broglie wavelengths

de Broglie's insight \Rightarrow you need a wave eqn to
represent a 'point' particle

wave fn $\rightarrow f(\vec{r}) = f(\vec{u})$ in one dimension



Restricting our choice of wave no.s (k)

$$k \in \left(k_0 - \frac{\Delta k}{2}, k_0 + \frac{\Delta k}{2} \right)$$

Schrödinger

→ Thorough knowledge of differential equations

→ de Broglie \Rightarrow particles show wave properties

Schrödinger \Rightarrow Every wave should have a wave eqn

↓
Always partial differential equations

$$\Psi(\vec{r}, t) \rightarrow \Psi(\vec{u}, t)$$

in one D

Considering a free particle

$$E = \frac{p^2}{2m}$$
 (classical mechanics)

$$(de\text{Broglie} \rightarrow p = \frac{h}{\lambda} = \hbar k)$$

$$\hbar\omega = \frac{\hbar^2 k^2}{2m}$$

(Planck)

∴ The wave eqn of this particle -

$$\Psi(u, t) = (\) e^{i(ku - \omega t)}$$

$$\hbar\omega = \frac{\hbar^2 k^2}{2m}$$

↓

$$\hbar\omega \Psi(u, t) = \frac{\hbar^2 k^2}{2m} \cdot \Psi(u, t)$$

\downarrow

$$i \frac{d\Psi}{dt}$$

$$\Rightarrow i\hbar \frac{d\Psi(u, t)}{dt} = -\frac{\hbar^2}{2m} \cdot \frac{d^2\Psi(u, t)}{du^2}$$

Schrödinger's wave eqn.

How we usually write it



$$i\hbar \frac{d\Psi}{dt} = \underline{H}\Psi$$



Hamiltonian

Special case when potential only depends on spatial coords, not on velocity or time



$$H = PE + KE$$

$$= \frac{|\vec{p}|^2}{2m} + V(\vec{r})$$

In general coordinates,

$$\vec{p} \longleftrightarrow -i\hbar \nabla$$

$$\vec{p} \cdot \vec{p} = p^2 \longleftrightarrow -\hbar^2 \nabla^2$$

General Schrödinger eqn

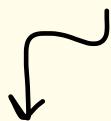


$$i\hbar \frac{\partial \Psi}{\partial t} = \left[-\frac{\hbar^2}{2m} \cdot \nabla^2 + V(r) \right] \Psi$$

Time dependent
Schrödinger Equation

* Suppose $V(r) = \frac{1}{2} m \omega^2 r^2$

$$= \frac{1}{2} k \hbar^2$$



$$\begin{aligned} E_n &= n \cancel{k} \cancel{\hbar^2} \\ &= n \cancel{k} \omega \end{aligned} \quad \left. \begin{array}{l} \text{Planck} \\ \text{ } \end{array} \right\}$$

$$E_n = \left(n + \frac{3}{2} \right) \hbar \omega$$

related to
angular momentum
quantisation

$\frac{3}{2} \hbar \omega$ = zero point energy

Time independent Schrödinger eqn.

$$i\hbar \frac{\partial \Psi(n,t)}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \Psi}{\partial n^2} + V(n) \cdot \Psi$$

Schrödinger solved

the problem for

$$V(r) = -\frac{e^2}{r}$$



Got same value for
Bohr's energy, but
didn't get angular
momentum quantisation



Got something
more intricate.

Schrödinger said that angular
momentum quantisation not
only dependent on n (principal
quantum no.), but on two more
numbers



l, m



* Uniqueness thm



for a diff eqn with
certain initial conditions,
there is a unique soln
to the eqn.

Assuming $\Psi(u,t) = T(t) \cdot \Psi(u)$

$$i\hbar \frac{d\Psi}{dt} = \left(-\frac{\hbar^2}{2m} \frac{d^2\Psi}{du^2} + V(u) \cdot \Psi \right) T(t)$$

$$i\hbar \frac{1}{T} \frac{dT}{dt} = \frac{1}{\Psi} \left(-\frac{\hbar^2}{2m} \frac{d^2\Psi}{du^2} + V(u) \Psi \right)$$

\downarrow
only
time
related
term

\downarrow
Only u
related terms

\downarrow
only
time
related
term

\downarrow

$$C(t) = f(u) \rightarrow C(t) = f(u) = \text{some const. } E$$

for all u & all t

\downarrow
only possible
solution

$$\therefore \frac{dT}{dt} = -\frac{i}{\hbar} \cdot T E \Rightarrow T(t) = T(0) e^{-\frac{iEt}{\hbar}}$$

(similar to radioactive decay law)

Moreover,

$$\left[-\frac{\hbar^2}{2m} \frac{d^2}{du^2} + V(u) \right] \Psi(u) = E \Psi(u)$$

$H\Psi = E\Psi \xrightarrow{\text{Energy of particle w/ wave fn } \Psi}$

* Eigenvalues & Eigenvectors

$M \rightarrow n \times n$ matrix

$\bar{v} \rightarrow n \times 1$ column matrix

$M\bar{v} = \bar{w} \rightarrow n \times 1 \Rightarrow \bar{v}$ is called an eigenvector
of M with eigenvalue λ if

'Matrix M multiplied'
by vector \bar{v}

$$\underline{M\bar{v}} = \underline{\lambda\bar{v}}$$

'Matrix M acts upon'
vector \bar{v}

Schrödinger's eqn is essentially an eigenvalue
problem for differential eqns instead of matrices

$$H\Psi = E\Psi$$

$\Psi(n)$ is the energy eigen for
of the operator H with
energy eigen value E .

* Sturm-Liouville Theory

↓
what kind of differential operators have
the solution of eigenvalue form.

$$\left. \begin{array}{l} Df = \lambda f \\ \vdash \end{array} \right.$$

→ Energy eigenvalues are
quadrised

Take certain free parameters → we get solns for certain integer values of the free parameters

Schrodinger eqn



$$i\hbar \frac{d\Psi(u,t)}{dt} = H\Psi(u,t) \Rightarrow \text{Time dependent Schrodinger eqn (TDSE)}$$

Time independent Schrodinger eqn (TISE)



$$H\Psi(u) = E\Psi(u) \Rightarrow \text{Eigen value eqn}$$

↓ ↓
 differential Number
 operator

For differential operators there is a set of functions such that the operators act on these functions in a way so as to reproduce the function in a scaled form.

Energy quantisation arises as $\Psi(\vec{r})$ is not well behaved unless E takes well-defined discrete values.

\therefore Quantisation arises not due to Schrodinger's equation, but as a property of differential equations.

Example of quantisation in classical physics



Vibration modes
of a stretched
string

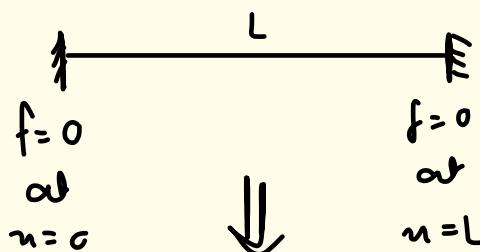
Defining a differential operator $D = \frac{d^2}{dx^2}$

$$Df = ()f \rightarrow f = A \cos kx + B \sin kx$$

k can take any value

$$Df = (k^2)f$$

Imposing conditions on f



This is the same problem we solve in quantum mechanics

We get well-behaved waveforms
only for certain discrete values
of k
↓
wave number
is quantised

But Schrödinger couldn't clarify the physical meaning behind $\Psi \rightarrow$ just a mathematical entity used to arrive at energy quantisation

Max Born proposed \rightarrow Probability interpretation

$|\Psi(u)|^2 du$ is probability of finding the particle
in du neighbourhood of u

↓

$|\Psi(\vec{r})|^2 dV =$ probability of finding particle in
 dV neighbourhood of \vec{r} .

$\int_{\text{universe}} |\Psi(\vec{r})|^2 dV = 1 \rightarrow$ for Ψ to be well behaved,

1. $\Psi(\vec{r})$ has to be finite everywhere

2. $\int |\Psi(\vec{r})|^2 dV \rightarrow$ finite for any volume

↓

$\Psi(\vec{r}) \rightarrow 0$ as $\vec{r} \rightarrow \infty$

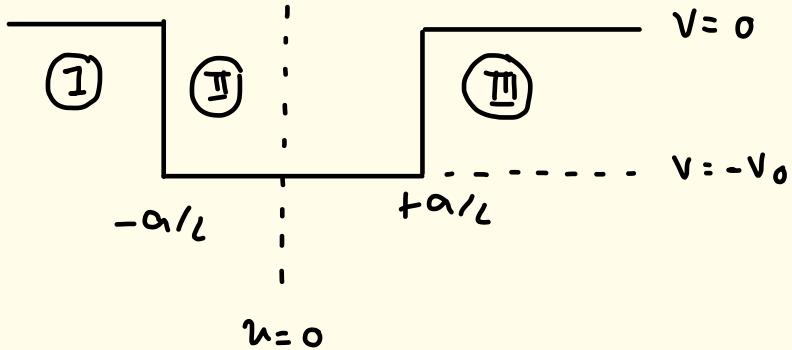
3. $\frac{d^2\Psi}{du^2}$ must exist

↓

Ψ & $\frac{d\Psi}{du}$ must
be continuous

Finite potential well

find the bound states
of this well



Three regions \textcircled{I} , \textcircled{II} , \textcircled{III} each with their own form of wave function of the particle Ψ_I , Ψ_{II} , Ψ_{III}

In a classical system,

$$\begin{array}{l} \text{lower bound} = -V_0 \\ \text{upper bound} = 0 \end{array} \quad \left\{ \rightarrow -V_0 < E < 0 \right.$$

E may have values in a continuous range from $-V_0$ to 0

However, in a quantum mechanical system,

$$-V_0 < E < 0 \rightarrow \text{still true,}$$

but there are a finite number of valid bound states

$$\textcircled{I}: -\frac{\hbar^2}{2m} \cdot \frac{d^2\Psi_I}{dx^2} = E \Psi_I = -|E| \Psi_I \quad \Rightarrow \text{form of wave function in region I}$$

$$\Rightarrow \frac{d^2\Psi_I}{dx^2} = \frac{2m|E|}{\hbar^2} \Psi_I \quad \Psi_I = \propto^z \Psi_I$$

$$\rightarrow \Psi_I = A e^{ikx} + B e^{-ikx}$$

As $x \rightarrow -\infty$

$\Psi_I \rightarrow 0$

$$\text{Similarly, } \Psi_{\text{III}} = F e^{-\alpha k u} + G e^{\alpha k u}$$

$$\text{(II)}: -\frac{\hbar^2}{2m} \cdot \frac{d^2 \Psi_{\text{II}}}{du^2} + (-V_0) \Psi_{\text{II}} = -|E| \Psi_{\text{II}}$$

$$\therefore \frac{d^2 \Psi_{\text{II}}}{du^2} = -\frac{2m}{\hbar^2} (V_0 - |E|) \Psi_{\text{II}} = -k^2 \Psi_{\text{II}}$$

$$\therefore \Psi_{\text{III}} = C \cos k u + D \sin k u$$

Ψ & $\frac{d\Psi}{du}$ both continuous at $u = \pi/2$

$$C \cos k \frac{a}{2} = A e^{-\alpha k a/2}$$

$$C k \sin k \frac{a}{2} = \alpha A e^{-\alpha k a/2}$$

↓

$$\tan k \frac{a}{2} = \frac{\alpha}{k}$$

$$\Rightarrow \boxed{\frac{ka}{2} \tan \frac{ka}{2} = \frac{\alpha a}{2}}$$

$$k^2 + \alpha^2 = \frac{2m}{\hbar^2} (V_0 - |E|) + \frac{2m|E|}{\hbar^2} \Rightarrow \boxed{\left(\frac{ka}{2}\right)^2 + \left(\frac{\alpha a}{2}\right)^2 = \frac{m V_0 a^2}{2 \hbar^2}}$$

Theorem → If potential has symmetry, then the lowest energy state shares the same symmetry

↓

As Ψ is symmetric

about y axis in regions beyond $\pm a/2$, so should be

the case in

the potential well.

↓

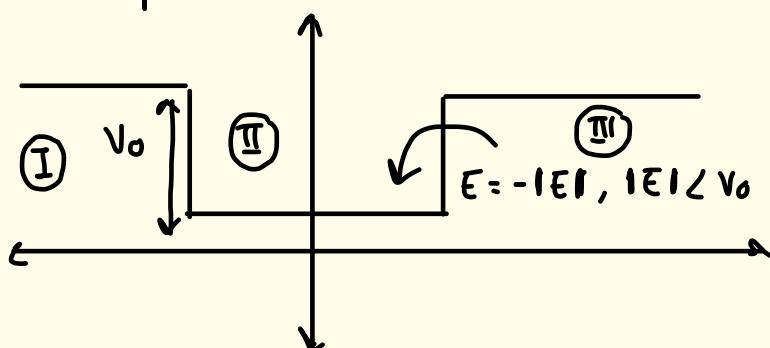
Ψ must be an even function, & sin is odd fn

Enforcing b'dary conditions gives us quantisation of energy



LEC-2S

Finite potential well



In region II, $E > -V_0$,

$$\therefore E > PE \Rightarrow KE > 0$$

But in regions I & III,
 $E < 0$, $PE = 0 \Rightarrow KE < 0$



Non-zero probability
of particle being located
here \Rightarrow K.E. - v.c.

We can't
observe the
particle in
these spaces,
but they
do exist

$$\int_{-\infty}^{\infty} |\Psi|^2 du = 1$$

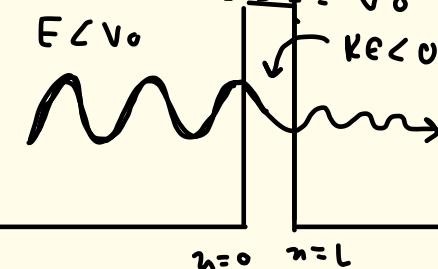
$|\Psi|^2 du$ = Prob of particle being
b/w in a width du

$$\frac{p^2}{2m} = -|E|$$

$$\Rightarrow p = \sqrt{-2m|E|}$$

So momentum is
imaginary?

"Quantum Quantum
Quantum"



\Rightarrow Non-zero
probability
of quantum
tunneling

Tunneling through
a barrier

George Gamow

→ Science

~50 years ago

- One, Two, Three infinity

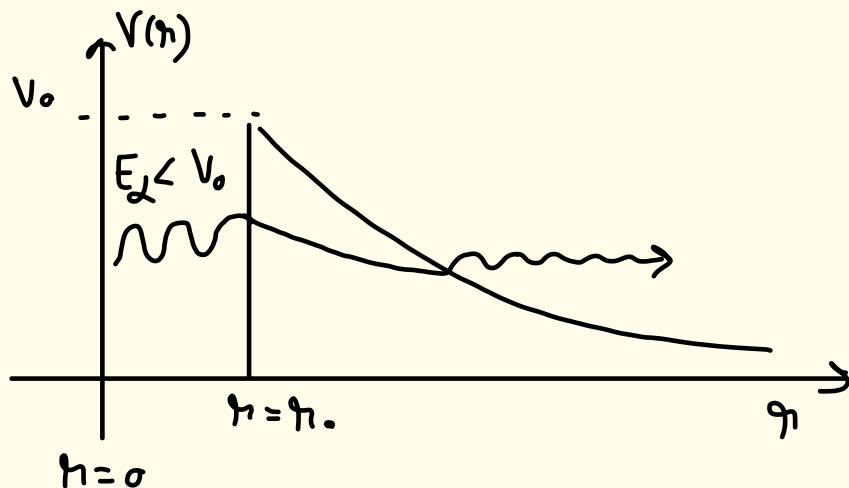
- Mr. Tompkins in Wonderland

Explained
concept of d.
decay using
quantum tunneling

$$_{\text{Pl}}^{239} \equiv \underbrace{\text{U}^{235}} + \alpha^4$$

loosely bound
together \rightarrow strong nuclear
force

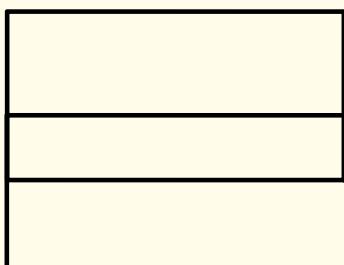
Bound = Potential well



Also explains properties of solids \rightarrow Electrical conduction
earlier theory \Rightarrow free electrons

using quantum mechanics \Rightarrow Band Theory of
solids

Conductors

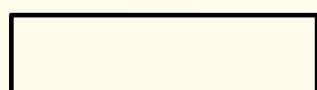


Conduction
Band

Valence
Band

Highest energy
band above which
there are no e⁻s

Allowed energy
levels form continuous
bands



Insulators



Conduction Band

Valence Band

Zener diode

Zener assumed the existence of materials where by modulating strength of electric field, the energy levels get affected

* Dirac → Relativistic version of Schrödinger's Equation
!!!

Dirac Equation

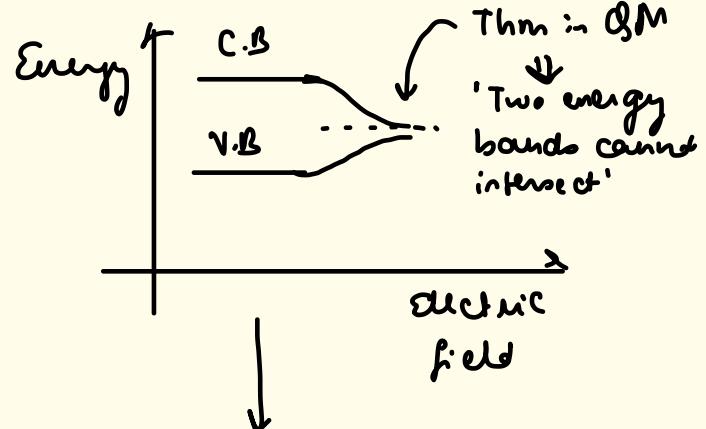
Suggestions of negative energy particles

Why don't we see -ve energy electrons around us?

Dirac: All -ve energy states are occupied

Everyone: ... wtf?

Dirac: All e's around us are excitations of -ve energy e's that we consider the ground state



Eventually such materials were discovered with diode properties imagined by Zener

$$\dots \dots \dots E = mc^2$$

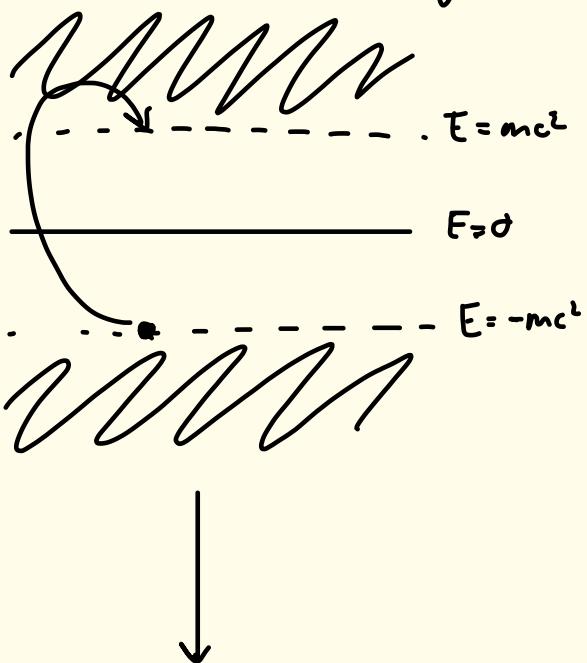
$$E=0$$

$$\dots \dots \dots E = -mc^2$$

$$\begin{aligned} \text{Dirac eqn} &\Rightarrow |E| > mc^2 \\ &\Rightarrow E > mc^2 \text{ or } E < -mc^2 \end{aligned}$$

Dirac eventually went on to mathematically prove this sh^2

Even more crazy shit

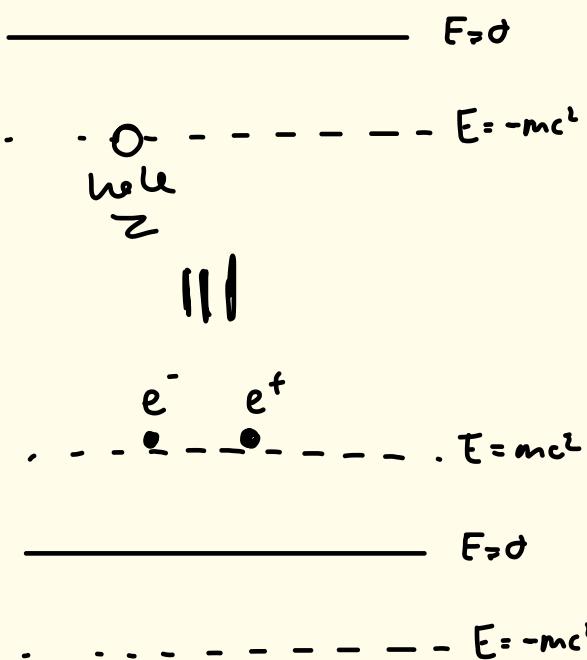


Suppose an e^- in the -ve energy levels was given enough energy to move into a +ve level



Formation of an e^- -hole pair

Now how do we comprehend this hole?



Hole = lack of -ve charge
+ve charge

However, this hole also

lack of -ve energy
+ve energy

Dirac eqn suggested the existence of a particle with the same physical properties as an electron but with +ve charge

→ Now called a positron → we eventually observed this particle
↓
Antimatter exists

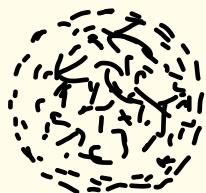
* Measurement in QM

Say we want to take the problem of e⁻'s bound in an atom

$\Psi(\vec{r})$ describes the 'state' of the electron

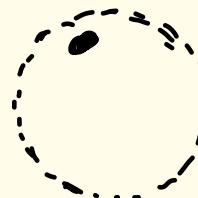
↓
Any physical
quantity we want
to measure

$|\Psi^2(\vec{r})| d^3\vec{r} \rightarrow$ Prob of finding an e⁻ in
 $d^3\vec{r}$ neighbourhood of \vec{r}



Ground state
energy eigenvalue
of the electron

Position measurement



Collapse of
wavefunction
into a small
region

LEC-26 → * Bonus - Prof Sasi Edition

Early era of Physics - How does the world work?

Modern era of Physics - Can we apply techniques of Physics to stuff very closely related to our daily lives?

Now, Quantum computers



I give you computer \Rightarrow How to give it ultimate performance
→ Hardware Constraints
→ Software Constraints

If the laws of Physics are true, and the Universe is causal, then

Bell's Inequality

Must be true

"Measurement outcomes are inherently random"

Quantum mechanics \Rightarrow Can't do "phase space physics"

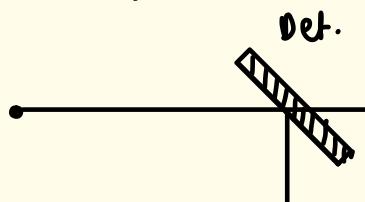
$$\nabla = m\ddot{a}$$

Physics involving momentum & stuff

If you send a single photon through



Detector will give you one click



output \Rightarrow inherently random

No amount of previous data can predict future data

'Inherently random'



It is impossible to predict future outcomes no matter the observer or previous state collected

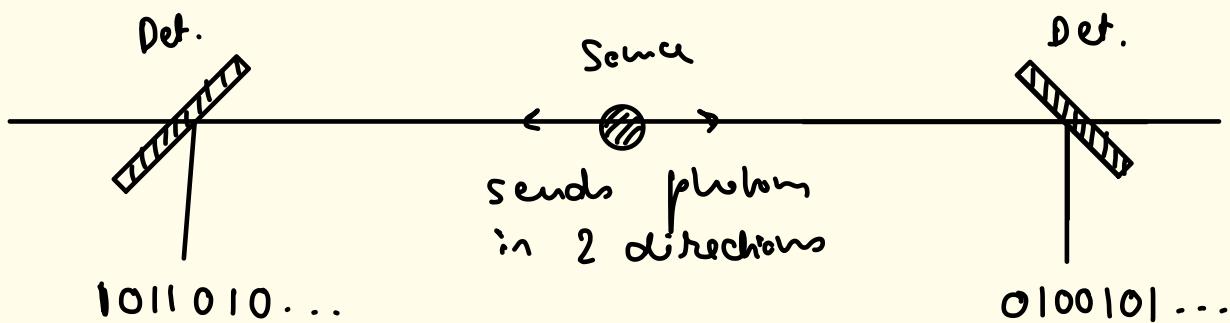
It is simply impossible to comprehend physics of small particles through lens of u & p

↓
Why?

Law of large numbers

↓
Statistical mech

Proof → EPR paradox



Observation → Results in each direction are
anti-correlated

* Tanger - Bertman's socks



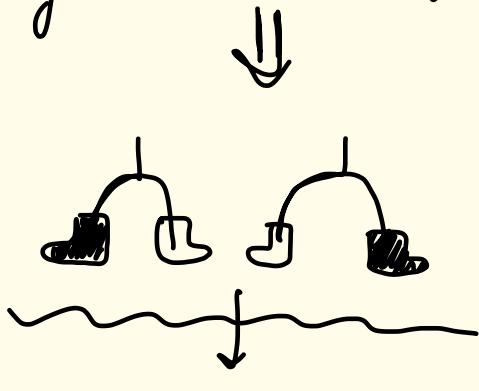
Used to wear different colored socks in each foot.

Say someone is a very good friend of Bertman's

↓
Has a pair
of blue socks

Bertman → pair of
red socks

They both exchange a sock



↓
Their socks
are 'anti-correlated'

Say Bertman is in
another galaxy
far far away → apparent
violation of causality

↑
No matter
how far Bertman
and his friend
are from each other

Now, say someone
hasn't met them
that day, but knows
of their peculiar
behaviour

↓
say the person meets
Bertman's friend.
As soon as they see
the friend's feet, they

instantaneously knew
what Bertman is
wearing

Differences

- Preparation uncertainty vs Idealized preparation
- In quantum mechanics, the colour of the socks isn't defined beforehand, it 'pops into existence' as soon as we make a measurement

Paradox lies in the fact that identical preparation leads to random outcomes AND anti-correlation still holds

According to EPR, maybe there is some 'hidden variable' which affects things

How does stuff hide?

→ Size

→ Lack of interaction strength

Einstein said that GM is nothing but CM with some hidden variables

↳ Philosophical + Causality reasons



'Bro eff all
this quantum BS
classical for the
win !'

Bell's intuition \rightarrow Classical mechanics generate probability theories
 \downarrow
Worked at CERN

If classical mechanics is correct

$$\langle A_0 B_0 \rangle + \langle A_0 B_1 \rangle + \langle A_1 B_0 \rangle - \langle A_1 B_1 \rangle < 2$$

Every result $\rightarrow \pm 1$

pre-determined reality

If Q.M is correct

$$\underbrace{\langle A_0 B_0 \rangle + \langle A_0 B_1 \rangle + \langle A_1 B_0 \rangle - \langle A_1 B_1 \rangle}_{< 2\sqrt{2}}$$

No pre-determined reality

You have to take a (?:) wide range of outcomes, which expands the range of possibilities.

