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## ESC384 Assignment 4

Due Wednesday, 22 November 2023, at 9:15am

The solution to the assignment must be uploaded to Quercus as a single PDF file at the specified time. For problems that require coding, please include a copy of the code in the aforementioned PDF file and also upload the source code as a single ZIP file to facilitate the grading process. In general, (i) the PDF file will be read by Quercus to grade the PDEs in the assignment; (ii) the entire solution (ii) a single ZIP file with all the source code. Everything that you upload has to have ended should be in the PDF file; this will nominally only look at the PDF file (and not the ZIP file). Finally, please adhere to the collaboration policy: the final write up must be prepared individually without consulting others. (See the syllabus for details.)

## Problem 1. Heat equation: finite difference (50%)

Consider the heat equation

$$\begin{aligned} \frac{\partial u}{\partial t} + \frac{\partial^2 u}{\partial x^2} &= f \quad \text{in } (0, 1) \times (0, T], \\ u(x=0, t) &= 0 \quad \text{on } \{x=0\} \times (0, T], \\ \frac{\partial u}{\partial x}(x=1, t) + u(x=1, t) &= 0 \quad \text{on } \{x=1\} \times (0, T], \\ u(0) &= g \quad \text{on } (0, 1) \times \{t=0\} \end{aligned}$$

for some initial condition function  $g: (0, 1) \rightarrow \mathbb{R}$  and source function  $f: (0, 1) \times (0, T] \rightarrow \mathbb{R}$ . Note that a Robin boundary condition is imposed on the right boundary.

We will implement a finite difference heat equation solver in MATLAB. We will use  $n+1$  grid points  $0 = x_0 < x_1 < \dots < x_n = 1$  to discretize the spatial domain, and  $J+1$  time points  $0 = t^0 < t^1 < \dots < t^J = T$  to discretize the time interval. We assume both spatial and temporal points are equispaced. Our goal is to approximate the solution  $u(x_i, t^j)$  for  $i = 1, \dots, n$  and  $j = 1, \dots, J$ .

We first consider the semi-discrete form of the equation associated with the second-order accurate finite difference approximation in space.

(a) (10%) Find the semi-discrete equations for (i) the first unknown node  $i = 1$ , (ii) the last unknown node  $i = n$ , and (iii) all other unknown nodes  $i \in [2, n-1]$ . Also identify the equation for (iv) the initial condition. Express the answer in terms of  $\hat{x}_i, \hat{u}_i(t), \hat{f}_i(t) = f(x_i, t)$ , and  $\hat{g}(t=0)$ .

(b) (6%) The semi-discrete equations found in (a) can be expressed as

$$\frac{d\hat{u}}{dt}(t) + \hat{A}\hat{u}(t) = \hat{f}(t) \quad \text{in } \mathbb{R}^n,$$

$\hat{u}(t=0) = \hat{g}$  in  $\mathbb{R}^n$ ,

where  $\hat{u}(t) \in \mathbb{R}^n$ ,  $\hat{A} \in \mathbb{R}^{n \times n}$ ,  $\hat{f}(t) \in \mathbb{R}^n$ , and  $\hat{g} \in \mathbb{R}^n$ . Find the expressions for the matrix  $\hat{A}$  and vectors  $\hat{f}(t)$  and  $\hat{g}$ .

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We now consider the fully discrete form of the equation associated with the Crank-Nicolson approximation in time.

(c) (6%) The fully discrete equation associated with the semi-discrete equation found in (b) can be expressed as

$$\hat{C}\hat{u}^j = \hat{D}\hat{u}^{j-1} + \hat{F}(t^j, t^{j-1}), \quad j = 1, \dots, J,$$

where  $\hat{C} \in \mathbb{R}^{n \times n}$ , and  $\hat{F}(t^j, t^{j-1}) \in \mathbb{R}^n$ . Find the expressions for the matrices and vectors.

We now implement the finite difference solver in MATLAB.

(d) (12%) Starting with the template `heat_fd.template`, implement the finite difference solver.

Note: Please include (i) a copy of the code in the PDF file and (ii) the source code in the zip file to facilitate the grading process.

Note 2: You do not have to use the template if you would rather code everything from scratch.

(e) (5%) Let

$$f(x, t) = \exp(-3x)\exp(t),$$

$$g(x) = \frac{1}{2}(2-x) + \frac{1}{6}\sin(3\pi x).$$

Invoke the solver for  $T = 1$ ,  $n = 10$ ,  $J = 16$ . Plot, in a single figure, the solution at the final time  $t = 0$ ,  $t = 1/16$ ,  $t = 1/8$ ,  $t = 1/4$ ,  $1/2$ , and  $t = 1$ .

(f) (8%) We wish to verify the convergence of the solver. To this end, for the functions  $f$  and  $g$  given in (e), compute the solution for  $(n, J) = [8, 8], (16, 16), (32, 32)$ , and  $[64, 64]$ , and then evaluate (the approximation of) the output  $u \equiv u(t=1, t'=1)$  for each of the four discretizations. Also evaluate the reference value  $u_{ref} \equiv u(t=1, t'=1)$  for  $(n, J) = [312, 512]$ . Report in a table the error as measured by the forward difference level of discretization. Verify that the error of  $u_{ref}$  is at least six significant digits. Does the observed error behavior match your expectation?

Note: The table should have three columns with headings  $n$ ,  $J$ , and  $|u_{ref} - u|$ . Please provide both the table and the value of  $u_{ref}$  in the hard copy of the assignment to facilitate the grading process.

## Problem 2. Laplace's equation in an annular domain (30%)

Let  $\Omega$  be an two-dimensional annular domain of the inner radius  $r_1$  and the outer radius of  $r_2$ : i.e.,

$$\Omega = \{x \in \mathbb{R}^2 \mid r_1 < |x| < r_2\}.$$

We consider Laplace's equation

$$-\Delta u = 0 \quad \text{in } \Omega,$$

$$u(r=r_1, \theta) = 0, \quad \theta \in [0, 2\pi],$$

$$u(r=r_2, \theta) = g(\theta), \quad \theta \in [0, 2\pi].$$

In words, the homogeneous Dirichlet boundary condition is imposed on the inner boundary, and a non-homogeneous Dirichlet boundary condition is imposed on the outer boundary. Answer the following questions:

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nagpalaaESC384 PDE's Assignment 4  
SOLUTIONS(a) PDE:

$$\frac{\partial u}{\partial t} - \frac{\partial^2 u}{\partial x^2} = f(x, t)$$

(i) Semi-discrete:

$$\frac{d\hat{u}_i(t)}{dt} - \frac{\hat{u}_{i+1}(t) - 2\hat{u}_i(t) + \hat{u}_{i-1}(t)}{\Delta x^2} = \hat{f}_i(t)$$

Left Dirichlet BC: want  $\hat{u}_0(t) = u(x=0, t) = 0$ 

$$\frac{d\hat{u}_1(t)}{dt} - \frac{1}{\Delta x^2} (\hat{u}_{2+} - 2\hat{u}_1(t) + \hat{u}_{1-}) = \hat{f}_1(t)$$

semi-discrete equation for node  $i=1$ .

$$\Rightarrow \frac{d\hat{u}_1(t)}{dt} - \frac{1}{\Delta x^2} (\hat{u}_2(t) - 2\hat{u}_1(t)) = \hat{f}_1(t), \text{ this is the semi-discrete equation for node } i=1.$$

(ii) Right Robin BC:

$$\frac{d\hat{u}_n(t)}{dt} = \frac{\partial u}{\partial x}(x=1, t) + \mathcal{N}(x=1, t) = 0$$

allows us to apply finite difference formula at  $i=n$ .

$$\Rightarrow \frac{\hat{u}_{n+1} - \hat{u}_{n-1}}{2\Delta x} + \hat{u}_{n+1} = 0 \Rightarrow \hat{u}_{n+1} = -\hat{u}_{n-1}(2\Delta x) + \hat{u}_{n-1}$$

Substitute into semi-discrete equation at  $i=n$ 

$$\frac{d\hat{u}_n}{dt} - \frac{1}{\Delta x^2} (\hat{u}_{n+1} - 2\hat{u}_n + \hat{u}_{n-1}) = \hat{f}_n(t)$$

$$\frac{d\hat{u}_n}{dt} - \frac{1}{\Delta x^2} ((\hat{u}_{n-1} - \hat{u}_n(2\Delta x)) - 2\hat{u}_n + \hat{u}_{n-1}) = \hat{f}_n(t)$$

semi-discrete equation for node  $i=n$ .this is the semi-discrete equation for node  $i=n$ .(iii)now we can apply finite difference formula at any  $i \in [2, n-1]$ 

$$\Rightarrow \frac{\hat{u}_{i+1} - \hat{u}_{i-1}}{2\Delta x} + \hat{u}_i = 0 \Rightarrow \hat{u}_{i+1} = -\hat{u}_i(2\Delta x) + \hat{u}_{i-1}$$

Substitute into semi-discrete equation for  $i \in [2, n-1]$ 

and we end with

$$\frac{d\hat{u}_i}{dt} - \frac{1}{\Delta x^2} (\hat{u}_{i+1} - 2\hat{u}_i + \hat{u}_{i-1}) = \hat{f}_i(t) \rightarrow \text{equation for unknown nodes } i \in [2, n-1]$$

this is the semi-discrete equation for all other unknown nodes  $i \in [2, n-1]$

(iv) Lastly, we can express the equation for the initial condition as:  $\hat{f}_i(t) = g(x_i)$ , but we are given that  $g(x_i) = \hat{g}_i$

where  $\hat{g}_i$  is the value of  $\hat{u}_i(t=0) = \hat{u}_i(0)$

so we finally have  $\hat{f}(t) = \hat{g} = \hat{u}(0) \text{ if } t \in [0, n]$   $\rightarrow$  equation for I.C.

(b) semi-discrete equation can be expressed as

$$\frac{d\hat{u}}{dt}(t) + \hat{A}\hat{u}(t) = \hat{f}(t) \text{ in } \mathbb{R}^n,$$

$$\hat{u}(t=0) = \hat{g} \text{ in } \mathbb{R}^n$$

so matrix  $\hat{A}$  can be written as,

$$\hat{A} = \frac{1}{\Delta x^2} \begin{bmatrix} 2 & -1 & 0 & \dots & 0 & 0 & 0 \\ -1 & 2 & -1 & 0 & \dots & 0 & 0 \\ 0 & -1 & 2 & -1 & \dots & 0 & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots & \vdots & \vdots \\ 0 & 0 & 0 & \dots & 2 & -1 & 0 \\ 0 & 0 & 0 & \dots & -1 & 2 & -1 \\ 0 & 0 & 0 & \dots & 0 & -2 & 2+\frac{1}{\Delta x} \end{bmatrix}$$

dependent on left B.C

-1 2 1 sequence shifts to right by 1 each time.

dependent on Right B.C  
From the  $\frac{1}{\Delta x}(2+\frac{1}{\Delta x})$  term.

now,  $\hat{f}(t)$  and  $\hat{g}$  vectors can be expressed as

$$\hat{f}(t) = \begin{bmatrix} \hat{f}_1(t) \\ \hat{f}_2(t) \\ \vdots \\ \hat{f}_n(t) \end{bmatrix}, \quad \hat{g} = \begin{bmatrix} \hat{g}_1 \\ \hat{g}_2 \\ \vdots \\ \hat{g}_n \end{bmatrix}$$

$\rightarrow$  expressions for matrix  $\hat{A}$ , vectors  $\hat{f}(t), \hat{g}$

(c) Crank Nicolson approximation intime for Fully discrete eqn

$$\Rightarrow \hat{u}^j = \Delta \hat{u}^{j-1} + \hat{F}(t^j, t^{j-1}), \quad j=0, 1, 2, 3, \dots, J, \text{ where } \hat{u}(t^j) = \hat{u}^j$$

In the C-N method, we can express the semi-discrete equation as

$$\frac{\hat{u}^j - \hat{u}^{j-1}}{\Delta t} = \frac{1}{2} (\hat{F}(\hat{u}^j, t^j) + \hat{F}(\hat{u}^{j-1}, t^{j-1})) \text{ in } \mathbb{R}^n,$$

↳ avg of the derivatives

Now, we know that  
 $\hat{F}(\hat{u}, t) = \frac{d\hat{u}}{dt}(t) + \hat{f}(t)$ , and so we can use the  $\hat{A}$  matrix to express this as done below

$$\Rightarrow \frac{\hat{u}^j - \hat{u}^{j-1}}{\Delta t} = \frac{1}{2} (-\hat{A}\hat{u}^j + \hat{f}(t^j) - \hat{A}\hat{u}^{j-1} + \hat{f}(t^{j-1}))$$

Rearranging and using the identity matrix 'I' we can write this as

$$\Rightarrow \left( \frac{I}{\Delta t} + \frac{1}{2}\hat{A} \right) \hat{u}^j = \left( \frac{I}{\Delta t} - \frac{1}{2}\hat{A} \right) \hat{u}^{j-1} + \frac{1}{2} (\hat{f}(t^j) + \hat{f}(t^{j-1}))$$

Now, if we compare this to the given fully discrete

$$\left( \frac{I}{\Delta t} + \frac{1}{2}\hat{A} \right) \hat{u}^j = \left( \frac{I}{\Delta t} - \frac{1}{2}\hat{A} \right) \hat{u}^{j-1} + \frac{1}{2} (\hat{f}(t^j) + \hat{f}(t^{j-1})) \Rightarrow \hat{C}\hat{u}^j = \Delta \hat{u}^{j-1} + \hat{F}(t^j, t^{j-1})$$

we have the expressions for matrices  $C, D$  and vector  $F$  by comparison as

$$\Rightarrow \hat{C} = \left( \frac{I}{\Delta t} + \frac{1}{2}\hat{A} \right), \quad \hat{D} = \left( \frac{I}{\Delta t} - \frac{1}{2}\hat{A} \right),$$

$\text{and } \hat{F}(t^j, t^{j-1}) = \frac{1}{2} (\hat{f}(t^j) + \hat{f}(t^{j-1}))$

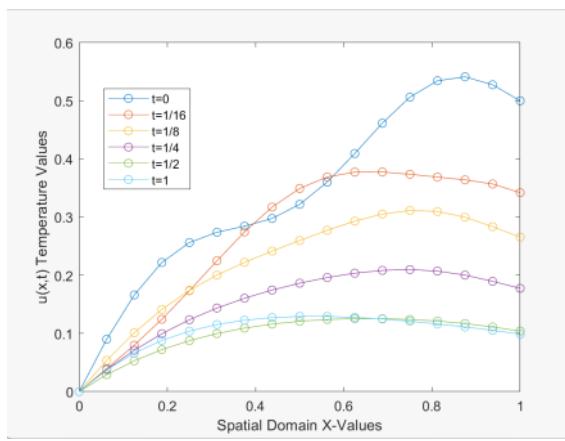
$\rightarrow$  expressions for  $C, D$  matrices and  $F$  vector

d) Code over here for heat\_fd\_temp



```
Question 1, part d heat_fd_temp code
function heat_fd_temp
% function to solve heat equation using finite difference method
gfun = @(x) 0.5*x.^2*x + 1/(6*pi)*sin(3*pi*x);
ffun = @(x,t) exp(-3*t)*exp(i*x);
% discretization parameters
n = 16;
J = 16;
% %Cover the time range (0,1) as given in the question
% setup spatial and temporal grids
xx = linspace(0,J+1);
tt = linspace(0,T+1);
%discretization
dt = J/2; %discretizing in the time domain over the given parameter
% assemble the A matrix
A = sparse(1:J,1:J,n);
A = sparse(J+1:J,J+1:n);
for i = 1:n
    if i == 1
        % left boundary
        A(i,i) = 2; %obtained from first element in first row of A matrix
        A(i,J+1) = -2; %obtained from second element in first row of A matrix
        else
            A(i,i) = n;
            A(i,J+1) = -n;
        end
    else
        % all other points
        % all these points are obtained from the 3 shifting elements in
        % each row of the A matrix except for the first and last rows which
        % represent the boundary conditions
        A(i,i) = 1;
        A(i,J+1) = -1;
        A(i,J+2) = -1;
        A(i,J+3) = 1;
    end
end
%In my expression for the A matrix I also have a 1/dx^2 term in front of
%dt, so I am adding that here for it to be consistent, once again using
%element wise division as A is a matrix
A=A*dt/2;
% iteration
for i = 1:J+1
    % assemble C and D matrices in terms of A, I, and dt
    %need to use element wise division and multiplication as they are vector
    C = (I*dt)/(12*A); %found in Q1c using C-N method
    D = (I*dt)/(12*A); %found in Q1c using C-N method
    % initial condition
    xx = xx(end); % excluding the left node
    U = gfun(xx);
    % record initial state in UU; note the first node is left as 0
    UU = zeros(n+1,1);
    UU(1) = U;
    for j = 1:J
        % evaluate t
        %P = heat_fd_temp(xx(j),tf(j)); %Using the expression we
        %found in Q1c but with indexing based on matlab formatting for the
        %timesteps in the function
        % solve linear system
        % expression involves U, C, D, and F. Do NOT use "inv".
        % Instead use "linsolve" or "mldivide" from Q1c, and also what was explained in
        %lecture and using \ instead of inv
        U = C\I\U + D\U(UU(end));
        % record solution in UU
        UU(2:end+1) = U;
    end
    % find time indices to plot
    tplot = [1/16, 1/8, 1/4, 1/2, 1];
    %P = heat_fd_temp(xx,tplot);
    % plot solution
    %P = heat_fd_temp(xx,tt);
    plot(xx,UU);
    xlabel("Spatial Domain X-Values");
    ylabel("Temperature Values");
    legend('t=0', 't=1/16', 't=1/8', 't=1/4', 't=1/2', 't=1', 'Position', [0.2 0.6 0.1 0.2]);
    % evaluate output
    s = UU(end,end);
end
```

e) solution representation for T=1, N=16, J=16 for various values of t using finite difference method



code over here for table

f) The observed error behaviour matches my expectation

as it follows the second order scheme discussed in lecture for both spatial and time domains. It can be noted that the error reduces by a factor of  $\approx 4$  for a doubling in the  $n, J$  values each time. This is observed and shown in the table below.

See table for error convergence values and confirmation.

s\_ref = 0.0997718155

n	J	s_ref-s
8	8	0.0047237
16	16	0.00078075
32	32	0.0001566
64	64	3.4994e-05

```
Question 1, part f heat_fd_solver and errorcalculator code
function s = heat_fd_solver(n,J)
% function to solve heat equation using finite difference method
gfun = @(x) 0.5*x.^2*x + 1/(6*pi)*sin(3*pi*x);
ffun = @(x,t) exp(-3*t)*exp(i*x);
% discretization parameters
n = 16;
J = 16;
% %Cover the time range (0,1) as given in the question
% setup spatial and temporal grids
xx = linspace(0,J+1);
tt = linspace(0,T+1);
%discretization
dt = J/2; %discretizing in the time domain over the given parameter
% assemble the A matrix
A = sparse(1:J,1:J,n);
A = sparse(J+1:J,J+1:n);
for i = 1:n
    if i == 1
        % left boundary
        A(i,i) = 2; %obtained from first element in first row of A matrix
        A(i,J+1) = -2; %obtained from second element in first row of A matrix
        else
            A(i,i) = n;
            A(i,J+1) = -n;
        end
    else
        % all other points
        % all these points are obtained from the 3 shifting elements in
        % each row of the A matrix except for the first and last rows which
        % represent the boundary conditions
        A(i,i) = 1;
        A(i,J+1) = -1;
        A(i,J+2) = -1;
        A(i,J+3) = 1;
    end
end
%In my expression for the A matrix I also have a 1/dx^2 term in front of
%dt, so I am adding that here for it to be consistent, once again using
%element wise division as A is a matrix
A=A*dt/2;
% iteration
for i = 1:J+1
    % assemble C and D matrices in terms of A, I, and dt
    %need to use element wise division and multiplication as they are vector
    C = (I*dt)/(12*A); %found in Q1c using C-N method
    D = (I*dt)/(12*A); %found in Q1c using C-N method
    % initial condition
    xx = xx(end); % excluding the left node
    U = gfun(xx);
    % record initial state in UU; note the first node is left as 0
    UU = zeros(n+1,1);
    UU(1) = U;
    for j = 1:J
        % evaluate t
        %P = heat_fd_temp(xx(j),tf(j)); %Using the expression we
        %found in Q1c but with indexing based on matlab formatting for the
        %timesteps in the function
        % solve linear system
        % expression involves U, C, D, and F. Do NOT use "inv".
        % Instead use "linsolve" or "mldivide" from Q1c, and also what was explained in
        %lecture and using \ instead of inv
        U = C\I\U + D\U(UU(end));
        % record solution in UU
        UU(2:end+1) = U;
    end
    % find time indices to plot
    tplot = [1/16, 1/8, 1/4, 1/2, 1];
    %P = heat_fd_temp(xx,tplot);
    % plot solution
    %P = heat_fd_temp(xx,tt);
    plot(xx,UU);
    xlabel("Spatial Domain X-Values");
    ylabel("Temperature Values");
    legend('t=0', 't=1/16', 't=1/8', 't=1/4', 't=1/2', 't=1', 'Position', [0.2 0.6 0.1 0.2]);
    % evaluate output
    s = UU(end,end);
end
```

```
%Now writing the code needed to get the error for the different levels of
%discretization mentioned in Q1f
function errorcalculator
%function to calculate the n, J values and reference levels of discretization
nJ_values = [8,16,32,64];
reference = [512,512];
ref_n = reference(1);
ref_J = reference(2);
s_ref = heat_fd_solver(ref_n, ref_J);
%creating an error vector with zeros that holds the error values at each level of
%discretization
error_values = zeros(length(nJ_values), size(nJ_values,1));
%obtaining the errors for each solution based on the parameters and storing
%them in an error vector we created earlier
for i = 1:length(nJ_values)
    s = heat_fd_solver(nJ_values(i), nJ_values(i));
    error_values(i)=abs(s-s_ref);
end
%Displaying the errors in the form of a table to show the errors at each
%level of discretization
T=table(nJ_values,nJ_values,error_values, VariableNames, {'n','J','s_ref-s'});
disp(T)
end
```

(Q2a) Given Laplace's Equation

$$\begin{aligned} -\Delta u &= 0 \text{ in } \Omega \\ u(r=r_1, \theta) &= 0, \theta \in [0, 2\pi] \\ u(r=r_2, \theta) &= g(\theta), \theta \in [0, 2\pi] \end{aligned}$$

Laplacian in Polar coordinates

$$\begin{aligned} \Delta u &= \nabla \cdot \nabla u = \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial u}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 u}{\partial \theta^2} \\ &= \frac{\partial^2 u}{\partial r^2} + \frac{1}{r} \frac{\partial u}{\partial r} + \frac{1}{r^2} \frac{\partial^2 u}{\partial \theta^2} \end{aligned}$$

We can use

Separation of variables

$$u_n(r, \theta) = R_n(r) \Theta_n(\theta) \rightarrow \text{Family of sols}$$

$$R_n'' \Theta_n + \frac{1}{r} R_n' \Theta_n + \frac{1}{r^2} R_n \Theta_n'' = 0$$

$$\text{Multiply by } r^2 \frac{1}{R_n \Theta_n}$$

$$\Rightarrow r^2 R_n'' + r R_n' + \frac{\Theta_n''}{\Theta_n} = 0$$

$$\underbrace{r^2 R_n''}_{\substack{\text{independent of} \\ \theta}} + \underbrace{r R_n'}_{\substack{\text{independent} \\ \text{of } r}} = - \underbrace{\frac{\Theta_n''}{\Theta_n}}_{\substack{\text{constant}}} = \alpha_n$$

Now, we can solve for

Eigenproblem for  $\Theta_n$

$$-\Theta_n'' = \alpha_n \Theta_n \text{ in } (0, 2\pi)$$

$$\Theta_n(0) = \Theta_n(2\pi) \rightarrow \text{periodic B.C. since we want continuous solutions.}$$

General soln is of the form,

$$\Rightarrow \Theta_n(\theta) = A_n \cos(n\theta) + B_n \sin(n\theta), n = 0, 1, 2, \dots$$

$$\alpha_n = n^2$$

Now,

$$\text{ODE for } R_n: r^2 R_n'' + r R_n' = n^2 R_n$$

$$r^2 R_n'' + r R_n' - n^2 R_n = 0 \quad (\text{Cauchy-Euler equation of order 2})$$

We can thus use the general soln for Cauchy-Euler equation of 2nd order

$$R_n(r) = \begin{cases} C_n + D_n \log(r), & n=0 \\ C_n r^n + D_n r^{-n}, & n=1, 2, 3, \dots \end{cases}$$

Now, we impose the boundary conditions

B.C. at  $r=r_1$ , and  $n=0$ , we have

$$R_n(r_1) = C_0 + D_0 \log(r_1) \stackrel{\text{want}}{=} 0,$$

$$\Rightarrow C_0 = -D_0 \log(r_1)$$

Plugging back in the general soln, we have

$$R_n(r) = -D_0 \log(r_1) + D_0 \log(r)$$

$$R_n(r) = D_0 \log(r/r_1)$$

B.C. at  $r=r_2$ , and  $n=1, 2, \dots$ , we have

$$R_n(r_2) = C_n r_2^n + D_n r_2^{-n} \stackrel{\text{want}}{=} 0$$

$$\Rightarrow C_n r_2^n = -D_n r_2^{-n}$$

$$\Rightarrow C_n = \frac{-D_n r_2^{-n}}{r_2^n} = -\frac{D_n}{r_2^{2n}}$$

Plugging back in the general soln, we have

$$R_n(r) = -\frac{D_n}{r_2^{2n}} r_2^n + D_n r^{-n}$$

$$R_n(r) = D_n \left( -\frac{r^n}{r_2^{2n}} + r^{-n} \right) = D_n (-r_2^{-2n} \cdot r^n + r^{-n})$$

Now using superposition and the general soln expression,  
we have that  $u_n(r, \theta) = R_n(r) \Theta_n(\theta)$

$$\Rightarrow u(r, \theta) = A_0 R_0(r) + \sum_{n=1}^{\infty} R_n(r) (A_n \cos(n\theta) + B_n \sin(n\theta))$$

we have that  $u_n(r, \theta) = R_n(r) \Theta_n(\theta)$

$$\Rightarrow u_n(r, \theta) = A_0 R_0(r) + \sum_{n=1}^{\infty} R_n(r)(A_n \cos(n\theta) + B_n \sin(n\theta))$$

$$u_n(r, \theta) = \begin{cases} A_0 \log(r/r_1), & n=0 \\ B_n (r_1^{-2n} r^n - r^{-n}) (\cos(n\theta) + \sin(n\theta)), & n=1, 2, 3 \dots \end{cases}$$

satisfies PDE,  
inner B.C.

family of functions satisfying  
PDE, inner B.C.

(b) Now, to find the solution for the PDE with the other B.C.

$$u(r=r_2, \theta) = g(\theta), \text{ where } g(\theta) = \cos(m\theta), m = \text{rc integer}$$

we can use the general solution expression found in part(a) and impose the B.C at  $r=r_2$  on it.

$$\text{i.e., } u(r=r_2, \theta) = A_0 \log(r_2/r_1) + \sum_{n=1}^{\infty} (r_1^{-2n} r_2^n - r_2^{-n}) (A_n \cos(n\theta) + B_n \sin(n\theta)) = \cos(m\theta)$$

The above expression in the RHS matches a Fourier series form, so deducing the coefficients for it.

$$\text{we know } g_0 = 2 \cdot \frac{1}{L} \int_{-L}^L g(\theta) d\theta$$

$\Rightarrow$  since the period in the PDE is  $[0, 2\pi]$ , we can use  $[-\pi, \pi]$ , instead as  $g(\theta) = \cos(m\theta)$ ,  $m > 0$  is even-periodic and the area under  $[-\pi, \pi]$  is the same as the area under  $[\pi, \pi]$  and the integrals remain unchanged otherwise.

$$\therefore \text{we have, } g_0 = 2 \cdot \frac{1}{\pi} \int_{-\pi}^{\pi} \cos(m\theta) d\theta \Rightarrow \frac{2}{\pi} [\sin(m\pi) - \sin(-m\pi)] = 0$$

comparing this to  $A_0 \log(r_1/r_2)$

$$\text{we have } A_0 = 0.$$

now, we take  $g_{n, \cos}$  and  $g_{n, \sin}$

$$g_{n, \cos} = \frac{1}{L} \int_{-L}^L \cos(n\theta) g(\theta) d\theta \Rightarrow \frac{1}{\pi} \int_{-\pi}^{\pi} \cos(n\theta) \cos(m\theta) d\theta = \frac{2}{\pi} \int_0^{\pi} \cos(n\theta) \cos(m\theta) d\theta$$

using the integral list we have,

$$\int \cos(\alpha x) \cos(\beta x) dx = \frac{\sin((\alpha-\beta)x)}{2(\alpha-\beta)} + \frac{\sin((\alpha+\beta)x)}{2(\alpha+\beta)}, \alpha \neq \beta$$

so we get

$$g_{n, \cos} = \frac{2}{\pi} \left[ \frac{\sin((n-m)\theta)}{2(\theta-m)} + \frac{\sin((n+m)\theta)}{2(n+m)} \right] \Big|_0^\pi$$

$$g_{n, \cos} \Rightarrow \begin{cases} 1, & n=m \\ 0, & n \neq m \end{cases}$$

now, we take  $g_{n, \sin}$

$$g_{n, \sin} = \frac{1}{L} \int_{-L}^L \sin(n\theta) g(\theta) d\theta \Rightarrow \frac{1}{\pi} \int_{-\pi}^{\pi} \sin(n\theta) \cos(m\theta) d\theta = \frac{2}{\pi} \int_0^{\pi} \sin(n\theta) \cos(m\theta) d\theta$$

using the integral list we have,

$$\int \sin(\alpha x) \cos(\beta x) dx = \frac{\cos((\alpha-\beta)x)}{2(\alpha-\beta)} + \frac{\cos((\alpha+\beta)x)}{2(\alpha+\beta)}, \alpha \neq \beta$$

so we get

$$g_{n, \sin} = \frac{2}{\pi} \left[ \frac{\cos((n-m)\theta)}{2(\theta-m)} + \frac{\cos((n+m)\theta)}{2(n+m)} \right] \Big|_0^\pi$$

$$\Rightarrow g_{n, \sin} = 0.$$

Now, if we compare coefficients with the expression, we have

$$g_{n, \cos} = \begin{cases} 0, & n \neq m \\ 1, & n=m \end{cases} = \left( -r_1^{-2n} \cdot r_2^n - r_2^{-n} \right) A_n \cos(n\theta)$$

$$\Rightarrow A_n \cos(n\theta) = \frac{1}{(-r_1^{-2n} \cdot r_2^n - r_2^{-n})} \quad \text{--- (1)}$$

$$\Rightarrow A_n \cos(n\theta) = \frac{1}{(-r_1^{-2n} r_2^n - r_2^{-n})} \quad \text{--- (1)}$$

$$B_n \sin(n\theta) = 0 \Rightarrow B_n \sin(n\theta) = 0$$

∴ plugging back into the original expression, we get

$$u(r, \theta) = \sum_{n=1}^{\infty} (r_1^{-2n} r_2^n - r_2^{-n}) \left( A_n \cos(n\theta) + B_n \sin(n\theta) \right) \xrightarrow{\text{plug in (1)}} g(\theta)$$

changing all 'n' subscripts to 'm' to match  $\cos(m\theta)$ , and shifting terms around we get

$$u(r, \theta) = \frac{\cos(m\theta) (r_1^{-2m} r^m - r_2^{-m})}{(r_1^{-2m} r_2^m - r_2^{-m})} \rightarrow \text{solution } u(r, \theta) \text{ for } g(\theta) = \cos(m\theta)$$

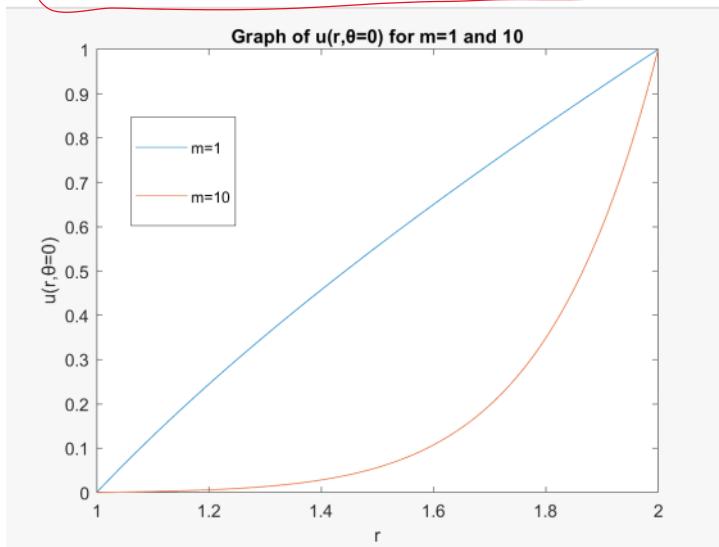
(Q3)(i) Yes, the graph obtained is consistent with my expectations as the solution with the higher mode ( $m=10$ ) decays a lot more rapidly from the outer boundary as it moves to the inner one as can be seen in the graph.

(ii) Yes, the maximum principle is also satisfied because the temperature value as seen in the graph between the two boundaries  $[r_1, r_2]$  does not exceed its max value as seen on the  $r_2$  boundary.

The graph for Laplace's solution at  $m=1, 10$  between the annular domain boundaries.

Code for (2c)

```
Question 2c, plot for m=1 and m=10 using the usolver function
%creating boundary values r_1, r_2 and also creating a vector with equally
%spaced points between them
r_1 = 1;
r_2 = 2;
r = linspace(r_1, 1, r_2);
%Creating a vector m for the two m values for which the solution will be
%graphed
m=[1,10];
%Plotting the solutions obtained by running the solver on the same plot
plot(r, usolver(r,m(1)), r, 1, 2);
hold on
plot(r, usolver(r,m(2)), r, 1, 2);
hold off
%Labelling
title('Graph of u(r,\theta=0) for m=1 and 10');
ylabel('u(r,\theta=0)');
title('Graph of u(r,\theta=0) for m=1 and 10');
legend('m=1', 'm=10', 'Position', [0.2 0.6 0.1 0.2]);
function u = usolver(r,m,r_1,r_2)
%Value is 0 in the question
th = 0;
% Using the value of the solution we obtained in Q2b of the assignment
u = cos(m*pi*(r_1-1)*(2*m)*r/m-(r_1-1)*(2*m)*r_2/m);
end
```



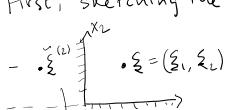
### (Q3a) Laplace's Equation

$$\begin{aligned} -\Delta u &= 0 \text{ on } \Omega \equiv R_{>0} \times R_{>0} \\ u &= g_B \text{ on } \Gamma_B \equiv R_{>0} \times \mathbb{R}_2 = 0 \\ u &= g_L \text{ on } \Gamma_L \equiv \mathbb{R}_1 \times R_{>0} \end{aligned}$$

To find Green's Function given  $\xi \in \Omega$ ,

$$\begin{aligned} -\Delta G(x, \xi) &= \delta(x - \xi) \forall x \in \Omega \\ G(x, \xi) &= 0 \forall x \in \Gamma_B \cup \Gamma_L \end{aligned}$$

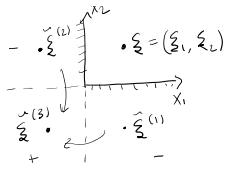
First, sketching the problem, using image points



Now we can express  $\hat{\xi}^{(1)}, \hat{\xi}^{(2)}, \hat{\xi}^{(3)}$  in terms of  $\xi_1, \xi_2$

$$\hat{\xi}^{(1)} = (\xi_1, -\xi_2), \hat{\xi}^{(2)} = (-\xi_1, \xi_2), \hat{\xi}^{(3)} = (-\xi_1, -\xi_2)$$

... ... ... Linear Ftn. fundamental solution. we have



Now we can express  $\xi_1, \xi_2$  in terms of  $x_1, x_2$   
 $\Rightarrow \xi^{(1)} = (\xi_1, -\xi_2), \xi^{(2)} = (-\xi_1, \xi_2), \xi^{(3)} = (-\xi_1, -\xi_2)$

writing out the solution in terms of the fundamental solution, we have

$$g(x-\xi) = \Phi(x-\xi) + \psi(x, \xi) \quad \text{where } \psi(x, \xi) = -\Phi(x-\xi) + x \cdot \xi$$

$$\therefore \text{here } \psi(x, \xi) = -\Phi(x-\xi^{(1)}) - \Phi(x-\xi^{(2)}) + \Phi(x-\xi^{(3)}),$$

and so  $\boxed{g(x-\xi) = \Phi(x-\xi) - \Phi(x-\xi^{(1)}) - \Phi(x-\xi^{(2)}) + \Phi(x-\xi^{(3)})}$  → green's function representation for the laplace equation.  
 where  $\|x-\xi\| = \|x-\xi^{(1)}\|$  on  $\Gamma_B$ ,  $\|x-\xi\| = \|x-\xi^{(2)}\|$  on  $\Gamma_L$ , and  
 $\|x-\xi\| = \|x-\xi^{(3)}\|$  on  $\Gamma_U$ .

(b)  $g_L = 0$ , we need to find integral expression of solution  $u(x=(x_1, x_2))$

we start by using the Green's representation formula

$$u(x) = - \int_{\Omega} \nabla \cdot \nabla_{\xi} G(x, \xi) g_L(\xi) d\xi$$

For  $u(x_1, x_2)$  in our case this can be expanded as follows

$$u(x_1, x_2) = - \int_{\Gamma_B} V_B \cdot \nabla_{\xi} G(x, \xi) g_B(\xi) d\xi - \int_{\Gamma_L} V_L \cdot \nabla_{\xi} G(x, \xi) g_L(\xi) d\xi$$

$\underbrace{\qquad\qquad\qquad}_{=0 \text{ here as given}}$

thus the entire integral vanishes.

⇒ Now plugging in the  $G(x, \xi)$  function we found in (a).

$$= - \int_{\Gamma_B} V_B \cdot \nabla_{\xi} (\Phi(x-\xi) - \Phi(x-\xi^{(1)}) - \Phi(x-\xi^{(2)}) + \Phi(x-\xi^{(3)})) g_B(\xi) d\xi \quad \text{--- (1)}$$

$$\text{Now, using the hint: we can compute } -\frac{\partial \Phi(x-\xi)}{\partial \xi_2} = \frac{\partial}{\partial \xi_2} \left( \frac{1}{2\pi} \log \left( \sqrt{(x_1-\xi_1)^2 + (x_2-\xi_2)^2} \right) \right)$$

Plugging the expressions in, we get

$$\begin{aligned} &= \frac{\partial}{\partial \xi_2} \left( \frac{1}{2\pi} \ln \left( (x_1-\xi_1)^2 + (x_2-\xi_2)^2 \right)^{1/2} \right) - \frac{\partial}{\partial \xi_2} \left( \frac{1}{2\pi} \ln \left( (x_1-\xi_1)^2 + (x_2+\xi_2)^2 \right)^{1/2} \right) \\ &- \frac{\partial}{\partial \xi_2} \left( \frac{1}{2\pi} \ln \left( (x_1+\xi_1)^2 + (x_2-\xi_2)^2 \right)^{1/2} \right) - \frac{\partial}{\partial \xi_2} \left( \frac{1}{2\pi} \ln \left( (x_1+\xi_1)^2 + (x_2+\xi_2)^2 \right)^{1/2} \right) \end{aligned}$$

plugging this expression into (1) after solving the partial derivatives, we get

$$\Rightarrow - \int_{\Gamma_B} \frac{1}{2\pi} \left( \frac{\xi_2-x_2}{(x_1-\xi_1)^2+(x_2-\xi_2)^2} - \frac{\xi_2+x_2}{(x_1-\xi_1)^2+(x_2+\xi_2)^2} - \frac{\xi_2-x_2}{(x_1+\xi_1)^2+(x_2-\xi_2)^2} + \frac{\xi_2+x_2}{(x_1+\xi_1)^2+(x_2+\xi_2)^2} \right) g_B(\xi) d\xi$$

we also know that on  $\Gamma_B$  the  $\xi_2$  term will be 0, so

$$\text{we get } \Rightarrow - \int_{\Gamma_B} \frac{1}{2\pi} \left( \frac{-x_2}{(x_1-\xi_1)^2+x_2^2} - \frac{x_2}{(x_1-\xi_1)^2+x_2^2} + \frac{x_2}{(x_1+\xi_1)^2+x_2^2} + \frac{x_2}{(x_1+\xi_1)^2+x_2^2} \right) g_B(\xi) d\xi$$

$$\Rightarrow \int_{\Gamma_B} \frac{1}{2\pi} \left( \frac{2x_2}{(x_1-\xi_1)^2+x_2^2} - \frac{2x_2}{(x_1+\xi_1)^2+x_2^2} \right) g_B(\xi) d\xi$$

but  $\Gamma_B$  goes from 0 to  $\infty$ , so replacing limits and moving  $\frac{1}{\pi}$  out  
 after cancelling the factor of 2

$$u(x_1, x_2) = \frac{1}{\pi} \int_0^\infty \left( \frac{x_2}{(x_1-\xi_1)^2+x_2^2} - \frac{x_2}{(x_1+\xi_1)^2+x_2^2} \right) g(\xi) d\xi$$

→ integral representation of the solution with  $g_L = 0$ .

(c)  $g_L = 0, g_B(x_1) = \begin{cases} 1, & x_1 \in (0, 1) \\ 0, & \text{o.w.} \end{cases}$

To find: integral representation of solution over finite limit  
 we can use the expression we found in (b) to find the solution,

so all we need to do is plug in over the new limit 0 to 1, and  
 evaluate  $g(\xi)$

We know  $g(\xi) = 1$  over the given limit of 0 to 1 based on the function  
 provided above,

so we have

$$u(x_1, x_2) = \frac{1}{\pi} \int_0^1 \left( \frac{x_2}{(x - \xi_1)^2 + x_2^2} - \frac{x_2}{(x_1 + \xi_1)^2 + x_2^2} \right) d\xi_1$$

↪ integral representation  
of solution given the  
 $gB(x)$  condition.