

Machine Learning for String Compactifications

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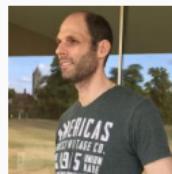
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Overview

Calabi–Yau metrics and hermitian Yang–Mills connections are crucial for string phenomenology

Numerical methods are the only way to access this data

Machine learning and neural networks provide a powerful set of tools to tackle geometric problems

Outline

Physics from geometry

Calabi–Yau metrics

Hermitian Yang–Mills connections

Machine learning and neural networks

Application: moduli dependence and the swampland

Physics from geometry

Motivation from physics

Does string theory describe our universe? Many semi-realistic MSSM-like string models from M-theory / F-theory / heterotic [...; Cole et al. '21; Abel et al. '21; Loges, Shiu '21, '22]

- Focus on models from **heterotic string** on Calabi–Yau

Coarse details: correct gauge group, matter spectrum, etc.

- **Topological** – do not need details of geometry

How many of these string vacua are physically reasonable?

- Predicted masses and couplings depend intricately on underlying geometry, i.e. **metric** and **gauge connection**
- No **explicitly** known (non-trivial) Calabi–Yau metrics or appropriate connections!

Calabi–Yau compactifications

Minimal supersymmetry on $\mathbb{R}^{1,3} \times X$ with $E_8 \times E_8$ bundle V [Candelas et al. '85]

- No H flux $\Rightarrow X$ equipped with **Calabi–Yau** metric g
- V admits **hermitian Yang–Mills** connection A
- Anomaly condition: $p_1(X) = p_1(V)$

Particle spectrum of low-energy theory determined by X and V

- e.g. standard embedding: $SU(3)$ bundle gives E_6 GUT gauge group in 4d with $\frac{1}{2}\chi(X)$ particle generations
- Most interesting MSSM examples from **non-standard embedding**, but not so simple... [...;Donagi et al. '98; Braun et al. '05; Anderson et al. '11;...]

Low-energy physics

Compactification on X leads to 4d $N = 1$ effective theory with gauge + chiral multiplets.

- Chiral multiplets split into moduli fields and **matter fields**

Particle content comes from topology of X and V , e.g.

- $SU(3)$ bundle V gives E_6 GUT group in 4d

$$E_8 \rightarrow E_6 \times SU(3)$$

$$\underline{248} \rightarrow \bigoplus_{\underline{r}, \underline{R}} (\underline{r}, \underline{R}) = (\underline{78}, \underline{1}) \oplus (\underline{1}, \underline{8}) \oplus (\underline{27}, \underline{3}) \oplus (\overline{\underline{27}}, \overline{\underline{3}})$$

- 4d multiplets transforming in \underline{r} come from $H^1(X, \underline{R})$, e.g. **matter fields** from $C^l \in H^1(X, \underline{3})$

Yukawa couplings

Yukawa terms in Standard Model include $\mathcal{L}_{\text{SM}} \supset \mathcal{L}_{\text{Yuk}} = Y_{ij}^d H Q^i d^j + \dots$

4d $N=1$ theory governed by **superpotential** and **Kähler potential**

$$W = \lambda_{IJK}(\phi) C^I C^J C^K + \dots \quad K = G_{IJ}(\phi, \bar{\phi}) C^I \bar{C}^J + \dots$$

- Holomorphic superpotential from triple overlap of zero modes on X

$$\lambda_{IJK} = \int_X \Omega \wedge \text{tr}(C^I \wedge C^J \wedge C^K)$$

- Matter field Kähler potential gives **normalisation** where C^I are **harmonic**

$$G_{IJ} = \int_X C^I \wedge \star_V \bar{C}^J$$

A string model wish list

MSSM spectrum, three families, etc. ✓

- Reduces to topology / algebraic methods

Superpotential couplings λ_{IJK} ✓

- Holomorphic – can use algebraic / differential methods

Harmonic modes and Kähler metric G_{IJ} on field space ✗

- Numerical methods

Supersymmetry breaking, moduli stabilisation, etc. ✗

- Soft masses and couplings c.f. $N = 1$ Kähler potential and normalised zero modes [Kaplunovsky, Louis '93; Blumenhagen et al. '09; ...]

The missing ingredients

How do we calculate Calabi–Yau metrics or hermitian
Yang–Mills connections?

Calabi–Yau metrics

Calabi–Yau geometry

Calabi–Yau manifolds are Kähler and admit **Ricci-flat** metrics

- **Existence** but no explicit constructions
- Kähler + $c_1(X) = 0 \Rightarrow$ there exists a Ricci-flat metric [Yau '77]

Kähler \Rightarrow **Kähler potential** K gives metric g and closed two-form $J = \partial\bar{\partial}K$

$$\text{vol}_g \equiv J \wedge J \wedge J$$

$c_1(X) = 0 \Rightarrow$ nowhere-vanishing holomorphic (3,0)-form Ω

$$\text{vol}_\Omega \equiv i\Omega \wedge \bar{\Omega}$$

Example: Fermat quintic

Calabi–Yau threefold is quintic hypersurface Q in \mathbb{P}^4

$$Q(z) \equiv z_0^5 + z_1^5 + z_2^5 + z_3^5 + z_4^5 = 0$$

(3,0)-form Ω determined by Q , e.g. in $z_0 = 1$ patch

$$\Omega = \frac{dz_2 \wedge dz_3 \wedge dz_4}{\partial Q / \partial z_1}$$

Metric g and Kähler form J determined by Kähler potential

$$g_{\bar{i}\bar{j}}(z, \bar{z}) = \partial_i \bar{\partial}_j K(z, \bar{z}).$$

How do we measure accuracy?

The Ricci-flat metric is given by a K that satisfies (c.f. Monge–Ampère)

$$\frac{\text{vol}_g}{\text{vol}_\Omega} \Big|_p = 1 \quad \Rightarrow \quad R_{\bar{i}\bar{j}} = 0$$

Define a **functional** of K [Douglas et al. '06]

$$\sigma(K) = \int_X \left| 1 - \frac{\text{vol}_g}{\text{vol}_\Omega} \right| \text{vol}_\Omega$$

The exact CY metric has $\sigma(K) = 0$

- Finding the Ricci-flat metric reduces to finding $K(z, \bar{z})$ that **minimises** σ

Algebraic metrics [Tian '90; Donaldson '05]

Natural Kähler metric on \mathbb{P}^4 given by

$$K_{FS} = \log \sum_{i=0}^4 z_i \bar{z}_i$$

Can generalise this with a **hermitian matrix** $h^{i\bar{j}}$

$$K(h) = \log \sum_{i,\bar{j}=0}^4 z_i h^{i\bar{j}} \bar{z}_{\bar{j}}$$

Restricting to $Q \subset \mathbb{P}^4$ gives Kähler metric but **not** Ricci-flat

- **25 real parameters** in $h^{i\bar{j}}$ that we can vary

Algebraic metrics [Tian '90; Donaldson '05]

Replace coordinates z_i with homogeneous monomials s_α of degree k

e.g. $k = 2 : s_\alpha = (z_0^2, z_0z_1, z_0z_2, \dots)$

Kähler potential is then

$$K(h) = \log \sum_{\alpha, \beta=1}^{15} s_\alpha h^{\alpha\bar{\beta}} \bar{s}_{\bar{\beta}}, \quad h^{\alpha\bar{\beta}} \sim 225 \text{ parameters}$$

At degree k have $\mathcal{O}(k^8)$ parameters, so can approximate the Ricci-flat metric to arbitrary precision

- Spectral method: $\frac{s_\alpha \bar{s}_{\bar{\beta}}}{(z_i \bar{z}_i)^k}$ gives basis for eigenspaces on $\mathbb{P}^4 = S^7 / U(1)$

How to fix K ?

Finding the “best” approximation to the Ricci-flat metric amounts to finding K so that σ is minimised

Three approaches:

- Iterative procedure for $h^{\alpha\bar{\beta}}$ [Donaldson ‘05; Douglas ‘06; Braun ‘07]
- Minimise σ directly as function of $h^{\alpha\bar{\beta}}$ [Headrick, Nassar ‘09; Anderson et al. ‘20]
- Find K or $g_{\bar{i}\bar{j}}$ **directly** by treating σ as a loss function for a neural network [Headrick, Wiseman ‘05; Douglas et al. 20; Anderson et al. ‘20; Jejjala et al. ‘20; Larfors et al. ‘21, ‘22]

In all cases, numerical integrals carried out by **Monte Carlo** [Shiffman, Zelditch ‘98]

Hermitian Yang–Mills connections

Hermitian Yang–Mills

A hermitian metric G on fibers of vector bundle V defines a connection and curvature

$$A_i = G^{-1} \partial_i G, \quad A_{\bar{i}} = 0 \quad \Rightarrow \quad F_{ij} = F_{\bar{i}\bar{j}} = 0, \quad F_{\bar{i}\bar{j}} = \partial_{\bar{j}}(G^{-1} \partial_i G)$$

We say A is **hermitian Yang–Mills** if

$$g^{\bar{i}\bar{j}} F_{\bar{i}\bar{j}} = \mu(V) \text{Id}$$

G is then known as a **Hermite–Einstein metric** on V

- Nonlinear PDE for G with no explicit solutions when X is Calabi–Yau
- HYM implies Yang–Mills: $d \star F = 0$
- **Supersymmetry** in 10d requires HYM with $\mu(V) = 0$

Existence and stability

Existence of HYM solutions [Donaldson '85; Uhlenbeck, Yau '86]

A holomorphic vector bundle V over a compact Kähler manifold (X, g) admits a Hermite–Einstein metric iff V is slope polystable

Slope of V

$$\mu(V) \equiv \int_X c_1(V) \wedge J^{n-1}$$

V is **stable** if $\mu(\mathcal{F}) < \mu(V)$ for all $\mathcal{F} \subset V$ (or **polystable** if sum of stable bundles with same slope)

- **Algebraic** condition (like $c_1(X) = 0$), but not constructive!

How do we measure accuracy?

Defining $F_g \equiv g^{\bar{i}\bar{j}} F_{\bar{i}\bar{j}}$, the HYM equation is $F_g = \mu(V) \text{Id}$

The **average** over the the Calabi–Yau is defined using the exact CY measure vol_Ω , e.g.

$$\langle \text{tr } F_g \rangle \equiv \int_X \text{vol}_\Omega \text{tr } F_g$$

Suitable choice of **accuracy measure** is

$$E[F, g] = \langle \text{tr } F_g^2 \rangle - \frac{1}{\text{rank } V} \langle \text{tr } F_g \rangle^2$$

$E[F, g]$ is positive semi-definite and vanishes on **HYM solutions**

$$F \text{ solves HYM} \iff E[F, g] = 0$$

Numerical approximations

Ansatz for **hermitian metric** [Wang '05; Douglas et al. '06; Anderson et al. '10]

$$(G^{-1})^{ab} = \sum_{\alpha, \beta}^{N_k} S_\alpha^a H^{\alpha\bar{\beta}} \bar{S}_\beta^b$$

where $S_\alpha^a \in H^0(X, V \otimes L^k)$ and $H^{\alpha\bar{\beta}}$ is a hermitian matrix of parameters

Varying $H^{\alpha\bar{\beta}}$, one can find an **approximate** HYM connection for $V \otimes L^k$ and then on V itself

Increasing k increases the number of sections N_k and hence the number of parameters in $H^{\alpha\bar{\beta}}$

The goal

Train a **neural network** to find solutions to the hermitian
Yang–Mills equation

Machine learning and neural networks

Overview

New era of **big data** in string theory

- Vacuum selection problem, huge number of CYs, larger number of flux vacua [Denef, Douglas '04; Taylor, Wang '15;...]

Many different types of machine learning

- **Supervised** – known inputs and outputs, e.g. recognise images, predict Hodge numbers [He '17; Bull et al. '18; Erbin, Finotello '20;...]
- **Unsupervised** – known inputs, e.g. looking for patterns or generate images
- **Self-supervised** – known inputs, output minimises a loss function, e.g. QM ground states, Ricci-flat metrics, **HYM connections**

Neural networks

Neural networks (NN) convert inputs to outputs: $\vec{x} \mapsto f(\vec{x}, \vec{w})$

- Network built from connected nodes called **neurons**
- **Weights** \vec{w} are parameters in network (strength of connections)
- Non-linear **activation functions**
- Training attempts to minimise a **loss function** computed from NN
- **Universal approximation theorem** for NNs [Cybenko '89]

More interested in the network itself than the actual values!

- NN gives a **variational ansatz** for some function you want to find, e.g. Hermite–Einstein metric G that solves HYM equation

Line bundles on CY manifolds

Line bundles crucial in many string models [Anderson, Gray, Lukas, Palti '11;...]

Holomorphic line bundle L determined by $c_1(L)$. Given a basis of divisors \mathcal{D}_I on X , denote by $\mathcal{O}_X(m^I)$ the line bundle with $c_1(L) = m^I \mathcal{D}_I$

Line bundles are **automatically** stable, so always admit a solution to HYM,
 $g^{i\bar{j}} F_{i\bar{j}} = \mu(L)$

We need the **functional** form of G to calculate harmonic representatives and the matter field Kähler metric

Bihomogenous networks [Douglas et al. '20]



$$\mathbb{C}^3 \rightarrow \mathbb{R}^9$$

$$z_i \mapsto (\operatorname{re} z_j \bar{z}_k, \operatorname{im} z_j \bar{z}_k)$$

$$\mathbb{R}^9 \rightarrow \mathbb{R}^{12}$$

$$\vec{x} \mapsto (W_1 \vec{x})^2$$

$$\mathbb{R}^{12} \rightarrow \mathbb{R}$$

$$\vec{y} \mapsto \log(W_2 \vec{y})$$

Parameters in W_1 and W_2 are **weights**, collectively denoted by \vec{w}

First implemented for CY metrics in TensorFlow by [Douglas et al. '20] at

<https://github.com/yidiq7/MLGeometry>

A loss function

Network output is treated as $\log G^{-1}$, which defines F

- Together with approximate CY metric g , this gives $F_g[\vec{w}]$ as a function of the network **weights** \vec{w}
- Mimics previous ansatz for hermitian structure

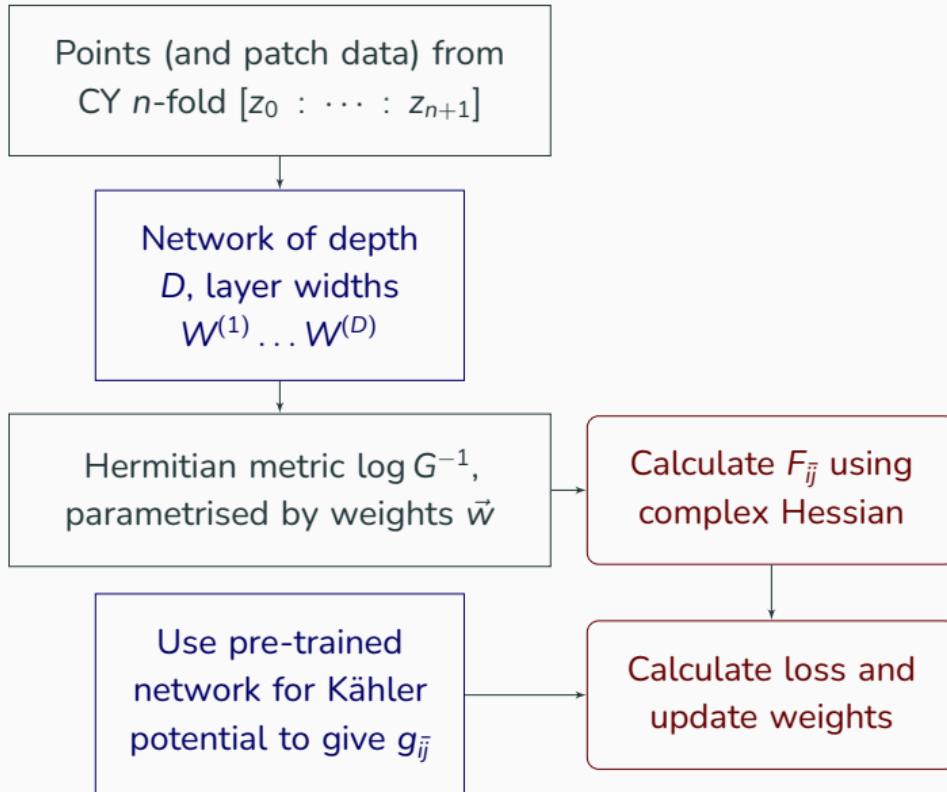
Loss function is

$$\text{Loss}[F, g] = E[F, g] \equiv \langle \text{tr } F_g^2 \rangle - \frac{1}{\text{rank } V} \langle \text{tr } F_g \rangle^2$$

After training, the network gives a **NN-based representation** of the HYM connection

- Effectively the **functional form** of G (plus A or F as can take derivatives, etc.)

General strategy



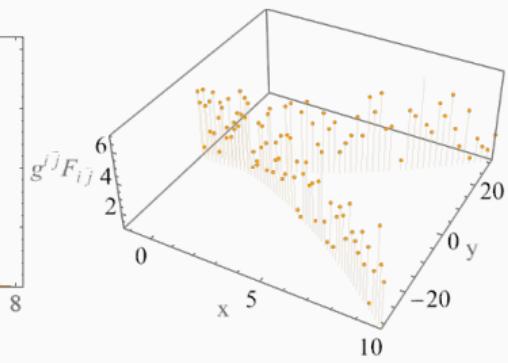
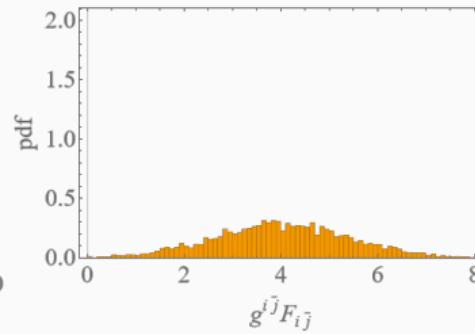
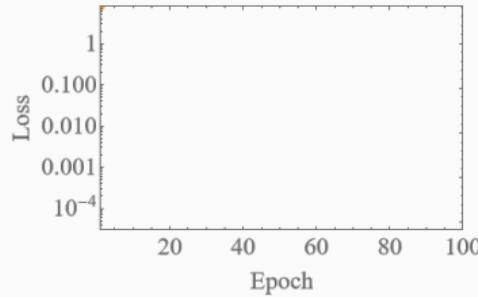
$\mathcal{O}_X(4)$ on elliptic curve

Line bundle $\mathcal{O}(4)$ over **elliptic curve** defined by

$$Q(z) \equiv z_1^3 - z_0^2 z_1 - z_0 z_2^2 + z_0^3 = 0 \quad \subset \mathbb{P}^2$$

- Solution to HYM should give $g^{i\bar{j}} F_{i\bar{j}} = 4$ **pointwise**

Evolution of loss, pdf of $g^{i\bar{j}} F_{i\bar{j}}$ and values of $g^{i\bar{j}} F_{i\bar{j}}$ on elliptic curve



$\mathcal{O}_X(4)$ on elliptic curve

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Evolution of loss, pdf of $g^{i\bar{j}} F_{i\bar{j}}$ and values of $g^{i\bar{j}} F_{i\bar{j}}$ on elliptic curve

$\mathcal{O}_X(m)$ on quintic threefold

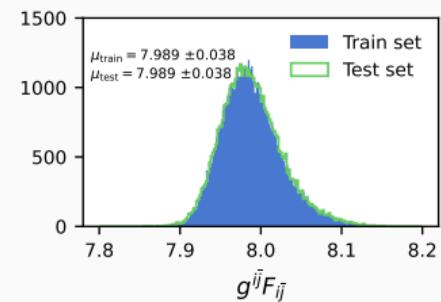
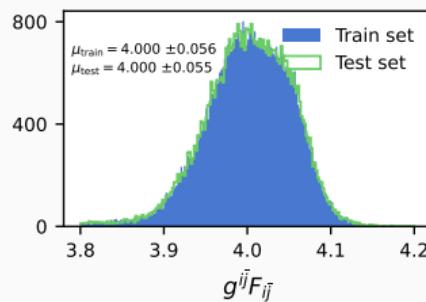
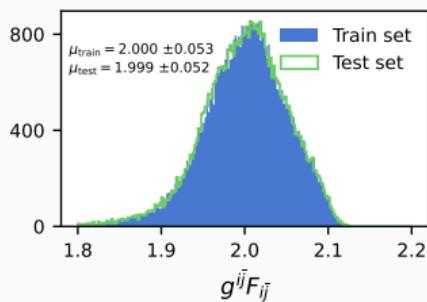
Dwork quintic defined by

$$Q(z) \equiv z_0^5 + \cdots + z_4^5 + \frac{1}{2}z_0z_1z_2z_3z_4 = 0 \quad \subset \mathbb{P}^4$$

Approximate CY metric computed with $\sigma = 0.001$

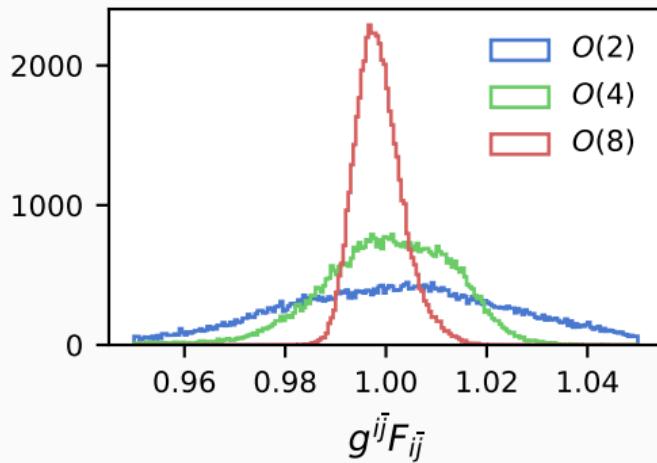
Neural networks of depth $D = 2, 3, 4$ with intermediate $W = 100$ layers

- Histogram of values of $g^{ij}\bar{F}_{ij}$ – should be **constant** over X



$\mathcal{O}_X(1)$ on quintic threefold

$D = 2, 3, 4$ networks give connections on $\mathcal{O}_X(2)$, $\mathcal{O}_X(4)$ and $\mathcal{O}_X(8)$ – **untwist** to give connections on $V = \mathcal{O}_X(1)$



Loss curves show that $D = 2$ network is **underparametrised**, but all still within 5% of expected result $g^{i\bar{j}}F_{i\bar{j}} = 1$

Application: moduli dependence and the swampland

The swampland distance conjecture

As an application of these numerical techniques, we can compute the spectrum of the Laplacian on a Calabi–Yau. This will teach us something about its KK spectrum and effective field theory.

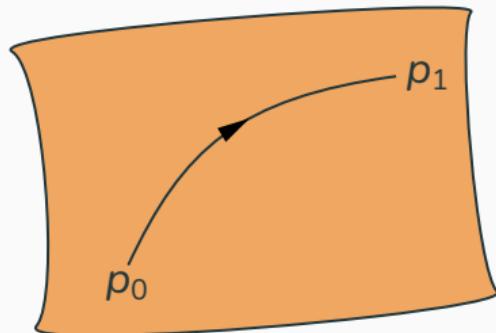
The Swampland Distance Conjecture [Ooguri, Vafa '06]

For an effective theory coupled to **gravity** with a moduli space parametrised by vevs of some fields, moving an **infinite distance** in moduli space brings down a tower of **massless states** that spoil the initial effective theory.

Finite version: there is a relation between the **mass** of a tower of states and the **distance** you move in moduli space

The swampland distance conjecture

Compare the effective theory at two points p_0 and p_1 in moduli space which are a **geodesic distance** $d(p_0, p_1)$ apart



Geodesic γ with $\gamma(\tau_i) = p_i$, distance along curve given by

$$d(p_0, p_1) = \int_{\tau_1}^{\tau_2} \sqrt{G_{ab} \dot{\gamma}^a \dot{\gamma}^b}$$

where G_{ab} is **metric** on moduli space

Conjecture implies that a tower of states becomes **light** on moving from p_0 to p_1 with masses

$$m(p_1) \sim m(p_0) e^{-\alpha d(p_0, p_1)}$$

where α is **order one** coefficient

The spectrum with varying complex structure [AA, Ruehle '21]

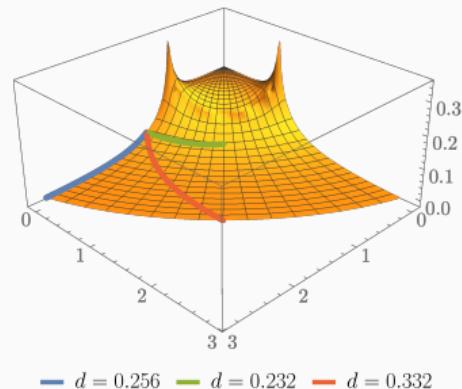
Focus on **complex structure** moduli space for quintic

$$Q_\psi \equiv z_0^5 + \cdots + z_4^5 - 5\psi z_0 \dots z_4$$

1. Compute the moduli space metric [Candelas et al. '98; Keller, Lukic '09]

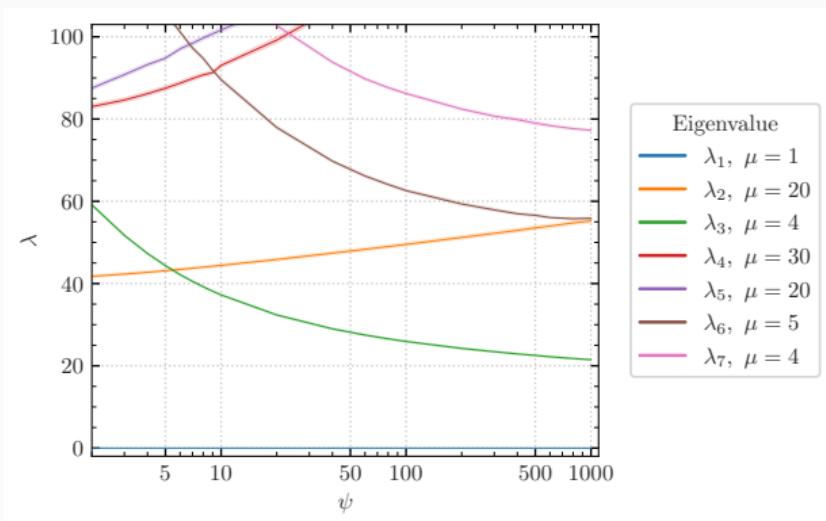
$$e^{-\mathcal{K}_{cs}} = i \int_{Q_\psi} \Omega \wedge \bar{\Omega}, \quad G_{a\bar{b}} = \partial_a \bar{\partial}_{\bar{b}} \mathcal{K}_{cs}$$

2. Compute geodesics and distances in moduli space for varying ψ
3. Compute the numerical Calabi–Yau metric and spectrum of the Laplacian Δ – scalars give **spin-two modes**



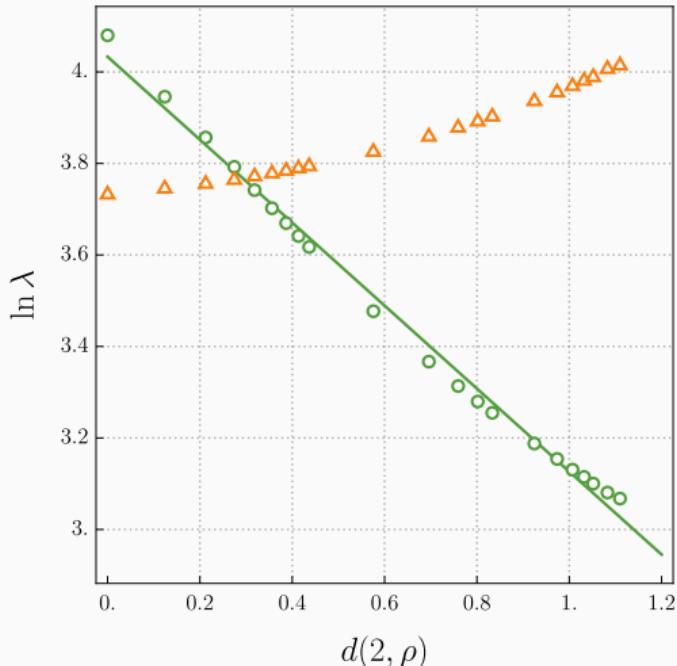
Scalar spectrum with varying ψ

Spectrum of $\Delta\phi = \lambda\phi$



1. Zero mode always present; massive modes appear with **multiplicities** given by irreps of $(S_5 \times \mathbb{Z}_2) \ltimes (\mathbb{Z}_5)^3$
2. Eigenvalues with small degeneracy become **lighter**
3. Local maximum in spectral gap for $\psi \approx 4$

A check of the distance conjecture



1. λ_3 falls exponentially as

$$\begin{aligned}\lambda_3 &= 56.4 e^{-(0.906 \pm 0.034) d(2, \rho)} \\ &\sim e^{-2\alpha d(p_0, p_1)}\end{aligned}$$

2. We see $\alpha \approx 0.45$ which is indeed order one
3. Almost saturates conjectured lower bound of $1/\sqrt{6} \approx 0.41$
[Andriot et al. '20]

Summary and outlook

Calabi–Yau metrics and HYM connections are accessible with **numerical methods** and **machine learning**

- Can extend to other bundles, compute (p, q) -form **Laplacian spectrum**
 - source of new non-BPS data [AA ‘20]

Ongoing work: **bundle-valued harmonic forms**, compute Yukawa couplings at a given point in moduli space

- SYZ conjecture? Non-Kähler metrics? G_2 metrics? Neural networks as general PDE solvers?
- 2d CFTs? [Afkhami-Jeddi, AA, Córdova ‘21] Input for **conformal bootstrap**?
[Lin et al. ‘15; Lin et al. ‘16]