Transverse energy analysis of relativistic heavy ion collisions through the use of identified particles spectra

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Table of Contents

List of Tables

List of Figures

Chapter 1

Introduction

One of the main focuses of current experimental and theoretical nuclear physics research is the study of the phase diagram of nuclear matter at a range of temperatures and baryon chemical potentials. In experiments involving the collisions of heavy ions at high and low energies, different regions of the phase diagram can be probed by varying the collision energy [3]. For instance, the high-baryon chemical potential regime corresponds to lower beam energies and higher temperatures correspond to higher beam energies. The results of these experiments and model calculations can be used to study the nature of transitions in the phase diagram.

Quantum chromodynamics (QCD) – the gauge theory of strong interaction [12, 20] – predicts a phase transition, at energy densities above 0.2-1 GeV/fm^3 [1] and around a critical temperature of about 200 MeV [14], of nuclear matter to a phase with quarks and gluons in thermal and chemical equilibrium representing the relevant degrees of freedom and behaving like an almost perfect quantum fluid [8]. This deconfined state of quarks and gluons is termed the quark-gluon plasma (QGP) in analogy to the quantum electrodynamical plasma phase of matter. The deconfinement is what the weakening of the strong interaction due to the polarization of the QCD vacuum is expected to lead to at high energies. The expectation of this phase transition also makes sense in terms of the chiral symmetry of the QCD Lagrangian, which is spontaneously broken at low temperatures, but restored at high temperatures, providing a sufficient condition for the deconfinement.

A schematic representing the QCD phase diagram on the temperature (T) and quark chemical potential (μ) plane is shown in Figure 1.1 [6]. A second-order transition

is predicted at low baryon chemical potentials (close to baryon-antibaryon symmetry) and high temperatures reminiscent of the early universe but within reach at modern facilities, specifically the Relativistic Heavy Ion Collider (RHIC) at the Brookhaven National Laboratory and the Large Hadron Collider (LHC) at CERN. At low temperatures and high chemical potentials, loose predictions have been made regarding the existence of exotic phases of high density matter, and programs, such as the Compressed Baryonic Matter experiment at the Facility for Antiproton and Ion Research in Germany, are being designed to study this region of the phase diagram.

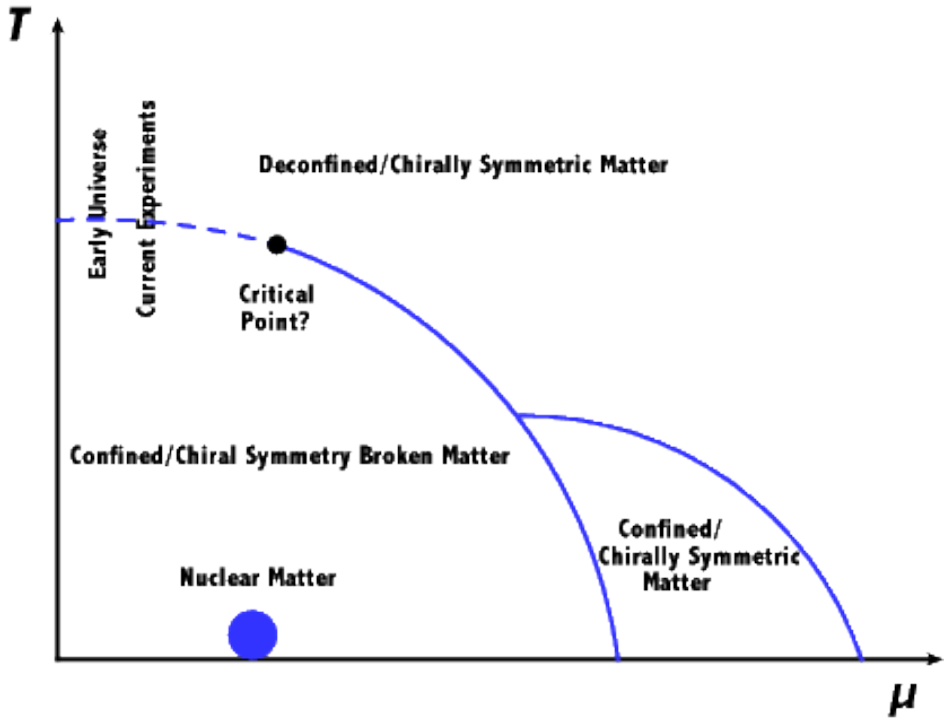

Figure 1.1: Schematic of the QCD phase diagram [6].

The existence and properties of the QGP in the aftermath of high-energy heavy-ion collisions can be probed using different techniques relevant to several theoretical characteristics of the phase. For instance, the interacting nuclei carry no net strangeness before colliding, and so a post-collision observation of strange and multi-strange particles can be a signal for an antecedent existence of deconfined quarks and gluons [10]. This signal,

when complemented with an observation of the suppression???????or enhancement of strange particles production, provides a strong hint of the formation of QGP. This can be further complemented with the estimate of the energy density and the temperature attained after the collision.

Analyses of experimental results have thus far provided signatures of the formation of matter with partonic degrees of freedom at the early stages of the collisions. Such signatures include suppression of high monentum hadrons, known as jet quenching, because the QGP is nearly opaque to colored probes, and large azimuthal anisotropies, indicating that the medium is a liquid of quarks and gluons [2]?????. Experiments also reveal the initial energy density of this matter to be about two orders of magnitude larger than that of low energy nuclear matter – comfortably more than the deconfinement phase transition critical density predicted by lattice QCD [11].

The state of the colliding nuclei before the collision at LHC and top RHIC energies has indications of being a Color Glass Condensate – strongly interacting, weakly coupled highly coherent gluonic matter [15]. The characteristics of the initial states of these nuclei affect the partonic distributions within the nuclei and ultimately the products of the collision. The collision products are also affected by variables such as the initial energy and entropy densities of the partonic matter [11].

Different observables can be used to study different aspects of heavy ion collisions. The charged particle multiplicity, $\langle N_{ch} \rangle$, is a global variable that relates to the entropy production during the collision (analysis note). The transverse energy, E_T, a global variable related to $\langle N_{ch} \rangle$, provides information about the conversion of the initial beam-direction kinetic energy into energy flowing in the transverse direction after the collision. Together, the studies of the fluctuation of the $\langle N_{ch} \rangle$ and the E_T pseudorapidity [footnote] density with respect to the beam energy and the collision centrality [footnote] help probe the characteristics of the initial conditions at the time of the collision. One can study, for instance, the distinctions between models based on quark participants against those based on nucleon participants [analysis note]. These quantities can also lead to the rough estimate of the initial energy

density through the use of the Bjorken formula [16]:

$$\epsilon \geq \frac{\frac{dE_T}{d\eta}}{\tau_0 \pi R^2} = \frac{3}{2}\langle\frac{E_T}{N}\rangle\frac{\frac{dN_{ch}}{d\eta}}{\tau_0 \pi R^2} \tag{1.1}$$

The transverse energy and the charged particle pseudorapidity densities have conventionally been calculated by using the transverse energy measurements obtained from calorimeters. This thesis details the use of particle spectra, reported as $\frac{d^2N}{dydp_T}$, from Au+Au collisions at RHIC to calculate the same global variables and serve as a method to cross check the ones involving calorimeters.

The organization of the thesis is as follows. Chapter II contains brief descriptions of different conventional methods used to estimate E_T as well as an elaboration of the method specific to this thesis.

Chapter 2

Method

In theory, E_T from a collision can be defined as the sum of the transverse masses, m_T, of all the particles produced in the collision, i.e.,

$$E_T \equiv \sum_i m_{T,i} \tag{2.1}$$

with

$$m_T \equiv \sqrt{p_T^2 + m^2} \tag{2.2}$$

where m is the rest mass of the particle and p_T is its transverse momentum. Using this definition to calculate the E_T requires perfect identification of all the particles. It has not been possible to do so in experiments, and so a more feasible, operational definition of E_T is fabricated. A commonly accepted definition in case of the feasibility of calorimetric measurements is [4, 8]:

$$E_T = \sum_i E_i \sin\theta_i, \tag{2.3}$$

$$\frac{dE_T}{d\eta} = \sin\theta \frac{dE}{d\eta}, \tag{2.4}$$

where the index i runs over all the particles going into a fixed solid angle for each event, θ is the polar angle, i.e, the angle with respect to the beam axis, η is the pseudorapidity defined as

$$\eta \equiv -\ln\tan\frac{\theta}{2}, \tag{2.5}$$

and E_i is the energy deposited in the calorimeter by the i^{th} particle. E_i is considered to be, by convention [5]???, the following

$$E_i = \begin{cases} E_i^{tot} - m_0 & \text{for baryons} \\ E_i^{tot} + m_0 & \text{for anti-baryons} \\ E_i^{tot} & \text{otherwise} \end{cases} \tag{2.6}$$

where E_i^{tot} is the total energy of the i^{th} particle defined canonically as

$$E^{tot} \equiv \sqrt{p^2 + m_0^2} \tag{2.7}$$

and m_0 is the particle's rest mass. In order to account for the portion of the emitted transverse energy not detected or overestimated by the calorimeters, corrections are made based on GEANT simulations. Transverse energy analysis can be done using tracking detectors as well if they are able to produce measurements of other physical quantities that implicitly contain information about the transverse energy. Specifically, the charged particle multiplicity distributions with respect to the transverse momenta can be used to calculate the particle's transverse energy pseudorapidity density. In fact, since the corrections related to the tracking detectors are very different from those related to the calorimeters, results from the two different methods can be used to test the assumptions involved in each.

The tracking detectors in experiments such as the STAR (Solenoidal Tracker At RHIC) experiment and ALICE (A Large Ion Collider Experiment) at CERN include Time Projection Chambers (TPCs) and Time-of-Flight (TOF) detectors that can give us the p_T spectra, yields and particle ratios of the identified charged hadrons [19, 2]. The TPCs provide measurements of particle trajectories – that can be used to determine the momenta for low-momentum particles – and of their specific energy loss,

$$\frac{dE}{dx}, \tag{2.8}$$

which can be used with the trajectories to make particle identifications using the Bethe-Bloch formula [7]. TOF detectors, on the other hand, cover the high-momentum part of

the measurements. In ALICE, the combination of the measurements of the TPC with those of the Inner Tracking System (ITS) effectively adds the tracking length, thereby improving the resolution of the measured p_T spectrum. Details about the particle identification and momentum determination capabilities of the detectors in ALICE can be found in [9].

In the STAR experiment, the TPC is the primary tracking detector. It is 4.2 m long and it cylindrically enshrouds the accelerator beam pipe from its outside, with an inner diameter of 1 m and an outer diameter of 4 m [17]. !!!!!!!!! more details about the TPC, then its limitation in high momentum resolution, then transition to TOF and some of its details !!!!!!!!!

The RHIC, in 2010, started a multi-phase Beam Energy Scan (BES) program to study the QCD phase diagram. The collider has the unique facility to collide nulclei at a range of center-of-mass energies per nucleon, $\sqrt{s_{NN}}$. It also has two different detectors that are currently operational, STAR and PHENIX (Pioneering High Energy Nuclear Interactions eXperiment), which facilitate the cross-checking of results. Between 2010 and 2011, under the exploratory phase I of the BES program, 7.7, 11.5 (not completed in PHENIX), 19.6, 27, and 39 GeV collisions were completed using pairs of Au nuclei. Together with the data formerly collected by the RHIC at higher collision energies, BES phase I data can scan the interval from 450 MeV to 20 MeV in μ_B space [18, 13]. One of the things that can be studied with the data associated with this region of the phase space is statedly the possibility of a "turn-off of new phenomena already established at higher RHIC energies" (https://drupal.star.bnl.gov/STAR/starnotes/public/sn0493). Results corresponding to the high-μ_B region might provide evidence of a first order phase transition, and possibly the critical point [13].

The manifestation of such phenomena would be in terms of the fluctuations in the properties of the post-collision system. One can, for instance, study the scaling of the transverse energy after the collision with the longidutional energy at the time of the collision, $\sqrt{s_{NN}}$. This can be done in multiple ways for a detector like STAR or PHENIX that is made up of sub-systems such as the TOF detectors, TPCs/Time Expansion Chambers, and calorimeters.

Adare et al. [3] use calorimetry in PHENIX to analyze the transverse energy corresponding to several different pairs of species colliding at a range of energies. They use the raw transverse energy measured by the EMCal, E_{TEMC}, to obtain the total hadronic E_T by making corrections in three different steps. They first scale the data by a constant factor calculated to account for the fiducial acceptance in azimuth and pseudorapidity. The second factor is calculated to adjust for the effects of the calorimeter towers that are disabled. The third factor, k, is computed as follows

$$k = k_{response} \times k_{inflow} \times k_{losses} \tag{2.9}$$

where $k_{response}$ corresponds to hadronic particles only depositing a fraction of their total energy while passing through the EMCal, k_{inflow} is attributable to the energy deposited by particles coming from outside the EMCal's fiducial aperture, and k_{losses} accounts for the energy not registered in the EMCal due to energy thresholds, edge effects, and more importantly due to the particles that make it into the fiducial aperture but decay into products outside the aperture.

Another method of transverse energy analysis, employed in this thesis, is to use the p_T spectra available from the tracking detectors. The TPCs and TOF detectors in STAR, for instance, can identify particles as well as their trajectories and ultimately their multiplicity distributions with respect to the momenta. Adamczyk et al. [2] report the results for the p_T spectra for six different identified hadrons, pi^+, pi^-, K^+, K^-, p, and $\bar{p}$, from the STAR experiment. The spectra come from Au+Au collisions – at $\sqrt{s_{NN}}$ = 7.7, 11.5, and 39 GeV in the year 2010 and at $\sqrt{s_{NN}}$ = 19.6 and 27 GeV in 2011 – under the BES Program. Figure 2.1 [2] shows the spectra corresponding to 39 GeV collisions categorized into seven different collision centralitiy classes. These spectra, and their counterparts for the rest of the energies, were used to calculate an estimate of the total transverse energy per event per particle species. This result was then used to estimate the total transverse energy due to all the collision products. The corrections applied by Adamczyk et al. [2] to the raw data to obtain the spectra, the reported systematic uncertainties in their results, and the

187 mathematics involved in calculating the transverse energy results using the p_T spectra are discussed below.

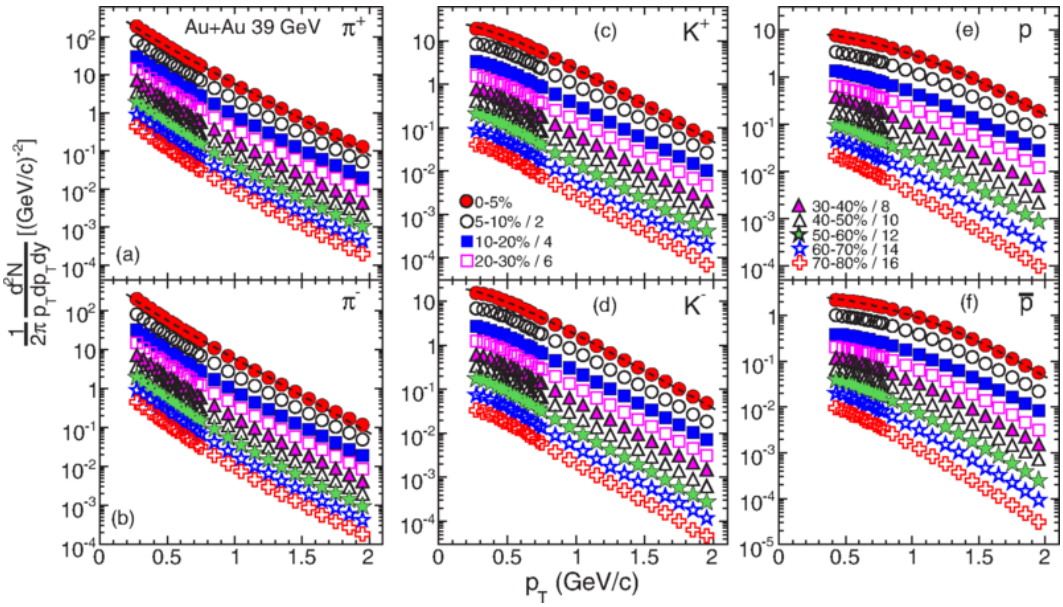

Figure 2.1: Transverse momentum spectra for pi^+, pi^-, K^+, K^-, p, and $\bar{p}$ at midrapidity (mod y ¡ 0.1) from 39 GeV Au+Au collisions at RHIC. The fitting curves on the 0-5% central collision spectra for pions, kaons, and protons/anti-protons represent, respectively, the Bose-Einstein, m_T-exponential, and double-exponential functions. [2].

188

189

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Appendices