

## Force

$$\text{Gravitational force } F = \frac{G * m_1 * m_2}{R^2} \text{ (N)}$$

$$\text{Electrical force } F = \frac{q_1 * q_2}{4 * \pi * \epsilon * R^2} * \hat{R} = E * q \text{ (N)}$$

$$\epsilon = 8.854 * 10^{-12}$$

$$\text{Gravitational field} = \frac{G * m_1}{R^2} * \hat{R} \text{ (N/kg)}$$

$$\text{Electrical field } E = \frac{q_1}{4 * \pi * \epsilon * R^2} \hat{R} \text{ (V/m)}$$

$$\mu = 4 * \pi * 10^{-7} \text{ H/m}$$

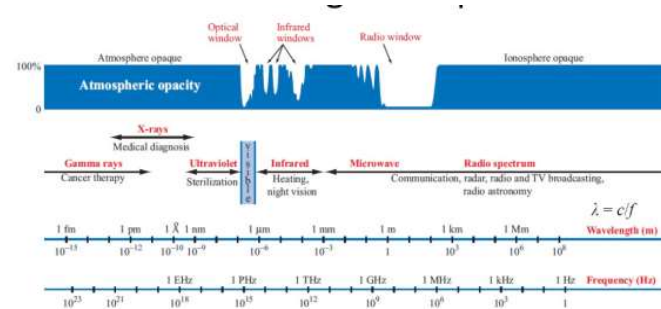
## Waves

$$y(x, t) = A * \cos\left(\frac{2 * \pi * t}{T} - \frac{2 * \pi * x}{\lambda} + \phi_0\right) \rightarrow y(x, t) = A * \cos(\omega t - \beta x)$$

<Opposite sign positive direction> <Same sign negative direction>

$$\text{Phase velocity } = u = \frac{\lambda}{T} = \lambda * f = \frac{\omega}{\beta}$$

$$\text{Lossy media } y(x, t) = A * e^{-\alpha x} * \cos(\omega t - \beta x + \phi)$$



## AC/DC circuit

$$\text{Real battery terminal voltage } \Delta V = V_{emf} - I * r$$

$$\text{Power } I * \Delta V = I * V_{emf} = I^2 * R = \frac{V^2}{R} \text{ (W)}$$

$$\sin(\theta) = \cos(90 - \theta) = \cos(\theta - 90)$$

$$\text{AC voltage } \Delta V = V_{max} * \sin(\omega t)$$

$$\text{Resistor in AC: } i = \frac{V_{max}}{R} * \sin(\omega t) \quad \text{no phase difference}$$

$$\text{Capacitor in AC: } I_C = \omega * C * \Delta V_{max} * \sin\left(\omega t + \frac{\pi}{2}\right) \quad \text{Current leads voltage by } 90^\circ$$

$$\text{Inductor in AC: } I_L = \frac{V_{max}}{\omega * L} * \sin\left(\omega t - \frac{\pi}{2}\right) \quad \text{Current lags voltage by } 90^\circ$$

Kirchhoff current node law and voltage loop rule

$$\text{Average power} = I_{rms}^2 * R$$

$$\cos(-\theta) = \cos(\theta)$$

$$\omega = 2\pi f \quad V_{rms} = \frac{V_{max}}{\sqrt{2}}$$

$$I_{max} = \frac{V_{max}}{R} \quad I_{rms} = \frac{I_{max}}{\sqrt{2}}$$

$$I_{max} = \omega * C * V_{max} \quad X_C = \frac{1}{\omega * C} = \frac{V_{max}}{I_{max}}$$

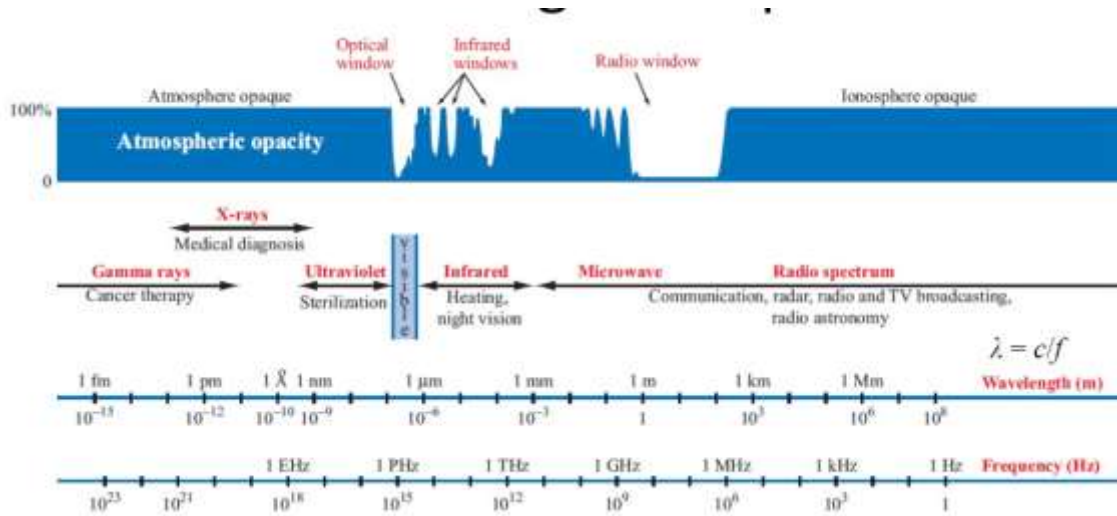
$$X_L = \omega * L * j$$

Euler's Identity: $e^{j\theta} = \cos \theta + j \sin \theta$	
$\sin \theta = \frac{e^{j\theta} - e^{-j\theta}}{2j}$	$\cos \theta = \frac{e^{j\theta} + e^{-j\theta}}{2}$
$\mathbf{z} = x + jy =  \mathbf{z} e^{j\theta}$	$\mathbf{z}^* = x - jy =  \mathbf{z} e^{-j\theta}$
$x = \Re(\mathbf{z}) =  \mathbf{z}  \cos \theta$	$ \mathbf{z}  = \sqrt{\mathbf{z}\mathbf{z}^*} = \sqrt{x^2 + y^2}$
$y = \Im(\mathbf{z}) =  \mathbf{z}  \sin \theta$	$\theta = \tan^{-1}(y/x)$
$\mathbf{z}^n =  \mathbf{z} ^n e^{jn\theta}$	$\mathbf{z}^{1/2} = \pm  \mathbf{z} ^{1/2} e^{j\theta/2}$
$\mathbf{z}_1 = x_1 + jy_1$	$\mathbf{z}_2 = x_2 + jy_2$
$\mathbf{z}_1 = \mathbf{z}_2$ iff $x_1 = x_2$ and $y_1 = y_2$	$\mathbf{z}_1 + \mathbf{z}_2 = (x_1 + x_2) + j(y_1 + y_2)$
$\mathbf{z}_1 \mathbf{z}_2 =  \mathbf{z}_1   \mathbf{z}_2  e^{j(\theta_1 + \theta_2)}$	$\frac{\mathbf{z}_1}{\mathbf{z}_2} = \frac{ \mathbf{z}_1 }{ \mathbf{z}_2 } e^{j(\theta_1 - \theta_2)}$
$-1 = e^{j\pi} = e^{-j\pi} = 1 \angle \pm 180^\circ$	
$j = e^{j\pi/2} = 1 \angle 90^\circ$	$-j = e^{-j\pi/2} = 1 \angle -90^\circ$
$\sqrt{j} = \pm e^{j\pi/4} = \pm \frac{(1 + j)}{\sqrt{2}}$	$\sqrt{-j} = \pm e^{-j\pi/4} = \pm \frac{(1 - j)}{\sqrt{2}}$

$z(t)$		$\tilde{z}$
$A \cos \omega t$	$\longleftrightarrow$	$A$
$A \cos(\omega t + \phi_0)$	$\longleftrightarrow$	$A e^{j\phi_0}$
$A \cos(\omega t + \beta x + \phi_0)$	$\longleftrightarrow$	$A e^{j(\beta x + \phi_0)}$
$A e^{-\alpha x} \cos(\omega t + \beta x + \phi_0)$	$\longleftrightarrow$	$A e^{-\alpha x} e^{j(\beta x + \phi_0)}$
$A \sin \omega t$	$\longleftrightarrow$	$A e^{-j\pi/2}$
$A \sin(\omega t + \phi_0)$	$\longleftrightarrow$	$A e^{j(\phi_0 - \pi/2)}$
$\frac{d}{dt}(z(t))$	$\longleftrightarrow$	$j\omega \tilde{z}$
$\frac{d}{dt}[A \cos(\omega t + \phi_0)]$	$\longleftrightarrow$	$j\omega A e^{j\phi_0}$
$\int z(t) dt$	$\longleftrightarrow$	$\frac{1}{j\omega} \tilde{z}$
$\int A \sin(\omega t + \phi_0) dt$	$\longleftrightarrow$	$\frac{1}{j\omega} A e^{j(\phi_0 - \pi/2)}$

$$\text{RLC circuit maximum frequency } \omega = \frac{1}{\sqrt{L * C}}$$

$$\text{Attenuating wave: } A e^{-\alpha x} \cos(\omega t - \beta x + \phi_0) \rightarrow A * e^{-\alpha x} * e^{j(-\beta x + \phi_0)}$$



## Transmission lines

Lines affects  $\phi = \frac{2\pi f l}{c} = \frac{2\pi l}{\lambda}$

$\lambda = \frac{c}{f}$

$R_s = \sqrt{\frac{\pi f \mu}{\sigma}} \rightarrow \text{surface resistance of conductor}$

$c = 3 \times 10^8 \text{ m/s}$

- When  $\frac{l}{\lambda}$  is very small ignore effects
- When  $\frac{l}{\lambda} > 0.01$ , need to account for phase delay and possibly reflection
- When  $\frac{l}{\lambda} > 0.25$ , definitely need to account for phase delay and possibly reflection

Parameter	Coaxial	Two-Wire	Parallel-Plate	Unit
$R'$	$\frac{R_s}{2\pi} \left( \frac{1}{a} + \frac{1}{b} \right)$	$\frac{2R_s}{\pi d}$	$\frac{2R_s}{w}$	$\Omega/\text{m}$
$L'$	$\frac{\mu}{2\pi} \ln(b/a)$	$\frac{\mu}{\pi} \ln \left[ (D/d) + \sqrt{(D/d)^2 - 1} \right]$	$\frac{\mu h}{w}$	$\text{H/m}$
$G'$	$\frac{2\pi \sigma}{\ln(b/a)}$	$\frac{\pi \sigma}{\ln \left[ (D/d) + \sqrt{(D/d)^2 - 1} \right]}$	$\frac{\sigma w}{h}$	$\text{S/m}$
$C'$	$\frac{2\pi \epsilon}{\ln(b/a)}$	$\frac{\pi \epsilon}{\ln \left[ (D/d) + \sqrt{(D/d)^2 - 1} \right]}$	$\frac{\epsilon w}{h}$	$\text{F/m}$

$L' * C' = \mu \epsilon$

$\frac{G'}{C'} = \frac{\sigma}{\epsilon}$

Air line:  $\epsilon = \epsilon_0 = 8.854 \times \frac{10^{-12} \text{ F}}{\text{m}}$

$\mu = \mu_0 = 4 * \pi * 10^{-7}$

$\sigma = 0$

$G' = 0$

Dispersion  $\rightarrow$  Distorts signals because different frequency components  $\rightarrow$  Proportional to the length of the transmission line

Telegraphers equations (time domain):

$$-\frac{\partial v(z,t)}{\partial z} = R' i(z,t) + L' \frac{\partial i(z,t)}{\partial t}$$

$$-\frac{\partial i(z,t)}{\partial z} = G' v(z,t) + C' \frac{\partial v(z,t)}{\partial t}$$

Telegraphers equations (phasor domain):

$$-\frac{d\tilde{V}(z)}{dz} = (R' + j\omega L') \tilde{I}(z)$$

$$-\frac{d\tilde{I}(z)}{dz} = (G' + j\omega C') \tilde{V}(z)$$

Derive the **wave equations** by separating variables

$$\frac{d^2 \tilde{V}}{dz^2} - \gamma^2 \tilde{V}(z) = 0$$

$$\frac{d^2 \tilde{I}}{dz^2} - \gamma^2 \tilde{I}(z) = 0$$

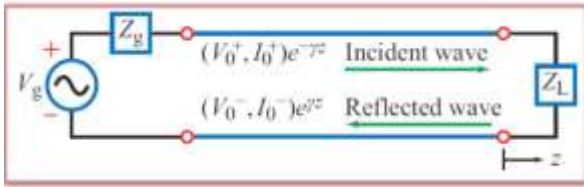
$$\text{Complex propagation constant} = \gamma = \sqrt{(R' + j\omega L')(G' + j\omega C')} = \alpha + j\beta$$

$$\text{Attenuation constant} = \alpha = \text{Re}(\gamma) \text{ Np/m}$$

$$\text{Phase constant} = \beta = \text{Im}(\gamma) \text{ rad/m}$$

$$\text{Characteristic Impedance} = Z_0 = \frac{R' + j\omega L'}{\gamma} = \frac{\sqrt{R' + j\omega L'}}{\sqrt{G' + j\omega C'}} = \frac{V_0^+}{I_0^+} = \frac{V_0^-}{I_0^-}$$

$$\text{Phase velocity} = u_p = \frac{\omega}{\beta} = f * \lambda$$



$$\tilde{V}(z) = V_0^+ e^{-\gamma z} + V_0^- e^{+\gamma z}$$

$$\tilde{I}(z) = I_0^+ e^{-\gamma z} + I_0^- e^{+\gamma z}$$

### Lossless transmission line

$$R' \text{ and } G' \text{ are negligible} \rightarrow \alpha = 0 \quad \beta = \omega * \sqrt{L' * C'} = \frac{\omega \sqrt{\epsilon_r}}{c} \quad \text{Phase velocity} = u_p = \frac{c}{\sqrt{\epsilon_r}} \quad Z_0 = \sqrt{\frac{L'}{C'}}$$

	Propagation Constant $\gamma = \alpha + j\beta$	Phase Velocity $u_p$	Characteristic Impedance $Z_0$
<b>General case</b>	$\gamma = \sqrt{(R' + j\omega L')(G' + j\omega C')}$	$u_p = \omega / \beta$	$Z_0 = \sqrt{\frac{(R' + j\omega L')}{(G' + j\omega C')}}$
<b>Lossless</b> ( $R' = G' = 0$ )	$\alpha = 0, \beta = \omega \sqrt{\epsilon_r} / c$	$u_p = c / \sqrt{\epsilon_r}$	$Z_0 = \sqrt{L' / C'}$
<b>Lossless coaxial</b>	$\alpha = 0, \beta = \omega \sqrt{\epsilon_r} / c$	$u_p = c / \sqrt{\epsilon_r}$	$Z_0 = (60 / \sqrt{\epsilon_r}) \ln(b/a)$
<b>Lossless two-wire</b>	$\alpha = 0, \beta = \omega \sqrt{\epsilon_r} / c$	$u_p = c / \sqrt{\epsilon_r}$	$Z_0 = \frac{120}{\sqrt{\epsilon_r}} \ln\left[\frac{D}{d}\right] + \sqrt{(D/d)^2 - 1}$ $Z_0 \simeq \frac{120}{\sqrt{\epsilon_r}} \ln(2D/d),$ if $D \gg d$
<b>Lossless parallel-plate</b>	$\alpha = 0, \beta = \omega \sqrt{\epsilon_r} / c$	$u_p = c / \sqrt{\epsilon_r}$	$Z_0 = \frac{120\pi}{\sqrt{\epsilon_r}} (h/w)$

$$\lambda = \frac{c}{f * \sqrt{\epsilon_r}}$$

$$\text{Normalized load impedance} = z_L = \frac{Z_L}{Z_0}$$

$$\text{voltage reflection coefficient} = \Gamma = \frac{Z_L - Z_0}{Z_L + Z_0} = \frac{V_0^+}{V_0^-} = -\frac{I_0^+}{I_0^-} = |\Gamma| * e^{j\theta_r}$$

$$\text{Wave impedance} = Z(d) = \frac{\tilde{V}(d)}{\tilde{I}(d)} \rightarrow \text{ratio of total voltage to total current}$$

$$\text{Wave impedance } Z(d) = Z_0 * \left[ \frac{1 + \Gamma_d}{1 - \Gamma_d} \right], \Gamma_d = \Gamma * e^{-j2\beta d}$$

$$\text{Input impedance} = Z_{in} = Z_0 * \left[ \frac{Z_L + j * \tan(\beta l)}{1 + j * Z_L * \tan(\beta l)} \right]$$

$$\text{Forward voltage} = V_0^+ = \left( \frac{\tilde{V}_g * Z_{in}}{Z_g + Z_{in}} \right) * \left[ \frac{1}{e^{j\beta l} + \Gamma * e^{-j\beta l}} \right] \rightarrow \text{this part may not be used}$$

Full final equation for **phasor voltage** and **phasor current** on the lossless line:

$$\begin{aligned} \tilde{V}(z) &= |V_0^+| e^{j\phi^+} [e^{-j\beta z} + |\Gamma| e^{j\theta_r} e^{j\beta z}] \\ \tilde{I}(z) &= \frac{|V_0^+|}{|Z_0|} e^{j\phi^+} [e^{-j\beta z} - |\Gamma| e^{j\theta_r} e^{j\beta z}] \end{aligned} \quad \begin{array}{l} \text{phasor} \\ \text{solutions} \end{array}$$

Full final equation for **instantaneous voltage** and **current** on the line:

$$\begin{aligned} v(z) &= \text{Re}\{\tilde{V}(z) e^{j\omega t}\} \\ &= |V_0^+| \{[\cos(\omega t - \beta z + \phi^+)] + |\Gamma| [\cos(\omega t + \beta z + \theta_r + \phi^-)]\} \\ i(z) &= \text{Re}\{\tilde{I}(z) e^{j\omega t}\} \\ &= \frac{|V_0^+|}{|Z_0|} \{[\cos(\omega t - \beta z + \phi^+ - \phi_z)] + |\Gamma| [\cos(\omega t + \beta z + \phi^- - \phi_z)]\} \end{aligned} \quad \begin{array}{l} \text{time domain} \\ \text{solutions} \end{array}$$

$$V_0^+ = \left( \frac{\tilde{V}_g Z_{in}}{Z_g + Z_{in}} \right) \left[ \frac{1}{e^{j\beta l} + \Gamma e^{-j\beta l}} \right]$$

## Vector analysis

Transformation	Coordinate Variables	Unit Vectors	Vector Components
<b>Cartesian to cylindrical</b>	$r = \sqrt{x^2 + y^2}$ $\phi = \tan^{-1}(y/x)$ $z = z$	$\hat{r} = \hat{x} \cos \phi + \hat{y} \sin \phi$ $\hat{\phi} = -\hat{x} \sin \phi + \hat{y} \cos \phi$ $\hat{z} = \hat{z}$	$A_r = A_x \cos \phi + A_y \sin \phi$ $A_\phi = -A_x \sin \phi + A_y \cos \phi$ $A_z = A_z$
<b>Cylindrical to Cartesian</b>	$x = r \cos \phi$ $y = r \sin \phi$ $z = z$	$\hat{x} = \hat{r} \cos \phi - \hat{\phi} \sin \phi$ $\hat{y} = \hat{r} \sin \phi + \hat{\phi} \cos \phi$ $\hat{z} = \hat{z}$	$A_x = A_r \cos \phi - A_\phi \sin \phi$ $A_y = A_r \sin \phi + A_\phi \cos \phi$ $A_z = A_z$
<b>Cartesian to spherical</b>	$R = \sqrt{x^2 + y^2 + z^2}$ $\theta = \tan^{-1}(\sqrt{x^2 + y^2}/z)$ $\phi = \tan^{-1}(y/x)$	$\hat{R} = \hat{x} \sin \theta \cos \phi + \hat{y} \sin \theta \sin \phi + \hat{z} \cos \theta$ $\hat{\theta} = \hat{x} \cos \theta \cos \phi + \hat{y} \cos \theta \sin \phi - \hat{z} \sin \theta$ $\hat{\phi} = -\hat{x} \sin \phi + \hat{y} \cos \phi$	$A_R = A_x \sin \theta \cos \phi + A_y \sin \theta \sin \phi + A_z \cos \theta$ $A_\theta = A_x \cos \theta \cos \phi + A_y \cos \theta \sin \phi - A_z \sin \theta$ $A_\phi = -A_x \sin \phi + A_y \cos \phi$
<b>Spherical to Cartesian</b>	$x = R \sin \theta \cos \phi$ $y = R \sin \theta \sin \phi$ $z = R \cos \theta$	$\hat{x} = \hat{R} \sin \theta \cos \phi + \hat{\theta} \cos \theta \cos \phi - \hat{\phi} \sin \phi$ $\hat{y} = \hat{R} \sin \theta \sin \phi + \hat{\theta} \cos \theta \sin \phi + \hat{\phi} \cos \phi$ $\hat{z} = \hat{R} \cos \theta - \hat{\theta} \sin \theta$	$A_x = A_R \sin \theta \cos \phi + A_\theta \cos \theta \cos \phi - A_\phi \sin \phi$ $A_y = A_R \sin \theta \sin \phi + A_\theta \cos \theta \sin \phi + A_\phi \cos \phi$ $A_z = A_R \cos \theta - A_\theta \sin \theta$
<b>Cylindrical to spherical</b>	$R = \sqrt{r^2 + z^2}$ $\theta = \tan^{-1}(r/z)$ $\phi = \phi$	$\hat{R} = \hat{r} \sin \theta + \hat{z} \cos \theta$ $\hat{\theta} = \hat{r} \cos \theta - \hat{z} \sin \theta$ $\hat{\phi} = \hat{\phi}$	$A_R = A_r \sin \theta + A_z \cos \theta$ $A_\theta = A_r \cos \theta - A_z \sin \theta$ $A_\phi = A_\phi$
<b>Spherical to cylindrical</b>	$r = R \sin \theta$ $\phi = \phi$ $z = R \cos \theta$	$\hat{r} = \hat{R} \sin \theta + \hat{\theta} \cos \theta$ $\hat{\phi} = \hat{\phi}$ $\hat{z} = \hat{R} \cos \theta - \hat{\theta} \sin \theta$	$A_r = A_R \sin \theta + A_\theta \cos \theta$ $A_\phi = A_\phi$ $A_z = A_R \cos \theta - A_\theta \sin \theta$

	Cartesian Coordinates	Cylindrical Coordinates	Spherical Coordinates
<b>Coordinate variables</b>	$x, y, z$	$r, \phi, z$	$R, \theta, \phi$
<b>Vector representation <math>\mathbf{A} =</math></b>	$\hat{x}A_x + \hat{y}A_y + \hat{z}A_z$	$\hat{r}A_r + \hat{\phi}A_\phi + \hat{z}A_z$	$\hat{R}A_R + \hat{\theta}A_\theta + \hat{\phi}A_\phi$
<b>Magnitude of <math>\mathbf{A}</math> <math> \mathbf{A}  =</math></b>	$\sqrt{A_x^2 + A_y^2 + A_z^2}$	$\sqrt{A_r^2 + A_\phi^2 + A_z^2}$	$\sqrt{A_R^2 + A_\theta^2 + A_\phi^2}$
<b>Position vector <math>\vec{OP}_1 =</math></b>	$\hat{x}x_1 + \hat{y}y_1 + \hat{z}z_1$ for $P = (x_1, y_1, z_1)$	$\hat{r}r_1 + \hat{z}z_1$ for $P = (r_1, \phi_1, z_1)$	$\hat{R}R_1$ for $P = (R_1, \theta_1, \phi_1)$
<b>Base vectors properties</b>	$\hat{x} \cdot \hat{x} = \hat{y} \cdot \hat{y} = \hat{z} \cdot \hat{z} = 1$ $\hat{x} \cdot \hat{y} = \hat{y} \cdot \hat{z} = \hat{z} \cdot \hat{x} = 0$ $\hat{x} \times \hat{y} = \hat{z}$ $\hat{y} \times \hat{z} = \hat{x}$ $\hat{z} \times \hat{x} = \hat{y}$	$\hat{r} \cdot \hat{r} = \hat{\phi} \cdot \hat{\phi} = \hat{z} \cdot \hat{z} = 1$ $\hat{r} \cdot \hat{\phi} = \hat{\phi} \cdot \hat{z} = \hat{z} \cdot \hat{r} = 0$ $\hat{r} \times \hat{\phi} = \hat{z}$ $\hat{\phi} \times \hat{z} = \hat{r}$ $\hat{z} \times \hat{r} = \hat{\phi}$	$\hat{R} \cdot \hat{R} = \hat{\theta} \cdot \hat{\theta} = \hat{\phi} \cdot \hat{\phi} = 1$ $\hat{R} \cdot \hat{\theta} = \hat{\theta} \cdot \hat{\phi} = \hat{\phi} \cdot \hat{R} = 0$ $\hat{R} \times \hat{\theta} = \hat{\phi}$ $\hat{\theta} \times \hat{\phi} = \hat{R}$ $\hat{\phi} \times \hat{R} = \hat{\theta}$
<b>Dot product <math>\mathbf{A} \cdot \mathbf{B} =</math></b>	$A_x B_x + A_y B_y + A_z B_z$	$A_r B_r + A_\phi B_\phi + A_z B_z$	$A_R B_R + A_\theta B_\theta + A_\phi B_\phi$
<b>Cross product <math>\mathbf{A} \times \mathbf{B} =</math></b>	$\begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ A_x & A_y & A_z \\ B_x & B_y & B_z \end{vmatrix}$	$\begin{vmatrix} \hat{r} & \hat{\phi} & \hat{z} \\ A_r & A_\phi & A_z \\ B_r & B_\phi & B_z \end{vmatrix}$	$\begin{vmatrix} \hat{R} & \hat{\theta} & \hat{\phi} \\ A_R & A_\theta & A_\phi \\ B_R & B_\theta & B_\phi \end{vmatrix}$
<b>Differential length <math>d\mathbf{l} =</math></b>	$\hat{x} dx + \hat{y} dy + \hat{z} dz$	$\hat{r} dr + \hat{\phi} r d\phi + \hat{z} dz$	$\hat{R} dR + \hat{\theta} R d\theta + \hat{\phi} R \sin \theta d\phi$
<b>Differential surface areas</b>	$ds_x = \hat{x} dy dz$ $ds_y = \hat{y} dx dz$ $ds_z = \hat{z} dx dy$	$ds_r = \hat{r} r d\phi dz$ $ds_\phi = \hat{\phi} dr dz$ $ds_z = \hat{z} 2r dr d\phi$	$ds_R = \hat{R} R^2 \sin \theta d\theta d\phi$ $ds_\theta = \hat{\theta} R \sin \theta dR d\phi$ $ds_\phi = \hat{\phi} R dR d\theta$
<b>Differential volume <math>dV =</math></b>	$dx dy dz$	$r dr d\phi dz$	$R^2 \sin \theta dR d\theta d\phi$

- Cylindrical  $\langle r, \phi, z \rangle$

- Spherical  $\langle R, \theta, \phi \rangle$

- Gradient ( $\nabla$ )

$$\nabla = \frac{\partial}{\partial x} \hat{x} + \frac{\partial}{\partial y} \hat{y} + \frac{\partial}{\partial z} \hat{z}$$

- Divergence ( $\nabla \cdot \vec{A}$ )

$$\nabla \cdot \vec{A} = \text{div} \vec{A} = \frac{\partial A_x}{\partial x} + \frac{\partial A_y}{\partial y} + \frac{\partial A_z}{\partial z}$$

- Curl ( $\nabla \times \vec{A}$ )

$$\nabla \times \vec{A} = \text{curl} \vec{A} = \left( \frac{\partial}{\partial x} \hat{x} + \frac{\partial}{\partial y} \hat{y} + \frac{\partial}{\partial z} \hat{z} \right) \times (A_x \hat{x} + A_y \hat{y} + A_z \hat{z})$$

- Stoke's theorem

$$\int \nabla \times \vec{B} \cdot d\vec{s} = \oint \vec{B} \cdot d\vec{l}$$

- Divergence theorem

$$\int_V \nabla \cdot \vec{E} dV = \int_S \vec{E} \cdot d\vec{s}$$

Gradient of cylindrical coordinate  $\rightarrow \frac{\partial A}{\partial r} \hat{r} + \frac{1}{r} \frac{\partial A}{\partial \phi} \hat{\phi} + \frac{\partial A}{\partial z} \hat{z}$

Gradient of Spherical coordinate  $\rightarrow \frac{\partial A}{\partial R} \hat{R} + \frac{1}{R} \frac{\partial A}{\partial \theta} \hat{\theta} + \frac{1}{R \sin \theta} \frac{\partial A}{\partial \phi} \hat{\phi}$

**Coulomb's Law** (Find electric field given charge)

$$\vec{E}(\vec{R}) = \frac{q}{4\pi\epsilon R^2} \hat{R} = \frac{q}{4\pi\epsilon |\vec{R}|^3} \vec{R} \quad (V/m) \rightarrow$$

*electric field at point P due to single charge*

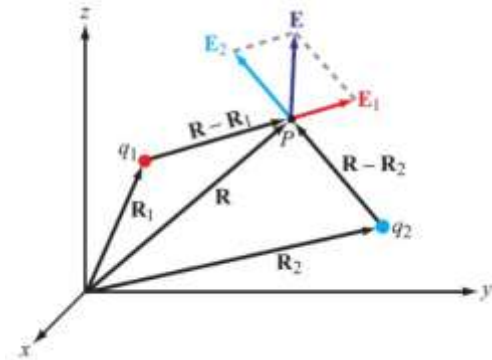
$\vec{F} = q' * \vec{E}(N) \rightarrow$  *electric force on a charge placed at P*

$$\vec{E}(\vec{R}) = \frac{1}{4\pi\epsilon} \sum_{i=0}^N \frac{q_i * (\vec{R} - \vec{R}_i)}{|\vec{R} - \vec{R}_i|^3} \quad (V/m)$$

$Q = \int_{V'} \rho_v(\vec{r}') dV' \rightarrow$  *total charge in a volume,  $\rho_v$  is charge density*

$$\vec{E}(\vec{R}) = \int_{V'} \frac{\rho_v(\vec{r}')(\vec{R} - \vec{r}')}{4\pi\epsilon |\vec{R} - \vec{r}'|^3} dV' \rightarrow$$
 *field at point P*

Infinite Plane (Disk) of charge  $\vec{E} = \pm \hat{z} \frac{\rho_v}{2\epsilon}$       Infinite line of charge  $E = \frac{\rho_v}{2\pi\epsilon r} \hat{r}$



$$\vec{E}(\vec{R}) = \int_{V'} \frac{\rho_v(\vec{r}')(\vec{R} - \vec{r}')}{4\pi\epsilon |\vec{R} - \vec{r}'|^3} dV' \quad \text{In a volume}$$

$$\vec{E}(\vec{R}) = \int_{s'} \frac{\rho_s(\vec{r}')(\vec{R} - \vec{r}')}{4\pi\epsilon |\vec{R} - \vec{r}'|^3} ds' \quad \text{Over a surface}$$

$$\vec{E}(\vec{R}) = \int_{l'} \frac{\rho_l(\vec{r}')(\vec{R} - \vec{r}')}{4\pi\epsilon |\vec{R} - \vec{r}'|^3} dl' \quad \text{On a line}$$

**Gauss's Law** (Find charge given a field)

$$\vec{D} = \epsilon \vec{E} \text{ (C/m}^2\text{)} \quad \epsilon = \epsilon_r * 8.854 * 10^{-12} \quad \vec{D} = \hat{R} \frac{q}{4 * \pi * R^2}$$

$$\oint_S \vec{D} \cdot d\vec{s}' = q = \int_{V'} \rho_v dV' \rightarrow \text{determine electric flux density } D \quad \nabla \cdot \vec{D} = \rho_v(x, y, z) \rightarrow \text{differential form}$$

## Electric Potential

$$V = - \int_{l'} \vec{E} \cdot d\vec{l}' \quad V_{21} = V_2 - V_1 = - \int_{P1}^{P2} \vec{E} \cdot d\vec{l}' \rightarrow \text{potential difference between } P1 \text{ and } P2$$

$$\oint \vec{E} \cdot d\vec{l}' = 0 \rightarrow \text{for any closed path} \quad V = - \int_{\infty}^P \vec{E} \cdot d\vec{l}' \rightarrow \text{zero reference at infinity (free space and material media)}$$

$$V(\vec{R}) = \frac{1}{4\pi\epsilon} \sum_{i=1}^N \frac{q_i}{|\vec{R} - \vec{R}_i|} \quad \vec{E} = \frac{V}{L} \quad \vec{E} = -\nabla V$$

$$V(\vec{R}) = \int_{V'} \frac{\rho_v(\vec{R}')}{4\pi\epsilon |\vec{R} - \vec{R}'|} dV' \quad \text{In a volume}$$

$$V(\vec{R}) = \int_{S'} \frac{\rho_s(\vec{R}')}{4\pi\epsilon |\vec{R} - \vec{R}'|} dS' \quad \text{Over a surface}$$

$$V(\vec{R}) = \int_{l'} \frac{\rho_l(\vec{R}')}{4\pi\epsilon |\vec{R} - \vec{R}'|} dl' \quad \text{On a line}$$

## Dielectrics

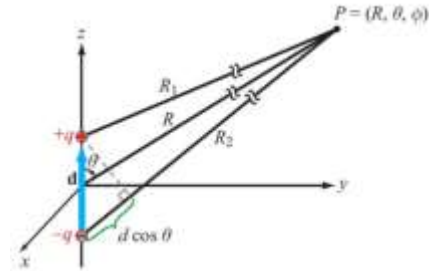
- An electric dipole consists of 2-point charges of equal magnitude but opposite polarity
  - o Applications: Dielectrics, molecular bonds, antennas

$$\vec{p} = q * \vec{d} \rightarrow \text{dipole moment} \quad V(\vec{R}) = \frac{\vec{p} \cdot \vec{R}}{4\pi\epsilon_0 |\vec{R} - \vec{R}_l|^2} = \frac{q * d * \cos\theta}{4\pi\epsilon_0 * R^2}$$

$$\vec{E}(\vec{R}) = \frac{qd}{4\pi\epsilon_0 |\vec{R} - \vec{R}_l|^3} (\hat{R} * 2 * \cos\theta + \hat{\theta} \sin\theta) \text{ V/m} \quad \text{only when } R \gg d$$

Types of dipoles in matter

- Permanent
  - o Molecule having atoms with different electronegativity
  - o Polar molecule  $\rightarrow$  water
- Instantaneous
  - o Electrons happen to concentrate in one place
- Induced
  - o A permanent dipole or applied electric field near another atom induces a dipole



In a dielectric material  $\vec{D} = \epsilon_0 \vec{E} + \vec{P} \rightarrow P$  is the electric flux density induced by applied field  $E$

$$\vec{P} = \epsilon_0 * X_e * \vec{E} \quad \vec{D} = \epsilon_0 \vec{E} + \epsilon_0 * X_e * \vec{E} = \epsilon_0 (1 + X_e) \vec{E} \quad (1 + X_e) = \text{relative permittivity} = \epsilon_r$$

## Conductors & Resistors

- Conductors are materials in which some of the electrons are free electrons
  - o Electrons can move relatively freely through the material
  - o Copper, aluminum, and silver
  - o Charge Carrier: A particle carrying charge that is free to move

$$\text{Total Current } I = \int_S \vec{J} \cdot d\vec{s} \rightarrow \vec{J} = \rho_v * \vec{u} \left( \frac{A}{m^2} \right) \rightarrow u \text{ is velocity} \rightarrow \rho_v = q * N, N \text{ is \# of charges per unit volume}$$

Drift velocity,  $u$ : Steady state average velocity of the electrons

Mobility  $\mu$ : Accounts for the effective mass of charged particle and the average distance before stopped by colliding

$$\vec{u}_e = -\mu_e \vec{E} \rightarrow \text{drift velocity of electrons (m/s)} \quad \vec{u}_h = \mu_h \vec{E} \rightarrow \text{drift velocity of holes (m/s)}$$

$$\vec{J} = \vec{J}_e + \vec{J}_h = \rho_e * \vec{u}_e + \rho_h * \vec{u}_h = (-\rho_e * \mu_e + \rho_h * \mu_h) \vec{E} \quad \vec{J} = \sigma \vec{E} \text{ (A/m}^2\text{)}$$

$$\text{Semiconductor / dielectric } \sigma = (-\rho_e * \mu_e + \rho_h * \mu_h) = -N_e q \mu_e + N_h q \mu_h \text{ (S/m)}$$

$$\text{Conductor } \sigma = -\rho_e * \mu_e = N_e e \mu_e \quad \text{For perfect dielectric: } N_e = 0, \sigma = 0, J = 0$$

$$\text{For perfect conductor } \mu_e = \infty, \sigma = \infty, E = D = 0$$

$$\text{For any conductor } R = \frac{V}{I} = \frac{- \int_{l'} \vec{E} \cdot d\vec{l}'}{\int_S \vec{J} \cdot d\vec{s}} = \frac{- \int_{l'} \vec{E} \cdot d\vec{l}'}{\int_S \sigma \vec{E} \cdot d\vec{s}} \quad R = \frac{l}{\sigma_1 * A_1 + \sigma_2 * A_2} \rightarrow \text{resistance coaxial cable}$$



## Electric Boundary Conditions

Field Component	Any Two Media	Medium 1 Dielectric $\epsilon_1$	Medium 2 Conductor
Tangential $\mathbf{E}$	$\mathbf{E}_{1t} = \mathbf{E}_{2t}$	$\mathbf{E}_{1t} = \mathbf{E}_{2t} = 0$	
Tangential $\mathbf{D}$	$\mathbf{D}_{1t}/\epsilon_1 = \mathbf{D}_{2t}/\epsilon_2$	$\mathbf{D}_{1t} = \mathbf{D}_{2t} = 0$	
Normal $\mathbf{E}$	$\epsilon_1 E_{1n} - \epsilon_2 E_{2n} = \rho_s$	$E_{1n} = \rho_s/\epsilon_1$	$E_{2n} = 0$
Normal $\mathbf{D}$	$D_{1n} - D_{2n} = \rho_s$	$D_{1n} = \rho_s$	$D_{2n} = 0$

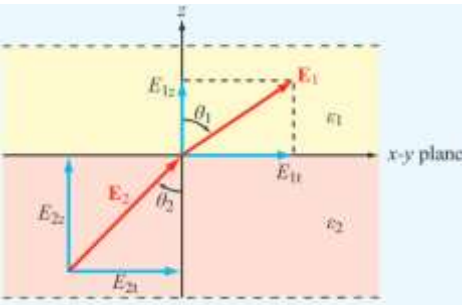
← General Boundary Conditions

$$\text{DD: } E_{1t} = E_{2t} \quad \& \quad E_{2n} = \frac{\epsilon_1}{\epsilon_2} E_{1n}$$

DC:  $E = D = 0$  in perfect conductor

$$D_{1t} = E_{1t} = 0 \quad \& \quad D_{1n} = \epsilon_1 E_{1n} = \rho_s$$

Notes: (1)  $\rho_s$  is the surface charge density at the boundary; (2) normal components of  $\mathbf{E}_1$ ,  $\mathbf{D}_1$ ,  $\mathbf{E}_2$ , and  $\mathbf{D}_2$  are along  $\hat{\mathbf{n}}_2$ , the outward normal unit vector of medium 2.



Dielectric – dielectric

$$\frac{\tan \theta_1}{\epsilon_1} = \frac{\tan \theta_2}{\epsilon_2}$$

Capacitor

$$C = \frac{Q}{V} = \frac{\int_S \epsilon \vec{E} \cdot d\vec{s}}{-\int_{l'} \vec{E} \cdot d\vec{l}} = \frac{\epsilon A}{d}$$

$$E = \frac{V}{d} \quad Q = \int_S \epsilon \vec{E} \cdot d\vec{s}$$

Magnetic Forces and torques

$$\text{Magnetic force} = \vec{F}_m = q\vec{u} \times \vec{B} \text{ (N)}$$

$$\text{Lorentz force} = q\vec{u} \times \vec{B} + q\vec{E}$$

$$dF_m = I d\vec{l} \times \vec{B} = dQ \vec{u} \times \vec{B}$$

Force on any closed current loop in a uniform magnetic field = 0

$$\text{Magnetic torque } T = \vec{m} \times \vec{B} \text{ (N} \cdot \text{m)}$$

$$|T| = N * I * A * B * \sin \theta$$

$$\vec{m} = \hat{n} N * I * A$$

Biot-Savart Law

$$\vec{B} = \mu_0 \vec{H} \text{ (T)}$$

$$d\vec{H} = \frac{I}{4\pi R^2} * d\vec{l} \times \hat{R} \left( \frac{A}{m} \right)$$

$$d\vec{B} = \frac{I \mu_0}{4\pi R^2} * d\vec{l} \times \hat{R} \text{ (T)}$$

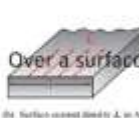
$$\vec{B}(\vec{R}) = \int_{V'} \frac{\mu \vec{J}(\vec{R}') \times (\vec{R} - \vec{R}')}{4\pi |\vec{R} - \vec{R}'|^3} dV' = \frac{\mu}{4\pi} \int_{V'} \frac{\vec{J} \times \vec{R}}{R^2} dV'$$



$$\text{Magnetic field of a loop } B = \frac{\mu I a^2}{2(a^2 + z^2)^{1.5}} \hat{z}$$

$$\text{At } z = 0 \quad B = \mu * \frac{I}{2 * a} \hat{z}$$

$$\vec{B}(\vec{R}) = \int_{S'} \frac{\mu \vec{J}_s(\vec{R}') \times (\vec{R} - \vec{R}')}{4\pi |\vec{R} - \vec{R}'|^3} ds' = \frac{\mu}{4\pi} \int_{S'} \frac{\vec{J}_s \times \vec{R}}{R^2} ds'$$



$$\text{At points far away } B = \frac{\mu I a^2}{2z^3} \hat{z}$$

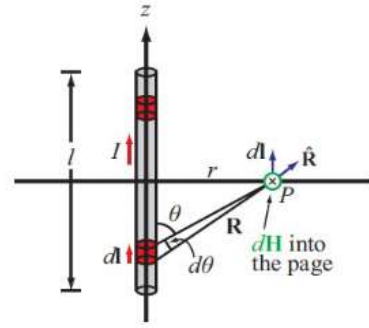
$$\vec{B}(\vec{R}) = \int_{l'} \frac{\mu I d\vec{l}' \times (\vec{R} - \vec{R}')}{4\pi |\vec{R} - \vec{R}'|^3} = \frac{\mu I}{4\pi} \int_{l'} \frac{d\vec{l}' \times \vec{R}}{R^2}$$

On a line

Magnetic field of a linear conductor  $B = \frac{\mu I a}{2\pi r \sqrt{4r^2 + a^2}} \hat{\phi}$

R is the distance from center to point P, a is length of wire segment

For an infinity long wire  $B = \frac{\mu I}{2\pi r} \hat{\phi}$



Ampere's law (Gauss's Law for magnetism) (net magnetic flux through a closed Gaussian surface is 0)

$$\oint_C \vec{H} \cdot d\vec{l} = I$$

$$\nabla \times \vec{H} = \vec{J}$$

H field for long wire:  $r_1 \leq a \rightarrow \vec{H} = \frac{r_1 * I}{2\pi a^2} \hat{\phi}$

$r_2 \geq a \rightarrow \vec{H} = \frac{I}{2\pi r} \hat{\phi}$

H field for toroidal coil:  $r < a \rightarrow \vec{H} = 0$

$r < a \rightarrow \vec{H} = 0$

$a < r < b \rightarrow \vec{H} = -\frac{NI}{2\pi r} \hat{\phi}$

H field inside long solenoid:  $r > a \rightarrow \vec{H} \approx 0$

$r < a \rightarrow \vec{H} = N * \frac{I}{L} \hat{z}$

H field of current sheet:  $z > 0 \rightarrow \vec{H} = -\frac{J}{2} \hat{y}$ , J is current density

$z < 0 \rightarrow \vec{H} = \frac{J}{2} \hat{y}$

Magnetic vector potential & Magnetic material

$\vec{B} = \nabla \times \vec{A}$  (Wb/m<sup>2</sup>)  $\rightarrow A$  is magnetic vector potential

$$\vec{A}(\vec{R}) = \int_{V'} \frac{\mu * \vec{J}(\vec{R}')}{4\pi |\vec{R} - \vec{R}'|} dV' \text{ (Wb/m}^2\text{)}$$

Spin Magnetic Moments:  $m = I * A = \frac{-e * v}{2\pi r} * \pi r^2 = -\frac{e * L}{2m}$

Angular momentum:  $L = m * v * r$

Magnetization:  $\vec{B} = \mu_0 \vec{H}$  in free space

$\vec{B} = \mu_0 \vec{H} + \mu_0 \vec{M}$  in magnetic material

$\vec{M} = X_m * \vec{H} \rightarrow X_m$  is magnetic susceptibility

M is sum of magnetic dipole moments in medium

$\vec{B} = \mu_0 \vec{H} + \mu_0 X_m * \vec{H} = \mu_0 (1 + X_m) \vec{H} \rightarrow (1 + X_m) = \mu_r$

	Diamagnetism	Paramagnetism	Ferrromagnetism
Permanent magnetic dipole moment	No	Yes, but weak	Yes, and strong
Primary magnetization mechanism	Electron orbital magnetic moment	Electron spin magnetic moment	Magnetized domains
Direction of induced magnetic field (relative to external field)	Opposite	Same	Hysteresis (see Fig. 5-22)
Common substances	Bismuth, copper, diamond, gold, lead, mercury, silver, silicon	Aluminum, calcium, chromium, magnesium, niobium, platinum, tungsten	Iron, nickel, cobalt
Typical value of $\chi_m$	$\approx -10^{-5}$	$\approx 10^{-5}$	$ \chi_m  \gg 1$ and hysteretic
Typical value of $\mu_r$	$\approx 1$	$\approx 1$	$ \mu_r  \gg 1$ and hysteretic

Inductor

$L = \frac{\text{Magnetic Flux}}{\text{Current}}$

$\phi_m = \int_S \vec{B} \cdot d\vec{s}$  (Wb)  $\rightarrow$  total magntic flux

Self inductance  $= L = \frac{\phi_m}{I} = \frac{\int_S \vec{B} \cdot d\vec{s}}{\oint_C \vec{H} \cdot d\vec{l}}$  (H or Wb/A)

Self-inductance in a solenoid:  $\phi_m = \frac{\mu * N * I}{L} * S \rightarrow S$  is area of one loop

$L = \frac{\mu * N^2}{L} * S$   $I = \frac{BL}{\mu N}$

Energy stored in solenoid:  $W = \frac{1}{2} * L * I^2 = \frac{1}{2} * \frac{B^2}{\mu}$

Energy density  $w = \frac{1}{2} \mu H^2$

Total energy in any volume  $W = \frac{1}{2} \int_V \mu H^2 dV$

Self-inductance of toroid:  $L = \frac{\mu * N^2}{2\pi r_m} * S \rightarrow r_m = \frac{a+b}{2}$

Mutual Inductance  $L_{12} = \frac{N_2}{I_1} \int_{S_2} \vec{B}_1 \cdot d\vec{s}$

Potential energy  $W = \int_0^I i * v * dt = L \int_0^I i * di = \frac{1}{2} LI^2$  (J)

## Faraday's Law

- A time varying magnetic field creates transformer emf  $V$
- A moving loop with time-varying surface area in static field  $B$  create motional emf
- A moving loop in time-varying field  $B$  is transformer emf + motional emf

$$V_{emf} = -N \frac{d\phi_m}{dt} = -N \frac{d}{dt} \int_S \vec{B} \cdot \vec{ds} \quad (V)$$

$$\text{Transformer } V_{emf} = -N \int_S \frac{d\vec{B}}{dt} \cdot \vec{ds}$$

$$\oint_C \vec{E} \cdot \vec{dl} = - \int_S \frac{\partial \vec{B}}{\partial t} \cdot \vec{ds}$$

$$\nabla \times \vec{E} = \frac{-\partial \vec{B}}{\partial t}$$

## Lenz's Law

- The current in the loop is always in a direction that opposes the change of magnetic flux that produced it

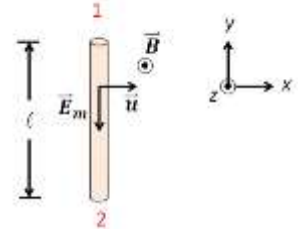
## Moving Conductor in a static magnetic field

$$\text{Motional EMF} = V_{12} = \int_2^1 (\vec{u} \times \vec{B}) \cdot \vec{dl}$$

$$\vec{u} \times \vec{B} = u\hat{x} \times B_0\hat{z} = -uB_0\hat{y}$$

$$V_{emf} = -uB_0L$$

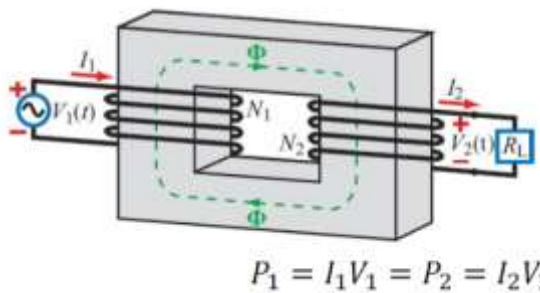
$$V_{emf} = - \int_S \frac{d\vec{B}}{dt} \cdot \vec{ds} + \int_2^1 (\vec{u} \times \vec{B}) \cdot \vec{dl}$$



## Transformers

$$\frac{V_1}{V_2} = \frac{N_1}{N_2} \rightarrow \text{turns ratio} \rightarrow N_2 > N_1 \text{ Step up transformer; } N_2 < N_1 \text{ "Step down transformer"}$$

$$\frac{I_1}{I_2} = \frac{N_2}{N_1}$$



$$R_{in} = \frac{V_1}{I_1} = \left(\frac{N_1}{N_2}\right)^2 * R_L$$

$$Z_{in} = \left(\frac{N_1}{N_2}\right)^2 * Z_L$$

$$V_{emf} = - \frac{d\phi_m}{dt} = A * w * B_0 * \sin(\omega t + C_0)$$

DC Generators: Same components as AC generator, main difference is contacts to the rotating loop are made using a split ring called commutator.

**Motor:** Electrical to mechanical energy      **Generators** works opposite

Motors are devices into which energy is transferred by electrical transmission while energy is transferred out by work

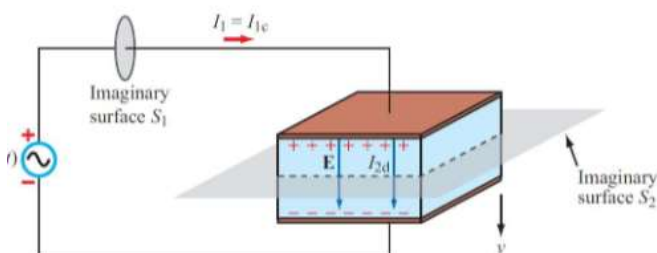
EM Motor: A current is supplied to the coil by a battery and the torque acting on the current carrying coil causes it to rotate

- Induced back emf, acts to reduce the current in the coil
- The back emf increases in magnitude as the rotational speed of coil increases
- $I = \frac{V_{app} - V_{emf}}{R}$

## The displacement current

$$\nabla \times \vec{H} = \vec{J} + \frac{\partial \vec{D}}{\partial t} \rightarrow \frac{\partial \vec{D}}{\partial t} = \vec{J}_d \rightarrow \text{displacement current density}$$

$$\oint_C \vec{H} \cdot \vec{dl} = \int_S \vec{J} \cdot \vec{ds}' + \int_S \frac{\partial \vec{D}}{\partial t} \cdot \vec{ds}' = I_c(\text{conduction current}) + I_d(\text{Displacement current}) = I$$



$$\text{In perfect conducting wire: } I_1 = I_{1c} + I_{1d} = -CV_0 \omega \sin(\omega t)$$

$$\text{In perfect conducting capacitor:}$$

$$I_2 = I_{2c} + I_{2d} = -CV_0 \omega \sin(\omega t) = I_1$$

Continuity of current flow through the circuit



- The displacement current behaves like a real current
- The displacement current accounts for polarization in the medium
- The perfect wire has infinite conductivity
  - o If it has finite conductivity, then  $D$  in the wire would be non-zero and  $I_1$  would consist of both conduction and displacement currents
- A magnetic field can be produced either by currents or by changing electric fields

$$\nabla \cdot (\nabla \times \mathbf{H}) = \nabla \cdot \mathbf{J} + \frac{\partial \nabla D}{\partial t} = 0$$

$$\text{Continuity equation: } \nabla \cdot \vec{J} = \frac{-\partial \rho}{\partial t}$$

$$\oint_S \vec{J} \cdot d\mathbf{s}' = 0 \text{ for steady currents} = \sum_i I_i = 0$$

- An electric field can be produced either by charges or changing magnetic fields