Induced Dirac-Schrödinger operators on S^1 -semi-free quotients

Disputation der Dissertation von **Juan Camilo Orduz Barrera** zur Erlangung des akademischen Grades Dr. rer. nat.

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The induced operator Dirac-Schrödinger operator

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Definition of \mathscr{D} Example revisited
Local description of the potential
Index computation (Witt case)
Index computation (non-Witt case)?

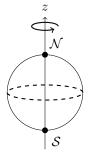
$$y = \frac{45}{1 + e^{-40(x - 0.1)}} + 2$$

$$x = -\frac{1}{40}\log\left(\frac{45}{y-2} - 1\right) + 0.1$$

Example: What is a semi-free action?

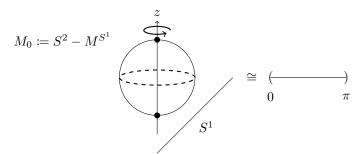
Let S^1 act on $M = S^2 \subset \mathbb{R}^3$ by rotations around the z-axis.

- $ightharpoonup M^{S^1} := \{\mathcal{N}, \mathcal{S}\} \text{ fixed point set.}$
- ▶ On the principal orbit $M_0 := S^2 M^{S^1}$, the action is free.

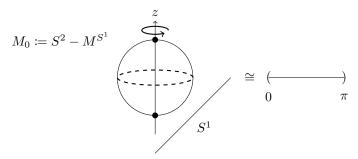


For a semi-free action the isotropy groups $S_x^1 := \{g \in S^1 \mid gx = x\}$ must be either $\{1\}$ or S^1 for all $x \in S^2$.

As the action on M_0 is free, the quotient space $M_0/S^1=(0,\pi)=:I$ is a smooth manifold.

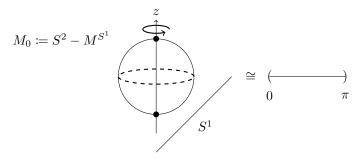


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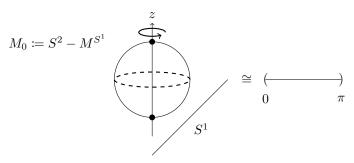
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- Equip S^2 with the round metric $g^{TS^2} = d\theta^2 + \sin^2\theta d\phi^2$.
- ▶ The quotient metric $g^{TI} := d\theta^2$ is incomplete.
- ► The Hodge de Rham operator

$$D_I := \left(\begin{array}{cc} 0 & -\partial_\theta \\ \partial_\theta & 0 \end{array} \right)$$

defined on $\Omega_c(I)$, is not essentially self-adjoint in $L^2(\wedge^*I)$.

- ▶ (M, g^{TM}) : 4k + 1-dimensional closed oriented Riemannian manifold on which S^1 acts by orientation preserving isometries.
- ▶ M^{S^1} : fixed point set and M_0 the principal orbit.
- ightharpoonup V: generating vector field of the action.

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- $\bullet \ \alpha := V^{\flat}/\|V\|^2 \in \Omega^1(M_0) \text{ satisfies } \alpha(V) = 1.$

Consider the differential complex

$$\Omega_{\mathrm{bas},c}(M_0) \coloneqq \{\omega \in \Omega_c(M_0) \mid L_V \omega = 0 \text{ and } \iota_V \omega = 0\}.$$

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Define $\sigma_{S^1}(M)$ to be the signature of the symmetric quadratic form,

$$H^{2k}_{\mathrm{bas},c}(M_0) \times H^{2k}_{\mathrm{bas},c}(M_0) \longrightarrow \mathbb{R}$$

$$([\omega], [\omega']) \longmapsto \int_M \alpha \wedge \omega \wedge \omega'.$$

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- ▶ It does not depend on the Riemannian metric.
- \blacktriangleright It is invariant under S^1 -homotopy equivalences.



$\sigma_{S^1}(M)$ formula for semi-free actions

Theorem (Lott, 00')

Suppose S^1 acts effectively and semi-freely on M, then

$$\sigma_{S^1}(M) = \int_{M_0/S^1} L(T(M_0/S^1), g^{T(M_0/S^1)}) + \eta(M^{S^1}).$$

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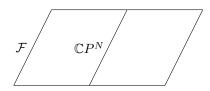
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If the codimension of ${\cal M}^{S^1}$ in ${\cal M}$ is divisible by four we call ${\cal M}/S^1$ a Witt space and

- $\eta(M^{S^1}) = 0.$
- ▶ $L(T(M_0/S^1), g^{T(M_0/S^1)})$ represents the homology L-class of M/S^1 .
- ▶ $\sigma_{S^1}(M)$ equals the intersection homology signature of M/S^1 .

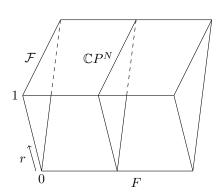
Strategy of Lott's proof

- ▶ $F \subset M^{S^1}$: connected closed 4k 2N 1 dimensional manifold.
- ▶ Let $NF \longrightarrow F$ be the normal bundle of F.
- ▶ SNF/S^1 is the total space of a Riemannian $\mathbb{C}P^N$ -fiber bundle \mathcal{F} .
- ▶ Model M/S^1 as the mapping cylinder $C(\pi_{\mathcal{F}}: \mathcal{F} \longrightarrow F)$.
- For r > 0 small enough $\sigma_{S^1}(M) = \sigma(M/S^1 N_r(F))$.
- ▶ Study the limit of the APS signature theorem as $r \longrightarrow 0$.
 - ▶ Use Dai's formula for the adiabatic limit of the eta invariant.
 - ▶ Prove that the form $\tilde{\eta}$ and Dai's tau invariant $\tau_{\mathcal{F}}$ vanish.



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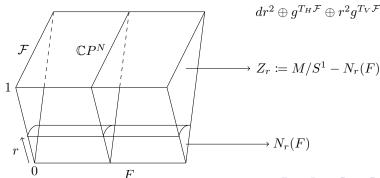


 $dr^2 \oplus q^{T_H \mathcal{F}} \oplus r^2 q^{T_V \mathcal{F}}$



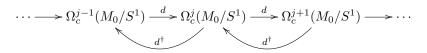
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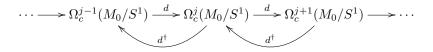
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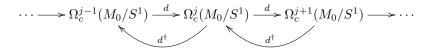
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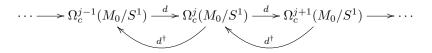


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- $\star: \wedge^j T^*(M_0/S^1) \longrightarrow \wedge^{4k-j} T^*(M_0/S^1)$ chirality operator.

$$\star^2 = 1$$
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 $ightharpoonup \operatorname{ind}(D^+) = ? \text{ where } D^+ : \Omega_c^+(M_0/S^1) \longrightarrow \Omega_c^-(M_0/S^1).$



Close to the fixed point set:

$$D = \gamma \left(\frac{\partial}{\partial r} + \bigstar \otimes A(r) \right) \Longrightarrow D^+ = \frac{\partial}{\partial r} + A(r),$$

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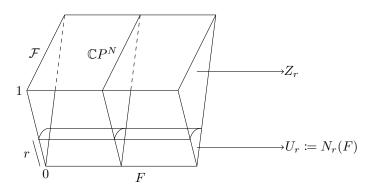
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- If N odd (Witt case) this is the case.
- ▶ If N even (non-Witt case) \Longrightarrow need boundary conditions.



Index formula for the Witt case (Brüning 09')



Use the Dirac-Systems formalism (Ballmann, Brüning & Carron, 08').

▶ In particular the gluing index formula:

$$\operatorname{ind}(D^+) = \operatorname{ind}\left(D^+_{Z_r, Q_{<}(A(r))(H)}\right) + \operatorname{ind}\left(D^+_{U_r, Q_{\geq}(A(r))(H)}\right).$$

1. Prove that for r > 0 small enough ind $\left(D_{U_r,Q_{\geq}(A(r))(H)}^+\right) = 0$. For the proof we require $A_H(r)A_V + A_VA_H(r)$ to be a first order vertical differential operator.

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- 2. Take the limit $r \longrightarrow 0^+$ of the signature of the manifold with boundary Z_r .

Note however that ind $\left(D_{Z_r,Q_<(A(r))(H)}^+\right)$ does not have the right APS boundary condition, $Q_<(A_0(r))(H)$, where

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The correction term is, as $r \longrightarrow 0^+$,

$$\begin{split} \operatorname{ind}(Q_{<}(A(r))(H), Q_{\geq}(A_0(r))(H)) &= \sigma_{(2)}(T_{\pi}) + \dim(\ker A_0(r))/2 \\ &= \tau_{\mathcal{F}} + \dim(\ker A_0(r))/2 \\ &= \dim(\ker A_0(r))/2. \end{split}$$

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3. Finally,

$$\lim_{r \to 0^+} \operatorname{ind} \left(D_{Z_r, Q_{<}(A_0(r))(H)}^+ \right) + \frac{1}{2} \dim(\ker A_0(r)) = \sigma_{S^1}(M)$$

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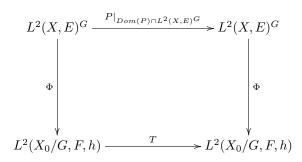
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Theorem (Brüning & Heintze, 79')



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Let $\kappa := -d \log(||V||) \in \Omega^1_{\text{bas}}(M_0)$ be the mean curvature form. Using Rummler's formula $\varphi_0 := d\chi + \kappa \wedge \chi \in \Omega^2_{\text{bas}}(M_0)$ one verifies

$$T = \left(\begin{array}{cc} D & 0 \\ 0 & D \end{array} \right) + \left(\begin{array}{cc} \iota_{\bar{\kappa}^\sharp} & \varepsilon(\bar{\varphi}_0 \wedge) \\ \varepsilon(\bar{\varphi}_0 \wedge)^\dagger & -\bar{\kappa} \wedge \end{array} \right)$$

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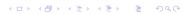
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We conjugate by multiplication by $U := h^{-1/2} = ||V||^{-1/2}$,

$$U^{-1}TU = \left(\begin{array}{cc} D & 0 \\ 0 & D \end{array} \right) + \left(\begin{array}{cc} \frac{1}{2}\hat{c}(\bar{\kappa}) & \varepsilon(\bar{\varphi}_0 \wedge) \\ \varepsilon(\bar{\varphi}_0 \wedge)^{\dagger} & -\frac{1}{2}\hat{c}(\bar{\kappa}) \end{array} \right)$$

where $\hat{c}(\bar{\kappa}) := \bar{\kappa} \wedge + \iota_{\bar{\kappa}^{\sharp}}$.



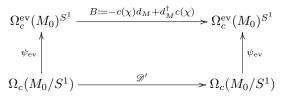


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- ▶ It is explicitly given by

$$\mathscr{D}' = D + \frac{1}{2}c(\bar{\kappa})\varepsilon - \frac{1}{2}\hat{c}(\bar{\varphi}_0)(1-\varepsilon)$$

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▶ It is enough to consider

$$\mathscr{D} \coloneqq D + \frac{1}{2} c(\bar{\kappa}) \varepsilon.$$



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$$= \gamma \left(\partial_\theta + \cot \theta \begin{pmatrix} -1/2 & 0 \\ 0 & 1/2 \end{pmatrix} \right), \text{ where } \gamma := \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}.$$

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▶ The spectrum of the cone coefficient satisfies

spec
$$\begin{pmatrix} -1/2 & 0 \\ 0 & 1/2 \end{pmatrix} \cap (-1/2, 1/2) = \emptyset.$$

Thus, \mathcal{D} is essentially self-adjoint.



Local description of the potential

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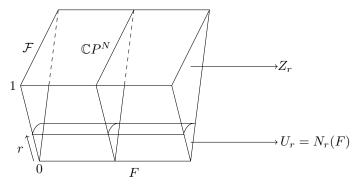
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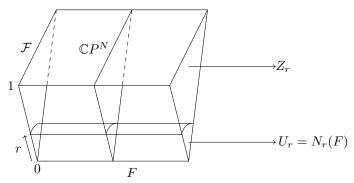
$$2j - N \pm \frac{1}{2} \not\in \left(-\frac{1}{2}, \frac{1}{2}\right),\,$$

thus we see \mathcal{D} is indeed essentially self-adjoint.

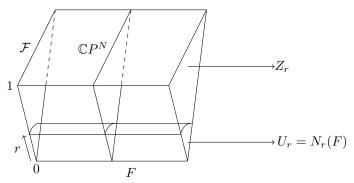




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- ▶ From a variation of Brüning's method we can prove

$$\lim_{r \to 0^+} \operatorname{ind} \left(\mathscr{D}^+_{Z_t, Q_{<}(A(r))(H)} \right) = \sigma_{S^1}(M).$$

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- ► For the Witt case:
 - As $r \to 0^+$, $\operatorname{ind}(Q_{<}(A(r))(H), Q_{\geq}(A_0(r))(H)) = \sigma_{(2)}(T_{\pi}) + \frac{1}{2} \dim(\ker A_0(r)).$
 - (Cheeger-Dai) $\sigma_{(2)}(\overline{T_{\pi}}) = \tau_{\mathcal{F}}$. For non-Witt spaces there is a similar formula developed by Hunsicker & Mazzeo for the L^2 -signature on the image $\mathcal{H}_{\rm rel} \longrightarrow \mathcal{H}_{\rm abs}$.

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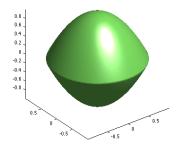
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- Study the case of the spin Dirac operator (Studied in the Witt case by Albin and Gell-Redman).





Thank you! Vielen Dank! Gracias! Gràcies!