

## Constantes

$$k_B = 1.381 \times 10^{-23} JK^{-1} = 8.62 \times 10^{-5} eVK^{-1}$$

$$m_e = 9.11 \times 10^{-31} kg = 0.511 MeVc^{-2}$$

$$m_p = 1.67 \times 10^{-27} kg = 938 MeVc^{-2}$$

$$\varepsilon_0 = \frac{1}{4\pi K} = 8.85 \times 10^{-12} Fm^{-1}$$

$$\hbar = 1.055 \times 10^{-34} Js = 6.58 \times 10^{-16} eVs$$

$$e = 1.602 \times 10^{-19} C$$

$$\text{Fermions: } e^-, p, n \quad (n_F(x) = \frac{1}{e^x + 1})$$

$$\text{Bosons: phonon, photon} \quad (n_B(x) = \frac{1}{e^x - 1})$$

$$n \sim 10^{22} cm^{-3}; \tau \sim 10^{-15} s; v \sim 10^{-5} \frac{m}{s}$$

## 1 Estructura cristalina

### 1.1 Redes de Bravais

<i>a</i>	triclínica
<i>m</i>	monoclínica
<i>o</i>	ortorómbica
<i>t</i>	tetragonal
<i>h</i>	hexagonal
<i>c</i>	cúbica

<i>P</i>	Primitiva
<i>S</i>	Centrada en una cara
<i>I</i>	Centrada en el cuerpo
<i>R</i>	Centrada romboidal
<i>F</i>	Centrada en las caras

14 posibles redes de Bravais

Tric.	Monoc.	Ortor.	Tetra.	Hex.	Cúbico
<i>aP</i>	<i>mP, mS</i>	<i>oP, oS, oF, oI</i>	<i>tP, tI</i>	<i>hP, hR</i>	<i>cP, cF, cI</i>

### 1.2 Cosas

Base dual y matriz métrica

$$a^* = \frac{b \times c}{V}, \quad b^* = \frac{c \times a}{V}, \quad c^* = \frac{a \times b}{V}, \quad V = \det(\bar{a}, \bar{b}, \bar{c})$$

$$(\bar{a}^*, \bar{b}^*, \bar{c}^*) = \left( \begin{pmatrix} \bar{a}^T \\ \bar{b}^T \\ \bar{c}^T \end{pmatrix} \right)^{-1}, \quad G = \begin{pmatrix} a \cdot a & a \cdot b & a \cdot c \\ b \cdot a & b \cdot b & b \cdot c \\ c \cdot a & c \cdot b & c \cdot c \end{pmatrix}, \quad G^* = G^{-1}$$

Cambio de base

$$(\bar{a}', \bar{b}', \bar{c}') = (\bar{a}, \bar{b}, \bar{c})P, \quad \begin{pmatrix} x' \\ y' \\ z' \end{pmatrix} = P^{-1} \begin{pmatrix} x \\ y \\ z \end{pmatrix}$$

$$(x, y, z) = (x^*, y^*, z^*)P, \quad \begin{pmatrix} a'^* \\ b'^* \\ z'^* \end{pmatrix} = P^{-1} \begin{pmatrix} a^* \\ b^* \\ c^* \end{pmatrix}$$

$$\text{Distancia interplanar } g_{hkl} = \frac{2}{d_{hkl}}; \quad g_{hkl}^2 = (hkl)G^* \begin{pmatrix} h \\ k \\ l \end{pmatrix}$$

$$\text{Transferencia de momento } Q = \frac{4\pi \sin \theta}{\lambda}$$

$$\text{Condiciones de Laue } \bar{Q} = 2\pi \bar{g}_{hkl}$$

$$\text{Ley de Bragg } g_{hkl} = \frac{2 \sin \theta_{hkl}}{\lambda}$$

$$\text{Módulo de Young } \nu_s = \sqrt{\frac{\gamma}{\rho}}$$

$$\text{Factor de estructura } F_{hkl} = \sum_p f_p e^{-i2\pi \bar{g}_{hkl} \cdot \bar{r}_p}; I \propto |F_{hkl}|^2$$

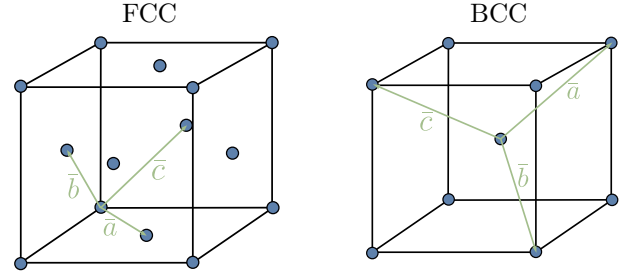
### 1.3 Estructuras comunes

FCC (primitiva volumen 1/4)

$$\begin{cases} \bar{a} = \frac{1}{2}(1 \ 1 \ 0) \\ \bar{b} = \frac{1}{2}(0 \ 1 \ 1) \\ \bar{c} = \frac{1}{2}(1 \ 0 \ 1) \end{cases} \quad \begin{cases} \bar{a}^* = (1 \ 1 \ -1) \\ \bar{b}^* = (-1 \ 1 \ 1) \\ \bar{c}^* = (1 \ -1 \ 1) \end{cases}$$

BCC (primitiva volumen 1/2)

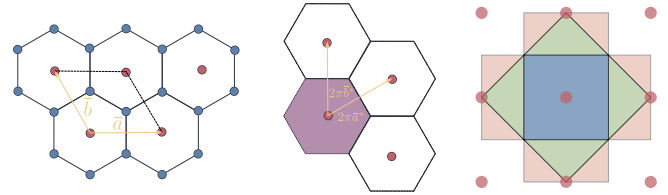
$$\begin{cases} \bar{a} = \frac{1}{2}(1 \ 1 \ -1) \\ \bar{b} = \frac{1}{2}(-1 \ 1 \ 1) \\ \bar{c} = \frac{1}{2}(1 \ -1 \ 1) \end{cases} \quad \begin{cases} \bar{a}^* = (1 \ 1 \ 0) \\ \bar{b}^* = (0 \ 1 \ 1) \\ \bar{c}^* = (1 \ 0 \ 1) \end{cases}$$



Hexagonal

$$\begin{cases} \bar{a} = (1, 0) \\ \bar{b} = (-\frac{1}{2}, \frac{\sqrt{3}}{2}) \end{cases} \quad \begin{cases} \bar{a}^* = \frac{2\sqrt{3}}{3}(\frac{\sqrt{3}}{2}, \frac{1}{2}) \\ \bar{b}^* = \frac{2\sqrt{3}}{3}(0, 1) \end{cases}$$

$$G = \begin{pmatrix} a^2 & -\frac{a^2}{2} & 0 \\ -\frac{a^2}{2} & a^2 & 0 \\ 0 & 0 & c^2 \end{pmatrix}, \quad G^* = \begin{pmatrix} \frac{4}{3a^2} & \frac{2}{3a^2} & 0 \\ \frac{2}{3a^2} & \frac{4}{3a^2} & 0 \\ 0 & 0 & \frac{1}{c^2} \end{pmatrix}$$



En una hcp  $c = 1.633a$

### 1.4 Grupos

$$m_{100} = \begin{pmatrix} -1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}; n_{001} = \begin{pmatrix} \cos\left(\frac{360}{n}\right) & -\sin\left(\frac{360}{n}\right) & 0 \\ \sin\left(\frac{360}{n}\right) & \cos\left(\frac{360}{n}\right) & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

Cambio de base a  $\mathcal{B} = \{\bar{u}, \bar{v}, \bar{w}\}$

$$M_{\mathcal{C}} = M_{\mathcal{B} \rightarrow \mathcal{C}} M_{\mathcal{B}}^{-1} M_{\mathcal{B} \rightarrow \mathcal{C}}, \quad M_{\mathcal{B} \rightarrow \mathcal{C}} = (\bar{u}, \bar{v}, \bar{w})$$

Reflexión vector director  $(a, b, c)$

$$M = \frac{1}{a^2 + b^2 + c^2} \begin{pmatrix} -a^2 + b^2 + c^2 & -2ab & -2ac \\ -2ab & a^2 - b^2 + c^2 & -2bc \\ -2ac & -2bc & a^2 + b^2 - c^2 \end{pmatrix}$$

Rotación respecto  $\hat{u} = (u_x, u_y, u_z)$  ( $c = \cos \theta$ ,  $s = \sin \theta$ ).  
 $R =$

$$\begin{pmatrix} c + u_x^2(1-c) & u_x u_y(1-c) - u_z s & u_x u_z(1-c) + u_y s \\ u_y u_x(1-c) + u_z s & c + u_y^2(1-c) & u_y u_z(1-c) - u_x s \\ u_z u_x(1-c) - u_y s & u_z u_y(1-c) + u_x s & c + u_z^2(1-c) \end{pmatrix}$$

Centrosimétricos  $(x, y, z) \rightarrow (-x, -y, -z)$  no tienen polarización espontánea

## 2 Dinámica de cristales

### 2.1 Densidad de estados

$$\bar{k} = \left( \frac{2\pi}{L}n \quad \frac{2\pi}{L}m \quad \frac{2\pi}{L}l \right) \quad \forall n, m, l \in \mathbb{Z}$$

Número de estados hasta  $k$

$$N(k) = \int_{(\frac{2\pi}{L})^2(n^2+m^2+l^2) \leq k^2} dV = \frac{L^3}{6\pi^2} k^3 = \frac{V}{6\pi^2} k^3$$

1, 2 y 3 dimensiones respectivamente (y se cumple  $\omega = \nu_s k$ )

$$\begin{cases} g(k) = \frac{L}{\pi} & \begin{cases} g(k) = \frac{L^2}{2\pi} k \\ g(\omega) = \frac{L}{\pi\nu} \end{cases} & \begin{cases} g(k) = \frac{V}{2\pi^2} k^2 \\ g(\omega) = \frac{V}{2\pi^2 \nu_s^3} \omega^2 \end{cases} \end{cases}$$

### 2.2 Dispersión

Oscilador con masa  $m$  y constante  $k_s$

$$F_n = m\ddot{x}_n = k_s(x_{n+1} + x_{n-1} - 2x_n)$$

$$\begin{aligned} -m\omega^2 A e^{i(kna - \omega t)} &= k_s A e^{i(kna - \omega t)} (e^{ika} + e^{-ika} - 2) = \\ &= -4k_s \sin^2\left(\frac{ka}{2}\right) \Rightarrow \boxed{\omega = 2\sqrt{\frac{k_s}{m}} \left| \sin\left(\frac{ka}{2}\right) \right|} \end{aligned}$$

Oscilador con masa  $m$  y constantes alternadas  $k_1, k_2$

$$\begin{cases} m\ddot{x}_n = k_1(y_{n-1} - x_n) + k_2(y_n - x_n) \\ m\ddot{y}_n = k_1(x_{n+1} - y_n) + k_2(x_n - y_n) \end{cases}$$

Ansatz

$$x_n = A e^{i(kna - \omega t)} \quad y_n = B e^{i(kna - \omega t)}$$

Ecuaciones

$$\begin{cases} -m\omega^2 A = -A(k_1 + k_2) + B(k_1 e^{ika} + k_2) \\ -m\omega^2 B = -A(k_1 e^{ika} + k_2) + B(-k_1 - k_2) \end{cases}$$

Forma matricial

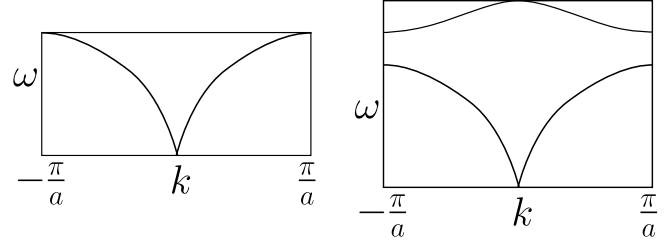
$$m\omega^2 \begin{pmatrix} A \\ B \end{pmatrix} = \begin{pmatrix} (k_1 + k_2) & -k_2 - k_1 e^{ika} \\ -k_2 - k_1 e^{ika} & (k_1 + k_2) \end{pmatrix} \begin{pmatrix} A \\ B \end{pmatrix} = K \begin{pmatrix} A \\ B \end{pmatrix}$$

$$0 = \det(K - m\omega^2 I) = |(k_1 + k_2) - m\omega^2|^2 - |k_2 + k_1 e^{ika}|^2$$

$$\boxed{\omega_{\pm}(k) = \sqrt{\frac{k_1 + k_2}{m}} \pm \frac{1}{m} \sqrt{(k_1 + k_2)^2 - 4k_1 k_2 \sin^2(ka/2)}}$$

Si  $m_1 \neq m_2$  y  $k_s$  es la misma, sea  $K_i = \frac{k}{m_i}$ , entonces

$$\omega_{\pm}(k) = \sqrt{(K_1 + K_2) \pm \sqrt{(K_1 + K_2)^2 - 4K_1 K_2 \sin^2(ka/2)}}$$



Si hay  $N$  átomos / celda:  $3N$  ramas:

- 3 acústicas (2 trans. < 1 long.)
- $3N - 3$  ópticas

### 2.3 Modelo de Einstein

$$E_n = \hbar\omega\left(n + \frac{1}{2}\right) \Rightarrow Z_1 = \frac{1}{2 \sinh\left(\frac{\beta\hbar\omega}{2}\right)}$$

$$\langle E_1 \rangle = -\frac{\partial}{\partial \beta} \ln Z_1 = \frac{\hbar\omega}{2} \coth\left(\frac{\beta\hbar\omega}{2}\right)$$

Energía y capacidad calorífica

$$\langle E \rangle = \frac{3}{2} N \hbar\omega \coth\left(\frac{\beta\hbar\omega}{2}\right)$$

$$C_v = \frac{\partial \langle E \rangle}{\partial T} = 3N k_B (\beta\hbar\omega)^2 \frac{e^{\beta\hbar\omega}}{(e^{\beta\hbar\omega} - 1)^2}$$

Definimos ahora  $T_E = \frac{\hbar\omega_E}{k_B}$ . En los límites

- Si  $T \gg T_E \Rightarrow C_v = 3N k_B$
- Si  $T \ll T_E \Rightarrow C_v = 3N k_B \left(\frac{T_E}{T}\right)^2 \frac{1}{\sinh^2\left(\frac{T_E}{2T}\right)}$

### 2.4 Modelo de Debye

Aproximamos la ecuación de dispersión para  $k$  baja como  $\omega = \nu k$

$$3N = \int_0^{\omega_D} 3g(\omega) d\omega = \frac{V}{2\pi^2 \nu^3} \omega_D^3 \Rightarrow \boxed{\omega_D = \sqrt[3]{\frac{6\pi^2 \nu^3 N}{V}}}$$

donde hemos contado cada partícula y cada estado 3 veces y hemos usado

$$\omega = \nu k, \quad g(k) = \frac{V}{2\pi^2} k^2, \quad g(\omega) = \frac{V}{2\pi^2 \nu^3} \omega^2 = 9N \frac{\omega^2}{\omega_D^3}$$

La energía y la capacidad calorífica

$$\begin{aligned} \langle E \rangle &= \int_0^{\omega_D} \hbar\omega 3g(\omega) \left( \frac{1}{e^{\beta\hbar\omega} - 1} + \frac{1}{2} \right) d\omega = \\ &= E_0 + \frac{3V\hbar}{2\pi^2 \nu^3} \int_0^{\omega_D} \frac{\hbar\omega^3}{e^{\beta\hbar\omega} - 1} d\omega \quad \left( x = \frac{\hbar\omega}{k_B T} \right) \end{aligned}$$

$$T_D := \frac{\hbar\omega_D}{k_B} \Rightarrow \boxed{\langle E \rangle = \frac{3V k_B^4 T^4}{2\pi^2 \nu^3 \hbar^3} \int_0^{\frac{T_D}{T}} \frac{x^3}{e^x - 1} dx}$$

La capacidad calorífica  $C_v = \frac{\partial \langle E \rangle}{\partial T}$  en los extremos:

- Si  $T \gg T_D \Rightarrow \langle E \rangle \sim 3Nk_B T \Rightarrow C_v \sim 3Nk_B$
- Si  $T \ll T_D \Rightarrow \langle E \rangle \sim \frac{3\pi^4 Nk_B T^4}{5T_D^3} \Rightarrow C_v \sim \frac{12\pi^4}{5} Nk_B \left(\frac{T}{T_D}\right)^3$

### 3 No se, cuanticocosas

#### 3.1 Drude model

$$n = \frac{N}{V}; \quad \frac{dp}{dt} = -\frac{p}{\tau} + F; \quad \bar{j} = -ne\bar{v} = \sigma \bar{E}$$

$$\sigma_0 = \frac{e^2 \tau n}{m}; \quad \text{si } F = e \text{Re}[E_0 e^{i\omega t}] \Rightarrow \sigma(\omega) = \frac{\sigma_0}{1 - i\omega\tau}$$

$$mv = p = -e\tau E; \quad R_H = \frac{-1}{ne} = \frac{\rho_{yx}}{|B|}$$

$$\bar{E} = \tilde{\rho}\bar{j}; \quad \rho_{xx} = \rho_{yy} = \rho_{zz} = \frac{m}{ne^2\tau}; \quad \frac{1}{2}mv_0^2 = \frac{3}{2}k_B T$$

(COSAS DEL EFECTO HALL PROBLEMA 4)

Hall resistivity  $\rho_{xy} = -\rho_{yx} = \frac{B}{ne} \quad (\bar{B} \propto \hat{z})$

Peltier coefficient  $\Pi = -\frac{k_B T}{2e} = \frac{-c_v T}{3e}$

Seebeck coefficient  $S = \frac{\Pi}{T}$

$$\langle v \rangle_{gasid.} = \sqrt{\frac{8k_B T}{\pi m}}; \quad \kappa = \frac{1}{3}nc\langle v \rangle^2\tau = \frac{4}{\pi} \frac{n\tau k_B^2 T}{m}$$

#### 3.2 Capacidad calorífica

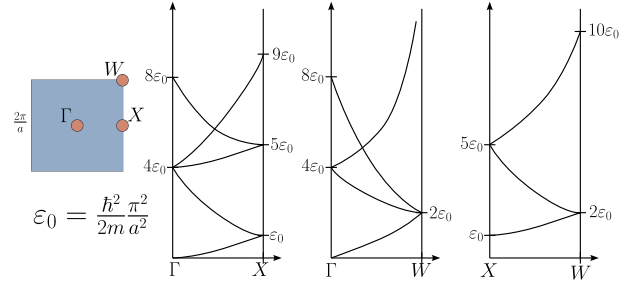
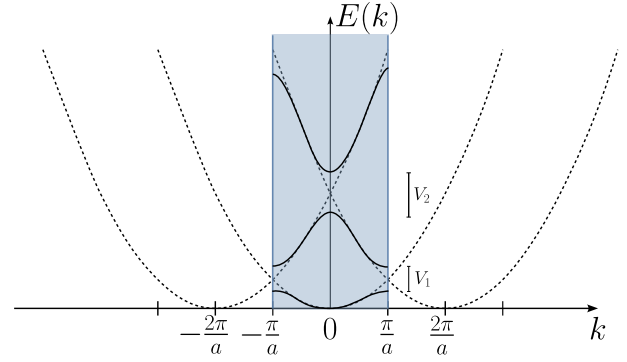
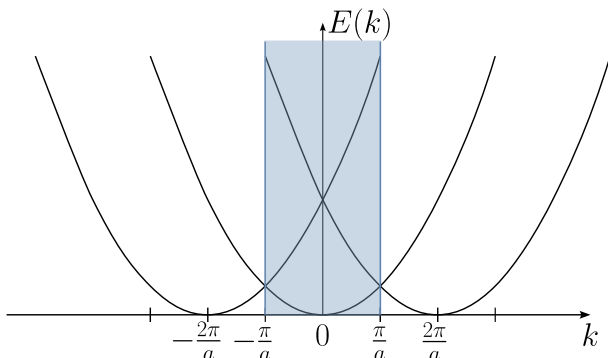
$g = \frac{\text{densidad de estados}}{V}$

$$g(\varepsilon) = \frac{3n}{2(E_F)^{\frac{3}{2}}} \varepsilon^{\frac{1}{2}} = \frac{(2m)^{\frac{3}{2}}}{2\pi^2 \hbar^3} \varepsilon^{\frac{1}{2}}, \quad k = \sqrt{\frac{2\varepsilon m}{\hbar^2}}$$

$$n = \int_0^\infty d\varepsilon g(\varepsilon) n_F(\beta(\varepsilon - \mu)), \quad \frac{E_T}{V} = \int_0^\infty d\varepsilon \varepsilon g(\varepsilon) n_F(\beta(\varepsilon - \mu))$$

$$C = \frac{\pi^2}{3} \left( \frac{3Nk_B}{2} \right) \left( \frac{T}{T_F} \right)$$

$$\bar{M} = g(E_F) \mu_B^2 \bar{B}; \quad \mu_B = 0.67 \left( \frac{K}{\text{Tesla}} \right) k_B$$



**Teorema de Bloch** ( $V(\bar{r})$  periódico)

$$\psi_{\bar{k}}(\bar{r}) = u_{\bar{k}}(\bar{r}) e^{i\bar{k} \cdot \bar{r}}, \quad E(\bar{k}) = E(\bar{k} + \bar{G})$$

(1D) Fourier del potencial de dos formas

$$V(x) = V_0 + \sum_{j=1}^{\infty} V_j \cos\left(\frac{2\pi j}{a} x\right) \quad \text{ó} \quad V(x) = \sum_{j=-\infty}^{\infty} V_{\frac{2\pi j}{a}} e^{i\frac{2\pi j}{a} x}$$

Con las relaciones  $V_j = 2V_{\frac{2\pi j}{a}}$ , y donde los coeficientes son

$$V_j = \frac{2}{a} \int_{-\frac{a}{2}}^{\frac{a}{2}} dx V(x) \cos\left(\frac{2\pi j}{a} x\right); \quad V_{\frac{2\pi j}{a}} = \frac{1}{a} \int_{-\frac{a}{2}}^{\frac{a}{2}} dx V(x) e^{-i\frac{2\pi j}{a} x}$$

Gas de electrones libres

$$\bar{k} = \frac{2\pi}{L}(n_x, n_y, n_z), \quad E(\bar{k}) = \frac{\hbar^2}{2m} |\bar{k}|^2, \quad n_F(x) = \frac{1}{e^x + 1}$$

$$N = 2 \sum_{\bar{k}} n_F(\beta(E(\bar{k}) - \mu)) = 2 \frac{V}{(2\pi)^3} \int d\bar{k} n_F(\beta(E(\bar{k}) - \mu))$$

Fermi energy ( $E_F = \mu(T \rightarrow 0)$ ) (d numero de dimensiones)

$$\varepsilon_F = \frac{\hbar^2 k_F^2}{2m} = k_B T_F; \quad p_F = \hbar k_F; \quad U_T = \frac{3}{5} \varepsilon_F N$$

$$N = 2 \frac{V}{(2\pi)^d} \int_{|k| < k_F} dk \Rightarrow k_F = (3\pi^2 n)^{\frac{1}{3}}, \quad \varepsilon_F = \frac{\hbar^2 (3\pi^2 n)^{\frac{2}{3}}}{2m}$$

Considerando los dos espines (multiplicamos por 2)

$$N_T = 2 \cdot \left( \frac{4}{3} \pi (n_x^2 + n_y^2 + n_z^2)^{3/2} \right) \Rightarrow k_{max}^2 = k_F^2 = (3n\pi^2)^{\frac{2}{3}}$$

Electrones casi-libres

$$\psi_+ \sim \cos\left(\pi \frac{x}{a}\right), \quad \psi_- \sim \sin\left(\pi \frac{x}{a}\right)$$

$$E^\pm = \frac{1}{2} (E_{\bar{k}-\bar{G}}^0 + E_{\bar{k}}^0) \pm \sqrt{\frac{1}{4} (E_{\bar{k}-\bar{G}}^0 - E_{\bar{k}}^0)^2 + |V_{\bar{G}}|^2}$$

Enlace fuerte, celda primitiva cúbica ( $B = \gamma, A = \beta$ )

$$E(\bar{k}) \approx E_i - A - 2B(\cos k_x a + \cos k_y a + \cos k_z a)$$

$$A = -\langle \varphi_{i,n} | v | \varphi_{i,n} \rangle, \quad B = -\langle \varphi_{i,m} | v | \varphi_{i,n} \rangle$$

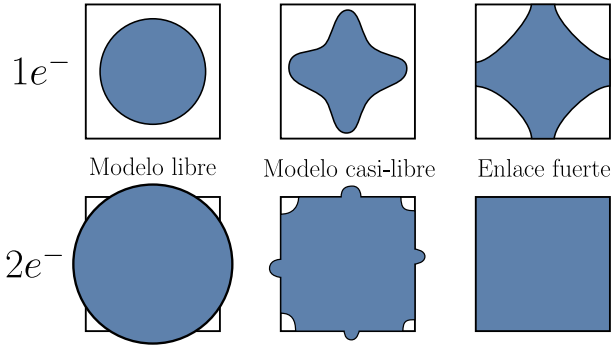
$$\bar{v} = \nabla_{\bar{k}} \omega(\bar{k}) = \frac{1}{\hbar} \nabla_{\bar{k}} E(\bar{k})$$

Carga de un campo  $\bar{\mathcal{E}}$

$$\dot{v}_i = \frac{1}{\hbar^2} \sum_j \frac{\partial^2 E}{\partial k_i \partial k_j} (-e \mathcal{E}_j), \quad \left( \frac{1}{m^*} \right)_{ij} = \frac{1}{\hbar^2} \frac{\partial^2 E(\bar{k})}{\partial k_i \partial k_j}$$

Caso totalmente degenerado

$$m^* = \frac{\hbar^2}{\left( \frac{d^2 E}{dk^2} \right)}, \quad E(\bar{k}) = E_0 + \frac{\hbar^2}{2m^*} |k|^2, \quad \sigma \simeq \frac{e^2 \tau (E_F) n}{m^*}$$

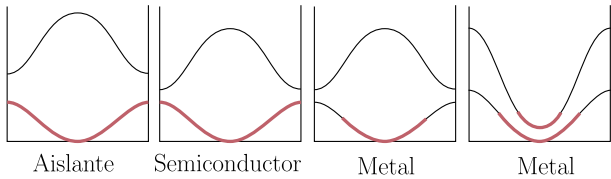


Tipos de materiales

**Aislante:** Banda llena ( $2e^-$ ).  $V_g > 4eV$

**Semiconductor** Banda llena ( $2e^-$ ).  $V_g < 4eV$ .

**Metal** Banda semillena ( $1e^2$  ó  $2e^-$  con bandas solapantes).



## 4 Semiconductores

extrínseco = dopado

Opacos si  $h\nu > E_g$

Nivel de Fermi  $E_F = \mu$

Densidad de estados ( $n$  electrones,  $p$  holes)

Energía de donadores / impurezas  $E_D$

$n$ (negativo)	$N_C$ (conducción)	$N_D$ (donadores)
$p$ (positivo)	$N_V$ (valencia)	$N_A$ (aceptores)

$$g_C(\varepsilon) = \frac{(2m_n^*)^{2/3}}{2\pi^2 \hbar^3} \sqrt{\varepsilon - \varepsilon_C}; \quad g_V(\varepsilon) = \frac{(2m_p^*)^{2/3}}{2\pi^2 \hbar^3} \sqrt{\varepsilon_V - \varepsilon}$$

$$n = \int_{\varepsilon_C}^{\infty} d\varepsilon g_C(\varepsilon) n_F(\beta(\varepsilon - \mu)) \approx \int_{\varepsilon_C}^{\infty} d\varepsilon g_C(\varepsilon) e^{\beta(\mu - \varepsilon)}$$

$$p = \int_{-\infty}^{\varepsilon_V} d\varepsilon g_V(\varepsilon) (1 - n_F(\beta(\varepsilon - \mu))) \approx \int_{-\infty}^{\varepsilon_V} d\varepsilon g_V(\varepsilon) e^{\beta(\varepsilon - \mu)}$$

$$n = \frac{1}{4} \left( \frac{2m_n^* k_B T}{\pi \hbar^2} \right)^{3/2} e^{\beta(\mu - \varepsilon_C)} = N_C e^{\beta(\mu - \varepsilon_C)}$$

$$p = \frac{1}{4} \left( \frac{2m_p^* k_B T}{\pi \hbar^2} \right)^{3/2} e^{\beta(\varepsilon_V - \mu)} = N_V e^{\beta(\varepsilon_V - \mu)}$$

$$np = N_C N_V e^{-\beta E_g} = 4 \left( \frac{k_B T}{2\pi \hbar^2} \right)^3 (m_n^* m_p^*)^{3/2} e^{-\beta E_g}$$

$$e^{2\beta\mu} = \frac{N_V}{N_C} e^{\beta(\varepsilon_V + \varepsilon_C)}, \quad \mu = \frac{\varepsilon_C + \varepsilon_V}{2} + \frac{3}{4} k_B T \ln \left( \frac{m_p^*}{m_n^*} \right)$$

$$\mu = \frac{e\tau}{m^*}, \quad \sigma = e(n\mu_n + p\mu_p), \quad E_g = \varepsilon_C - \varepsilon_V$$

$$n = p = \sqrt{N_C N_V} e^{-\frac{\beta E_g}{2}} \quad \text{si intrínseco}$$

Semiconductores dopados ( $n = 1$  ionización)  $\varepsilon = \varepsilon_0 \varepsilon_r$

$$E_n = \frac{m^* e^4}{2(4\pi \varepsilon \hbar)^2} \frac{1}{n^2}, \quad r_n = \varepsilon \frac{4\pi \hbar^2}{m^* e^2} n^2$$

$$n_n \approx \frac{2N_D}{1 + \sqrt{1 + 4 \frac{N_D}{N_C} e^{\beta E_d}}}$$

Unión p-n

$$n_n = N_C e^{\beta(\mu - \varepsilon_C^0)}; \quad p_p = N_V e^{\beta(\varepsilon_V^0 - \mu)}$$

$$d_n^0 = \sqrt{\frac{2\varepsilon V_D}{e} \frac{N_A/N_D}{N_A + N_D}}; \quad d_p^0 = \sqrt{\frac{2\varepsilon V_D}{e} \frac{N_D/N_A}{N_A + N_D}}$$

$$d_n(V) = d_n^0 \sqrt{\frac{V_D - V}{V_D}}; \quad d_p(V) = d_p^0 \sqrt{\frac{V_D - V}{V_D}}$$

ancho de zona de carga espacial =  $d_n(V) + d_p(V)$

$$eV_D = k_B T \ln \left( \frac{n_n p_p}{n_i^2} \right); \quad I(V) = (I_n^{gen} + I_p^{gen}) (e^{\beta eV} - 1)$$

p-Dopado ( $N_A \gg N_D$ ):

- $T$  bajas  $\Rightarrow p \approx N_V e^{-\beta E_A}$
- $T$  intermedias  $\Rightarrow p \approx N_A - N_D$
- $T$  alta  $\Rightarrow p = n = \sqrt{N_C N_V} e^{-\beta \frac{E_g}{2}}$

## 5 Mates

$$\sin^2 \left( \frac{x}{2} \right) = \frac{1 - \cos x}{2}$$

$$\int_0^{\infty} \frac{1}{e^x - 1} dx = +\infty, \quad \int_0^{\infty} \frac{1}{e^x + 1} dx = \ln(2)$$

$$\int_0^{\infty} \frac{x}{e^x - 1} dx = \frac{\pi^2}{6}, \quad \int_0^{\infty} \frac{x}{e^x + 1} dx = \frac{\pi^2}{12}$$

$$\int_0^{\infty} \frac{x^2}{e^x - 1} dx = 2\zeta(3), \quad \int_0^{\infty} \frac{x^2}{e^x + 1} dx = \frac{3}{2}\zeta(3)$$

$$\int_0^{\infty} \frac{x^3}{e^x - 1} dx = \frac{\pi^4}{15}, \quad \int_0^{\infty} \frac{x^3}{e^x + 1} dx = \frac{7\pi^4}{120}$$