

# ACCQ 206 - Entanglement

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# Recap

## Quantum states

$n$ -qubit state  $|\psi\rangle = \begin{pmatrix} \alpha_0 \\ \vdots \\ \alpha_{2^n} \end{pmatrix} \in \mathbb{C}^{2^n}$   
 $\sum_i |\alpha_i|^2 = 1$

- position:  $\{|x\rangle\}$   
- momentum:  $\{|p\rangle\}$

$|0\rangle / |1\rangle$   
position

## Evolution

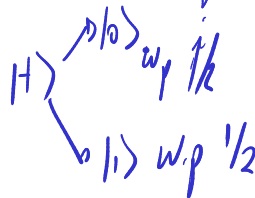
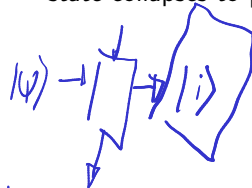
Unitary matrices  
 $U : UU^\dagger = U(U^*)^T = I$

$|\psi\rangle \rightarrow |\phi\rangle$

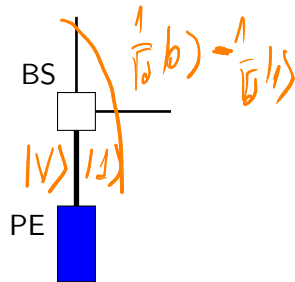
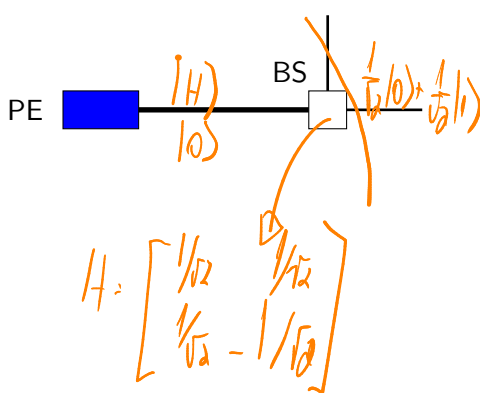
$|+\rangle / |-\rangle$   
momentum

## Measurement in comp. basis

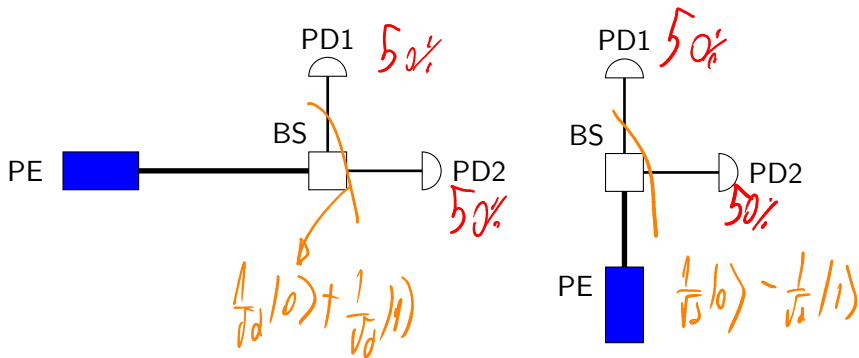
output  $i$  w.p.  $|\alpha_i|^2$   
state collapses to  $|i\rangle$



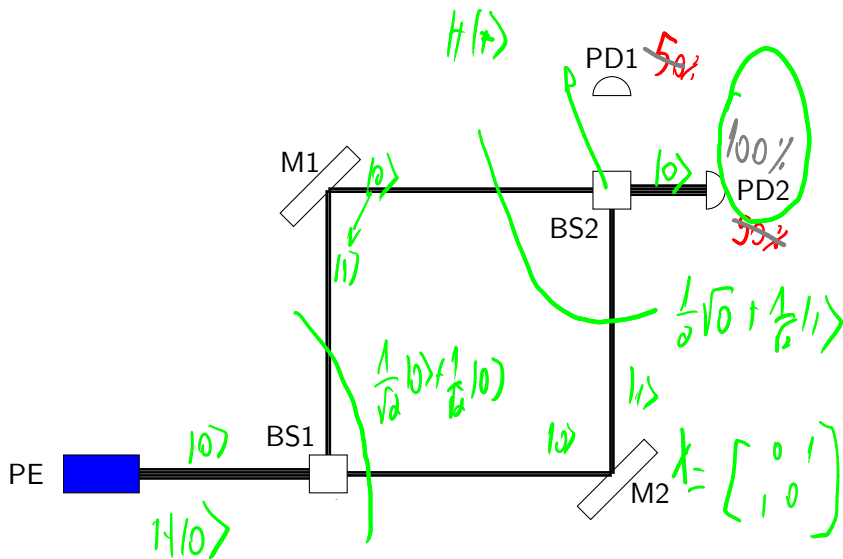
# Beam splitter

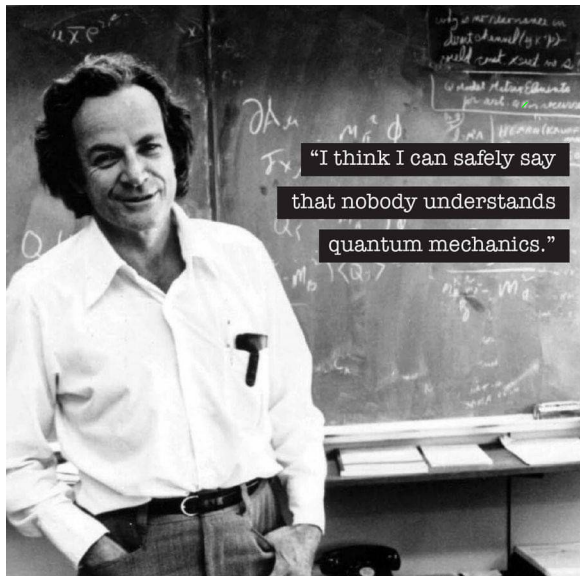


## Beam splitter



# Interferometer

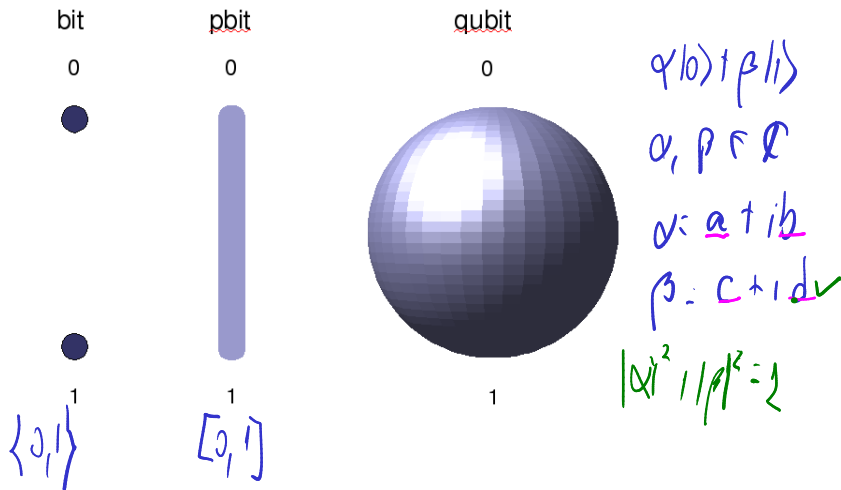




Quantum mechanics is very worthy of regard. But an inner voice tells me that this is not yet the right track. The theory yields much, but it hardly brings us closer to the Old One's secrets. I, in any case, am convinced that He does not play dice.

Albert Einstein

# Bit vs. pbit vs qubit



# Observables

- Hermitian operator  $A = A^\dagger = (A^T)^\dagger$

- ▶ Real values  $a_1, \dots, a_k$  and corresponding projectors,  $P_1, \dots, P_k$
- ▶ We have  $A = \sum_{i \in [k]} a_i P_i$
- ▶ Output  $a_i$  with probability  $q_i := \langle \psi | P_i | \psi \rangle = \|P_i |\psi\rangle\|^2$
- ▶ State collapses to

$$|\psi_i\rangle = \frac{P_i |\psi\rangle}{\sqrt{q_i}}$$

$P_i$  is a projector  $P_i^2 = P_i$

$$P_i^2 = P_i$$

$$\| |\psi_i\rangle \|^2 = 1$$

- Average value of the outcome is

$|\psi\rangle \rightarrow$  measure wrt  $A$   
 $E[a_i]$

$$\langle A \rangle = \langle \psi | A | \psi \rangle = \sum_i \langle \psi | P_i | \psi \rangle a_i$$

$\rightarrow \sum_i q_i a_i$



Global phases:  $|\psi\rangle$  vs.  $e^{i\theta} |\psi\rangle$   
 $\hookrightarrow$  global phase

Fix arb.  $A = \sum a_i P_i$   $\nearrow \cos\theta + i \sin\theta$

$|\psi\rangle$   
 $\rightarrow$   $a_i$  w.p.  $q_i = \langle \psi | P_i | \psi \rangle$   
 $\rightarrow$  post. meas. state is  $P_i |\psi\rangle / \sqrt{q_i}$

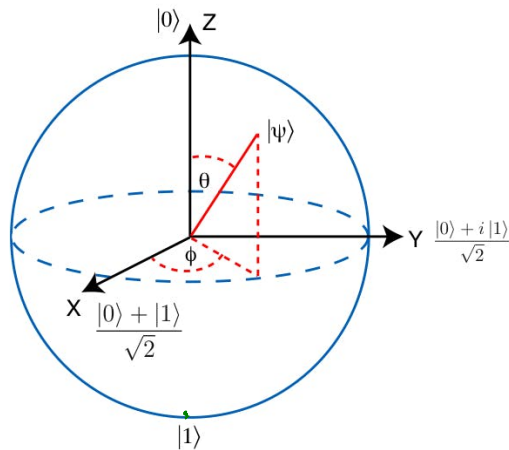
$e^{i\theta} |\psi\rangle$   
 $\rightarrow$   $a_i$  w.p.  $e^{-i\theta} \langle \psi | P_i e^{i\theta} |\psi\rangle = \langle \psi | P_i |\psi\rangle$   
 $\rightarrow$  post. meas. state  $\frac{e^{i\theta} P_i |\psi\rangle}{\sqrt{q_i}}$

Global phases  
 $\alpha |0\rangle + \beta |1\rangle$   
 $\alpha, \beta \in \mathbb{R}$

are not important!  
 $\alpha = \sqrt{p}$   
 $\beta = \sqrt{1-p}$   
 $|\alpha|^2 + |\beta|^2 = 1$

# Bloch sphere

Geometrical representation of 1 qubit



Polarization of Photons

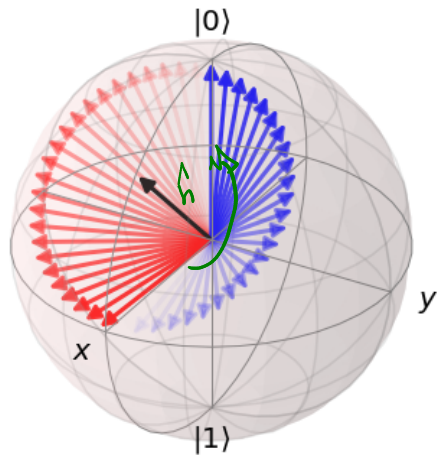


$$|\psi\rangle = \cos\theta |0\rangle + e^{i\phi} \sin\theta |1\rangle$$

$$X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \Rightarrow 180^\circ \text{ degrees around the } x \text{ axis}$$

$$Z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \Rightarrow 180^\circ \text{ degrees around the } z \text{ axis}$$

## Bloch sphere



$$H = \begin{bmatrix} 1/\sqrt{2} & 1/\sqrt{2} \\ 1/\sqrt{2} & -1/\sqrt{2} \end{bmatrix}$$

$$R_{\hat{n}}(\theta)$$

## Product states

A two qubit state  $|\psi\rangle$  is said to be a product state if there exists  $|\psi_1\rangle$  and  $|\psi_2\rangle$  such that

$$|\psi\rangle = \alpha_{00}|00\rangle + \alpha_{01}|01\rangle + \alpha_{10}|10\rangle + \alpha_{11}|11\rangle$$

$$|\psi\rangle = |\psi_1\rangle \otimes |\psi_2\rangle = |\psi_1\rangle |\psi_2\rangle$$

$\begin{matrix} \alpha_{00} & \alpha_{01} & \alpha_{10} & \alpha_{11} \\ \beta_{00} + \beta'_{01} & \gamma_{00} + \gamma'_{01} & \beta_{10} + \beta'_{11} & \gamma_{10} + \gamma'_{11} \end{matrix}$

Not all states are product

$$|EPR\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle) \rightarrow \text{EPR}$$

Proof Let us assume  $\exists |\psi_1\rangle, |\psi_2\rangle$  st  ~~$|EPR\rangle = |\psi_1\rangle \otimes |\psi_2\rangle$~~

$$|EPR\rangle = (\beta|0\rangle + \beta'|1\rangle) \otimes (\gamma|0\rangle + \gamma'|1\rangle) = \beta\gamma|00\rangle + \beta\gamma'|01\rangle + \beta'\gamma|10\rangle + \beta'\gamma'|11\rangle$$

$$\begin{aligned} (1) \beta\gamma &= \beta'\gamma' = \frac{1}{\sqrt{2}} \\ (2) \beta\gamma' &= \beta'\gamma = 0 \end{aligned}$$

$$\begin{aligned} (1) &\Rightarrow \beta, \gamma, \beta', \gamma' \neq 0 \\ (2) &\Rightarrow \begin{matrix} \beta & \text{or} & \gamma' \\ \beta' & \text{or} & \gamma \end{matrix} = 0 \end{aligned}$$

} contradiction!

# Entangled states

**Entangled states:** quantum states that are not product

$|FPR\rangle$

## Examples

$$GHZ_n = \frac{1}{\sqrt{2}}(|\underbrace{0\dots 0}_n\rangle + |\underbrace{1\dots 1}_n\rangle)$$

$$GHZ_3 = \frac{1}{\sqrt{2}}(|0\rangle_A |00\rangle_B + |1\rangle_A |11\rangle)$$

$\begin{matrix} |\psi_1\rangle \\ |\psi_2\rangle \end{matrix}$  on 1 qubit  
on 2 qubits

$$|GHZ_3\rangle = |\psi_1\rangle \otimes |\psi_2\rangle$$

$$W_n = \frac{1}{\sqrt{n}}(|10\dots 0\rangle + |01\dots 0\rangle + |00\dots 1\rangle)$$

$$W_3 = \frac{1}{\sqrt{3}}(|100\rangle + |010\rangle + |001\rangle)$$

Bell basis - basis for  $\mathbb{C}^4$

Computational  
Canonical basis

$$\begin{aligned} Z|0\rangle &= |0\rangle \\ Z|1\rangle &= -|1\rangle \end{aligned}$$

every vector in the basis is entangled

Bell basis

$$\left\{ \begin{array}{l} |00\rangle \\ |01\rangle \\ |10\rangle \\ |11\rangle \end{array} \right\} \begin{array}{l} |0\rangle \otimes |0\rangle \\ |0\rangle \otimes |1\rangle \\ \dots \end{array}$$

$$\begin{aligned} X &= \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} & Z &= \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \\ X^\dagger &= X & Z^\dagger &= Z \\ X^0 &= I & Z^0 &= I \end{aligned}$$

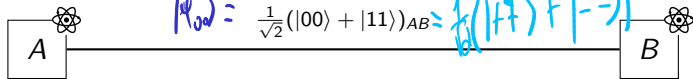
$$\begin{aligned} \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle) &= |\psi_{00}\rangle \\ \frac{1}{\sqrt{2}}(|00\rangle - |11\rangle) &= |\psi_{01}\rangle \\ \frac{1}{\sqrt{2}}(|01\rangle + |10\rangle) &= |\psi_{10}\rangle \\ \frac{1}{\sqrt{2}}(|01\rangle - |10\rangle) &= |\psi_{11}\rangle \end{aligned}$$

$$\begin{aligned} |\psi_{00}\rangle &= \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle) \\ |\psi_{01}\rangle &= \frac{1}{\sqrt{2}}(|00\rangle - |11\rangle) \end{aligned}$$

Lemma:  $|\psi_{00}\rangle$  is entangled

$$|\psi_{ab}\rangle = (X^a Z^b \otimes I) |\psi_{00}\rangle = (I \otimes X^a Z^b) |\psi_{00}\rangle$$

# Source of quantum weirdness



$$|\psi\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)_{AB} = \frac{1}{\sqrt{2}}(|++\rangle + |--\rangle)$$

exercise

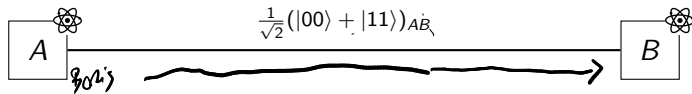
$$\begin{aligned} |00\rangle &= \\ &= |0\rangle\left(\frac{1}{\sqrt{2}}|+\rangle + \frac{1}{\sqrt{2}}|-\rangle\right) = \\ &= \left(\frac{1}{2}|0\rangle|+\rangle + \frac{1}{2}|0\rangle|-\rangle\right) \end{aligned}$$

- what happens if A measures her share on comp basis
  - outcome is 0 w.p.  $\frac{1}{2}$ 
    - Bob measures outcome is 0 w.p. 1
  - outcome is 1 w.p.  $\frac{1}{2}$ 
    - Bob measures outcome is 1 w.p. 1

Hadamard basis

$$\begin{aligned} \frac{1}{\sqrt{2}}|00\rangle &= \frac{1}{2}|00\rangle + \frac{1}{2}|01\rangle + \frac{1}{2}|10\rangle + \frac{1}{2}|11\rangle \\ &= \frac{1}{2}|00\rangle + \frac{1}{2}|01\rangle + \frac{1}{2}|10\rangle + \frac{1}{2}|11\rangle \end{aligned}$$

## Source of quantum weirdness



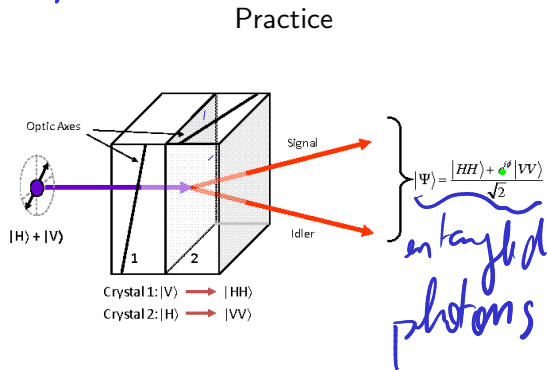
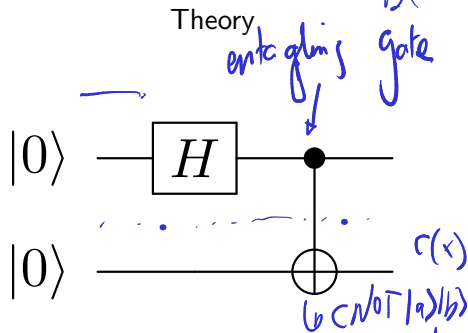
### Violation of relativity theory?

"I cannot seriously believe in it [quantum theory] because the theory cannot be reconciled with the idea that physics should represent a reality in time and space, free from spooky actions at a distance" Albert Einstein

- non-local games
- quantum teleportation: how to transmit quantum states via classical channels
- superdense coding: 2 bits of information w/ 1 qubit



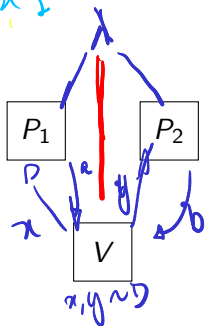
# Generating entangled states $\frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$



$$\begin{aligned} \text{CNOT} \left( H \frac{1}{\sqrt{2}} |00\rangle \right) &= \text{CNOT} \left( \frac{1}{\sqrt{2}} (|00\rangle + |10\rangle) \right) \\ &= \frac{1}{\sqrt{2}} (\text{CNOT} |00\rangle + \text{CNOT} |10\rangle) \\ &= \frac{1}{\sqrt{2}} (|00\rangle + |11\rangle) = |\text{EPR}\rangle \end{aligned}$$

# Non-local games

phase 1



$D$

- 1)  $P_1$  and  $P_2$  share some strategy
- 2) Not allowed to communicate after
- 3)  $V$  pick questions to  $P_1$  and  $P_2$
- 4)  $x, y \sim D$   
 $P_1$  and  $P_2$  answer
- 5)  $V(a, b | x, y) = 0/1$

Classical value  $\omega(G)$

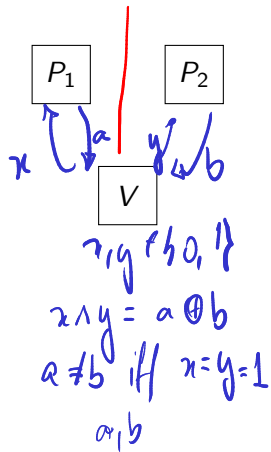
$\lambda$  is a bit string  
 $\omega(G) = \max_{\text{classical strategy}} \Pr[V(a, b | x, y) = 1]$

Quantum value  $\omega^*(G)$

$\lambda$  is a quantum state  $|\psi\rangle_{AB}$   
 $\omega^*(G) = \max_{\text{quantum strategy}} \Pr[V(a, b | x, y) = 1]$

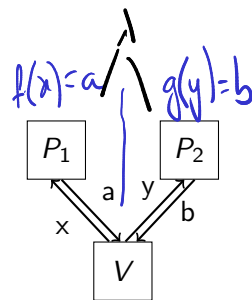
Goal: find  $G$  st.  $\omega^*(G) > \omega(G)$

## CHSH game



$x, y$  are random bits

# CHSH game - classical strategies



$x, y \in_R \{0, 1\}$   
 $x \oplus y = a \oplus b$   
 $a \neq b \text{ iff } x = y = 1$   
 $a, b \in \{0, 1\}$

Example  $f$  in  $f$  or  $g$

$f(x) = 0$   
 $g(y) = 0$

$x$	$y$	$a$	$b$	Win?
0	0	0	0	✓
0	1	0	0	✓
1	0	0	0	✓
1	1	0	0	✗

Win  $\frac{3}{4}$

$x \neq y = 0$   
 $x \neq y = 0$

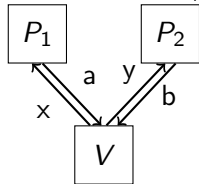
$a \oplus b = 0$

Lim:  $w(G) = \frac{3}{4}$

$V(a, b | x, y) = 1 \text{ iff } x \oplus y = a \oplus b$   
 $x = y = 0 \Rightarrow 0 = 0$

# CHSH game - quantum strategies

1. quantum state.
2. quantum operations



$$x, y \in_R \{0, 1\}$$

$$x \cdot y = a \oplus b$$

$$a \neq b \text{ iff } x = y = 1$$

$$\omega^*(\text{CHSH}) \approx \cos^2 \pi/8 \approx 0.85$$

$$1. \quad \frac{1}{\sqrt{2}}(|00\rangle - |11\rangle) = |\psi_{01}\rangle = Z^b |\psi_{00}\rangle$$

$$2. \quad R(\theta) = \begin{bmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{bmatrix}$$

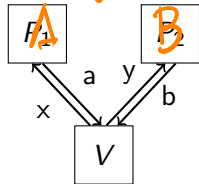
$$\text{str. of } P_1: \quad \begin{array}{ll} \text{if } x=0 & R(-\pi/16) \\ \text{if } x=1 & R(3\pi/16) \end{array}$$

then measure in comp basis

$$P_2: \quad \begin{array}{ll} \text{if } y=0 & R(-\pi/16) \\ \text{if } y=1 & R(3\pi/16) \end{array}$$

then measure in comp. basis

# CHSH game - quantum strategies



$$x, y \in_R \{0, 1\}$$

$$x \cdot y = a \oplus b$$

$$a \neq b \text{ iff } x = y = 1$$

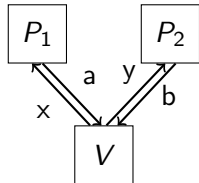
if  $x \wedge y = 0$

$$\Rightarrow \theta_A + \theta_B = \pm \frac{\pi}{8}$$

if  $x \wedge y = 1$   $\theta_A + \theta_B = \frac{3\pi}{8}$

$$\begin{aligned} & \frac{1}{\sqrt{2}} (R(\theta_A) \otimes R(\theta_B)) (\underline{100} - \underline{111}) = \\ & \frac{1}{\sqrt{2}} (\cos(\theta_A) |0\rangle - \sin(\theta_A) |1\rangle) (\cos(\theta_B) |0\rangle - \sin(\theta_B) |1\rangle) \\ & - (\sin(\theta_A) |0\rangle + \cos(\theta_A) |1\rangle) (\sin(\theta_B) |0\rangle + \cos(\theta_B) |1\rangle) \\ & = \frac{1}{\sqrt{2}} (\cos(\theta_A) \cos(\theta_B) - \sin(\theta_A) \sin(\theta_B)) (100 - 111) \\ & - (\cos(\theta_A) \sin(\theta_B) + \sin(\theta_A) \cos(\theta_B)) (101 + 110) \\ & = \frac{1}{\sqrt{2}} (\cos(\theta_A + \theta_B) (100 - 111) + \sin(\theta_A + \theta_B) (101 + 110)) \end{aligned}$$

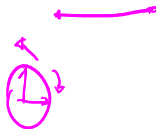
## CHSH game - quantum strategies



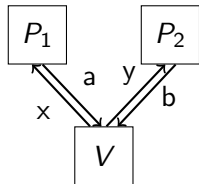
$$x, y \in_R \{0, 1\}$$

$$x \cdot y = a \oplus b$$

$$a \neq b \text{ iff } x = y = 1$$



## CHSH game - quantum strategies



$$x, y \in_R \{0, 1\}$$

$$x \cdot y = a \oplus b$$

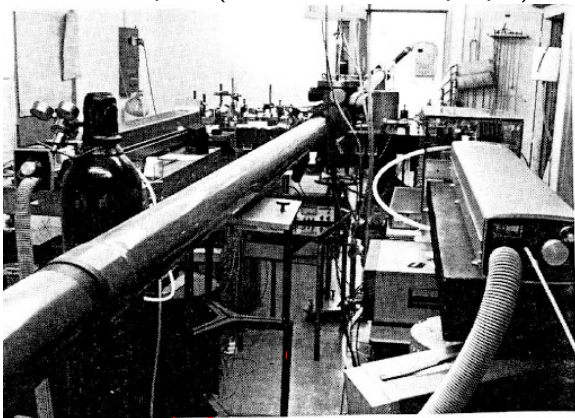
$$a \neq b \text{ iff } x = y = 1$$



# Bell inequalities

Bell inequalities give us a way for testing if Nature is not classical  $\Rightarrow$  implement the quantum experiments in the lab and if the acceptance probability is strictly larger than

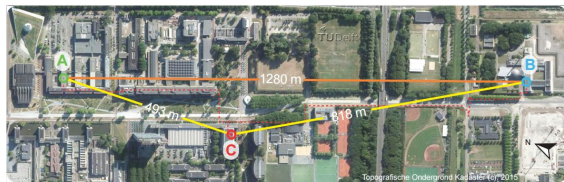
Alain Aspect (1982, Institut d'Optique)



$> 3/4$

loophole experiment

QuTech group (2015, TU Delft)



$> 3/4$

## Further notions

We have seen a quantum strategy that achieves  $\omega^*(CHSH) = \cos^2 \frac{\pi}{8} \approx 0.85$

Is this strategy optimal?  $\checkmark$

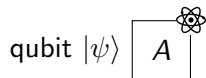
Tsirelson's inequalities  
to bound the quantum  
value of games

Is this <sup>maximally</sup> optimal strategy unique?  $\checkmark$

if achieve  $\cos^2 \frac{\pi}{8}$   
then one can  
characterize your  
q. devices

↳ DI crypto  
↳ disproof of Connes conj.

# Quantum teleportation



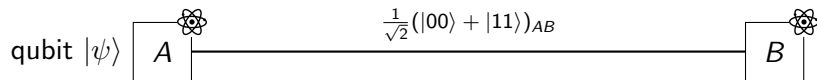
**Problem:** Alice wants to send a qubit  $|\psi\rangle$  to Bob.

**If they have a quantum channel to communicate:**

**If they have a classical channel to communicate:**

**If they have a classical channel to communicate + pre-shared quantum state:**

# Quantum teleportation



## Super-dense coding

**Alice receives two random bits  $a$  and  $b$  and she wants Bob to learn both of them.**

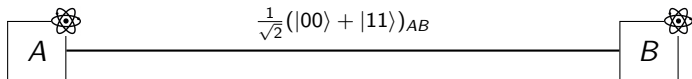
If Alice sends a single classical bit:

If they have shared randomness and Alice sends a single classical bit:

If they have shared quantum state and Alice sends a single classical bit:

If they have shared quantum state and Alice sends a single qubit:

# Super-dense coding



# Mixed states

- Mixed states: probabilistic distribution of quantum states

## Examples

$$((\frac{1}{2}, |0\rangle), (\frac{1}{2}, |+\rangle)), ((\frac{1}{3}, |0\rangle), (\frac{2}{3}, |1\rangle))$$

- Density matrices: mathematical representation of mixed states

$$(1, |\psi\rangle)$$

$$((p_1, |\psi_1\rangle), (p_2, |\psi_2\rangle), (p_k, |\psi_k\rangle))$$

- Properties of density matrices

- ▶ Its trace is 1 (The trace of a square matrix is  $\sum_i a_{i,i}$ )
- ▶ Positive semi-definitive (all its eigenvalues are non-negative)

- Definition of evolution and measurements can be extended to density matrices

## “Parts” of quantum states

- Trace-out: ignore qubits of a larger quantum state