UPSC PHYSICS PYQ SOLUTION

Quantum Mechanics - Part 2

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41 A beam of particles of energy 9 eV is incident on a potential step 8 eV high from the left. What percentage of particles will reflect back?

Introduction:

In quantum mechanics, the reflection and transmission of particles at a potential step is a fundamental problem. When a particle encounters a potential step, part of the wave function is reflected, and part is transmitted. The reflection coefficient (R) gives the probability of the particle being reflected.

Solution:

The energy of the incident particles is $E = 9\,\text{eV}$, and the height of the potential step is $V_0 = 8\,\text{eV}$. The reflection coefficient (R) is given by:

$$R = \left(\frac{k_1 - k_2}{k_1 + k_2}\right)^2$$

where k_1 and k_2 are the wave numbers of the particle in the regions before and after the potential step, respectively.

The wave number k is related to the energy E and the potential V by:

$$k = \sqrt{\frac{2m(E - V)}{\hbar^2}}$$

For the region before the potential step ($E = 9\,\text{eV}$ and $V = 0\,\text{eV}$):

$$k_1 = \sqrt{\frac{2m(9\,\text{eV})}{\hbar^2}}$$

For the region after the potential step ($E = 9\,\text{eV}$ and $V = 8\,\text{eV}$):

$$k_2 = \sqrt{\frac{2m(9\,\text{eV} - 8\,\text{eV})}{\hbar^2}} = \sqrt{\frac{2m(1\,\text{eV})}{\hbar^2}}$$

The ratio of the wave numbers is:

$$\frac{k_1}{k_2} = \frac{\sqrt{9\,\text{eV}}}{\sqrt{1\,\text{eV}}} = 3$$

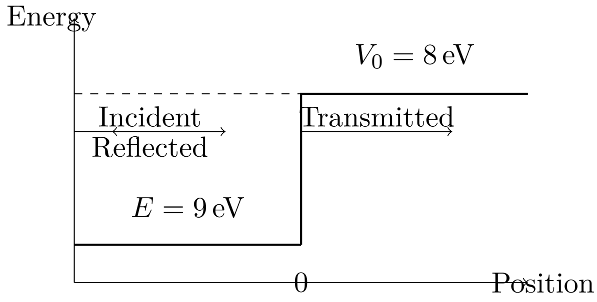
Substituting into the reflection coefficient formula:

$$R = \left(\frac{3 - 1}{3 + 1}\right)^2 = \left(\frac{2}{4}\right)^2 = \left(\frac{1}{2}\right)^2 = 0.25$$

Therefore, the reflection percentage is:

$$R \times 100\% = 0.25 \times 100\% = 25\%$$

The following diagram illustrates the potential step and the wave function behavior:



Conclusion:

The reflection coefficient indicates that 25% of the particles will reflect back when a beam of particles with energy 9 eV encounters a potential step of 8 eV. This result highlights the wave nature of particles, where partial reflection and transmission occur due to quantum mechanical effects.

Applications:

1. **Tunneling in Semiconductors**: Quantum tunneling is crucial in the operation of semiconductor devices such as diodes and transistors.
2. **Scanning Tunneling Microscopy (STM)**: STM relies on the quantum tunneling of electrons to image surfaces at the atomic level.
3. **Nuclear Fusion**: Quantum tunneling allows particles to overcome the Coulomb barrier, facilitating nuclear reactions in stars and experimental fusion reactors.

42 Estimate the size of hydrogen atom and the ground state energy from the uncertainty principle.

Introduction: The problem requires an estimation of the characteristic size (Bohr radius) and ground state energy of a hydrogen atom using the Heisenberg uncertainty principle. The hydrogen atom consists of an electron bound to a proton via Coulomb attraction. We aim to estimate:

- The approximate radius r of the hydrogen atom,
- The ground state energy E of the electron.

We assume a non-relativistic quantum mechanical model and apply the uncertainty relation $\Delta x \Delta p \sim \hbar$.

Solution:

Let the electron be confined within a region of size r, so the uncertainty in position is $\Delta x \sim r$. Then the uncertainty in momentum is:

$$\Delta p \sim \frac{\hbar}{r}\text{order of magnitude is satisfied even when we don't take } 1/2 \text{ as a factor}$$

The kinetic energy of the electron can be approximated using:

$$T \sim \frac{(\Delta p)^2}{2m} = \frac{\hbar^2}{2mr^2},$$

where m is the mass of the electron.

The potential energy due to Coulomb attraction between the proton and the electron is:

$$V \sim -\frac{e^2}{4\pi\varepsilon_0 r},$$

where e is the elementary charge and ε_0 is the vacuum permittivity.

The total energy of the electron is approximately:

$$E(r) = T + V \sim \frac{\hbar^2}{2mr^2} - \frac{e^2}{4\pi\varepsilon_0 r}.$$

To find the equilibrium (ground state), we minimize $E(r)$ with respect to r:

$$\frac{dE}{dr} = -\frac{\hbar^2}{mr^3} + \frac{e^2}{4\pi\varepsilon_0 r^2} = 0.$$

The second derivative greater than 0, confirms a minimum.

Solving for r:

$$\frac{\hbar^2}{mr^3} = \frac{e^2}{4\pi\varepsilon_0 r^2} \quad \Rightarrow \quad r = \frac{4\pi\varepsilon_0\hbar^2}{me^2}.$$

This is the Bohr radius:

$$a_0 = \frac{4\pi\varepsilon_0\hbar^2}{me^2} \approx 5.29 \times 10^{-11}\,\text{m}.$$

Substitute $r = a_0$ into the expression for energy:

$$E = \frac{\hbar^2}{2ma_0^2} - \frac{e^2}{4\pi\varepsilon_0 a_0}.$$

This gives the ground state energy:

$$E_0 = -13.6\,\text{eV}.$$

Conclusion: By applying the uncertainty principle, we estimate the size of the hydrogen atom to be approximately $a_0 = 5.29 \times 10^{-11}\,\text{m}$, known as the Bohr radius. The ground state energy is approximately $E_0 = -13.6\,\text{eV}$, consistent with experimental results and Bohr's model.

43 Write down the Hamiltonian operator for a linear harmonic oscillator. Show that the energy eigenvalue of the same can be given by $E_n = \left(n + \frac{1}{2}\right)\hbar\omega_0$ at energy state n with ω_0 being the natural frequency of vibration of the linear oscillator. Prove that $n = 0$ energy state has a wave function of typical Gaussian form.

Introduction: The problem involves analyzing the quantum harmonic oscillator. We are asked to:

- Write the Hamiltonian operator for a linear harmonic oscillator.
- Derive the energy eigenvalues, demonstrating the quantized form $E_n = \left(n + \frac{1}{2}\right)\hbar\omega_0$.
- Show that the ground state wavefunction ($n = 0$) has a Gaussian form.

Assumptions include a one-dimensional oscillator and standard canonical quantization with position operator $\hat{x}$ and momentum operator $\hat{p}$ satisfying $[\hat{x}, \hat{p}] = i\hbar$.

Solution:

The Hamiltonian for a one-dimensional quantum harmonic oscillator is

$$\hat{H} = \frac{\hat{p}^2}{2m} + \frac{1}{2}m\omega_0^2\hat{x}^2.$$

This represents the total energy of the system the sum of kinetic and potential energies in quantum mechanical form.

To simplify the problem and reveal its underlying algebraic structure, we introduce ladder (creation and annihilation) operators:

$$\begin{aligned}
\hat{a} &= \sqrt{\frac{m\omega_0}{2\hbar}}\left(\hat{x} + \frac{i}{m\omega_0}\hat{p}\right), \\
\hat{a}^\dagger &= \sqrt{\frac{m\omega_0}{2\hbar}}\left(\hat{x} - \frac{i}{m\omega_0}\hat{p}\right).
\end{aligned}$$

These satisfy the commutation relation:

$$[\hat{a}, \hat{a}^\dagger] = 1.$$

Ladder operators provide an elegant way to analyze the harmonic oscillator because they allow us to raise or lower the energy levels of the system in discrete steps, corresponding to the quantized nature of energy in quantum mechanics.

In terms of these operators, the Hamiltonian becomes:

$$\hat{H} = \hbar\omega_0\left(\hat{a}^\dagger\hat{a} + \frac{1}{2}\right).$$

The number operator is defined as $\hat{n} = \hat{a}^\dagger\hat{a}$, and its eigenstates $|n\rangle$ satisfy:

$$\hat{n}\,|n\rangle = n\,|n\rangle\,, \quad n = 0, 1, 2, \ldots$$

Hence, the energy eigenvalues are:

$$E_n = \hbar\omega_0 \left(n + \frac{1}{2}\right).$$

This result shows that the energy levels are quantized and equally spaced, with a minimum energy of $\frac{1}{2}\hbar\omega_0$, known as the zero-point energy. This non-zero minimum energy reflects the Heisenberg uncertainty principle: even in the ground state, the particle cannot have both definite position and momentum.

Now consider the ground state $|0\rangle$, which satisfies:

$$\hat{a}\,|0\rangle = 0.$$

Using the coordinate representation, we have:

$$\hat{x} = x,$$
$$\hat{p} = -i\hbar\frac{d}{dx}.$$

Thus the annihilation operator becomes:

$$\hat{a} = \sqrt{\frac{m\omega_0}{2\hbar}} \left(x + \frac{\hbar}{m\omega_0}\frac{d}{dx}\right).$$

Apply $\hat{a}$ to the ground state wavefunction $\psi_0(x)$:

$$\hat{a}\psi_0(x) = 0 \quad \Rightarrow \quad \left(x + \frac{\hbar}{m\omega_0}\frac{d}{dx}\right)\psi_0(x) = 0.$$

Solving this differential equation:

$$\frac{d\psi_0}{dx} = -\frac{m\omega_0}{\hbar}x\psi_0(x).$$

This is a separable differential equation. Integrating both sides:

$$\int \frac{1}{\psi_0}\,d\psi_0 = -\frac{m\omega_0}{\hbar}\int x\,dx,$$
$$\ln\psi_0 = -\frac{m\omega_0}{2\hbar}x^2 + C,$$
$$\psi_0(x) = Ae^{-\frac{m\omega_0}{2\hbar}x^2},$$

where $A = e^C$ is the normalization constant.

To normalize, we impose:

$$\int_{-\infty}^{\infty} |\psi_0(x)|^2 dx = 1 \Rightarrow A = \left(\frac{m\omega_0}{\pi\hbar}\right)^{1/4}.$$

Conclusion: The Hamiltonian operator for a linear harmonic oscillator is $\hat{H} = \frac{\hat{p}^2}{2m} + \frac{1}{2}m\omega_0^2\hat{x}^2$. Its energy eigenvalues are quantized as $E_n = \left(n + \frac{1}{2}\right)\hbar\omega_0$, reflecting the discrete and equally spaced energy levels characteristic of quantum oscillators. The ground state ($n = 0$) wavefunction is of Gaussian form:

$$\psi_0(x) = \left(\frac{m\omega_0}{\pi\hbar}\right)^{1/4} e^{-\frac{m\omega_0}{2\hbar}x^2}.$$

This confirms both the quantized energy spectrum and the Gaussian nature of the ground state in quantum harmonic oscillators. Moreover, excited states can be generated by applying the creation operator repeatedly on the ground state.

44 Prove that Bohr hydrogen atom approaches classical conditions, when n becomes very large and small quantum jumps are involved.

Introduction:

The correspondence principle, formulated by Niels Bohr, states that the behavior of systems described by quantum mechanics replicates classical physics in the limit of large quantum numbers. For the Bohr model of the hydrogen atom, this principle can be demonstrated by showing that the energy levels become closely spaced and the frequency of radiation approaches the classical orbital frequency as n becomes very large.

Solution:

1. **Energy Levels in Bohr Model:**

The energy levels of a hydrogen atom in the Bohr model are given by:

$$E_n = -\frac{13.6\,\text{eV}}{n^2}$$

where n is the principal quantum number.

2. **Frequency of Radiation:**

When an electron transitions from a higher energy level n_i to a lower energy level n_f, the frequency of the emitted photon is:

$$f = \frac{E_i - E_f}{h}$$

Substituting the energy levels:

$$f = \frac{-\frac{13.6\,\text{eV}}{n_i^2} + \frac{13.6\,\text{eV}}{n_f^2}}{h}$$

Let $n_i = n$ and $n_f = n - \Delta n$ where Δn is small compared to n. Then,

$$f = \frac{13.6\,\text{eV}}{h}\left(\frac{1}{(n-\Delta n)^2} - \frac{1}{n^2}\right)$$

For large n and small Δn, we can use the binomial approximation:

$$(n-\Delta n)^2 \approx n^2 - 2n\Delta n$$

So,

$$\frac{1}{(n-\Delta n)^2} \approx \frac{1}{n^2}\left(1 + \frac{2\Delta n}{n}\right)$$

Therefore,

$$f \approx \frac{13.6\,\text{eV}}{h}\left(\frac{1}{n^2} - \frac{1}{n^2}\left(1 + \frac{2\Delta n}{n}\right)\right) = \frac{13.6\,\text{eV}}{h}\frac{2\Delta n}{n^3}$$

3. **Classical Orbital Frequency:**

The classical orbital frequency $f_{\text{classical}}$ of an electron in the nth orbit is given by:

$$f_{\text{classical}} = \frac{\omega}{2\pi} = \frac{v}{2\pi r}$$

Using Bohr's model, $v = \frac{e^2}{2\epsilon_0 h}\frac{1}{n}$ and $r = \frac{4\pi\epsilon_0 h^2 n^2}{e^2 m}$, we get:

$$f_{\text{classical}} = \frac{\left(\frac{e^2}{2\epsilon_0 h}\frac{1}{n}\right)}{2\pi\left(\frac{4\pi\epsilon_0 h^2 n^2}{e^2 m}\right)} = \frac{e^4 m}{16\pi^3\epsilon_0^2 h^3}\frac{1}{n^3}$$

4. **Comparison and Conclusion:**

For large n,

$$f \approx \frac{13.6\,\text{eV}}{h}\frac{2\Delta n}{n^3} = f_{\text{classical}}\Delta n$$

Thus, the frequency of the radiation approaches the classical orbital frequency when the quantum number n is very large, confirming Bohr's correspondence principle.

Conclusion:

As the quantum number n becomes very large, the energy levels of the Bohr hydrogen atom become closely spaced, and the frequency of emitted radiation for small quantum jumps approaches the classical orbital frequency. This demonstrates that the Bohr model converges to classical physics in the limit of large quantum numbers, highlighting the correspondence principle. Practical applications include understanding atomic spectra and transitions in high-energy physics and astrophysics.

45 Find the probability current density for the wave function $\Psi(x,t) = \left[Ae^{ipx/\hbar} + Be^{-ipx/\hbar}\right] e^{-ip^2t/2m\hbar}$. Interpret the result physically.

Introduction:

In quantum mechanics, the probability current density $\mathbf{j}(x,t)$ represents the flow of probability associated with the wave function $\Psi(x,t)$. It is defined as:

$$\mathbf{j}(x,t) = \frac{\hbar}{2mi}\left(\Psi^*\frac{\partial\Psi}{\partial x} - \Psi\frac{\partial\Psi^*}{\partial x}\right)$$

where Ψ^* is the complex conjugate of Ψ.

Solution:

Given the wave function:

$$\Psi(x,t) = \left[Ae^{ipx/\hbar} + Be^{-ipx/\hbar}\right] e^{-ip^2t/2m\hbar}$$

First, find the partial derivatives of Ψ and Ψ^* with respect to x.

1. **Partial Derivative of Ψ with Respect to x:**

$$\frac{\partial\Psi}{\partial x} = \frac{\partial}{\partial x}\left[\left(Ae^{ipx/\hbar} + Be^{-ipx/\hbar}\right) e^{-ip^2t/2m\hbar}\right]$$

Since $e^{-ip^2t/2m\hbar}$ is a constant with respect to x:

$$\frac{\partial\Psi}{\partial x} = e^{-ip^2t/2m\hbar}\left[\frac{\partial}{\partial x}\left(Ae^{ipx/\hbar} + Be^{-ipx/\hbar}\right)\right]$$

$$= e^{-ip^2t/2m\hbar}\left[\frac{ip}{\hbar}Ae^{ipx/\hbar} - \frac{ip}{\hbar}Be^{-ipx/\hbar}\right]$$

$$= \frac{ip}{\hbar}e^{-ip^2t/2m\hbar}\left[Ae^{ipx/\hbar} - Be^{-ipx/\hbar}\right]$$

2. **Partial Derivative of Ψ^* with Respect to x:**

The complex conjugate of $\Psi(x,t)$ is:

$$\Psi^*(x,t) = \left[A^*e^{-ipx/\hbar} + B^*e^{ipx/\hbar}\right] e^{ip^2t/2m\hbar}$$

$$\frac{\partial\Psi^*}{\partial x} = e^{ip^2t/2m\hbar}\left[\frac{\partial}{\partial x}\left(A^*e^{-ipx/\hbar} + B^*e^{ipx/\hbar}\right)\right]$$

$$= e^{ip^2t/2m\hbar}\left[-\frac{ip}{\hbar}A^*e^{-ipx/\hbar} + \frac{ip}{\hbar}B^*e^{ipx/\hbar}\right]$$

$$= \frac{ip}{\hbar}e^{ip^2t/2m\hbar}\left[-A^*e^{-ipx/\hbar} + B^*e^{ipx/\hbar}\right]$$

3. **Probability Current Density:**

Using the definition of $\mathbf{j}(x,t)$:

$$\mathbf{j}(x,t) = \frac{\hbar}{2mi}\left(\Psi^*\frac{\partial\Psi}{\partial x} - \Psi\frac{\partial\Psi^*}{\partial x}\right)$$

Substitute Ψ, Ψ^*, and their derivatives:

$$\Psi^*\frac{\partial\Psi}{\partial x} = \left[A^*e^{-ipx/\hbar} + B^*e^{ipx/\hbar}\right]e^{ip^2t/2m\hbar}\cdot\frac{ip}{\hbar}e^{-ip^2t/2m\hbar}\left[Ae^{ipx/\hbar} - Be^{-ipx/\hbar}\right]$$

$$= \frac{ip}{\hbar}\left[A^*e^{-ipx/\hbar}Ae^{ipx/\hbar} - A^*e^{-ipx/\hbar}Be^{-ipx/\hbar} + B^*e^{ipx/\hbar}Ae^{ipx/\hbar} - B^*e^{ipx/\hbar}Be^{-ipx/\hbar}\right]$$

$$= \frac{ip}{\hbar}\left[A^*A - A^*Be^{-2ipx/\hbar} + B^*Ae^{2ipx/\hbar} - B^*B\right]$$

Similarly,

$$\Psi\frac{\partial\Psi^*}{\partial x} = \left[Ae^{ipx/\hbar} + Be^{-ipx/\hbar}\right]e^{-ip^2t/2m\hbar}\cdot\frac{ip}{\hbar}e^{ip^2t/2m\hbar}\left[-A^*e^{-ipx/\hbar} + B^*e^{ipx/\hbar}\right]$$

$$= \frac{ip}{\hbar}\left[Ae^{ipx/\hbar}(-A^*e^{-ipx/\hbar}) + Ae^{ipx/\hbar}B^*e^{ipx/\hbar} + Be^{-ipx/\hbar}(-A^*e^{-ipx/\hbar}) + Be^{-ipx/\hbar}B^*e^{ipx/\hbar}\right]$$

$$= \frac{ip}{\hbar}\left[-AA^* + AB^*e^{2ipx/\hbar} - BA^*e^{-2ipx/\hbar} + BB^*\right]$$

Therefore, the probability current density is:

$$\mathbf{j}(x,t) = \frac{\hbar}{2mi}\left[\frac{ip}{\hbar}\left(A^*A - A^*Be^{-2ipx/\hbar} + B^*Ae^{2ipx/\hbar} - B^*B\right) - \frac{ip}{\hbar}\left(-AA^* + AB^*e^{2ipx/\hbar} - BA^*e^{-2ipx\right.\right.$$

$$= \frac{\hbar}{2mi}\cdot\frac{2ip}{\hbar}\left(A^*A - B^*B\right)$$

$$\mathbf{j}(x,t) = \frac{p}{m}\left(A^*A - B^*B\right)$$

Conclusion:

The probability current density for the given wave function is:

$$\mathbf{j}(x,t) = \frac{p}{m}\left(A^*A - B^*B\right)$$

Interpretation:

This result indicates that the probability current density depends on the coefficients A and B. If $|A|^2 = |B|^2$, the probability current density $\mathbf{j}(x,t)$ is zero, implying no net flow of probability. If $|A|^2 \neq |B|^2$, there is a net flow of probability in the direction of the momentum p.

This reflects the physical interpretation that the probability current density represents the flow of probability for a particle described by the wave function $\Psi(x,t)$. The terms $|A|^2$ and $|B|^2$ represent the probabilities of the particle moving in positive and negative directions, respectively. The difference between these probabilities determines the net flow of probability in the system.

46 A particle is described by the wave function $\Psi(x) = \left(\frac{\pi}{2}\right)^{-1/4} e^{-ax^2/2}$. Calculate Δx and Δp for the particle, and verify the uncertainty relation $\Delta x \Delta p = \frac{\hbar}{2}$.

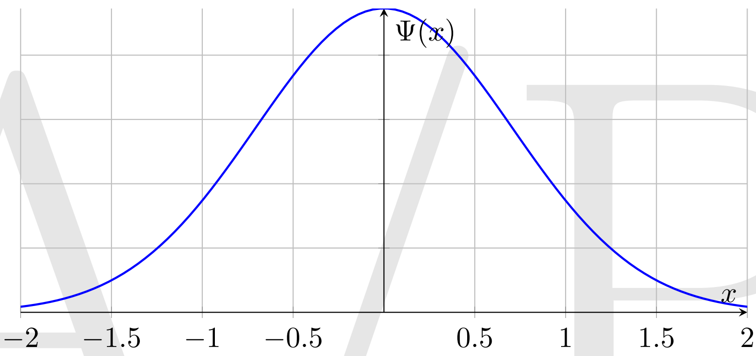
Introduction:

The uncertainty principle is a fundamental concept in quantum mechanics, indicating that the position and momentum of a particle cannot both be precisely determined simultaneously. For a Gaussian wave packet, the product of the uncertainties in position and momentum is minimized, reaching the value $\Delta x \Delta p = \frac{\hbar}{2}$.

The given wave function is a Gaussian wave packet, which can be represented as:

$$\Psi(x) = \left(\frac{\pi}{2}\right)^{-1/4} e^{-ax^2/2}$$

The graph of this wave function is shown below:



Solution:

Given the wave function:

$$\Psi(x) = \left(\frac{\pi}{2}\right)^{-1/4} e^{-ax^2/2}$$

1. **Normalization:**

First, we verify that $\Psi(x)$ is normalized:

$$\int_{-\infty}^{\infty} |\Psi(x)|^2 \, dx = 1$$

$$\int_{-\infty}^{\infty} \left(\frac{\pi}{2}\right)^{-1/2} e^{-ax^2} \, dx = 1$$

Using the Gaussian integral $\int_{-\infty}^{\infty} e^{-ax^2} \, dx = \sqrt{\frac{\pi}{a}}$:

$$\left(\frac{\pi}{2}\right)^{-1/2} \cdot \sqrt{\frac{\pi}{a}} = 1 \implies a = 2$$

So, the normalized wave function is:

$$\Psi(x) = \left(\frac{2}{\pi}\right)^{1/4} e^{-x^2}$$

2. **Expectation Value of x:**

$$\langle x \rangle = \int_{-\infty}^{\infty} x|\Psi(x)|^2 \, dx$$

Since $|\Psi(x)|^2$ is an even function and x is an odd function, the integrand is an odd function, and thus:

$$\langle x \rangle = 0$$

3. **Expectation Value of x^2:**

$$\langle x^2 \rangle = \int_{-\infty}^{\infty} x^2|\Psi(x)|^2 \, dx$$

$$\langle x^2 \rangle = \left(\frac{2}{\pi}\right)^{1/2} \int_{-\infty}^{\infty} x^2 e^{-2x^2} \, dx$$

Using the integral $\int_{-\infty}^{\infty} x^2 e^{-2x^2} \, dx = \frac{1}{4}\sqrt{\frac{\pi}{2}}$:

$$\langle x^2 \rangle = \left(\frac{2}{\pi}\right)^{1/2} \cdot \frac{1}{4}\sqrt{\frac{\pi}{2}} = \frac{1}{8}$$

4. **Uncertainty in x:**

$$\Delta x = \sqrt{\langle x^2 \rangle - \langle x \rangle^2} = \sqrt{\frac{1}{8}} = \frac{1}{2\sqrt{2}}$$

5. **Expectation Value of p:**

$$\langle p \rangle = \int_{-\infty}^{\infty} \Psi^*(x) \left(-i\hbar \frac{\partial}{\partial x}\right) \Psi(x) \, dx$$

$$\langle p \rangle = 0 \quad \text{(since the wave function is real and even)}$$

6. **Expectation Value of p^2:**

$$\langle p^2 \rangle = \int_{-\infty}^{\infty} \Psi^*(x) \left(-\hbar^2 \frac{\partial^2}{\partial x^2}\right) \Psi(x) \, dx$$

$$\frac{\partial \Psi}{\partial x} = -2x \left(\frac{2}{\pi}\right)^{1/4} e^{-x^2}$$

$$\frac{\partial^2 \Psi}{\partial x^2} = \left(4x^2 - 2\right) \left(\frac{2}{\pi}\right)^{1/4} e^{-x^2}$$

$$\langle p^2 \rangle = \hbar^2 \left(\frac{2}{\pi}\right)^{1/2} \int_{-\infty}^{\infty} \left(4x^2 - 2\right) e^{-2x^2} \, dx$$

Using $\int_{-\infty}^{\infty} e^{-2x^2}\,dx = \sqrt{\frac{\pi}{2}}$ and $\int_{-\infty}^{\infty} x^2 e^{-2x^2}\,dx = \frac{1}{4}\sqrt{\frac{\pi}{2}}$:

$$\langle p^2 \rangle = \hbar^2 \left(\frac{2}{\pi}\right)^{1/2} \left(4 \cdot \frac{1}{4}\sqrt{\frac{\pi}{2}} - 2 \cdot \sqrt{\frac{\pi}{2}}\right) = \hbar^2$$

7. **Uncertainty in p:**

$$\Delta p = \sqrt{\langle p^2 \rangle - \langle p \rangle^2} = \hbar$$

8. **Uncertainty Relation:**

$$\Delta x \Delta p = \frac{1}{2\sqrt{2}} \cdot \hbar = \frac{\hbar}{2}$$

Conclusion:

We have calculated the uncertainties in position Δx and momentum Δp for the given wave function. The product $\Delta x \Delta p = \frac{\hbar}{2}$ verifies the Heisenberg uncertainty principle. This result illustrates that for a Gaussian wave packet, the uncertainties are minimized, reaching the lower bound of the uncertainty principle, $\Delta x \Delta p = \frac{\hbar}{2}$.

47 A beam of 12eV electron is incident on a potential barrier of height 25eV and width 0.05 nm. Calculate the transmission coefficient.

Introduction: The transmission coefficient (T) in quantum mechanics describes the probability of a particle tunneling through a potential barrier. It is a key concept in the study of quantum tunneling and is particularly relevant in scenarios involving potential barriers that are higher than the energy of the incident particle.

Solution: Given:

1. Energy of the electron, $E = 12\,\text{eV}$
2. Height of the potential barrier, $V_0 = 25\,\text{eV}$
3. Width of the potential barrier, $a = 0.05\,\text{nm} = 0.05 \times 10^{-9}\,\text{m}$

The transmission coefficient for a rectangular potential barrier can be calculated using the formula:

$$T = \exp\left(-2\kappa a\right)$$

where

$$\kappa = \frac{\sqrt{2m(V_0 - E)}}{\hbar}$$

Here,

1. m is the mass of the electron ($9.11 \times 10^{-31}\,\text{kg}$)
2. $\hbar$ is the reduced Planck's constant ($1.055 \times 10^{-34}\,\text{Js}$)

First, calculate κ:

$$\kappa = \frac{\sqrt{2 \times 9.11 \times 10^{-31}\,\text{kg} \times (25\,\text{eV} - 12\,\text{eV}) \times 1.602 \times 10^{-19}\,\text{J/eV}}}{1.055 \times 10^{-34}\,\text{Js}}$$

Simplify inside the square root:

$$\kappa = \frac{\sqrt{2 \times 9.11 \times 10^{-31} \times 13 \times 1.602 \times 10^{-19}}}{1.055 \times 10^{-34}}$$

Calculate the values:

$$\kappa = \frac{\sqrt{2 \times 9.11 \times 10^{-31} \times 13 \times 1.602 \times 10^{-19}}}{1.055 \times 10^{-34}} = \frac{\sqrt{3.785 \times 10^{-30}}}{1.055 \times 10^{-34}} \approx 5.98 \times 10^{9}\,\text{m}^{-1}$$

Now, calculate the transmission coefficient T:

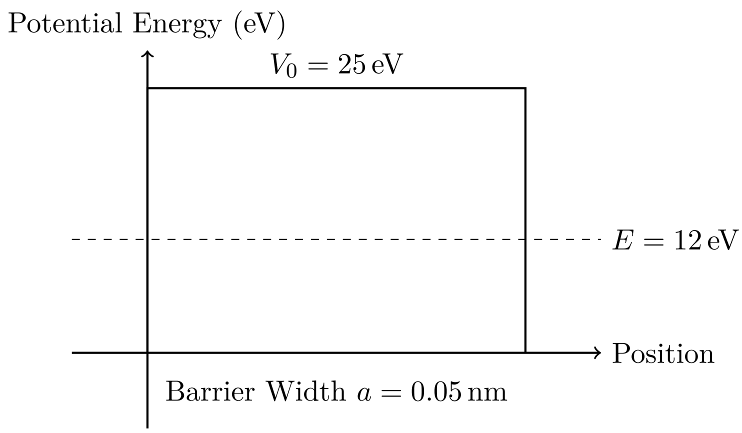
$$T = \exp\left(-2 \times 5.98 \times 10^{9} \times 0.05 \times 10^{-9}\right)$$

Simplify the exponent:

$$T = \exp\left(-2 \times 5.98 \times 0.05\right) = \exp\left(-0.598\right)$$

Finally, calculate T:

$$T \approx 0.55$$



Conclusion: The transmission coefficient T for the given potential barrier is approximately 0.55. This means that there is a 55% probability for the electron to tunnel through the potential barrier. Quantum tunneling has significant applications in modern physics, including the functioning of tunnel diodes and the nuclear fusion process in stars.

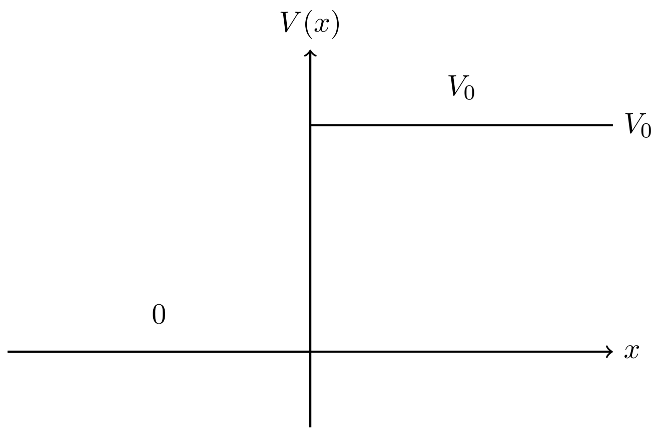
48 Solve the Schrödinger equation for a step potential and calculate the transmission and reflection coefficients for the case when the kinetic energy of the particle E_0 is greater than the potential energy V (i.e., $E_0 > V$).

Introduction:
The step potential is a fundamental problem in quantum mechanics that illustrates the behavior of a particle encountering a sudden change in potential energy. This problem is essential for understanding phenomena such as quantum tunneling and reflection.

Consider a particle encountering a step potential:

$$V(x) = \begin{cases} 0 & \text{for } x < 0 \\ V_0 & \text{for } x \geq 0 \end{cases}$$

Below is a diagram illustrating the step potential:



Solution:

Consider a particle encountering a step potential:

$$V(x) = \begin{cases} 0 & \text{for } x < 0 \\ V_0 & \text{for } x \geq 0 \end{cases}$$

The Schrödinger equation in regions where $V(x)$ is constant is:

$$-\frac{\hbar^2}{2m}\frac{d^2\psi(x)}{dx^2} + V(x)\psi(x) = E\psi(x)$$

For $x < 0$ (Region I), where $V(x) = 0$:

$$-\frac{\hbar^2}{2m}\frac{d^2\psi(x)}{dx^2} = E_0\psi(x)$$

The general solution is:

$$\psi_I(x) = Ae^{ik_1x} + Be^{-ik_1x}$$

where:

$$k_1 = \sqrt{\frac{2mE_0}{\hbar^2}}$$

For $x \geq 0$ (Region II), where $V(x) = V_0$:

$$-\frac{\hbar^2}{2m}\frac{d^2\psi(x)}{dx^2} + V_0\psi(x) = E_0\psi(x)$$

This simplifies to:

$$\frac{d^2\psi(x)}{dx^2} = k_2^2\psi(x)$$

where:

$$k_2 = \sqrt{\frac{2m(E_0 - V_0)}{\hbar^2}}$$

The general solution is:

$$\psi_{II}(x) = Ce^{ik_2x}$$

Since we consider the particle coming from the left and moving to the right, there will be no wave traveling to the left in Region II ($D = 0$):

$$\psi_{II}(x) = Ce^{ik_2x}$$

Boundary Conditions:

At $x = 0$, the wavefunctions and their first derivatives must be continuous:

$$\psi_I(0) = \psi_{II}(0)$$

$$\left.\frac{d\psi_I}{dx}\right|_{x=0} = \left.\frac{d\psi_{II}}{dx}\right|_{x=0}$$

Applying these conditions:

1. Continuity of wavefunction:

$$A + B = C$$

2. Continuity of derivative:

$$ik_1 A - ik_1 B = ik_2 C$$

Solving these equations for A, B, and C:

From the first equation:

$$C = A + B$$

Substituting into the second equation:

$$ik_1 A - ik_1 B = ik_2(A + B)$$

Rearranging:

$$k_1 A - k_1 B = k_2 A + k_2 B$$

$$(k_1 - k_2)A = (k_1 + k_2)B$$

$$\frac{A}{B} = \frac{k_1 + k_2}{k_1 - k_2}$$

Therefore, the reflection coefficient R is:

$$R = \left|\frac{B}{A}\right|^2 = \left|\frac{k_1 - k_2}{k_1 + k_2}\right|^2$$

The transmission coefficient T is given by:

$$T = \left|\frac{C}{A}\right|^2 = \left|\frac{2k_1}{k_1 + k_2}\right|^2$$

Conclusion:
For a particle encountering a step potential with $E_0 > V_0$, the transmission and reflection coefficients are given by:

$$R = \left|\frac{k_1 - k_2}{k_1 + k_2}\right|^2, \quad T = \left|\frac{2k_1}{k_1 + k_2}\right|^2$$

These coefficients describe the probability of the particle being reflected or transmitted at the potential step.

An application of the step potential is seen in the behavior of electrons in semiconductor devices, where they encounter potential barriers at junctions, leading to phenomena like tunneling and reflection that are crucial for the operation of diodes and transistors.

49 Write the wave functions for a particle on both sides of a step potential, for $E > V_0$:

$$V(x) = \begin{cases} V_0, & x > 0 \\ 0, & x < 0 \end{cases}$$

Interpret the results physically.

Introduction: This problem involves a quantum particle encountering a 1D step potential. The potential energy function is piecewise constant, and the total energy of the particle satisfies $E > V_0$. Our goal is to determine the wavefunctions in both regions and interpret the behavior of the particle, including any reflection or transmission effects due to the step.

Solution:

The time-independent Schrödinger equation is given by:

$$-\frac{\hbar^2}{2m}\frac{d^2\psi(x)}{dx^2} + V(x)\psi(x) = E\psi(x).$$

We solve this separately in regions I ($x < 0$) and II ($x > 0$).

Region I: $x < 0$ (where $V(x) = 0$)

The Schrödinger equation becomes:

$$-\frac{\hbar^2}{2m}\frac{d^2\psi}{dx^2} = E\psi,$$

which simplifies to:

$$\frac{d^2\psi}{dx^2} + k_1^2\psi = 0, \quad \text{where } k_1 = \frac{\sqrt{2mE}}{\hbar}.$$

General solution:

$$\psi_I(x) = Ae^{ik_1x} + Be^{-ik_1x}.$$

Here, Ae^{ik_1x} represents the incident wave, and Be^{-ik_1x} is the reflected wave.

Region II: $x > 0$ (where $V(x) = V_0$)

The Schrödinger equation becomes:

$$-\frac{\hbar^2}{2m}\frac{d^2\psi}{dx^2} + V_0\psi = E\psi \quad \Rightarrow \quad \frac{d^2\psi}{dx^2} + k_2^2\psi = 0,$$

where

$$k_2 = \frac{\sqrt{2m(E - V_0)}}{\hbar}.$$

General solution:

$$\psi_{II}(x) = Ce^{ik_2x}.$$

We exclude the term De^{-ik_2x} because it would represent a wave incoming from $x \to \infty$, which contradicts the physical setup of a wave incident from the left.

Boundary Conditions:

Continuity of the wavefunction and its derivative at $x = 0$:

$$\psi_I(0) = \psi_{II}(0) \Rightarrow A + B = C,$$
$$\psi'_I(0) = \psi'_{II}(0) \Rightarrow ik_1(A - B) = ik_2 C.$$

Solving this system:

$$A + B = C$$
$$k_1(A - B) = k_2(A + B)$$

Solving for B/A and C/A:

$$\frac{B}{A} = \frac{k_1 - k_2}{k_1 + k_2},$$
$$\frac{C}{A} = \frac{2k_1}{k_1 + k_2}.$$

Interpretation:

Even though the energy $E > V_0$, the particle has a finite probability of being reflected. The reflection coefficient R and transmission coefficient T are given by:

$$R = \left|\frac{B}{A}\right|^2 = \left(\frac{k_1 - k_2}{k_1 + k_2}\right)^2,$$
$$T = \frac{k_2}{k_1}\left|\frac{C}{A}\right|^2 = \frac{4k_1 k_2}{(k_1 + k_2)^2}.$$

Note that $R + T = 1$, as required by probability conservation.

Physically, even when the particle has enough energy to surpass the potential step, there is a non-zero probability of reflection due to the abrupt change in potential, a purely quantum mechanical phenomenon with no classical analog.

Conclusion: The wavefunctions in each region are:

$$\psi_I(x) = Ae^{ik_1 x} + Be^{-ik_1 x}, \quad x < 0,$$
$$\psi_{II}(x) = Ce^{ik_2 x}, \quad x > 0.$$

Despite having energy $E > V_0$, the particle experiences partial reflection and transmission due to the discontinuity in potential. This highlights the wave nature of particles in quantum mechanics and the non-classical behavior at potential boundaries.

50 Normalize the 1s state of the hydrogen atom in the ground state and calculate the expectation value of position

Introduction:

The wave function for the hydrogen atom in the 1s state is given by:

$$\psi_{100}(r) = Ae^{-r/a_0}$$

where $a_0 = \frac{\hbar^2}{me^2}$ is the Bohr radius. To normalize this wave function, we need to find the constant A such that the total probability of finding the electron anywhere in space is equal to 1:

$$\int |\psi_{100}(r)|^2 d^3r = 1$$

After normalizing the wave function, we will calculate the expectation value of the position $\langle r \rangle$ in this state.

Solution:

1. Normalization of the wave function:

The wave function must satisfy the normalization condition:

$$\int_0^{2\pi} \int_0^{\pi} \int_0^{\infty} |\psi_{100}(r)|^2 r^2 \sin\theta \, dr \, d\theta \, d\phi = 1$$

Substituting $\psi_{100}(r) = Ae^{-r/a_0}$, we have:

$$\int_0^{2\pi} d\phi \int_0^{\pi} \sin\theta \, d\theta \int_0^{\infty} |A|^2 e^{-2r/a_0} r^2 \, dr = 1$$

The angular integrals are straightforward:

$$\int_0^{2\pi} d\phi = 2\pi, \quad \int_0^{\pi} \sin\theta \, d\theta = 2$$

Thus, the normalization condition simplifies to:

$$4\pi |A|^2 \int_0^{\infty} e^{-2r/a_0} r^2 \, dr = 1$$

To solve the radial integral, we use the integral formula:

$$\int_0^{\infty} x^2 e^{-ax} dx = \frac{2}{a^3} \quad \text{for } a > 0$$

Setting $x = r/a_0$ and $a = 2/a_0$, the integral becomes:

$$\int_0^{\infty} e^{-2r/a_0} r^2 \, dr = \frac{a_0^3}{4}$$

Substituting this into the normalization condition:

$$4\pi |A|^2 \cdot \frac{a_0^3}{4} = 1$$

$$|A|^2 \cdot \pi a_0^3 = 1 \quad \Rightarrow \quad |A| = \frac{1}{\sqrt{\pi a_0^3}}$$

Thus, the normalized wave function is:

$$\psi_{100}(r) = \frac{1}{\sqrt{\pi a_0^3}} e^{-r/a_0}$$

2. Expectation value of position $\langle r \rangle$:

The expectation value of the position is given by:

$$\langle r \rangle = \int_0^{2\pi} \int_0^{\pi} \int_0^{\infty} \psi_{100}^*(r) r \psi_{100}(r) r^2 \sin\theta \, dr \, d\theta \, d\phi$$

Substituting the normalized wave function:

$$\langle r \rangle = \frac{1}{\pi a_0^3} \int_0^{2\pi} d\phi \int_0^{\pi} \sin\theta \, d\theta \int_0^{\infty} r^3 e^{-2r/a_0} \, dr$$

The angular integrals again give:

$$\int_0^{2\pi} d\phi = 2\pi, \quad \int_0^{\pi} \sin\theta \, d\theta = 2$$

Thus:

$$\langle r \rangle = \frac{2}{a_0^3} \int_0^{\infty} r^3 e^{-2r/a_0} \, dr$$

To solve the radial integral, we use the integral formula:

$$\int_0^{\infty} x^3 e^{-ax} dx = \frac{6}{a^4} \quad \text{for } a > 0$$

Setting $x = r/a_0$ and $a = 2/a_0$, the integral becomes:

$$\int_0^{\infty} r^3 e^{-2r/a_0} \, dr = \frac{3a_0^4}{8}$$

Substituting this into the expectation value:

$$\langle r \rangle = \frac{2}{a_0^3} \cdot \frac{3a_0^4}{8} = \frac{3a_0}{4}$$

Conclusion:

The normalized wave function for the hydrogen atom in the 1s state is:

$$\psi_{100}(r) = \frac{1}{\sqrt{\pi a_0^3}} e^{-r/a_0}$$

The expectation value of the position $\langle r \rangle$ in this state is:

$$\langle r \rangle = \frac{3a_0}{2}$$

This result indicates that the average distance of the electron from the nucleus in the 1s state is 1.5 times the Bohr radius, highlighting the probabilistic nature of quantum mechanics where the electron does not occupy a fixed orbit but has a distribution of possible positions.