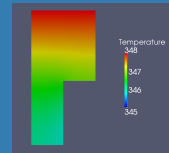
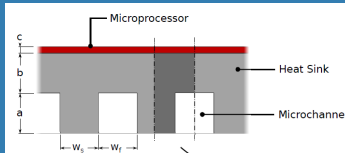




German Research School
for Simulation Sciences

Unstructured FEM: 2D Unsteady Heat Diffusion



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Contents

Objective

Finite Element Method

FEM Formulation

Implementation

Verification

Optimization

Parallel Computing

Application

Summary



Objective

- To develop a parallel finite element solver that uses unstructured grid to solve a 2D transient heat diffusion problem.
- Application to a realistic engineering heat transfer problem, Heat diffusion in Microchannels in the heat sink of microprocessors.



Contents

Objective

Finite Element Method

FEM Formulation

Implementation

Verification

Optimization

Parallel Computing

Application

Summary



General Steps

1. Preprocessing

- Define a domain of interest
- Discretization of domain into finite elements

2. Solution

- Define shape functions to represent the behaviour of elements
- Construct local stiffness matrix and then ultimately global stiffness matrix
- Apply boundary condition, initial conditions and forcing function
- Solving linear system of algebraic equations

3. Postprocessing

- Analyse the solution



Contents

Objective

Finite Element Method

FEM Formulation

Implementation

Verification

Optimization

Parallel Computing

Application

Summary



2D Unsteady Heat Diffusion Equation

$$\frac{\partial T}{\partial t} - \alpha \nabla^2 T = f, \text{ in } \Omega \in \mathbb{R}^2$$

where α is diffusion coefficient.

$$\alpha = \frac{k}{\rho c_p}$$

k is heat conductivity, ρ is the density and c_p is the specific heat capacity of the material.

$$f = \frac{\dot{q}}{\rho c_p}$$

where \dot{q} is heat generation per unit volume.

Initial and Boundary Conditions

- Initial condition: $T(x, y, 0) = T_0(x, y)$
- Boundary conditions
 - Dirichlet:
 $T(x, 0, t) = T_1, T(x, L, t) = T_2, T(0, y, t) = T_3,$
 $T(L, y, t) = T_4$
 - Neumann:
 $k \left(n_x \frac{\partial T}{\partial x} + n_y \frac{\partial T}{\partial y} \right) = q_0$
 - Robin (Mixed):
 $h(T - T_{\text{inf}}) = -k \left(n_x \frac{\partial T}{\partial x} + n_y \frac{\partial T}{\partial y} \right)$

Weighted Residual Method

$$\int_{\Omega} (w \frac{\partial T}{\partial t} + \alpha w \nabla^2 T) d\Omega = \int_{\Omega} w f d\Omega$$

where w is a weighing function or test function.
Integrating by parts,

$$\int_{\Omega} (w \frac{\partial T}{\partial t} + \alpha \nabla w \cdot \nabla T) d\Omega = \int_{\Omega} w f d\Omega + \frac{1}{\rho c_p} \int_{\Gamma} w \underbrace{(k \hat{n} \cdot \nabla T)}_{k(n_x \frac{\partial T}{\partial x} + n_y \frac{\partial T}{\partial y})} d\Gamma$$

Approximation

- Approximate the temperature as a piecewise linear combination of shape functions $S_j(x, y)$.
- These shape functions remain constant over time. The coefficients $T_j(t)$ are time dependent.

$$T(x, y, t) = \sum_{j=1}^{nn} T_j(t) S_j(x, y)$$

- Galerkin Finite Element Method:

$$w = S_i(x)$$

Approximation [contd.]

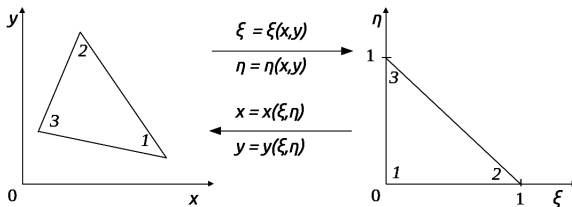
$$\underbrace{\sum_{j=1}^{nn} \int_{\Omega} (S_i S_j) d\Omega \frac{\partial T_j}{\partial t}}_{[M][\dot{T}]} - \underbrace{\sum_{j=1}^{nn} \left[\int_{\Omega} \alpha \left(\frac{\partial S_i}{\partial x} \frac{\partial S_j}{\partial x} + \frac{\partial S_i}{\partial y} \frac{\partial S_j}{\partial y} \right) d\Omega \right] T_j}_{[K][T]}$$

$$= \underbrace{\int_{\Omega} S_i f d\Omega}_{[F]} + \underbrace{\frac{1}{\rho c_p} \int_{\Gamma} S_i (k \hat{n} \cdot \nabla T) d\Gamma}_{[B]}$$

$$i, j = 1, 2, \dots, nn$$

$$[K]_{ij} = \int_{\Omega} \alpha \left(\frac{\partial S_i}{\partial x} \frac{\partial S_j}{\partial x} + \frac{\partial S_i}{\partial y} \frac{\partial S_j}{\partial y} \right) d\Omega$$

Triangular Elements



Co-ordinate transformation

Shape functions:

$$S_1(\xi, \eta) = 1 - \xi - \eta, \quad S_2(\xi, \eta) = \xi, \quad S_3(\xi, \eta) = \eta$$

$$\frac{\partial S_1}{\partial \xi} = -1, \frac{\partial S_2}{\partial \xi} = 1, \frac{\partial S_3}{\partial \xi} = 0, \frac{\partial S_1}{\partial \eta} = -1, \frac{\partial S_2}{\partial \eta} = 0, \frac{\partial S_3}{\partial \eta} = 1$$

$$\frac{\partial S_i}{\partial \xi} = \frac{\partial S_i}{\partial x} \frac{\partial x}{\partial \xi} + \frac{\partial S_i}{\partial y} \frac{\partial y}{\partial \xi}, \frac{\partial S_i}{\partial \eta} = \frac{\partial S_i}{\partial x} \frac{\partial x}{\partial \eta} + \frac{\partial S_i}{\partial y} \frac{\partial y}{\partial \eta}$$

$$\begin{bmatrix} \frac{\partial S_i}{\partial \xi} \\ \frac{\partial S_i}{\partial \eta} \end{bmatrix} = \underbrace{\begin{bmatrix} \frac{\partial x}{\partial \xi} & \frac{\partial y}{\partial \xi} \\ \frac{\partial x}{\partial \eta} & \frac{\partial y}{\partial \eta} \end{bmatrix}}_{\text{Jacobian}} \begin{bmatrix} \frac{\partial S_i}{\partial x} \\ \frac{\partial S_i}{\partial y} \end{bmatrix}$$

Isoparametric formulation:

$$x(\xi, \eta) = \sum_{j=1}^{nen} x_j^e S_j(\xi, \eta), \quad y(\xi, \eta) = \sum_{j=1}^{nen} y_j^e S_j(\xi, \eta)$$

Jacobian:

$$[J] = \begin{bmatrix} \frac{\partial S_1}{\partial \xi} & \frac{\partial S_2}{\partial \xi} & \frac{\partial S_3}{\partial \xi} \\ \frac{\partial S_1}{\partial \eta} & \frac{\partial S_2}{\partial \eta} & \frac{\partial S_3}{\partial \eta} \end{bmatrix} \begin{bmatrix} x_1^e & y_1^e \\ x_2^e & y_2^e \\ x_3^e & y_3^e \end{bmatrix} = \begin{bmatrix} \frac{\partial x}{\partial \xi} & \frac{\partial y}{\partial \xi} \\ \frac{\partial x}{\partial \eta} & \frac{\partial y}{\partial \eta} \end{bmatrix}$$

But we need $[J]^{-1}$ which is the inverse mapping.

$$\begin{bmatrix} \frac{\partial S_i}{\partial x} \\ \frac{\partial S_i}{\partial y} \end{bmatrix} = \underbrace{\begin{bmatrix} \frac{\partial \xi}{\partial x} & \frac{\partial \eta}{\partial x} \\ \frac{\partial \xi}{\partial y} & \frac{\partial \eta}{\partial y} \end{bmatrix}}_{[J]^{-1}} \underbrace{\begin{bmatrix} \frac{\partial S_i}{\partial \xi} \\ \frac{\partial S_i}{\partial \eta} \end{bmatrix}}_{\text{known}}$$

Calculation of element level stiffness matrix:

$$K_{ij}^e = \int_{\Omega_e} \alpha \left(\frac{\partial S_i}{\partial x} \frac{\partial S_j}{\partial x} + \frac{\partial S_i}{\partial y} \frac{\partial S_j}{\partial y} \right) dx dy \quad i, j = 1, 2, \dots, n_{en}$$

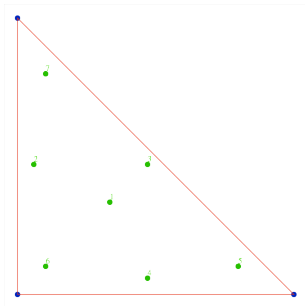
Integral transformation to master element:

$$K_{ij}^e = \int_0^1 \int_0^{1-\xi} \alpha[f_k(\xi, \eta)] |J| d\xi d\eta$$

$$f_k(\xi, \eta) = \left(\frac{\partial S_i}{\partial \xi} \frac{\partial \xi}{\partial x} + \frac{\partial S_i}{\partial \eta} \frac{\partial \eta}{\partial x} \right) \left(\frac{\partial S_j}{\partial \xi} \frac{\partial \xi}{\partial x} + \frac{\partial S_j}{\partial \eta} \frac{\partial \eta}{\partial x} \right) \\ + \left(\frac{\partial S_i}{\partial \xi} \frac{\partial \xi}{\partial y} + \frac{\partial S_i}{\partial \eta} \frac{\partial \eta}{\partial y} \right) \left(\frac{\partial S_j}{\partial \xi} \frac{\partial \xi}{\partial y} + \frac{\partial S_j}{\partial \eta} \frac{\partial \eta}{\partial y} \right)$$

$$i, j = 1, 2, \dots, n_{en}$$

Numerical Integration using 7 point Gauss Quadrature rule:



$$K_{ij}^e \simeq \sum_{l=1}^{nGQP} f_k(\xi_l, \eta_l) w_l$$

Solution

In matrix form : $[M][\dot{T}] + [K][T] = [F] + [B]$
Forward Euler explicit differencing in time:

$$\frac{dT}{dt} \Big|_s = \frac{T^{s+1} - T^s}{\Delta t} + O(\Delta t)$$

$$[M] \frac{[T]^{s+1} - [T]^s}{\Delta t} + [K][T]^s = [F] + [B]$$

$$[M][T]^{s+1} = [M][T]^s + \Delta t([F] + [B] - [K][T]^s)$$

$$[M][T]^{s+1} = [RHS]^s$$

Lumping of Mass Matrix

The inversion of Mass matrix can be avoided if we do a physical approximation called lumping of mass matrix.

Consistent element mass matrix: $[M_{cij}]$

Lumped element mass matrix: $[M_{Lij}]$

$$M_{Lii} = M_{cii} \frac{\sum_{i=1}^n \sum_{j=1}^n M_{cij}}{\sum_{j=1}^n M_{cjj}}$$

This is correct only for linear elements.

$[M_{Lij}]$ is a diagonal matrix.

Boundary Integral

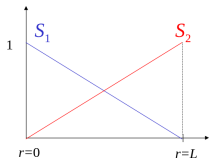
- Dirichlet: $[B]_j = T_w$, $[K]_{jj} = 1$, $[K]_{ij} = 0 \forall i \neq j$
For example, if for some element, node temperatures T_1 and T_3 are known then $[B]$ and $[K]$ for that element can be written as:

$$[B] = \begin{bmatrix} T_w \\ B_2 \\ T_w \end{bmatrix} \text{ and } [K] = \begin{bmatrix} 1 & 0 & 0 \\ * & * & * \\ 0 & 0 & 1 \end{bmatrix}$$

* indicate non-zero values.

Boundary Integral [contd.]

- Neumann (Flux Variable(FV) = q_{in}):



$$B_1^e = \frac{1}{\rho c_p} \int_{\Gamma} S_1(FV) d\Gamma = \frac{1}{\rho c_p} \int_0^L \left(1 - \frac{r}{L}\right) q_{in} dr = \frac{q_{in} L}{2\rho c_p}$$

$$B_2^e = \frac{1}{\rho c_p} \int_{\Gamma} S_2(FV) d\Gamma = \frac{1}{\rho c_p} \int_0^L \frac{r}{L} q_{in} dr = \frac{q_{in} L}{2\rho c_p}$$

$$[B] = \begin{bmatrix} \frac{q_{in} L}{2\rho c_p} \\ \frac{q_{in} L}{2\rho c_p} \\ 0 \end{bmatrix}$$

Boundary Integral [contd.]

- Robin(Mixed type): Flux variable (FV) can be expressed as

$$-k\hat{n}.\nabla T = h(T - T_{\infty})$$

$$FV = aPV + b \text{ where } a = -h \text{ and } b = hT_{\infty}$$

$$B_1^e = \frac{1}{\rho c_p} \int_{\Gamma} S_1(FV) d\Gamma = \frac{1}{\rho c_p} \int_0^L \left(1 - \frac{r}{L}\right) (aPV + b) dr$$

$$B_2^e = \frac{1}{\rho c_p} \int_{\Gamma} S_2(FV) d\Gamma = \frac{1}{\rho c_p} \int_0^L \frac{r}{L} (aPV + b) dr$$

Primary variable (PV) temperature:

$$PV = S_1 T_1 + S_2 T_2$$

Boundary Integral [contd.]

$$\begin{aligned} B_1^e &= \frac{1}{\rho c_p} \int_0^L \left(1 - \frac{r}{L}\right) \left[a \left(\left(1 - \frac{r}{L}\right) T_1 + \frac{r}{L} T_2 \right) + b \right] dr \\ &= \frac{L}{6\rho c_p} (3b + 2aT_1 + aT_2) \end{aligned}$$

$$\begin{aligned} B_2^e &= \frac{1}{\rho c_p} \int_0^L \frac{r}{L} \left[a \left(\left(1 - \frac{r}{L}\right) T_1 + \frac{r}{L} T_2 \right) + b \right] dr \\ &= \frac{L}{6\rho c_p} (3b + aT_1 + 2aT_2) \end{aligned}$$

Boundary Integral [contd.]

Element boundary vector can be written as:

$$[B] = \begin{bmatrix} \frac{L}{6\rho c_p} (3b + 2aT_1 + aT_2) \\ \frac{L}{6\rho c_p} (3b + aT_1 + 2aT_2) \\ 0 \end{bmatrix}$$

Coefficients of T_1 and T_2 transferred to K. Modified $[B]$ and corresponding $[K]$ can be written as:

$$[B] = \begin{bmatrix} \frac{Lb}{2\rho c_p} \\ \frac{Lb}{2\rho c_p} \\ 0 \end{bmatrix} \quad \text{and} \quad [K] = \begin{bmatrix} \left(* - \frac{aL}{3\rho c_p}\right) & \left(* - \frac{aL}{6\rho c_p}\right) & * \\ \left(* - \frac{aL}{6\rho c_p}\right) & \left(* - \frac{aL}{3\rho c_p}\right) & * \\ * & * & * \end{bmatrix}$$



Contents

Objective

Finite Element Method

FEM Formulation

Implementation

Verification

Optimization

Parallel Computing

Application

Summary

Program Flow Chart

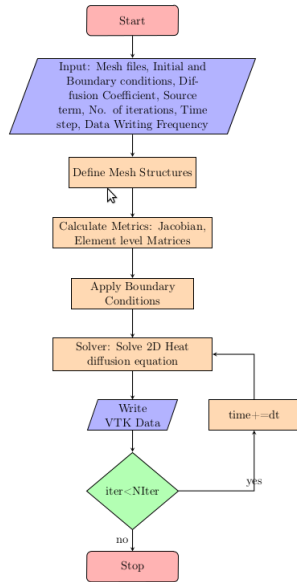
1 Preprocessing Stage:

- Reading settings file
- Reading mesh data

2 Solution Stage:

- Calculating Jacobian matrix for all elements
- Calculating element stiffness matrix for all elements
- Applying boundary conditions
- Explicit Solver

3 Post Processing Stage





Contents

Objective

Finite Element Method

FEM Formulation

Implementation

Verification

Optimization

Parallel Computing

Application

Summary

Case 1: 1D

A simple 1D case is taken as an example. Domain: $x \in [0, L]$
 $L = 2$.

Initial condition: $T_i = T(x, 0) = 1000K$

Boundary Conditions: $T_L = T(0, t) = 0$, $T_R = T(L, t) = 0$,
 T_{top} and T_{bottom} are insulated

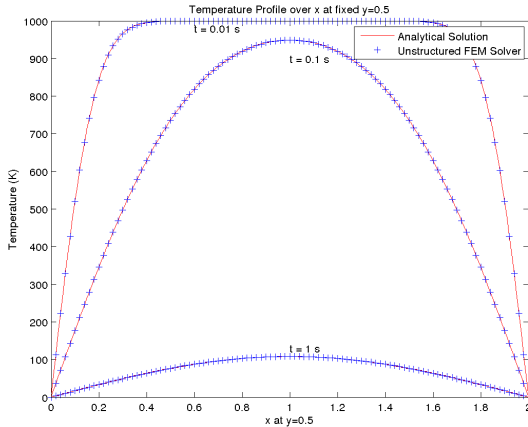
Analytical Solution:

$$T(x, t) = \sum_{n=1}^{\infty} B_n \sin\left(\frac{n\pi x}{L}\right) e^{-\frac{n^2 \pi^2 \alpha t}{L^2}}$$

where α is diffusion coefficient and

$$B_n = -T_i \frac{2(-1 + (-1)^n)}{n\pi}$$

Comparison with Analytical solution



Case 2: 2D Quenching of a Billet[1]

Domain: A long billet of rectangular cross section $2a \times 2b$ ($a=2, b=1$). Because of symmetry, analyze one quarter of the cross section

Shifted temperature: $\theta(x, y, t) = T(x, y, t) - T_\infty$

Material properties: $\alpha = 1, \rho = 1, c_p = 1, h = 1$

Initial condition: $\theta(x, y, 0) = T_i - T_\infty = \theta_i, T_i = 1000,$
 $T_\infty = 300$

Boundary Conditions:

$$x = 0: \frac{\partial \theta}{\partial x} = 0$$

$$x = a: \frac{\partial \theta}{\partial x} = -\frac{h}{k}\theta(a, y, t)$$

$$y = 0: \frac{\partial \theta}{\partial y} = 0$$

$$y = b: \frac{\partial \theta}{\partial y} = -\frac{h}{k}\theta(x, b, t)$$



Analytical solution:

$$\frac{\theta}{\theta_i} = 4 \sum_{n=1}^{\infty} \sum_{m=1}^{\infty} e^{(-\alpha(\lambda_n^2 + \beta_m^2)t)} \times$$

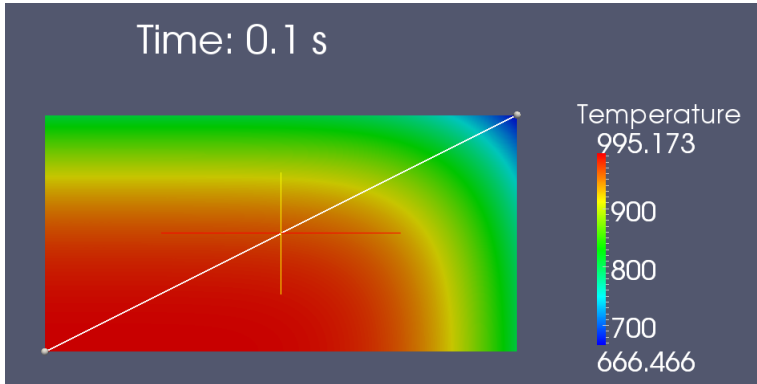
$$\frac{\sin(\lambda_n a) \cos(\lambda_n x) \sin(\beta_m b) \cos(\beta_m y)}{[\lambda_n a + \sin(\lambda_n a) \cos(\lambda_n a)][\beta_m b + \sin(\beta_m b) \cos(\beta_m b)]}$$

where α is diffusion coefficient and the eigenvalues λ_n and β_m are the roots of transcendental equations.

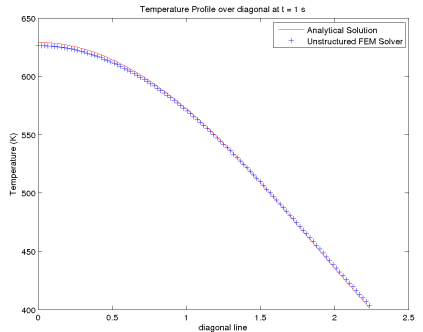
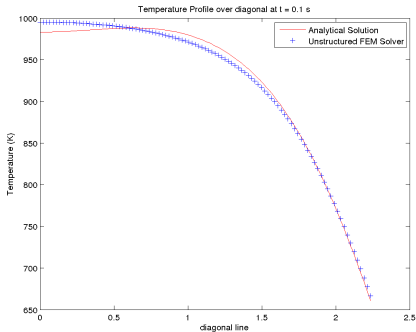
$$\lambda_n \tan(\lambda_n a) = \frac{h}{k} \qquad \beta_m \tan(\beta_m b) = \frac{h}{k}$$

where k is conductivity of the billet material.

Results



Results [contd.]





Contents

Objective

Finite Element Method

FEM Formulation

Implementation

Verification

Optimization

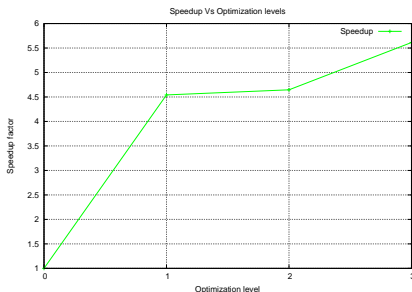
Parallel Computing

Application

Summary

Optimization

- Manual Optimization
 - Local arrays and variables
 - Loop unrolling
- Automatic Optimization
 - O1
 - O2
 - O3





Contents

Objective

Finite Element Method

FEM Formulation

Implementation

Verification

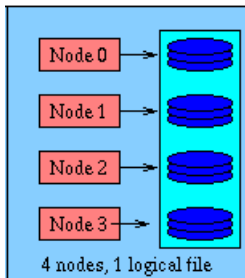
Optimization

Parallel Computing

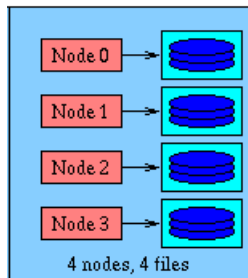
Application

Summary

Parallelizing



Parallel I/O 1 [3]



Parallel I/O 2 [3]

Structure of Mesh files (Illustration)

mxyz

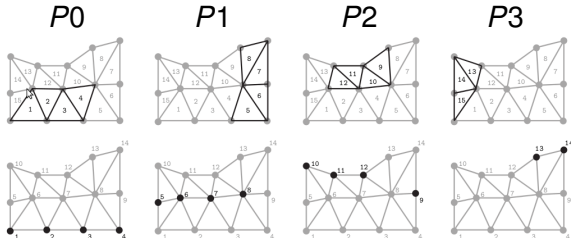
0	0
0.5	0
1	0
2	0
2	0.25
2	0.5
1.5	0.75
1	0.75
0.5	0.75
0	0.5
0	0.25
0.5	0.25

mien

8	11	10
9	2	8
1	4	7
3	5	8
5	11	10
4	7	9
6	8	3
10	3	9
7	1	2
5	4	8
11	9	7

First approach

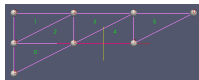
- All the node information not present on individual processors. Mesh data needs to be communicated before the actual FEM code starts. This adds to communication overhead.



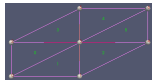
Division of elements and nodes [4]

Second approach

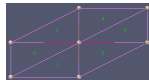
- All mesh data present on the same processor.



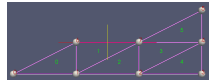
P0



P1



P2

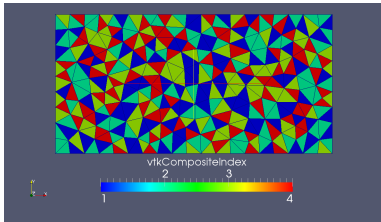


P3

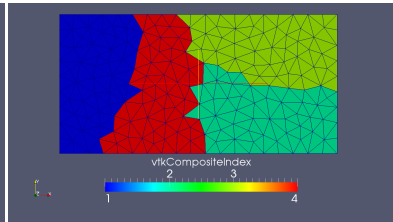
- A preprocessor is needed to generate the separate domains for each processor. Once the mesh is divided it can be reused in further simulations.

Mesh Partitioning

- Mesh partitioning tool needs to be used.



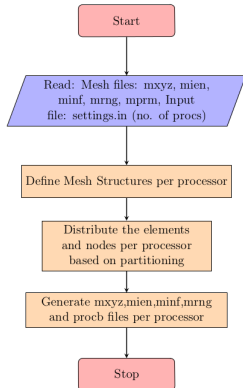
Before Mesh partitioning



After Mesh partitioning

Several public tools available, e.g. SCOTCH, CHACO, HARP, JOSTLE, METIS, etc. We used METIS.

Parallelizing with second approach



Contents of a sample probc file

0	0	-1	-1	-1	-1
1	0	-1	-1	-1	-1
2	0	-1	-1	-1	-1
3	1	1	-1	-1	-1
13	1	1	-1	-1	-1
14	1	1	-1	-1	-1
15	2	1	2	-1	-1
16	2	1	2	-1	-1

Preprocessor Flow Chart

Program Flow Chart

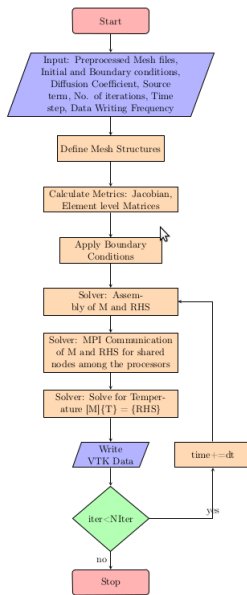
1 Preprocessing Stage:

- Reading settings and mesh data

2 Solution Stage:

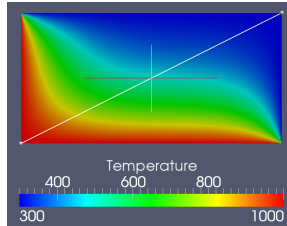
- Calculating mesh metrics
- Applying boundary conditions
- Assembly of $[M]$ and $\{RHS\}$
- Communicate $[M]$ and $\{RHS\}$ for boundary nodes
- Solve $[M]\{T\} = \{RHS\}$

3 Post Processing Stage



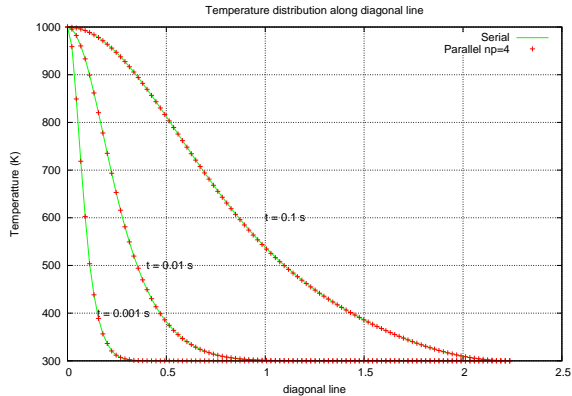
Verification of Parallel code

A rectangular domain is considered for heat diffusion problem. Left and bottom walls are kept at 1000 K while top and right walls are maintained at 300 K and the temperature profile is allowed to develop over time. The snapshots are taken at $t=0.001\text{s}$, $t=0.01\text{s}$ and $t=0.1\text{s}$.

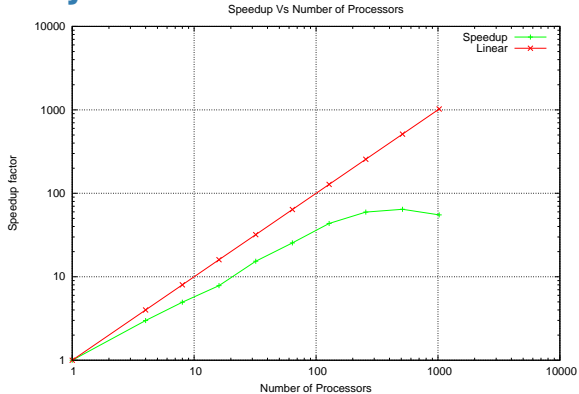


$t = 0.1\text{ s}$

Verification of Parallel code

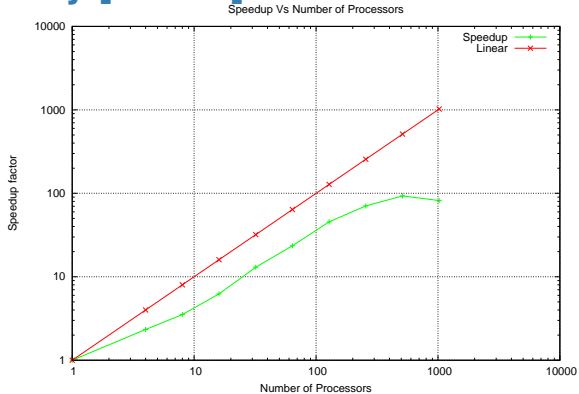


Scalability



Finestmesh: $ne = 516160$, $nn = 259082$

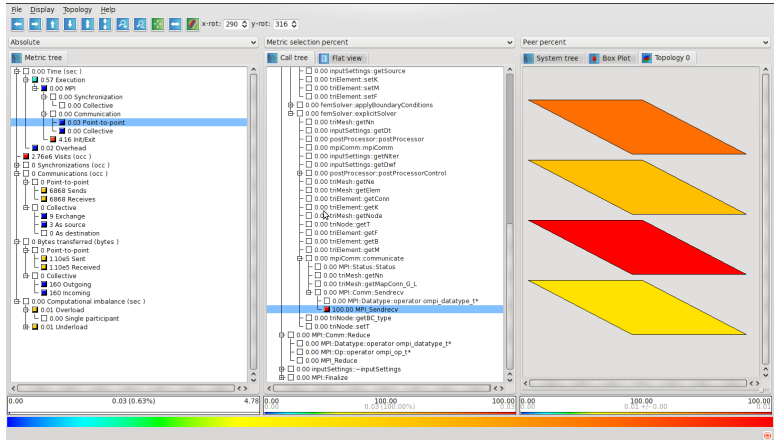
Scalability [contd.]



Superfine: ne = 945746, nn = 474074



Analysis using Scalasca





Contents

Objective

Finite Element Method

FEM Formulation

Implementation

Verification

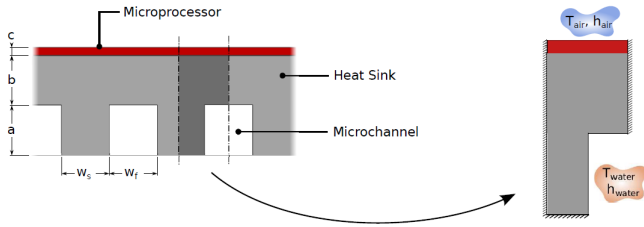
Optimization

Parallel Computing

Application

Summary

Application

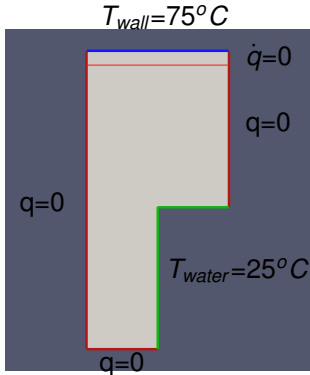


A heat sink for cooling computer chip, which is fabricated from copper with machined microchannels. Within these microchannels, water flows and carries away the heat dissipated by the chips.

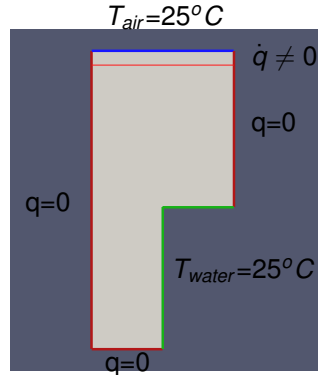
Physical properties and dimensions:

- Dimensions of the heat sink: $a = b = w_s = w_f = 0.2 \text{ mm}$, $c = 0.02 \text{ mm}$
- Thermal diffusivity of copper: $\alpha_{copper} = 1.11 \times 10^{-4} \text{ m}^2/\text{s}$
- Density of copper: $\rho = 8940 \text{ kg}/\text{m}^3$
- Specific heat capacity: $C_p = 385 \text{ J}/\text{kg}/\text{K}$
- Properties of cooling fluid: $T_{water} = 25^\circ \text{ C}$, $h_{water} = 30000 \text{ W}/\text{m}^2/\text{K}$
- Properties of ambient air: $T_{air} = 25^\circ \text{ C}$, $h_{air} = 2 \text{ W}/\text{m}^2/\text{K}$

Case 1: Steady state
temperature distribution

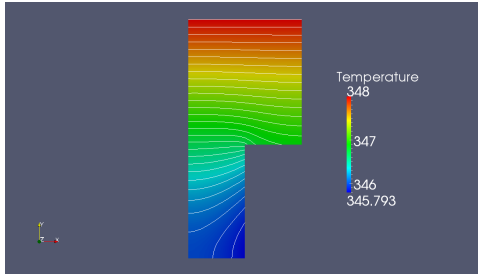


Case 2: Find \dot{q}_{max} for
 $T_{max} < 75^{\circ}C$



Case 1

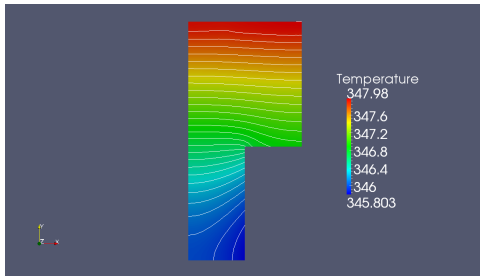
Steady state criteria: Maximum rate of change of temperature in the domain $\left(\frac{dT}{dt}\right)_{max} \leq 0.001$



Time required to reach steady state = **8.4237117 ms**

Case 2

Use the result of case 1 as the initial condition. Maximum allowable temperature in the domain 75°C or 348 K



Maximum heat generation rate $\approx 9.6 \times 10^{10} \text{ W/m}^3$. If we consider a chip of dimensions 10 mm \times 10 mm \times 0.02 mm, it is allowed to generate maximum **192 W**.



Contents

Objective

Finite Element Method

FEM Formulation

Implementation

Verification

Optimization

Parallel Computing

Application

Summary



Summary

- A two dimensional transient heat diffusion equation is solved numerically using finite element method on unstructured mesh.
- A serial object oriented C++ code is developed based on the FE formulation.
- The code is verified against analytical solutions found in [1].
- The code is parallelized to be able to solve larger problems faster.
- The parallel code shows sufficient amount of scalability.
- Finally, the capability of the software to be applied to a realistic problem has been demonstrated.



References



Jordan Wall.

Class materials for Conduction Heat Transfer.
pages 5-8, 21.



Class materials for Parallel Programming I.



<https://computing.llnl.gov/LCdocs/ioguide/>.



Class materials for Parallel Computing for Continuum
Mechanics.



Questions?