

Lecture # 3

Euler-Lagrange(EL)
eqⁿ.

$u: \Omega \rightarrow \mathbb{R}$

$\Omega \subset \mathbb{R}^n$

For ($n=1$):

$$\frac{\partial L}{\partial u} - \frac{d}{dx} \left(\frac{\partial L}{\partial u'} \right) = 0$$

For ($n > 1$)

$$\frac{\partial L}{\partial u} - \nabla \cdot \left(\frac{\partial L}{\partial u'} \right) = 0$$

In general, $L = L(x, u, u')$

Spl. case: $n=1$, suppose L does NOT explicitly depend on x

Then EL eqⁿ for $n=1$ could be simplified:

$$\frac{\partial L}{\partial u} = \frac{d}{dx} \left(\frac{\partial L}{\partial u'} \right) \quad [\text{EL eqⁿ for } n=1]$$

$$\Rightarrow \boxed{u' \frac{\partial L}{\partial u} = u' \frac{d}{dx} \left(\frac{\partial L}{\partial u'} \right)} \quad [\text{multiplying both sides by } u']$$

On the other hand, by chain rule: $\frac{dL}{dx} = \frac{\partial L}{\partial u} u' + \frac{\partial L}{\partial u'} u'' + \frac{\partial L}{\partial x}$

Combining the red & blue box :

$$\frac{dL}{dx} - \frac{\partial L}{\partial u'} u'' = u' \frac{d}{dx} \left(\frac{\partial L}{\partial u'} \right)$$

(bring everything in LHS)

$$\Rightarrow \frac{d}{dx} \left(L - u' \frac{\partial L}{\partial u'} \right) = 0$$

$$\Rightarrow \boxed{L - u' \frac{\partial L}{\partial u'} = \text{constant}}$$

(Beltrami identity)

Generalization of EL eqn for higher order derivatives :

We do for 1D :

EL eqn (If $L(x, u, u')$) : $\frac{\partial L}{\partial u} = \frac{d}{dx} \left(\frac{\partial L}{\partial u'} \right)$

$$\Leftrightarrow L_u = (L_{u'})'$$

EL eqn (If $L(x, u, u', u'', u''', u''', \dots, u^{(n)})$)

$$\text{Then } L_u = (L_{u'})' - (L_{u''})'' + (L_{u'''})''' - \dots = \sum_{n=1}^{\infty} (-1)^{n+1} (L_{u^{(n)}})^{(n)}$$

$$\Rightarrow \boxed{L_u = \sum_{r=1}^n (-1)^{r+1} (L_{u^{(r)}})^{(r)}}$$

The sol^{*} of EL eq^m u(.) may NOT be $C^2(\mathbb{R})$

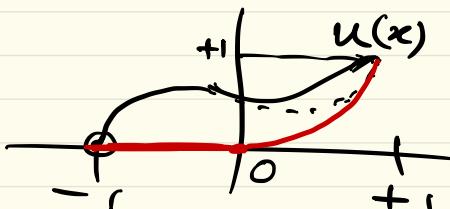
Example: (1D) [$u''(.)$ may NOT exist]

$$\min_{\substack{u(. \in C^1([-1, +1]) \\ u(-1) = 0, \quad u(1) = 1}} I(u) = \int_{-1}^{+1} u^2 (2x - u')^2 dx$$

$$\text{s. t. } u(-1) = 0, \quad u(1) = 1$$

Verify that sol^{*} of
EL eq^m:

$$u^*(x) = \begin{cases} 0 & \text{if } x \in [-1, 0] \\ x^2 & \text{if } x \in (0, 1] \end{cases}$$

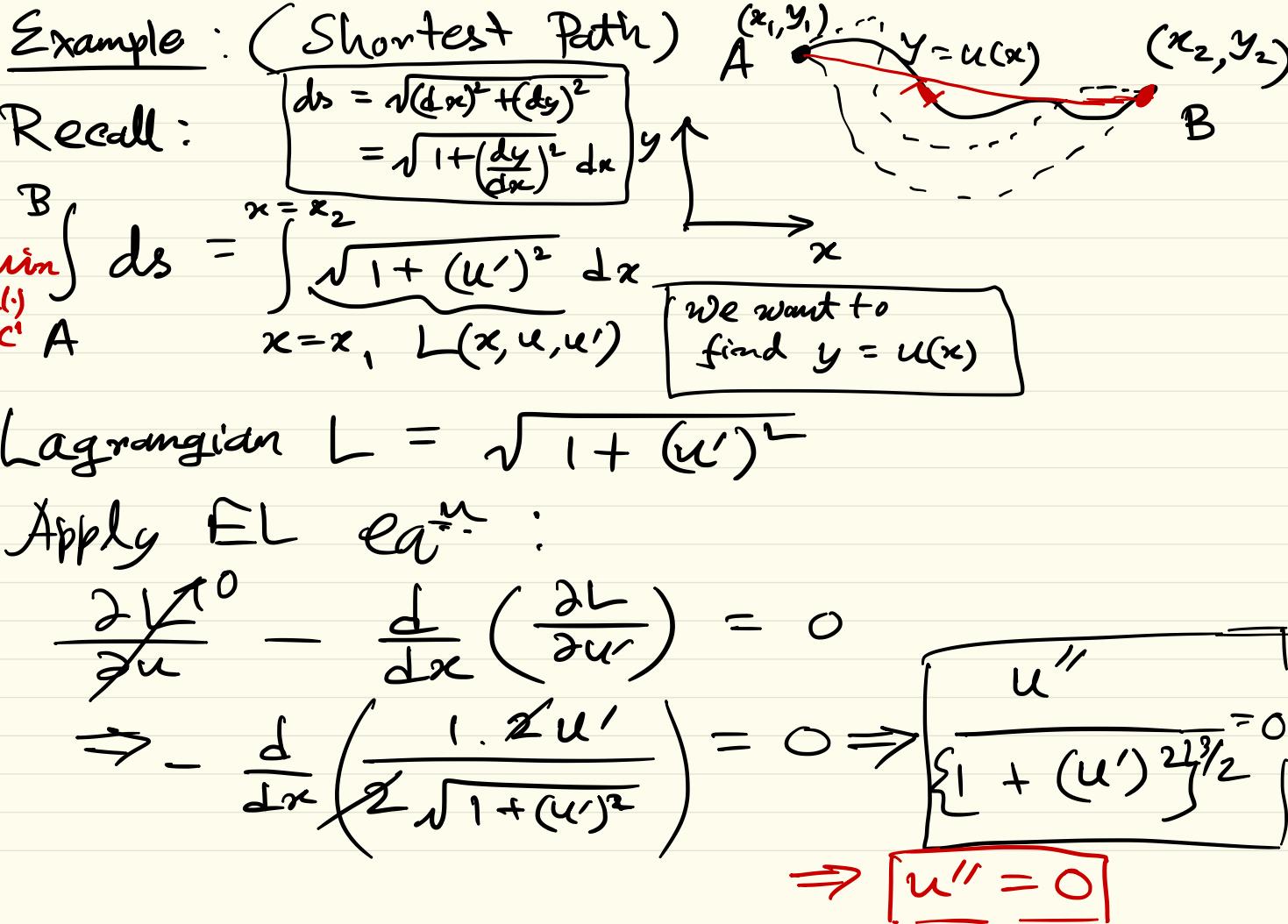


When $u^*(-)$ $\in C^2(\Omega)$

Hilbert's Thm.

If $\frac{\partial^2 L}{\partial u'^2} \neq 0$ in the entire $\text{dom}(u) \equiv \Omega$,
then the extremal $u^*(-)$ $\in C^2(\Omega)$, and
is called nonsingular.

Corollary: Suppose $u^*(-)$ is nonsingular
and L is $C^3(\Omega)$ w.r.t. u' ,
then $u^*(-)$ is the unique extremal.



$$\Rightarrow u'' = 0 \Rightarrow u' = C_1 \Rightarrow u(x) = C_1 x + C_2$$

$$y_1 = u(x_1) = C_1 x_1 + C_2$$

$$y_2 = u(x_2) = C_1 x_2 + C_2$$

Subtract:

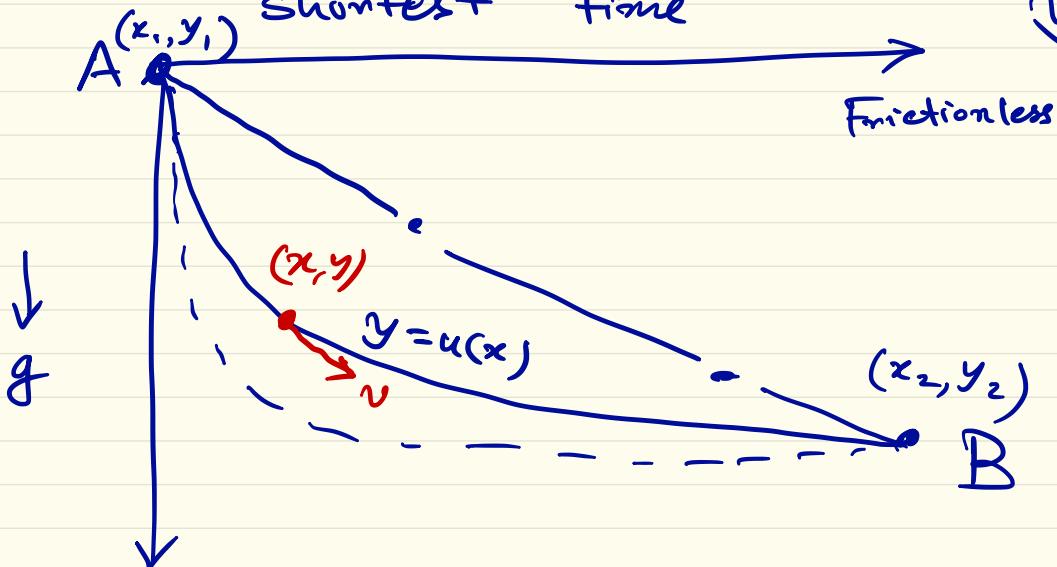
$$C_1 = \frac{y_1 - y_2}{x_1 - x_2}$$

$$\& C_2 = \frac{x_1 y_2 - x_2 y_1}{x_1 - x_2}$$

Example (Brachistochrone)

shortest time

$$y = u(x)$$



$$\min_{u(\cdot) \in C^1(\mathbb{R})} I(u) = \int_A^B dt$$

$$v = \frac{ds}{dt}$$

$$\Rightarrow dt = \frac{ds}{v} = \frac{\sqrt{1+(u')^2} dx}{\sqrt{2g} \sqrt{u}}$$

$$\therefore L_{(x, u, v)} = \frac{1}{\sqrt{2g}} \sqrt{\frac{1+(u')^2}{u}}$$

Beltrami identity:

$$L - u' \frac{\partial L}{\partial u'} = c \quad \text{constant}$$

$$\Rightarrow \sqrt{\frac{1+(u')^2}{u}} - u' \frac{\cancel{u'}}{\sqrt{u} \cdot \cancel{\sqrt{1+(u')^2}}} = c$$

$$\Rightarrow \frac{\sqrt{1+(u')^2}}{(u')^2} - \frac{(u')^2}{\cancel{u'} \cancel{\sqrt{1+(u')^2}}} = c$$

$$\Rightarrow \frac{\sqrt{u}}{1+(u')^2} - \frac{1}{\sqrt{u} \sqrt{1+(u')^2}} = c \Rightarrow \frac{1}{\sqrt{u} \sqrt{1+(u')^2}} = c$$

use Energy (E) conservation

$$E = \underbrace{T}_{\text{(Kinetic energy)}} + \underbrace{V}_{\text{(Potential energy)}}$$

$$= \boxed{\frac{1}{2} m v^2 + (-m g y)}$$

$$@A, E|_A = \frac{1}{2} m \cdot (0)^2 + (-m g \cdot 0)$$

$$= 0$$

$$\therefore E = E|_A$$

$$\Rightarrow \frac{1}{2} m v^2 = m g y \Rightarrow v = \sqrt{2g y}$$

$$= \sqrt{2g u(x)}$$

$$\Rightarrow (u')^2 = \frac{1}{c^2 u} - 1$$

$$\Rightarrow u' = \sqrt{\frac{1 - c^2 u}{c^2 u}} = \sqrt{\frac{1/c^2 - u}{u}}$$

$$\text{Let } k := \frac{1}{c^2} = \text{constant}$$

$$\Rightarrow u' = \sqrt{\frac{k-u}{u}} \quad \left. \begin{array}{l} \text{1st order} \\ \text{nonlinear} \\ \text{ODE} \end{array} \right\}$$

$$\text{Substitute: } u = k \sin^2 \phi$$

$$\Rightarrow \frac{du}{dx} = \sqrt{\frac{k-u}{u}} = \cot(\phi) \longrightarrow \cot \phi$$

$$\text{By chain rule: } \frac{d\phi}{dx} = \frac{d\phi}{du} \frac{du}{dx} = \frac{1}{2k \sin \phi \cos \phi}$$

$$= \frac{1}{2k \sin^2(\phi)}$$

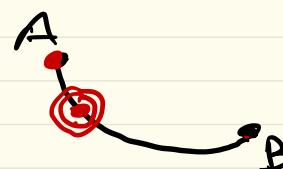
$$\Rightarrow dx = 2k \sin^2 \phi d\phi$$

$$\Rightarrow x = k \int (1 - \cos 2\phi) d\phi = k\phi - \frac{k \sin(2\phi)}{2} + C,$$

$$\Rightarrow (x, u(x)) = \left(k\phi - \frac{k}{2} \sin(2\phi) + C_1, \frac{k}{2} (1 - \cos(2\phi)) \right)$$

$$@ \text{pt. } A(0,0) : u = k \sin^2 \phi \Rightarrow \phi = 0$$

$$\cancel{x_0} = k \cancel{\phi} - \frac{k}{2} \sin(2\phi) + C_1$$



 A point A is at the origin (0,0). A curve starts at A and ends at point B. The curve is labeled with a red circle containing a dot.

$$\Rightarrow \boxed{C_1 = 0}$$

Introducing $a := \frac{k}{2}$, and $\theta := 2\phi$, we get

$$x = a(\theta - \sin \theta)$$

$$y \equiv u(x) = a(1 - \cos \theta)$$

These are parametric eqⁿs of Cycloid

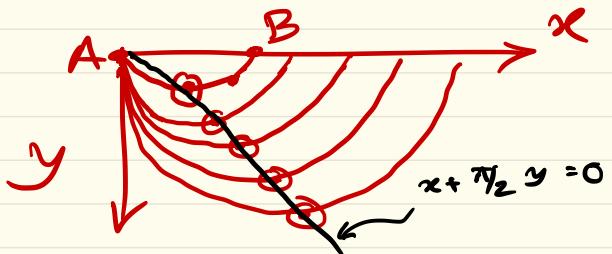
Exercise: Show that $t^* = \min_{\Gamma} t = \min_{\Gamma} \int_A^B dt$

is independent of (x_1, y_1)

\Leftrightarrow Cycloid curve is an isochrone/tautochrone

Exercise: Show that the straight line

$x + \frac{\pi}{2}y = 0$ passes through
the minima of all Brachistochrone curves



If $x_2 + \frac{\pi}{2}y_2 > 0$, then slide down,
then climb

If $x_2 + \frac{\pi}{2}y_2 \leq 0$, then slide downhill

Brachistochrone with variable g :

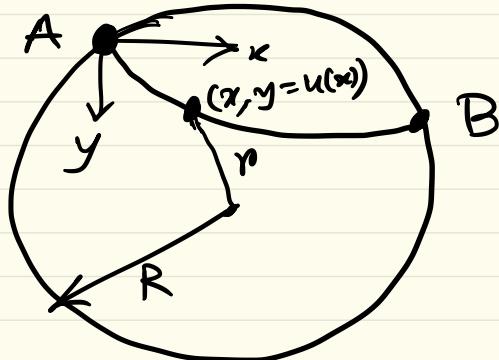
$$E|_{@A} = \frac{1}{2} m v^2 + V(r)$$

uniform
interior
density

$$= \frac{G M m}{2 R^3}$$

$$= \frac{G M m}{2 R}$$

$$E|_{@A} = E|_{@(x, u(x))}$$



Let $g := \underbrace{\frac{G M}{R^2}}_{\text{act. on surface}}$

$$\Rightarrow \frac{G M m}{2 R} = \frac{1}{2} m v^2 + \frac{G M m}{2 R^3} r$$

$$\Rightarrow \frac{v}{r} = \sqrt{\frac{G M (R^2 - r^2)}{R^3}} = \sqrt{\frac{g}{R}} \sqrt{R^2 - x^2 - u^2}$$

$$dt = \frac{ds}{v} = \sqrt{\frac{R}{g}} \frac{\sqrt{1 + (u')^2} dx}{\sqrt{R^2 - x^2 - u^2}}$$

$$\therefore \min_{u(\cdot) \in C^1(A \rightarrow B)} \int_A^B dt = \sqrt{\frac{R}{g}} \int_{x_1}^{x_2} \sqrt{\frac{1 + (u')^2}{R^2 - x^2 - u^2}} dx$$

Optimal of \mathcal{L} $\stackrel{\text{def}}{=} \mathcal{L}(x, u, u')$

$$y = u(x)$$

$$\begin{cases} x(\theta) = R \left[(1-b) \cos \theta + b \cos \left(\frac{1-b}{b} \theta \right) \right] \\ y(\theta) = R \left[(1-b) \sin \theta - b \sin \left(\frac{1-b}{b} \theta \right) \right] \end{cases}$$

\Rightarrow Hypocycloid

$$b \in [0, 1]$$

Example $u: \Omega \mapsto \mathbb{R}$, $\Omega \subset \mathbb{R}^n$. Dirichlet energy

$$\min_{u \in C^1(\Omega)} I(u) = \int_{\Omega} \frac{1}{2} \|\nabla u\|_2^2 dx$$

s.t. $u(x) = g(x)$ for $x \in \partial\Omega$.

Euler-Lagrange Eq: $\cancel{\frac{\partial L}{\partial u}} - \nabla \cdot \left(\frac{\partial L}{\partial \nabla u} \right) = 0$

$$\Rightarrow -\nabla \cdot \nabla u = 0$$

$\boxed{\Delta u = 0 \text{ with B.C. } u=g @ x \in \partial\Omega}$

Laplacian

$$\Delta = \nabla \cdot \nabla$$

$$= \left(\frac{\partial}{\partial x_1}, \frac{\partial}{\partial x_2}, \dots, \frac{\partial}{\partial x_n} \right) \begin{pmatrix} \frac{\partial^2}{\partial x_1^2} \\ \vdots \\ \frac{\partial^2}{\partial x_n^2} \end{pmatrix} = \sum_{i=1}^n \frac{\partial^2}{\partial x_i^2}$$

Example: $\min_{u \in C^4(\Omega)} I(u) = \int_{\Omega} \left\{ \frac{1}{2} \|\nabla u\|^2 - f(x)u \right\} dx$

$$\frac{\partial L}{\partial u} - \nabla \cdot \frac{\partial L}{\partial \nabla u} = 0$$

$$\Rightarrow -f(x) - \nabla \cdot \nabla u = 0$$

$$\Rightarrow \boxed{\Delta u = -f(x)}$$

(Linear Poisson eq⁴)

Example: $\min_{u \in C^1(\Omega)} I(u) = \int_{\Omega} \left\{ \frac{1}{2} \|\nabla u\|^2 - f(u)u \right\} dx$

$$\Rightarrow \frac{\partial L}{\partial u} - \nabla \cdot \frac{\partial L}{\partial \nabla u} = 0$$

$$\Rightarrow \Delta u = -f'(u) \xrightarrow{\text{Nonlinear Poisson eq}^5} \boxed{\Delta u = -\phi(u)}$$