

Lamperti Transform For Power Systems

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We consider the system of 2nd-order SDEs given by

$$m_i \ddot{\theta}_i + \gamma_i \dot{\theta}_i = P_i - \sum_{j=1}^n k_{ij} \sin(\theta_i - \theta_j) + \sigma_i \times \text{stochastic forcing } i = 1, \dots, n.$$

We can transform the system into

$$\begin{aligned} d\boldsymbol{\theta} &= \boldsymbol{\omega} dt \\ d\boldsymbol{\omega} &= -\text{diag}(\boldsymbol{\gamma} \oslash \boldsymbol{m}) \boldsymbol{\omega} - \nabla_{\boldsymbol{\theta}} V(\boldsymbol{\theta}) + \text{diag}(\boldsymbol{\sigma} \oslash \boldsymbol{m}) d\boldsymbol{w} \end{aligned} \quad (1)$$

where $\boldsymbol{\theta} = (\theta_1, \theta_2, \dots, \theta_n)$ and $\boldsymbol{\omega} = (\dot{\theta}_1, \dot{\theta}_2, \dots, \dot{\theta}_n)$ and the potential function is given by

$$V(\boldsymbol{\theta}) := \sum_{i=1}^n \frac{1}{m_i} P_i \theta_i + \sum_{(i,j) \in \mathcal{E}} \frac{1}{m_i} k_{ij} (1 - \cos(\theta_i - \theta_j)) \quad (2)$$

Theorem 1 (Lamperti Transform). *Let X_t be the solution to the (Ito) SDE*

$$d\mathbf{X}_t = \mathbf{f}(\mathbf{X}_t, t) dt + \boldsymbol{\beta}(\mathbf{X}_t, t) \mathbf{R}(t) d\mathbf{W}_t \quad (3)$$

where $\mathbf{X}_t, w_t \in \mathbb{R}^d$, $\mathbf{R}(t) \in \mathbb{R}^{d \times d}$ is any matrix function of t and $\boldsymbol{\beta} \in \mathbb{R}^{d \times d}$ is a diagonal matrix whose diagonal elements which we denote

$$\beta_{i,i}(\mathbf{X}_t, t) = \beta_i(X_{i,t}, t),$$

i.e., each diagonal element depend only on the i th component of \mathbf{X}_t . Then, the transformation defined by

$$Z_{i,t} = \psi_i(X_{i,t}, t) = \int \frac{1}{\beta_i(x, t)} dx \Big|_{x=X_{i,t}} \quad (4)$$

will result in the diffusion process

$$dZ_{i,t} = \left(\frac{\partial}{\partial t} \psi_i(x, t) \Big|_{x=X_{i,y}} + \frac{f_i(\psi^{-1}(\mathbf{Z}, t), t)}{\beta_i(\psi_i^{-1}(Z_{i,t}, t), t)} + \frac{1}{2} \frac{\partial}{\partial x} \beta_i(\psi_i^{-1}(Z_{i,t}, t), t) \right) dt + \sum_{j=1}^d r_{i,j}(t) dw_{j,t} \quad (5)$$

where $r_{i,j}(t)$ are the elements of $\mathbf{R}(t)$ and $\mathbf{X}_t = \psi^{-1}(\mathbf{Z}, t)$.

Proposition 2. Let $d = 2n$ and set

$$\mathbf{X}_t = \begin{bmatrix} \boldsymbol{\theta} \\ \boldsymbol{\omega} \end{bmatrix}, \quad \mathbf{Z}_t := \begin{bmatrix} \boldsymbol{\xi} \\ \boldsymbol{\eta} \end{bmatrix}. \quad (6)$$

Let

$$\boldsymbol{\beta}(\mathbf{X}, t) \equiv \begin{bmatrix} \mathbf{I} & 0 \\ 0 & \text{diag}(\boldsymbol{\sigma} \otimes \mathbf{m}) \end{bmatrix}, \quad \mathbf{R}(t) \equiv \begin{bmatrix} 0 & 0 \\ 0 & \mathbf{I} \end{bmatrix} \quad (7)$$

then applying Theorem 1 to (1) results in the diffusion process

$$d\mathbf{z}_t \quad (8)$$

Proof. Notice that performing the change of variable (4) results in

$$\mathbf{Z}_t = \boldsymbol{\psi}(\mathbf{X}, t) = \begin{bmatrix} \mathbf{I} & 0 \\ 0 & \text{diag}(\mathbf{m} \otimes \boldsymbol{\sigma}) \end{bmatrix} \mathbf{X}_t \quad (9)$$

so that

$$\mathbf{X}_t = \boldsymbol{\psi}^{-1}(\mathbf{Z}, t) = \begin{bmatrix} \mathbf{I} & 0 \\ 0 & \text{diag}(\boldsymbol{\sigma} \otimes \mathbf{m}) \end{bmatrix} \mathbf{Z}_t. \quad (10)$$

Using (6), the inverse transformation can be written component wise as

$$\theta_i = \xi_i, \quad \omega_i = \frac{\sigma_i}{m_i} \eta_i \quad (11)$$

Notice that since $\boldsymbol{\psi}$ and $\boldsymbol{\beta}$ are independent of t and x , then the first and third term in the drift of (5) vanish and we are left with the second term. Notice that we have

$$\frac{f_i(\mathbf{X}_t, t)}{\beta_i(\mathbf{X}_{i,t}, t)} = \begin{cases} \omega_i & 1 \leq i \leq n \\ -\frac{\gamma_i}{m_i} \frac{m_i}{\sigma_i} \omega_i - \frac{m_i}{\sigma_i} \frac{P_i}{m_i} - \sum_{j=1}^n \frac{m_i}{\sigma_i} \frac{1}{m_i} k_{ij} \sin(\theta_i - \theta_j) & n+1 \leq i \leq 2n \end{cases} \quad (12)$$

which means that

$$\frac{f_i(\psi(\mathbf{Z}, t), t)}{\beta_i(\psi(Z_i, t, t), t)} = \begin{cases} \frac{\sigma_i}{m_i} \eta_i & 1 \leq i \leq n \\ -\frac{\gamma_i}{m_i} \eta_i - \frac{P_i}{\sigma_i} - \sum_{j=1}^n \frac{1}{\sigma_i} k_{i,j} \sin(\xi_i - \xi_j) & n+1 \leq i \leq 2n \end{cases} \quad (13)$$

This results in the diffusion process

$$d \begin{bmatrix} \boldsymbol{\xi} \\ \boldsymbol{\eta} \end{bmatrix} = \begin{bmatrix} \text{diag}(\boldsymbol{\sigma} \otimes \mathbf{m}) \boldsymbol{\eta} \\ -\text{diag}(\boldsymbol{\gamma} \otimes \mathbf{m}) \boldsymbol{\eta} - \nabla_{\boldsymbol{\xi}} U(\boldsymbol{\xi}) \end{bmatrix} dt + \begin{bmatrix} 0 & 0 \\ 0 & \mathbf{I} \end{bmatrix} d\mathbf{W}_t \quad (14)$$

$$:= \begin{bmatrix} \mathbf{G}_1 \boldsymbol{\eta} \\ -\nabla_{\boldsymbol{\eta}} F(\boldsymbol{\eta}) - \nabla_{\boldsymbol{\xi}} U(\boldsymbol{\xi}) \end{bmatrix} dt + \begin{bmatrix} 0 & 0 \\ 0 & \mathbf{I} \end{bmatrix} d\mathbf{W}_t \quad (15)$$

where $F(\boldsymbol{\eta}) = \frac{1}{2} \boldsymbol{\eta}^\top \mathbf{G}_2 \boldsymbol{\eta}$ where the potential function U is given by

$$U(\boldsymbol{\xi}) := \sum_{i=1}^n \frac{P_i}{\sigma_i} \xi_i - \sum_{i,j \in \mathcal{E}} \frac{1}{\sigma_i} k_{i,j} (1 - \cos(\xi_i - \xi_j)) \quad (16)$$

□

The Fokker-Planck equation of this SDE is

$$\frac{\partial \rho}{\partial t} = -\langle \mathbf{G}_1 \boldsymbol{\eta}, \nabla_{\boldsymbol{\xi}} \rho \rangle + \nabla_{\boldsymbol{\eta}} \cdot ((\nabla_{\boldsymbol{\eta}} F(\boldsymbol{\eta}) + \nabla_{\boldsymbol{\xi}} U(\boldsymbol{\xi})) \rho) + \Delta_{\boldsymbol{\eta}} \rho \quad (17)$$

0.1 Stationary Solution

To compute the stationary solution, we set

$$0 = \langle \mathbf{G}_1 \boldsymbol{\eta}, \nabla_{\boldsymbol{\xi}} \rho \rangle + \nabla_{\boldsymbol{\eta}} \cdot ((\nabla_{\boldsymbol{\eta}} F(\boldsymbol{\eta}) + \nabla_{\boldsymbol{\xi}} U(\boldsymbol{\xi})) \rho) + \Delta_{\boldsymbol{\eta}} \rho \quad (18)$$

and we have the ansatz

$$\rho_\infty \propto \exp(-(F(\boldsymbol{\eta}) + U(\boldsymbol{\xi}))). \quad (19)$$

From this, we get the derivatives

$$\nabla_{\boldsymbol{\xi}} \rho_\infty = -\rho_\infty \nabla_{\boldsymbol{\xi}} U(\boldsymbol{\xi}), \quad \nabla_{\boldsymbol{\eta}} \rho_\infty = -\rho_\infty \nabla_{\boldsymbol{\eta}} F(\boldsymbol{\eta}). \quad (20)$$

Using (20), we obtain

$$\Delta_{\boldsymbol{\eta}} \rho_\infty = \rho_\infty \|\mathbf{G}_2 \boldsymbol{\eta}\|^2 - \rho_\infty \text{tr}(\mathbf{G}_2) \quad (21)$$

and

$$\nabla_{\boldsymbol{\eta}} \cdot ((\nabla_{\boldsymbol{\eta}} F(\boldsymbol{\eta}) + \nabla_{\boldsymbol{\xi}} U(\boldsymbol{\xi}))) = -\rho_\infty \|\mathbf{G}_2 \boldsymbol{\eta}\|^2 - \rho_\infty \langle \mathbf{G}_2 \boldsymbol{\eta}, \nabla_{\boldsymbol{\xi}} U(\boldsymbol{\xi}) \rangle + \rho_\infty \text{tr}(\mathbf{G}_2) \quad (22)$$

which implies that (18) reduces to

$$0 = \rho_\infty \langle (\mathbf{G}_2 - \mathbf{G}_1) \boldsymbol{\eta}, \nabla_{\boldsymbol{\xi}} U(\boldsymbol{\xi}) \rangle. \quad (23)$$

In order for this to hold for all $\boldsymbol{\eta}, \boldsymbol{\xi}$ we must have $\mathbf{G}_2 = \mathbf{G}_1$ which means that $\boldsymbol{\sigma} = \boldsymbol{\gamma}$, i.e., $\sigma_i = \gamma_i$ for all i .

0.2 Free Energy

Consider the functional

$$\begin{aligned}\Phi(\rho) &= \int \rho(\boldsymbol{\eta}, \boldsymbol{\xi}) \log \rho(\boldsymbol{\eta}, \boldsymbol{\xi}) \, d\boldsymbol{\eta} d\boldsymbol{\xi} + \beta \int \rho(\boldsymbol{\eta}, \boldsymbol{\xi}) U(\boldsymbol{\xi}) \, d\boldsymbol{\eta} d\boldsymbol{\xi} + \beta \int \rho(\boldsymbol{\eta}, \boldsymbol{\xi}) F(\boldsymbol{\eta}) \, d\boldsymbol{\eta} d\boldsymbol{\xi} \\ &= \beta \tilde{\Phi}(\rho)\end{aligned}\tag{24}$$

We compute

$$\begin{aligned}\frac{d\Phi}{dt} &= \int \frac{\partial \rho}{\partial t} (\log \rho + \beta F + \beta U) \, d\boldsymbol{\eta} d\boldsymbol{\xi} \, d\boldsymbol{\eta} d\boldsymbol{\xi} \\ &= \int \nabla_{\boldsymbol{\eta}} \cdot (\rho (\nabla_{\boldsymbol{\eta}} \log \rho + \nabla_{\boldsymbol{\eta}} F + \nabla_{\boldsymbol{\xi}} U)) (\log \rho + \beta F + \beta U) \, d\boldsymbol{\eta} d\boldsymbol{\xi} \\ &\quad - \int \langle G_1 \boldsymbol{\eta}, \nabla_{\boldsymbol{\xi}} \rho \rangle (\log \rho + \beta F + \beta U) \, d\boldsymbol{\eta} d\boldsymbol{\xi} \\ &= \int -|\nabla_{\boldsymbol{\eta}} \log \rho|^2 \rho - (\beta + 1) \langle \nabla_{\boldsymbol{\eta}} \log \rho, \nabla_{\boldsymbol{\eta}} F \rangle \rho - \beta |\nabla_{\boldsymbol{\eta}} F|^2 \rho \, d\boldsymbol{\eta} d\boldsymbol{\xi} \\ &\quad \int -\langle \nabla_{\boldsymbol{\xi}} U, \nabla_{\boldsymbol{\eta}} \log \rho \rangle \rho - \beta \langle \nabla_{\boldsymbol{\xi}} U, \nabla_{\boldsymbol{\eta}} F \rangle \rho + \langle G_1 \boldsymbol{\eta}, \nabla_{\boldsymbol{\xi}} \log \rho \rangle \rho + \beta \langle G_1 \boldsymbol{\eta}, \nabla_{\boldsymbol{\xi}} U \rangle \rho \, d\boldsymbol{\eta} d\boldsymbol{\xi} \\ &= \int -|\nabla_{\boldsymbol{\eta}} \log \rho|^2 \rho - (\beta + 1) \langle \nabla_{\boldsymbol{\eta}} \log \rho, \nabla_{\boldsymbol{\eta}} F \rangle \rho - \beta |\nabla_{\boldsymbol{\eta}} F|^2 \rho \, d\boldsymbol{\eta} d\boldsymbol{\xi} \\ &\quad + t \int -\beta \langle \nabla_{\boldsymbol{\xi}} U, \nabla_{\boldsymbol{\eta}} F \rangle \rho + \beta \langle G_1 \boldsymbol{\eta}, \nabla_{\boldsymbol{\xi}} U \rangle \rho.\end{aligned}$$

A sufficient condition for $\frac{d\Phi}{dt} \leq 0$ is to set $G_1 \boldsymbol{\eta} = \nabla_{\boldsymbol{\eta}} F = G_2 \boldsymbol{\eta} \Rightarrow \boldsymbol{\sigma} = \boldsymbol{\gamma} \Rightarrow \sigma_i = \gamma_i$ for all i .