

Linearised NSE.

NSE

$$\frac{\partial \bar{u}}{\partial t} + (\bar{u} \cdot \nabla) \bar{u} = -\nabla p + \frac{1}{Re} \nabla^2 \bar{u} \quad \text{--- (1)}$$

$$\nabla \cdot \bar{u} = 0 \quad \text{--- (2)}$$

$$\bar{u}(x, y, z, t) = \bar{U}(y) + u'(x, y, z, t) \quad \text{--- (3)}$$

$$p(x, y, z, t) = P(y) + p'(x, y, z, t) \quad \text{--- (4)}$$

Using (3) & (4) in (1)

$$\frac{\partial \bar{u}}{\partial t} = \frac{\partial u'}{\partial t} \quad \therefore \frac{\partial \bar{U}}{\partial t} = 0$$

* For this derivation
I am using
 $\gamma = Re^{-1}$

$$\begin{aligned} (\bar{u} \cdot \nabla) \bar{u} &= (u' + \bar{U}) \cdot \nabla (u' + \bar{U}) \\ &= u' \cdot \nabla u' + \bar{U} \cdot \nabla u' + u' \cdot \nabla \bar{U} + \bar{U} \cdot \nabla \bar{U} \end{aligned}$$

$$\nabla p = \nabla p' + \nabla P$$

$$\gamma \nabla^2 \bar{u} = \gamma \nabla^2 u' + \gamma \nabla^2 \bar{U}$$

$$\begin{aligned} \therefore \frac{\partial u'}{\partial t} + (u' \cdot \nabla) u' + \bar{U} \cdot \nabla u' + u' \cdot \nabla \bar{U} + \bar{U} \cdot \nabla \bar{U} &= -\nabla p' - \nabla P \\ &\quad + \gamma \nabla^2 u' + \gamma \nabla^2 \bar{U} \end{aligned} \quad \text{--- (5)}$$

Eqⁿ for the steady flow

$$(\bar{U} \cdot \nabla) \bar{U} = -\nabla P + \gamma \nabla^2 \bar{U} \quad \text{--- (6)}$$

Eqⁿ (5) - Eqⁿ (6) and dropping $O(u'^2)$ term.

$$\frac{\partial \bar{u}'}{\partial t} + \bar{u}' \cdot \nabla \bar{U} + \bar{U} \cdot \nabla \bar{u}' = -\nabla \phi' + \gamma \nabla^2 \bar{u}' \quad \text{--- (7)}$$

Using (3) in (2)

$$\nabla \cdot (\bar{u}' + \bar{U}) = 0$$

$$\Rightarrow \nabla \cdot \bar{u}' + \nabla \cdot \bar{U} = 0 \quad \text{--- (8)}$$

for steady flow $\nabla \cdot \bar{U} = 0$ --- (9)

(8) - (9),

$$\nabla \cdot \bar{u}' = 0 \quad \text{--- (10)}$$

Orr - Sommerfeld eq.

For parallel flow

$$\bar{U} = U(y) \hat{e}_x \quad \text{--- (11)}$$

Using (11) in (7)

$$\frac{\partial \bar{u}'}{\partial t} + \underbrace{\bar{u}' \cdot \nabla \bar{U}}_{\text{"}} + \bar{U} \cdot \nabla \bar{u}' = -\sigma \bar{p}' + \nu \nabla^2 \bar{u}'$$

$$\left(\underbrace{u' \frac{\partial}{\partial x} + v' \frac{\partial}{\partial y} + w' \frac{\partial}{\partial z}}_{\text{"}} \right) U(y) \hat{e}_x$$

$$\underbrace{v' \frac{dU}{dy}}_{\text{"}} \hat{e}_x$$

$$v' U^{(1)} \hat{e}_x$$

$$\text{let } U^{(1)} = \frac{dU}{dy}$$

$$(U \cdot \nabla) \bar{u}' = \left(\underbrace{U \frac{\partial}{\partial x}}_0 + \underbrace{v \frac{\partial}{\partial y}}_0 + \underbrace{w \frac{\partial}{\partial z}}_0 \right) \bar{u}' = U \frac{\partial}{\partial x} \bar{u}'$$

$$\therefore \left[\frac{\partial \bar{u}'}{\partial t} + v' u^{(1)} \hat{e}_x + v \frac{\partial \bar{u}'}{\partial x} = -\nabla p' + \gamma \nabla^2 u' \right] \cdot \hat{e}_y \quad (12)$$

$$\Rightarrow \nabla^2 \left[\frac{\partial v'}{\partial t} + v \frac{\partial v'}{\partial x} = -\frac{\partial p'}{\partial y} + \gamma \nabla^2 v' \right]$$

$$\Rightarrow \frac{\partial}{\partial t} \nabla^2 v' + \nabla^2 v \frac{\partial v'}{\partial x} + \frac{\partial}{\partial y} \nabla^2 p' = \gamma \nabla^4 v' \quad (12)$$

$\nabla \cdot \text{eq}^n (12)$

$$\frac{\partial}{\partial t} \underbrace{\nabla \cdot \bar{u}'}_0 + \nabla \cdot (v' u^{(1)} \hat{e}_x) + \nabla \cdot \left(v \frac{\partial \bar{u}'}{\partial x} \right) = -\nabla^2 p' + \gamma \nabla^2 (\underbrace{\nabla \cdot \bar{u}'}_0)$$

$$\therefore \nabla^2 p' = - \underbrace{\nabla \cdot (v' u^{(1)} \hat{e}_x)}_{\partial_x (v' u^{(1)})} - \underbrace{\nabla \cdot \left(v \frac{\partial \bar{u}'}{\partial x} \right)}_{\partial_i (v \partial_x \bar{u}'_i)}$$

$$\nabla^2 p' = - \underbrace{u^{(1)} \partial_x v'}_0 - \underbrace{\partial_x v \partial_x u'}_0 - \underbrace{\partial_z v \partial_x v'}_{u^{(1)} \partial_x v'} - \underbrace{\partial_x v \partial_x \omega'}_0$$

$$\therefore \nabla^2 \phi' = -2 U^{(1)} \frac{\partial \psi'}{\partial x}$$

$$\therefore \frac{\partial}{\partial y} \nabla^2 \phi' = -2 \left(\frac{\partial}{\partial y} U^{(1)} \frac{\partial \psi'}{\partial x} + U^{(1)} \frac{\partial}{\partial y} \frac{\partial \psi'}{\partial x} \right)$$

$$\text{let } U^{(2)} = \frac{d^2 U}{dy^2}$$

$$\therefore \frac{\partial}{\partial y} \nabla^2 \phi' = -2 \left(U^{(2)} \frac{\partial \psi'}{\partial x} + U^{(1)} \frac{\partial}{\partial y} \frac{\partial \psi'}{\partial x} \right) \quad \text{--- (14)}$$

$$\frac{\partial}{\partial t} \nabla^2 \psi' + \nabla^2 U \frac{\partial \psi'}{\partial x} + \frac{\partial}{\partial y} \nabla^2 \phi' = \nabla^4 \psi' \quad \text{--- (12)}$$

$$\nabla^2 \left(U \frac{\partial \psi'}{\partial x} \right) = \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) U \frac{\partial \psi'}{\partial x}$$

$$= \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) U \frac{\partial \psi'}{\partial x} + \frac{\partial^2}{\partial y^2} U \frac{\partial \psi'}{\partial x}$$

$$= U \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) \frac{\partial \psi'}{\partial x} + \frac{\partial}{\partial y} \left\{ \frac{\partial}{\partial y} \left(U \frac{\partial \psi'}{\partial x} \right) \right\}$$

$$\underbrace{U \frac{\partial}{\partial x} \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) \psi'}_{\text{"}} \quad \left\{ U^{(1)} \frac{\partial \psi'}{\partial x} + U \frac{\partial}{\partial x} \frac{\partial \psi'}{\partial y} \right\}$$

$$= \underbrace{U \frac{\partial}{\partial x} \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right)}_{\text{green}} \psi' + \frac{\partial}{\partial y} \left\{ \underbrace{U^{(1)} \frac{\partial \psi'}{\partial x} + U \partial_n \partial_y \psi'}_{\text{blue}} \right\}$$

$$\begin{aligned} & U^{(2)} \frac{\partial \psi'}{\partial x} + \underbrace{U^{(1)} \partial_n \partial_y \psi'}_{\text{pink}} \\ & + \underbrace{U^{(1)} \partial_n \partial_y \psi'}_{\text{pink}} + \underbrace{U \partial_n \partial_y^2 \psi'}_{\text{green}} \end{aligned}$$

$$= U \frac{\partial}{\partial x} \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) \psi' + U^{(2)} \frac{\partial \psi'}{\partial x} + 2 U^{(1)} \partial_n \partial_y \psi'$$

$$\nabla^2 \left(U \frac{\partial \psi'}{\partial x} \right) = U \frac{\partial}{\partial x} \nabla^2 \psi' + U^{(2)} \frac{\partial \psi'}{\partial x} + 2 U^{(1)} \partial_n \partial_y \psi' \quad \text{--- (15)}$$

Using (14) & (15) in (13)

$$\frac{\partial}{\partial t} \nabla^2 \psi' + \nabla^2 U \frac{\partial \psi'}{\partial x} + \frac{\partial}{\partial y} \nabla^2 \psi' = \nabla^4 \psi' \quad \text{--- (12)}$$

$$\frac{\partial}{\partial t} \nabla^2 \psi' + U \frac{\partial}{\partial x} \nabla^2 \psi' + U^{(2)} \frac{\partial \psi'}{\partial x} + 2 U^{(1)} \cancel{\partial_n \partial_y \psi'}$$

$$- 2 \left(U^{(2)} \frac{\partial \psi'}{\partial x} + U^{(1)} \cancel{\partial_y \partial_n \psi'} \right)$$

$$= \nabla^4 \psi'$$

$$\Rightarrow \left(\frac{\partial}{\partial t} + v \frac{\partial}{\partial x} \right) \nabla^2 \psi' - v^{(2)} \frac{\partial \psi'}{\partial x} = \gamma \nabla^4 \psi'$$

$$\therefore \left[\left(\frac{\partial}{\partial t} + v \frac{\partial}{\partial x} \right) \nabla^2 - v^{(2)} \frac{\partial}{\partial x} \right] \psi' = \gamma \nabla^4 \psi'$$

———— (16)

$$\psi'(x, y, z, t) = \hat{v}(y) e^{i(\alpha x + \beta z)} e^{-i\omega t} \quad \text{--- (17)}$$

$$\begin{aligned} \nabla^2 &\equiv \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} = \frac{\partial^2}{\partial y^2} - \alpha^2 - \beta^2 \\ &= D^2 - (\alpha^2 + \beta^2) \\ &= D^2 - k^2 \end{aligned}$$

$$\frac{\partial}{\partial t} \equiv -i\omega$$

$$\frac{\partial}{\partial x} = i\alpha$$

Here $k^2 = \alpha^2 + \beta^2$

$$D \equiv \frac{d}{dy}$$

$$D^2 \equiv \frac{d^2}{dy^2}$$

$$= \left[(-i\omega + v i\alpha)(D^2 - k^2) - v^{(2)} i\alpha \right] \hat{\psi}(y) \\ = \gamma (D^2 - k^2)^2 \hat{\psi}(y)$$

$$\Rightarrow (D^4 + k^4 - 2D^2 k^2) \hat{\psi} = \frac{1}{\gamma} \left\{ \underbrace{\left(-i\frac{\omega}{\alpha} \alpha + v i\alpha \right)}_c (D^2 - k^2) - v^{(2)} i\alpha \right\} \hat{\psi}$$

$$\text{let } \frac{\omega}{\alpha} = c$$

$$\Rightarrow \hat{\psi}^{(4)} - 2k^2 \hat{\psi}^{(2)} + k^4 \hat{\psi} = \frac{i\alpha}{\gamma} \left\{ (v - c)(\hat{\psi}^{(2)} - k^2 \hat{\psi}) - v^{(2)} \hat{\psi} \right\}$$

$$\Rightarrow \hat{\psi}^{(4)} - 2k^2 \hat{\psi}^{(2)} + k^4 \hat{\psi} = i\alpha k c \left\{ (v - c)(\hat{\psi}^{(2)} - k^2 \hat{\psi}) - v^{(2)} \hat{\psi} \right\}$$

— (18)

Above is the Orr-Sommerfeld Eqⁿ.

PPF

Solution of the
base flow.

$$y = +1$$

$$y = -1$$

$$\text{B.C. } U(y = \pm 1) = 0$$

Base flow eq.

$$(\bar{U} \cdot \nabla) \bar{U} = -\nabla P + \frac{1}{Re} \nabla^2 \bar{U}$$

$$\text{For parallel flow } \bar{U} = U(y) \hat{e}_x$$

$$\therefore (\bar{U} \cdot \nabla) \bar{U} = \left(\underbrace{U}_{=0} \frac{\partial}{\partial x} + \underbrace{V}_{=0} \frac{\partial}{\partial y} + \underbrace{W}_{=0} \frac{\partial}{\partial z} \right) U(y) \hat{e}_x$$

$$= 0$$

$$-\frac{\partial P}{\partial x} + \frac{1}{Re} \nabla^2 U = 0$$

$$\text{Let } G = -\frac{\partial P}{\partial x} = \text{constant pressure gradient.}$$

$$\therefore \frac{d^2 U}{dy^2} = -G Re$$

$$\Rightarrow \frac{dU}{dy} = -G Re y + A$$

$$= U(y) = -\frac{C_1 k_e}{2} y^2 + Ay + B$$

At $y=1$

$$0 = -\frac{C_1 k_e}{2} + A + B \quad \text{--- (i)}$$

At $y=-1$

$$0 = -\frac{C_1 k_e}{2} - A + B \quad \text{--- (ii)}$$

(i) - (ii)

$$0 = 0 + 2A + 0 \Rightarrow A = 0$$

(i) + (ii)

$$0 = -C_1 k_e + 2B$$

$$\therefore B = \frac{C_1 k_e}{2}$$

$$\therefore U(y) = -\frac{C_1 k_e}{2} y^2 + \frac{C_1 k_e}{2}$$

$$\therefore U(y) = \frac{C_1 k_e}{2} (1 - y^2)$$

3f $C_1 = \frac{2}{k_e}$

$$\therefore \boxed{U(y) = 1 - y^2}$$

Solution of OS eq.

B.C. on perturbation

$$u' = v' = \omega' = 0 \quad \text{at } y = \pm 1$$

$$\hat{v}^{(4)} - 2k^2 \hat{v}^{(2)} + k^4 \hat{v} = i\alpha k e \{ (v-c)(\hat{v}^{(2)} - k^2 \hat{v}) - v^{(2)} \hat{v} \} \quad (18)$$

We need to solve eq (18) using above B.C.

Above B.C. implies.

$$\hat{u}(y) = \hat{v}(y) = \hat{\omega}(y) = 0 \quad \text{at } y = \pm 1$$

$$\text{or } \hat{v}(y) = 0 \quad \text{At } y = \pm 1$$

Eq (18) is 4th order eq., therefore we need two more condition to solve OS.

$$\therefore \nabla \cdot \bar{u}' = 0$$

$$\Rightarrow \frac{\partial u'}{\partial x} + \frac{\partial v'}{\partial y} + \frac{\partial \omega'}{\partial z} = 0$$

$$\Rightarrow i\alpha \hat{u} + \frac{d\hat{v}}{dy} + i\beta \hat{\omega} = 0$$

$$\therefore \hat{u}(y) = 0 \text{ \& } \hat{v}(y) = 0 \quad \text{at } y = \pm 1$$

$$\Rightarrow \frac{d\hat{V}}{dy} = 0 \quad \text{at} \quad y = \pm 1$$

Thus B.C. are.

$$\hat{V}(y) = 0 \quad \frac{d\hat{V}}{dy} = 0 \quad \text{at} \quad y = \pm 1$$

Matrix Eigenvalue problem.

$$A \hat{V} = \omega B \hat{V}$$

$$\hat{V}^{(4)} - 2k^2 \hat{V}^{(2)} + k^4 \hat{V} = i\alpha \text{Re} \left\{ (U-C) (\hat{V}^{(2)} - k^2 \hat{V}) - U^{(2)} \hat{V} \right\} \quad \text{--- (18)}$$

From eq (18),

Using $C = \frac{\omega}{\alpha}$ and multiplying both sides by $\frac{i}{\text{Re}}$

$$\begin{aligned} \therefore \frac{i}{\text{Re}} \left[\hat{V}^{(4)} - 2k^2 \hat{V}^{(2)} + k^4 \hat{V} \right] &= -\alpha \left\{ \left(U - \frac{\omega}{\alpha} \right) (\hat{V}^{(2)} - k^2 \hat{V}) - U^{(2)} \hat{V} \right\} \\ &= \alpha U^{(2)} \hat{V} - (\alpha U - \omega) (\hat{V}^{(2)} - k^2 \hat{V}) \\ &= \alpha U^{(2)} \hat{V} - \left\{ \alpha U (\hat{V}^{(2)} - k^2 \hat{V}) - \omega (\hat{V}^{(2)} - k^2 \hat{V}) \right\} \\ &= \alpha U^{(2)} \hat{V} + \omega (\hat{V}^{(2)} - k^2 \hat{V}) - \alpha U (\hat{V}^{(2)} - k^2 \hat{V}) \end{aligned}$$

$$\Rightarrow \frac{i}{\hbar c} \left[\hat{U}^{(4)} - 2k^2 \hat{U}^{(2)} + k^4 \hat{U} \right] + \alpha U (\hat{U}^{(2)} - k^2 \hat{U}) - \alpha U^{(2)} \hat{U} \\ = \omega (\hat{U}^{(2)} - k^2 \hat{U})$$

$$\therefore A = \frac{i}{\hbar c} (\mathcal{D}^4 - 2k^2 \mathcal{D}^2 + k^4 \mathbb{I}) + \alpha U (\mathcal{D}^2 - k^2 \mathbb{I}) \\ - \alpha U^{(2)} \mathbb{I}$$

$$B = \mathcal{D}^2 - k^2 \mathbb{I}$$

* How to find \hat{u} from \hat{V} in 2D.

From continuity eqⁿ $\nabla \cdot \vec{u}' = 0$

$$\Rightarrow i\alpha \hat{u} + \frac{d\hat{V}}{dy} + i\beta \hat{w} = 0$$

For 2D $\beta = 0$.

$$i\alpha \hat{u} = - \frac{d\hat{V}}{dy}$$

$$\Rightarrow i\alpha \hat{u} = -i \frac{d\hat{V}}{dy}$$

$$\therefore \hat{u} = \frac{1}{\alpha} \frac{d\hat{V}}{dy}$$

* How to obtain u' from $\hat{V}(y)$

$$u'(x, y) = \hat{V}(y) e^{i\alpha x}$$

$$u'(x, y) = \hat{V}(y) (\cos(\alpha x) + i \sin(\alpha x))$$

$$\Rightarrow u' = \text{Real} \left\{ \hat{V} (\cos \alpha x + i \sin \alpha x) \right\}$$

$$= \text{Real} \left\{ (\hat{V}_R + i \hat{V}_I) (\cos \alpha x + i \sin \alpha x) \right\}$$

$$V' = \hat{V}_p \cos \alpha - \hat{V}_T \sin \alpha$$