

Galerkin Technique

AS6320

08th April, 2024

$$\gamma M \frac{\partial u'}{\partial t} + \frac{\partial p'}{\partial x} = 0$$

$$\frac{\partial p'}{\partial t} + \gamma M \frac{\partial u'}{\partial x} = K \left[\sqrt{\left| \frac{1}{3} + u'_f(t - \tau) \right|} - \sqrt{\frac{1}{3}} \right] \delta(x - x_f)$$

Galerkin Technique or Modal Expansion

Change PDE into ODE

ODEs are much easier
to solve than PDEs

1. Expand variables (u', p') in terms of basis functions
2. Project PDEs onto the basis functions
3. Individual equations for all the terms

Projecting along basis functions??

$$\vec{F} = m \vec{a} = m \frac{d^2 \vec{x}}{dt^2}$$

$$F_x \vec{i} + F_y \vec{j} + F_z \vec{k} = m \left(\frac{d^2 x}{dt^2} \vec{i} + \frac{d^2 y}{dt^2} \vec{j} + \frac{d^2 z}{dt^2} \vec{k} \right)$$

$$\vec{F} = F_x \vec{i} + F_y \vec{j} + F_z \vec{k}$$

$$F_x = m \frac{d^2 x}{dt^2} \quad F_y = m \frac{d^2 y}{dt^2} \quad F_z = m \frac{d^2 z}{dt^2}$$

$$\vec{X} = x \vec{i} + y \vec{j} + z \vec{k}$$

How does one “mathematically” arrive at this statement ?

Projecting along basis functions??

$$F_x \vec{i} + F_y \vec{j} + F_z \vec{k} = m \left(\frac{d^2x}{dt^2} \vec{i} + \frac{d^2y}{dt^2} \vec{j} + \frac{d^2z}{dt^2} \vec{k} \right)$$

$$\vec{i} \cdot \left[F_x \vec{i} + F_y \vec{j} + F_z \vec{k} = m \left(\frac{d^2x}{dt^2} \vec{i} + \frac{d^2y}{dt^2} \vec{j} + \frac{d^2z}{dt^2} \vec{k} \right) \right]$$

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$$F_x = m \frac{d^2x}{dt^2}$$

Projecting along basis functions??

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$$F_y = m \frac{d^2y}{dt^2}$$

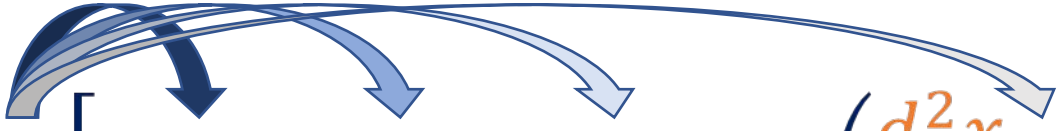
Projecting along basis functions??

$$F_x \vec{i} + F_y \vec{j} + F_z \vec{k} = m \left(\frac{d^2x}{dt^2} \vec{i} + \frac{d^2y}{dt^2} \vec{j} + \frac{d^2z}{dt^2} \vec{k} \right)$$

$$\vec{k} \cdot \left[F_x \vec{i} + F_y \vec{j} + F_z \vec{k} = m \left(\frac{d^2x}{dt^2} \vec{i} + \frac{d^2y}{dt^2} \vec{j} + \frac{d^2z}{dt^2} \vec{k} \right) \right]$$

$$F_z = m \frac{d^2z}{dt^2}$$

Projecting along basis functions



The diagram illustrates the projection of a vector \vec{k} onto the components of a force vector. A dark blue vector \vec{k} originates from the left. Three curved arrows originate from its tip: a dark blue arrow pointing to $F_x \vec{i}$, a medium blue arrow pointing to $F_y \vec{j}$, and a light blue arrow pointing to $F_z \vec{k}$. These arrows represent the projection of \vec{k} onto the respective basis vectors.

$$\vec{k} \cdot \left[F_x \vec{i} + F_y \vec{j} + F_z \vec{k} \right] = m \left(\frac{d^2 x}{dt^2} \vec{i} + \frac{d^2 y}{dt^2} \vec{j} + \frac{d^2 z}{dt^2} \vec{k} \right)$$

$$\gamma M \frac{\partial u'}{\partial t} + \frac{\partial p'}{\partial x} = 0 \qquad \frac{\partial p'}{\partial t} + \gamma M \frac{\partial u'}{\partial x} = K \left[\sqrt{\left| \frac{1}{3} + u'_f(t - \tau) \right|} - \sqrt{\frac{1}{3}} \right] \delta(x - x_f)$$

Duct open at both ends: $p'(x = 0) = 0, p'(x = 1) = 0$

$$p'(x, t) = \sum_{j=1}^N a_j(t) \sin(j\pi x)$$

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$$a_j(t) = -\frac{\gamma M}{j\pi} \dot{\eta}_j(t)$$

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$$\frac{\partial p'(x, t)}{\partial t} = - \sum_{j=1}^N \frac{\gamma M}{j\pi} \frac{d}{dt} (\dot{\eta}_j) \sin(j\pi x)$$

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Multiply both sides by $\frac{j\pi}{\gamma M}$

$$\sum_{j=1}^N \left[\frac{d \dot{\eta}_j}{dt} + (j\pi)^2 \eta_j \right] \sin(j\pi x) = - \frac{j\pi K}{\gamma M} \left[\sqrt{\left| \frac{1}{3} + u'_f(t - \tau) \right|} - \sqrt{\frac{1}{3}} \right] \delta(x - x_f)$$

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$k_j = j\pi =$ nondimensionalized wave number

Projecting along the basis function

$$\sum_{j=1}^N \left[\frac{d \dot{\eta}_j}{dt} + k_j^2 \eta_j \right] \sin(j\pi x) = -\frac{j\pi K}{\gamma M} \left[\sqrt{\left| \frac{1}{3} + u'_f(t - \tau) \right|} - \sqrt{\frac{1}{3}} \right] \delta(x - x_f)$$

Multiply both sides by $\sin(j\pi x)$ and integrate along the domain $[0, 1]$

$$\int_{x=0}^{x=1} \sum_{j=1}^N \left[\frac{d \dot{\eta}_j}{dt} + k_j^2 \eta_j \right] \sin(j\pi x) \sin(j\pi x) dx = \int_{x=0}^{x=1} -\frac{j\pi K}{\gamma M} \left[\sqrt{\left| \frac{1}{3} + u'_f(t - \tau) \right|} - \sqrt{\frac{1}{3}} \right] \delta(x - x_f) \sin(j\pi x) dx$$

$$\int_{x=0}^{x=1} \sum_{j=1}^N \left[\frac{d \dot{\eta}_j}{dt} + k_j^2 \eta_j \right] \sin(j\pi x) \sin(j\pi x) \, dx = \int_{x=0}^{x=1} -\frac{j\pi K}{\gamma M} \left[\sqrt{\left| \frac{1}{3} + u'_f(t - \tau) \right|} - \sqrt{\frac{1}{3}} \right] \delta(x - x_f) \sin(j\pi x) \, dx$$

$$\int_0^1 \sin(j\pi x) \sin(n\pi x) dx = \frac{1}{2} \delta_{jn}$$

$$\int_0^1 f(x) \delta(x - x_f) dx = f(x_f)$$

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$$\frac{d \dot{\eta}_j}{dt} + k_j^2 \eta_j = -\frac{2j\pi K}{\gamma M} \left[\sqrt{\left| \frac{1}{3} + u'_f(t - \tau) \right|} - \sqrt{\frac{1}{3}} \right] \sin(j\pi x_f)$$

$$\frac{d \eta_j}{dt} = \dot{\eta}_j$$

Equation 12,
Balasubramanian
and Sujith, Physics
of Fluids, 2008

$$\frac{d \dot{\eta}_j}{dt} + 2\zeta_j \omega_j \dot{\eta}_j + k_j^2 \eta_j = -\frac{2j\pi K}{\gamma M} \left[\sqrt{\left| \frac{1}{3} + u'_f(t - \tau) \right|} - \sqrt{\frac{1}{3}} \right] \sin(j\pi x_f)$$

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$$\zeta_j = \text{frequency dependent damping} = \frac{1}{2\pi} \left[c_1 \frac{\omega_j}{\omega_1} + c_2 \sqrt{\frac{\omega_1}{\omega_j}} \right]$$

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Numerical solution of ODEs using RK4 scheme

$$\frac{d \, \eta_j}{dt} = \, \dot{\eta}_j$$

$$\frac{d \, \dot{\eta}_j}{dt} + 2\zeta_j \omega_j \dot{\eta}_j \, + k_j^2 \eta_j \, = - \frac{2j\pi K}{\gamma M} \left[\sqrt{\left| \frac{1}{3} + u'_f(t - \tau) \right|} - \sqrt{\frac{1}{3}} \right] \sin(j\pi x_f)$$

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$$\frac{d \, \dot{\eta}_j}{dt} = -2\zeta_j\omega_j\dot{\eta}_j - k_j^2\eta_j - \frac{2j\pi K}{\gamma M} \left[\sqrt{\left| \frac{1}{3} + u'_f(t-\tau) \right|} - \sqrt{\frac{1}{3}} \right] \sin(j\pi x_f)$$

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To simplify, we have the following system of equations:

$$\frac{d \eta_j}{dt} = \dot{\eta}_j$$

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To simplify, we have the following system of equations:

$$\frac{d \eta_j}{dt} = f(\eta_j, \dot{\eta}_j)$$

for $j = 1, 2, \dots, N$ Galerkin modes

$$\frac{d \dot{\eta}_j}{dt} = g(\eta_j, \dot{\eta}_j)$$

For a general ordinary first-order differential equation,

$$\frac{dy(t)}{dx} = f(t, y), \quad y(t_0) = y_0$$

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Thus, from expanding $y(t)$ in the vicinity of t_0 through Taylor's expansion, we get,

$$y(t_0 + h) = y(t_0) + h y'(t_0) + \frac{h^2}{2!} y''(t_0) + \dots$$

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$$y(t_0 + h) = y(t_0) + h y'(t_0)$$

$$y_1 = y_0 + h y'(t_0) = y_0 + h f(x_0, y_0)$$

$$y_{n+1} = y_n + h y'(x_n, y_n)$$

If we seek higher order accuracy, we must retain higher order terms in the Taylor Expansion. The **Runge-Kutta** method is an extension of the Euler Method where we retain terms up to h^5 . Thus, we have:

$$y_{n+1} = y_n + \frac{1}{6}(k_0 + 2k_1 + 2k_2 + k_3)$$

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where,

$$\begin{aligned}k_0 &= hf(t_n, y_n) \\k_1 &= hf\left[t_n + \frac{h}{2}, y_n + \frac{k_0}{2}\right] \\k_2 &= hf\left[t_n + \frac{h}{2}, y_n + \frac{k_1}{2}\right] \\k_3 &= hf[t_n + h, y_n + k_2]\end{aligned}$$

Runge-Kutta for the Rijke tube model

In the Rijke tube model, we have:

$$\frac{d \eta_j}{dt} = f(t, \eta_j, \dot{\eta}_j) \quad \frac{d \dot{\eta}_j}{dt} = g(t, \eta_j, \dot{\eta}_j) \quad \text{for } j = 1, 2, \dots, N \text{ Galerkin modes}$$

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The solution will proceed along the lines of

$$\eta_j^{n+1} = \eta_j^n + \frac{1}{6} [k_0 + 2k_1 + 2k_2 + k_3]$$

$$\dot{\eta}_j^{n+1} = \dot{\eta}_j^n + \frac{1}{6} [l_0 + 2l_1 + 2l_2 + l_3]$$

Runge-Kutta for the Rijke tube model

In the Rijke tube model, we have:

$$\frac{d \eta_j}{dt} = f(t, \eta_j, \dot{\eta}_j) \quad \frac{d \dot{\eta}_j}{dt} = g(t, \eta_j, \dot{\eta}_j) \quad \text{for } j = 1, 2, \dots, N \text{ Galerkin modes}$$

These two equations need to be solved simultaneously.

The solution will proceed along the lines of

$$\eta_j^{n+1} = \eta_j^n + \frac{1}{6} [k_0 + 2k_1 + 2k_2 + k_3]$$

$$\dot{\eta}_j^{n+1} = \dot{\eta}_j^n + \frac{1}{6} [l_0 + 2l_1 + 2l_2 + l_3]$$

Tackling time-delay τ

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You will still have to calculate u'_f for the time steps $t = 0, \dots, \tau$ and use these values to find the solution to $\dot{\eta}(t = \tau + 1), \dot{\eta}(t = \tau + 2), \dots$ and so on.

$u_f'(t - \tau)$ inside the square root term

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$$u_f'(t - \tau) \approx u_f'(t) - \tau \frac{\partial u_f'(t)}{\partial t} \qquad u'(x_f, t) = \sum_{j=1}^N \eta_j(t) \cos(j\pi x_f)$$

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Conclusion

References

Contact

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