

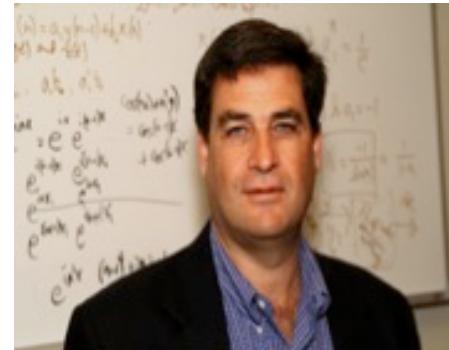
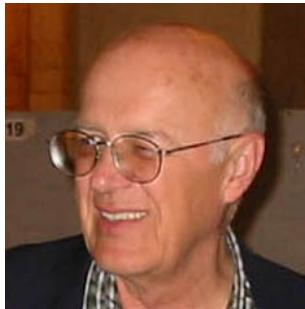


21st International Congress on Acoustics
165th Meeting of the Acoustical Society of America
52nd Meeting of the Canadian Acoustical Association

2-7 June 2013
Montréal, Canada

Buckling as a source of sound, with application to the modeling of cicada sound generation

Allan D. Pierce, Derke R. Hughes, Kossi Edoh, Richard A. Katz, Robert M. Koch

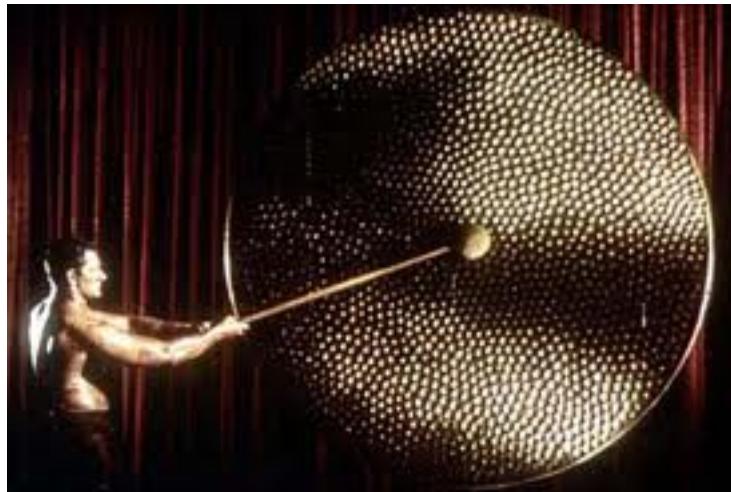
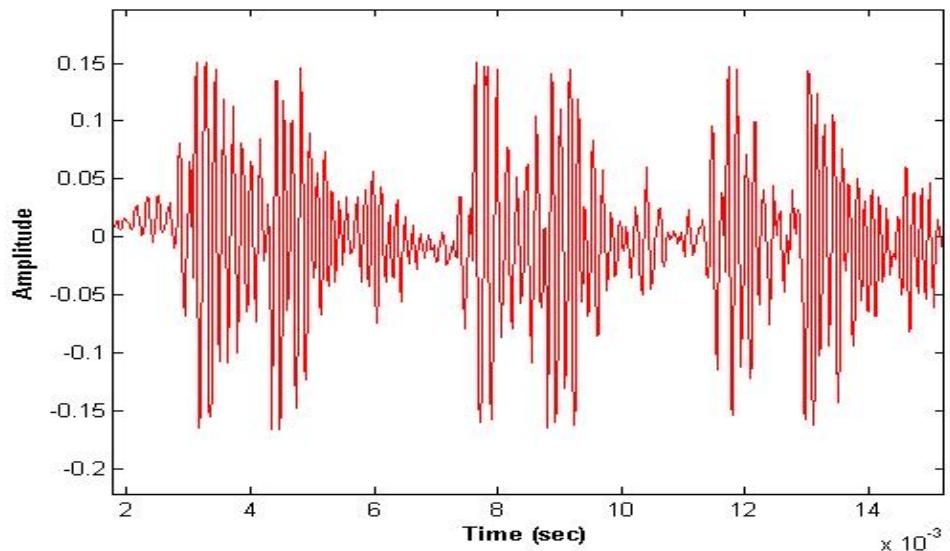




Noisy little (not necessarily) bugs

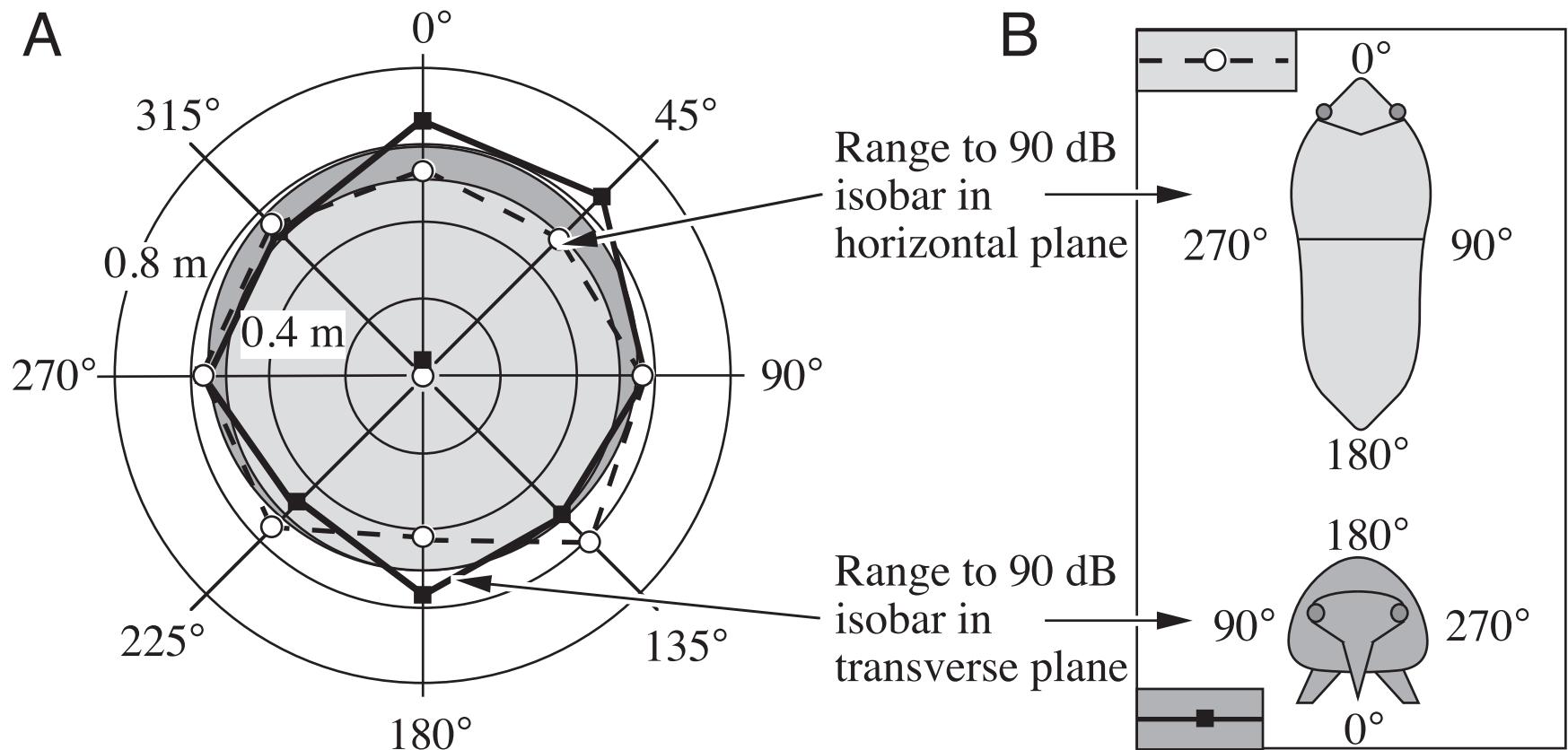


Source: Massachusetts Cicada



cymbals

90 dB in air at a distance of 0.8 m



Bennet-Clark and Daws, J. Exper. Biol. (1999)

A PHYSIOLOGICAL ANALYSIS OF CICADA SONG

By J. W. S. PRINGLE

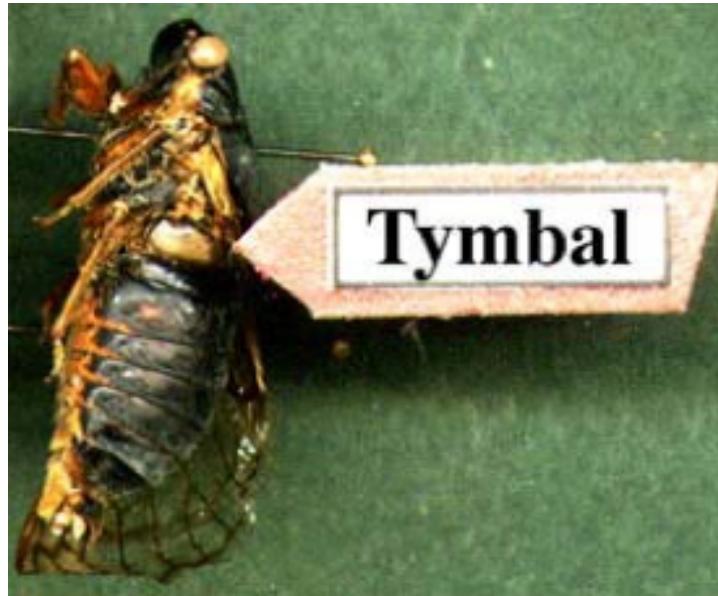
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of Physiology, University of Ceylon, Colombo*

(Received 10 February 1954)

John William Sutton Pringle, FRS (1912-1982)

Pringle was the first to explicitly mention the word “buckling”

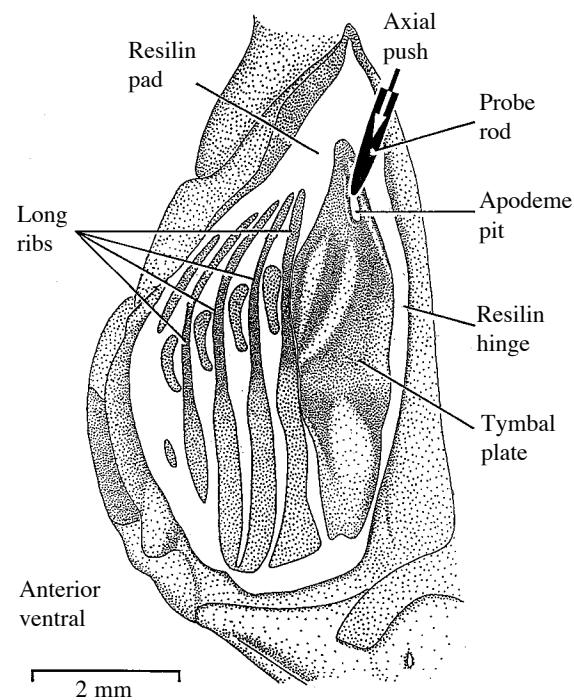
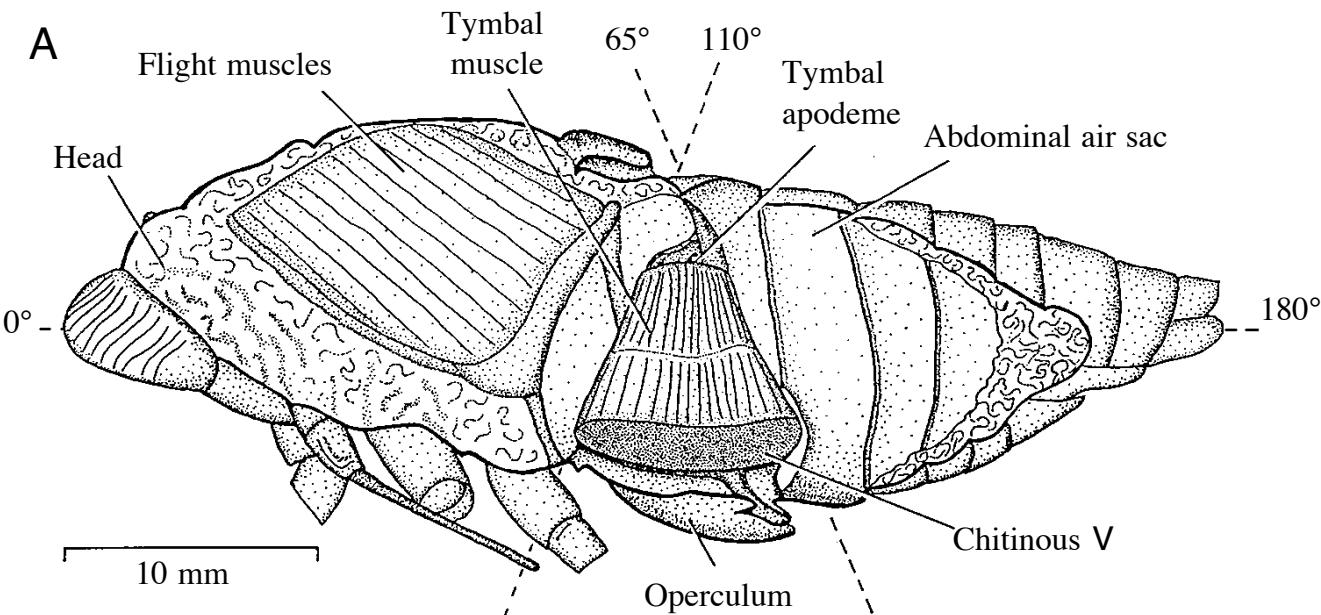
A pulse of sound is emitted when the tymbal suddenly buckles or is restored to its resting position by its natural elasticity; in the song of some species both movements are effective. The tymbal muscles, which are responsible for the buckling, have a myogenic rhythm of activity, initiated, but only slightly controlled in frequency, by impulses in the single nerve fibre supplying each muscle. The two tymbals normally act together.



The word **tymbal** is defined in Merriam-Webster Unabridged Dictionary as the **“vibrating membrane in the shrilling organ of a cicada”**--- so it is a term invented by biologists to describe part of the anatomy of cicadas. OED attributes its origin to someone named Bushman in 1854



cymbals



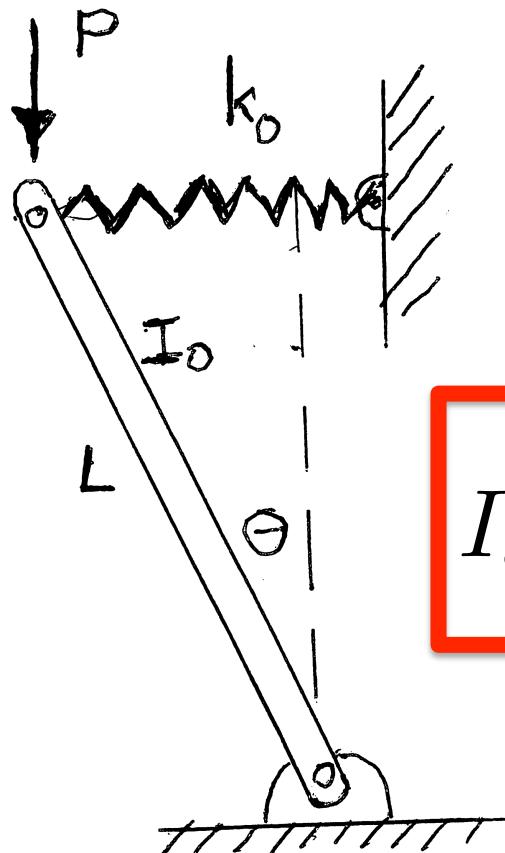
*Long ribs do the
actual buckling*

What is buckling?

Buckling caused the recent Bangladesh building collapse



Exposition to explain how buckling generates sound



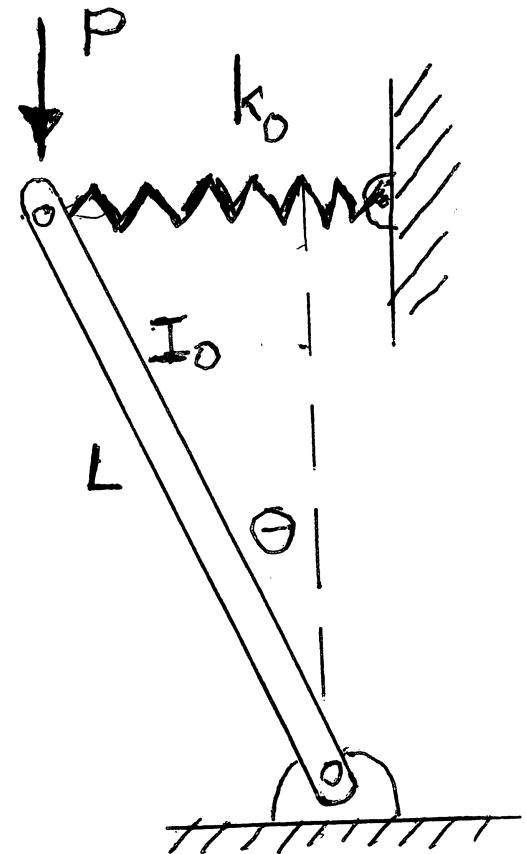
$$I_o \frac{d^2\theta}{dt^2} = PL \sin \theta - k_0 L^2 \sin \theta.$$

approximates to

$$I_o \frac{d^2\theta}{dt^2} + (k_0 L^2 - PL) \theta = 0$$

Buckling (critical) load

$$P_{\text{cr}} = k_o L$$



$$\frac{d^2\theta}{dt^2} + \frac{L}{I_o} (P_{\text{cr}} - P) \theta = 0$$

If (muscle) applied compressive force is ***below the buckling load***:

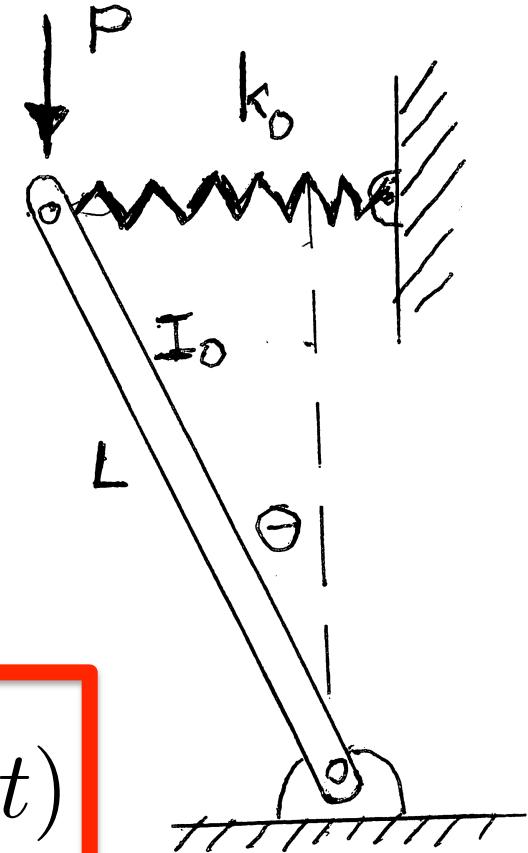
$$P < P_{\text{cr}}$$

The rod (rib) ***oscillates*** back and forth

$$\theta = A \cos(\omega_o t) + B \sin(\omega_o t)$$

with the angular frequency

$$\omega_o = \left(\frac{L}{I_o} \right)^{1/2} (P_{\text{cr}} - P)^{1/2}$$



If (muscle) applied compressive force is ***above the buckling load***:

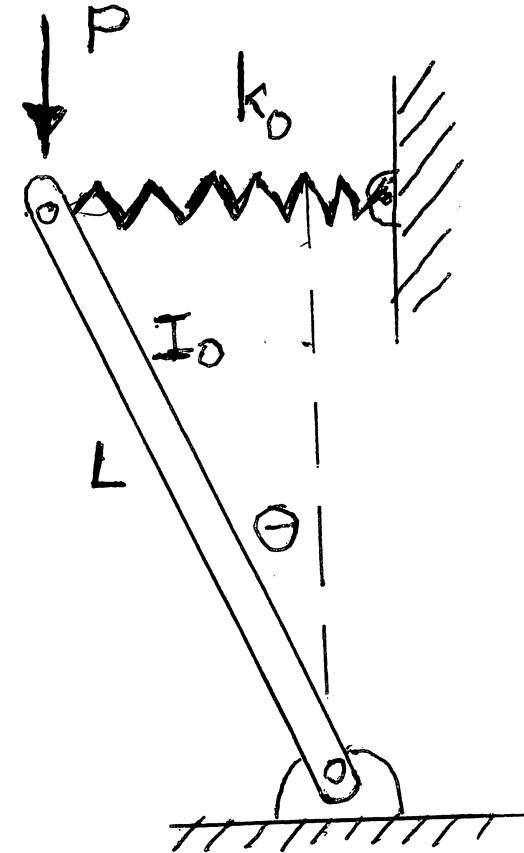
$$P > P_{\text{cr}}$$

The displacement of the rod (rib) grows exponentially

$$\theta = A e^{\alpha_o t} + B e^{-\alpha_o t}$$

with the exponential growth rate

$$\alpha_o = \left(\frac{L}{I_o} \right)^{1/2} (P - P_{\text{cr}})^{1/2}$$



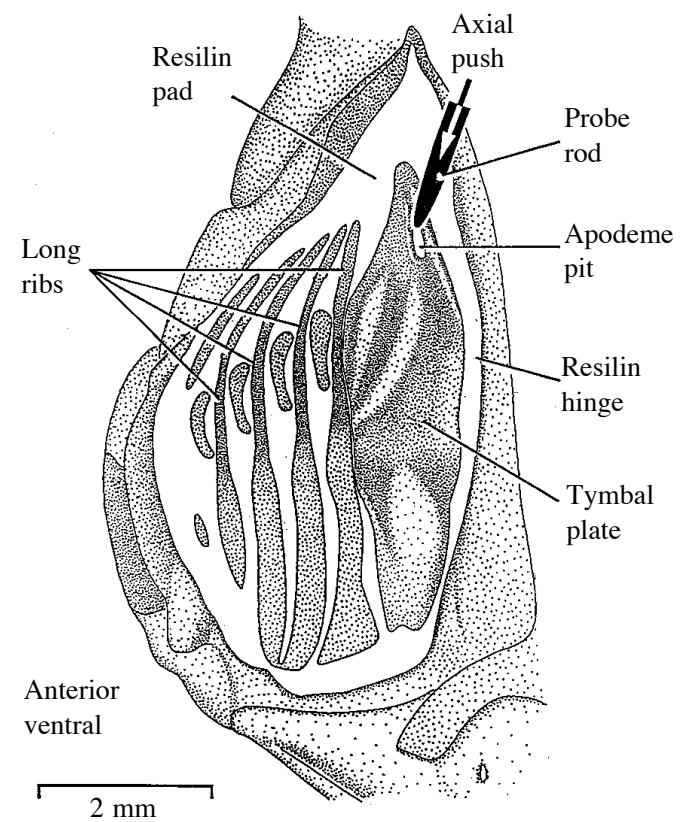
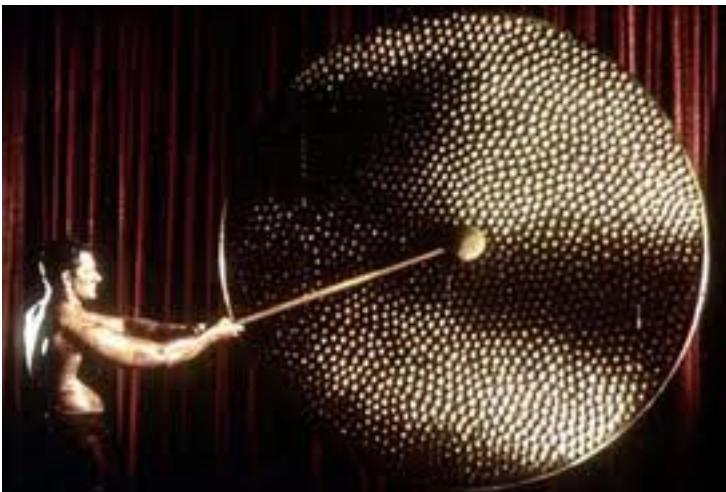
If this exponential growth isn't arrested ***then the rib will break.***



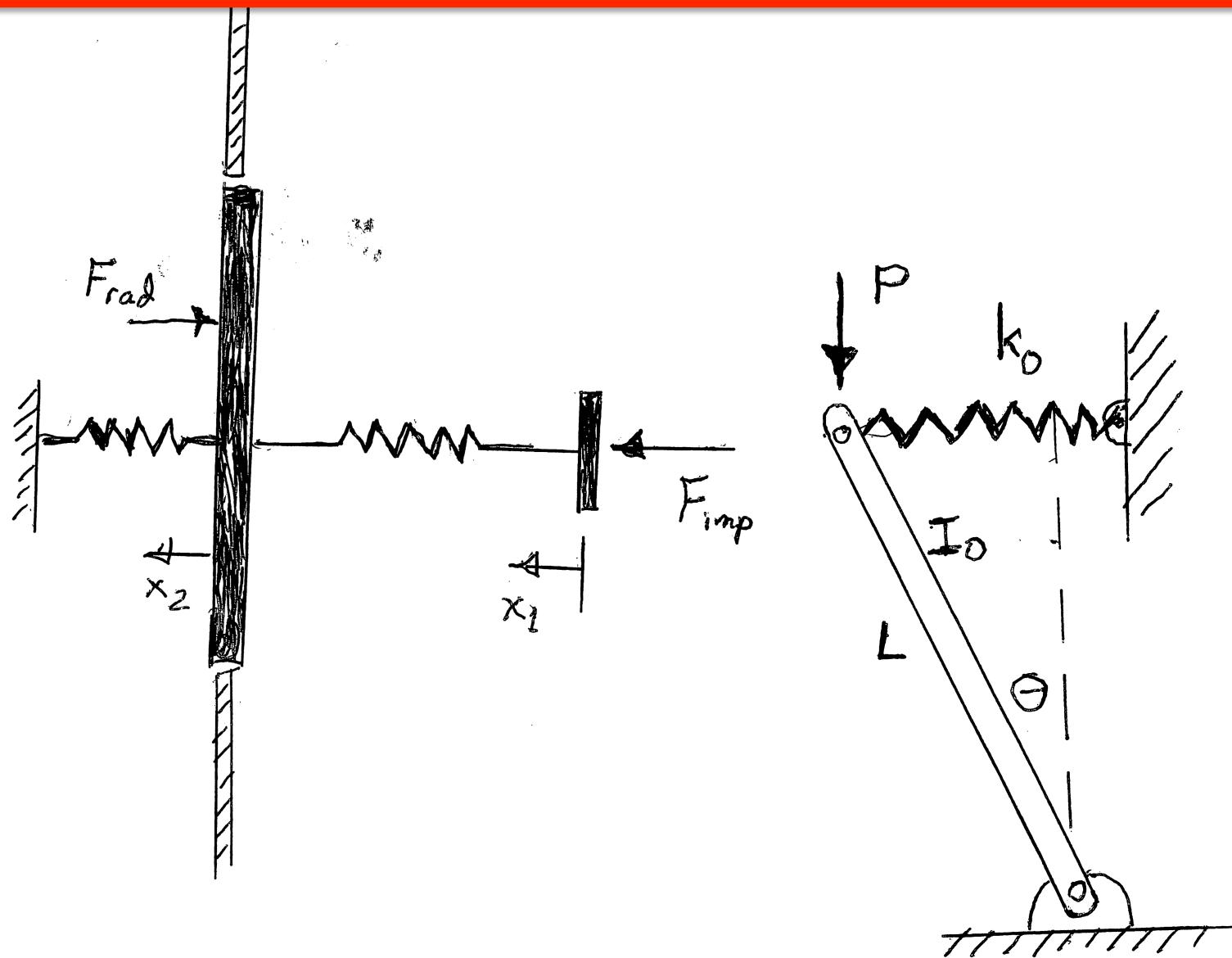
Something like what happened in Bangladesh.

In the cicada the exponential growth of the buckling rib is **arrested**.

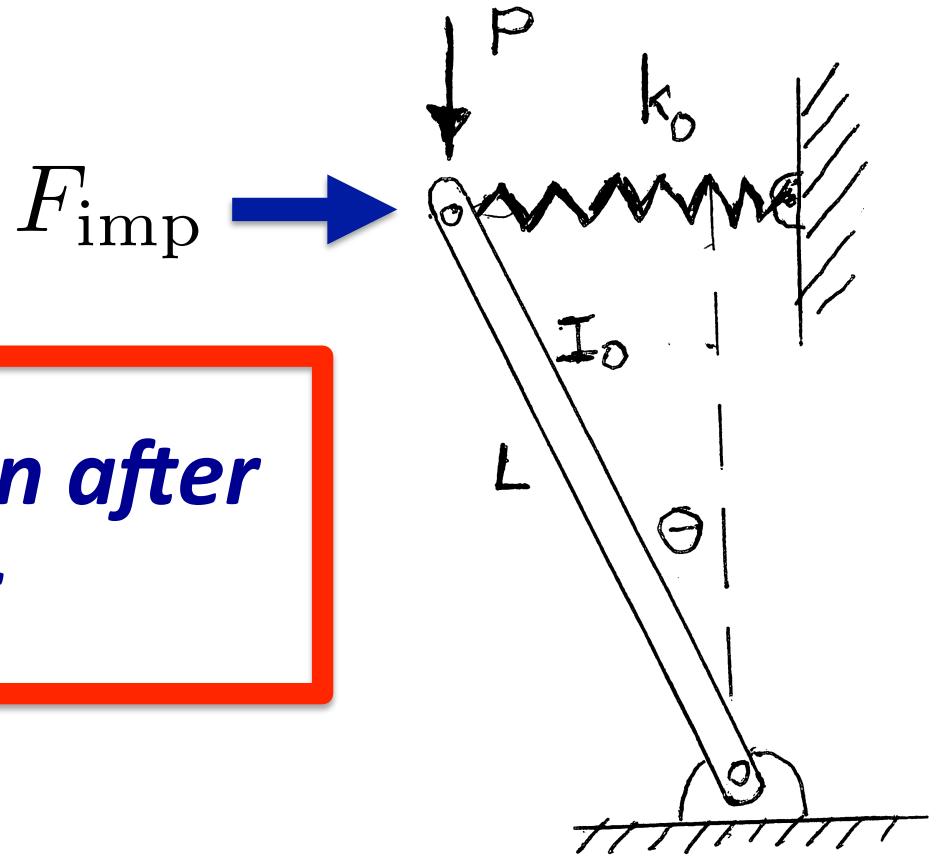
Result is analogous to a hammer striking a gong.



Sketch of a model for an arresting mechanism

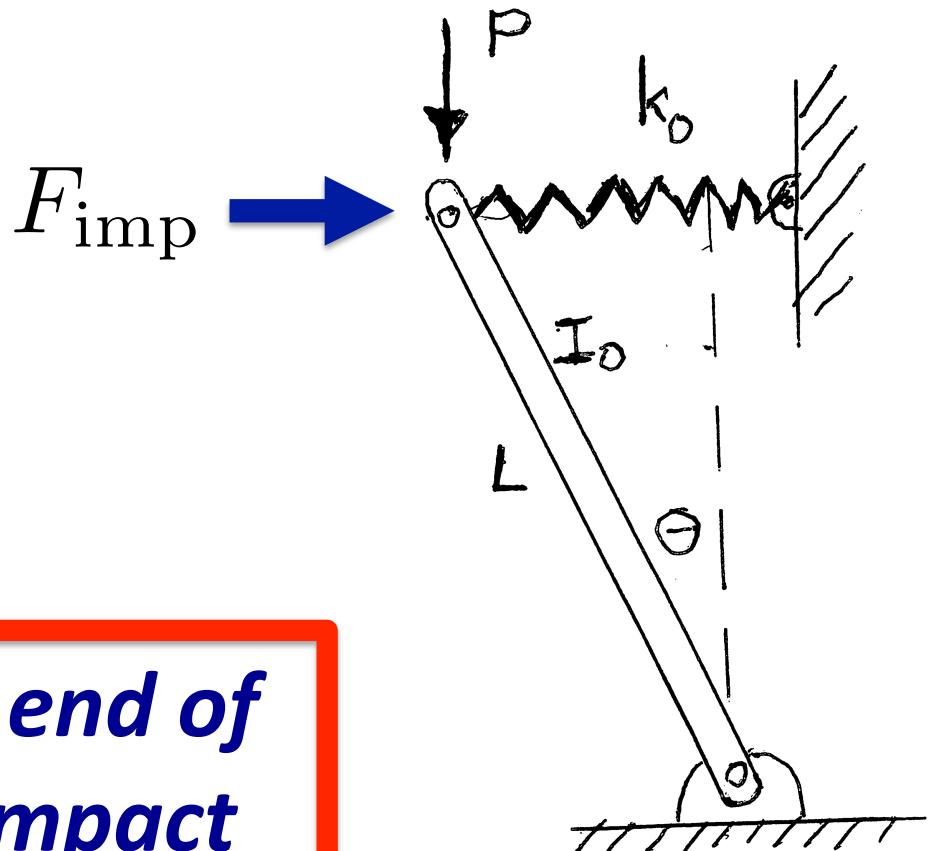


Motion of rod motion after impact with arrester



$$I_o \frac{d^2\theta}{dt^2} + (k_0 L^2 - PL) \theta = -F_{\text{imp}} L$$

$$x_1 = L (\theta - \theta_{\text{ar}})$$



Displacement of the end of the rod after initial impact with the arrester

Coupled differential equations

$$\frac{I_o}{L^2} \frac{d^2x_1}{dt^2} + \left(k_o - \frac{P}{L} \right) (x_1 + L\theta_{\text{ar}}) + k_1 (x_1 - x_2) = 0$$

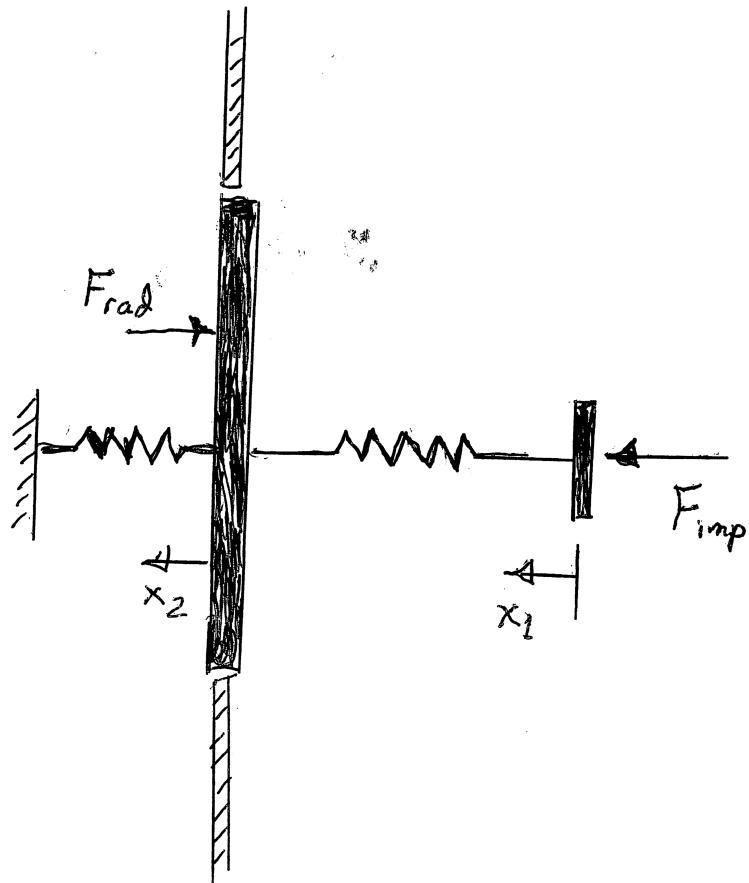
$$M \frac{d^2x_2}{dt^2} + k_2 x_2 + k_1 (x_2 - x_1) = -F_{\text{rad}}$$

The second displacement variable can be regarded as a representative tymbal displacement

Initial conditions at time of impact

$$x_1 = 0$$

$$\frac{dx_1}{dt} = V_0 \quad \text{at} \quad t = 0$$



After the impact, the second mass (M) is quickly accelerated to the velocity

$$V_{\text{begin}} = \frac{(I_o/L^2)}{M + (I_o/L^2)} V_o$$

so the second mass (tymbal) starts to oscillate with an initial energy

$$E_{\text{begin}} = \frac{1}{2} \frac{M(I_o/L^2)^2}{[M + (I_o/L^2)]^2} V_o^2$$

Radiated energy is maximized if tymbal mass is

$$M_{\text{opt}} = \frac{I_o}{L^2}$$

Shortly after impact the tymbal is in damped oscillation

$$M \frac{d^2 x_2}{dt^2} + k_2 x_2 = -F_{\text{rad}}$$

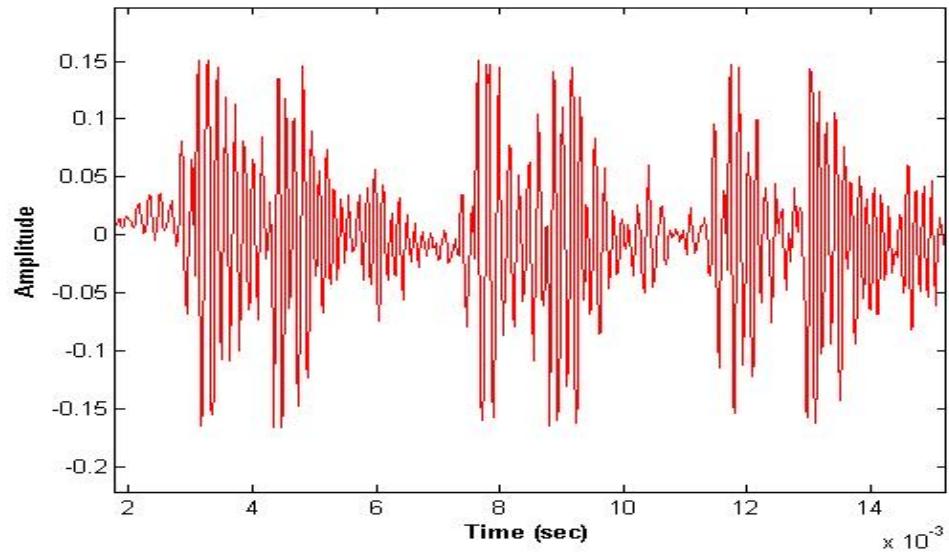
$$F_{\text{rad}} \approx -\frac{\rho \pi a^4}{2c} \frac{d^3 x_2}{dt^3}$$

Solution for tymbal motion

$$x_2 = \left(\frac{V_{\text{begin}}}{\omega_{\text{vib}}} \right) G(t) \sin(\omega_{\text{vib}} t)$$

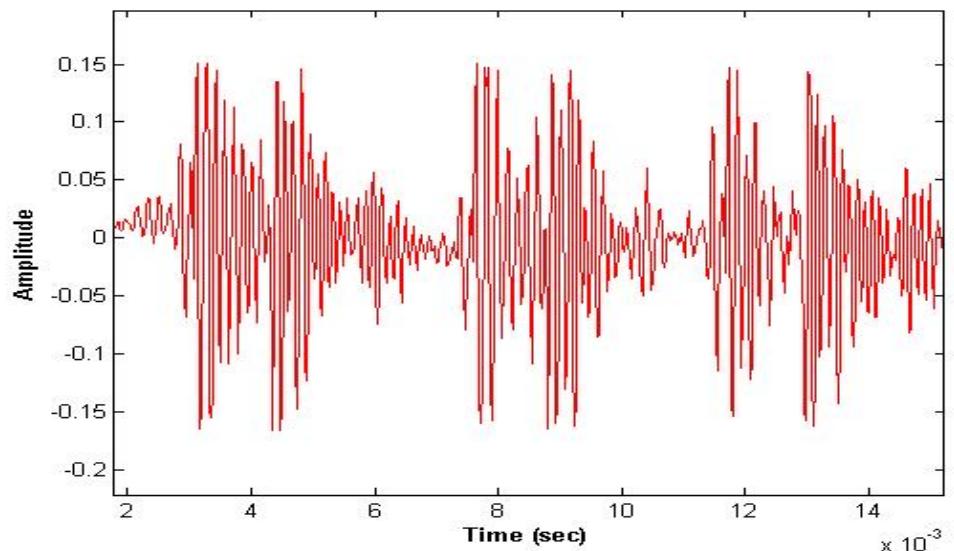
$$\omega_{\text{vib}} = \left(\frac{k_2}{M} \right)^{1/2}$$

$$G(t) = e^{-\alpha_{\text{rad}} t} H(t)$$



Radiated pressure

$$p = -\omega_{\text{vib}} V_{\text{begin}} \frac{\rho A}{2\pi r} G \left(t - \frac{r}{c} \right) \sin \left[\omega_{\text{vib}} \left(t - \frac{r}{c} \right) \right]$$



Radiated power

$$P_{\text{rad}} = \frac{\omega_{\text{vib}}^2 \rho A^2}{4\pi c} V_{\text{begin}}^2 G^2(t)$$

Total energy radiated

$$E_{\text{rad}} = \frac{1}{2\alpha_{\text{rad}}} \frac{\omega_{\text{vib}}^2 \rho A^2}{4\pi c} V_{\text{begin}}^2$$

Conservation of energy

$$E_{\text{rad}} = E_{\text{rad}}$$

Concluding remarks

- *Energy initially stored in muscle transferred to rib motion*
- *Energy in rib motion transferred to tymbal motion*
- *Tymbal radiates nearly as a baffled piston*
- *Energy in vibrating piston transferred to acoustic energy*