

Stars :

Radius of the Sun $\approx 7 \times 10^5$ km : Sun is a normal star.

Normal
Chemical (1)
or
Main Sequence
star →

Classical physics is mostly enough to explain the structure, stability and existence of Sun

Main Sequence in the H-R diagram.

(2) Quantum Stars : white dwarf, Neutron stars,
→ Black holes

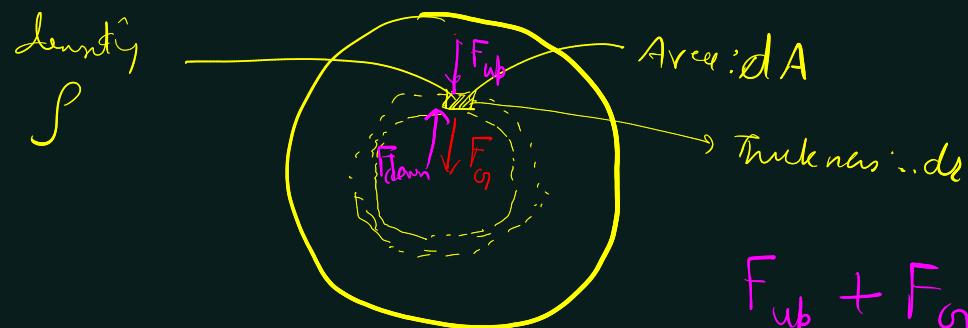
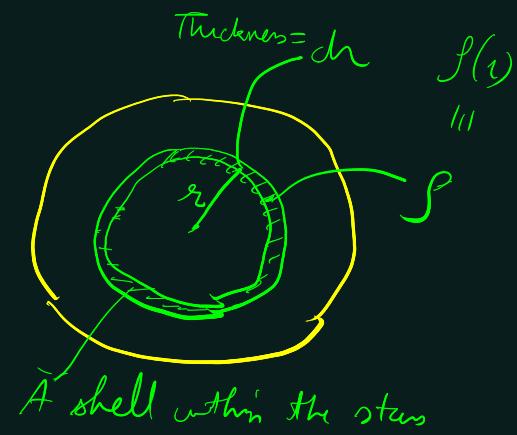
Astrophysics Books:

1. Astronomy: A Physical perspective, M. Kutner
2. Foundations of Astrophysics, B. Ryden
3. Physical Universe, F. Shu
4. An Introduction to Modern Astrophysics, Carroll & Ostlie

Stability of stars: Basic Equation

Mass of the shell: $dM = 4\pi r^2 \rho dr$

$$\rightarrow \boxed{\frac{dM}{dr} = 4\pi r^2 \rho} \quad (1)$$



$$F_{up} + F_{in} = F_{down}$$

$$(P + dP) dA + \frac{GM}{r^2} dr = P dA$$

$$\rightarrow (P + dP) dA + \frac{GM \int dA dr}{r^2} = P dA$$

$$\Rightarrow \left\{ \begin{array}{l} \frac{dP}{dr} = - \frac{GM}{r^2} \rho \\ \frac{dM}{dr} = 4\pi r^2 \rho \end{array} \right. \quad - (2) \quad \begin{array}{l} \text{Mass} \\ \uparrow \\ M, P, \rho \end{array} \quad \begin{array}{l} \text{pressure} \\ \uparrow \\ r \end{array} \quad \begin{array}{l} \text{density} \\ \uparrow \\ \rho \end{array}$$

are unknowns

student answer : Ideal gas law : $P \propto ST \rightarrow P = \frac{\rho k_B T}{m_{\text{atom}}}$

But this introduce one more variable which is T

we need to consider some general way

Let's assume $P = K \rho^\gamma$

$K = constant$

A general form
law
Polytropic
relation
through index
 γ

Let's solve analytically :

Scaling (Grav) way

$$\frac{dM}{dr} = 4\pi r^2 \rho \quad \rightarrow \quad \frac{M}{R} = R^2 \frac{M}{R^3} = \frac{M}{R}$$

$$\frac{dp}{dr} = -\frac{GM}{r^2} \rho \quad \rightarrow \quad \frac{P}{R} = \frac{M}{R^2} \frac{M}{R^3} \Rightarrow P = \frac{GM^2}{R^4}$$

$$P = K \frac{M^\gamma}{R^{3\gamma}}$$

$$\Rightarrow K \frac{M^\gamma}{R^{3\gamma}} = \frac{GM^2}{R^4} \Rightarrow M^{\gamma-2} = \frac{G}{K} R^{3\gamma-4}$$

for $\gamma = \frac{5}{3}$:

$$M^{\frac{5}{3}-2} = \frac{G}{K} R^{3 \cdot \frac{5}{3} - 4}$$

$$\Rightarrow M^{-1/3} = \frac{G}{K} R$$

$$\Rightarrow \boxed{R \propto \frac{1}{M^{1/3}}}$$

$$\text{For } \gamma = \frac{4}{3} : \Rightarrow M^{\frac{4}{3}-2} = \frac{C}{k} R^{3 \cdot \frac{4}{3}-4}$$

$$\Rightarrow M^{-2/3} = ? = \frac{C}{k}$$

Equation of state:

Boltzmann's ratio = $\frac{N_2}{N_1} = e^{-h\nu/k_B T}$

where $\nu = \nu_2 - \nu_1$

Kinetic Theory

Pressure: $P \propto n (\text{KE})$: Classical Ideal gas
 $P \propto n k_B T$

Now what will happen for a quantum gas:

Let's see

Imagine Sun contracted to very small size and it is a plasma of e^- and p^+

Then what is K.E. ?

Let's calculate:

Imagine we have N electrons packed densely in volume V

$$\text{then } (\Delta v) \text{ volume available to each electron} = \frac{V}{N} = \frac{1}{n}$$

$$\text{the space available to each } n \quad \Delta x \approx (\Delta v)^{1/3} \approx \frac{1}{n^{1/3}}$$

$$\Rightarrow \Delta x \Delta p \approx \hbar$$

Therefore roughly the momentum of such electrons is of the order $p \propto \frac{\hbar}{\Delta x}$

$$\Rightarrow \boxed{p_e \approx \hbar n^{1/3}}$$

$$\text{Kinetic energy : } K.E. = \frac{p_e^2}{2m_e} = \frac{\hbar^2 n^{2/3}}{2m_e}$$

$$\Rightarrow \text{Pressure} = P = n(K.E.) = \frac{\hbar^2 n^{5/3}}{2m_e}$$

$$\Rightarrow \rho = \frac{N m_e + N m_p}{V} \approx n m_p \quad \text{because } m_e \ll m_p$$

And $KE = \frac{p_e^2}{2m_e} + \frac{p_p^2}{2m_p} = \frac{\hbar^2 n^{2/3}}{2m_e} + \frac{\hbar^2 n^{2/3}}{2m_p} \approx \frac{\hbar^2 n^{2/3}}{2m_e}$

$$\Rightarrow n = \rho/m_p$$

$$\Rightarrow P = \frac{\hbar^2}{2m_e m_p^{5/3}} \rho^{5/3} \Rightarrow P = K \rho^\gamma$$

where $K = \frac{\hbar^2}{2m_e m_p^{5/3}}$ and $\gamma = \frac{5}{3}$

Stellar structure:

$$\frac{M}{R}$$

$$M^{-1/3} = \frac{G}{K} R \Rightarrow R = \frac{\hbar^2}{2(Gm_e m_p^{5/3})} M^{-1/3}$$

we can drop 2

Find out R for $M = 2 \times 10^{30} \text{ kg}$

only taking physics

However $K.E. = \frac{p^2}{2m}$ is only valid for non-relativistic particles

what if the particles are relativistic

The $K.E. = \sqrt{p^2 c^2 + m_0^2 c^4} - m_0 c^2$

\downarrow
rest mass energy

Relativistic case $p c \gg m_0 c^2 \Rightarrow K.E. \approx p c$

Non-relativistic case $p c \ll m_0 c^2 \Rightarrow K.E. = m_0 c^2 \left(1 + \frac{p^2}{m_0^2 c^2}\right)^{1/2} - m_0 c^2$

$$\Rightarrow K.E. \approx \cancel{m_0^2 c^2} + m_0 c^2 \frac{p^2}{2m_0^2 c^2} - \cancel{m_0 c^2}$$

$$\Rightarrow K.E. \approx \frac{p^2}{2m_0}$$

For electrostatic case: $K.E. = \frac{1}{2}pc$

$$\Rightarrow \text{Prism: } P = n(K.E.) = n \frac{1}{2}pc$$

but $p_e = \frac{1}{3}n^{4/3}$ (electrons) as derived above

$$\Rightarrow P = \frac{1}{3}n^{4/3}c = \frac{\frac{1}{3}c}{m_p^{4/3}} \rho^{4/3} = K \rho^\gamma$$

where $K = \frac{\frac{1}{3}c}{m_p^{4/3}}$ and $\gamma = 4/3$

From stellar structure equation: in this case: $M^{-2/3} = \frac{G}{K}$

$$\Rightarrow M = \left(\frac{K}{G} \right)^{3/2} = \left(\frac{\frac{1}{3}c}{G m_p^{4/3}} \right)^{3/2}$$

$$\approx M = \sqrt[3]{\frac{\frac{1}{3}c^3}{G^3 m_p^4}}$$

A constant mass for which
radius $R = 0$

$$\approx M_c = \sqrt{\frac{G^2 c^3}{G^3 m_p^4}}$$

$\hbar = 1.05 \times 10^{-34}$
 $c = 3 \times 10^8$
 $G = 6.67 \times 10^{-11}$
 $m_p = 1.67 \times 10^{-24}$

S.I.
units

Remember mass of Sun = 2×10^{30} kg = M_\odot

$$M_c \approx 2 M_\odot \quad \left\{ \text{we get this but verify} \right.$$

very close to actual

$$\underline{\underline{M_{ch} \approx 1.4 M_\odot}}$$

This is the Chandrasekhar limit for white stars.

Let's look at the speed of moving electron (degenerate electrons) in a white dwarf star.

$$\Delta x \Delta p \approx \hbar \quad \text{when} \quad v_e \approx \frac{p_e}{m_e} \approx \frac{\Delta p_e}{m_e} \approx \frac{\hbar}{\Delta x} \propto \frac{\hbar n^{1/3}}{m_e}$$

ΔV = Volume available to each electron : $\frac{V}{N} = \frac{1}{n}$

$$\Delta x \approx (\Delta V)^{1/3} \approx \frac{1}{n^{1/3}}$$

Therefore $v_e = \frac{\hbar s^{1/3}}{m_e m_p^{1/3}}$ $\Rightarrow v_e \propto s^{1/3}$

$$s \propto v_e^3$$

$$v_e = \frac{\hbar}{m_e m_p^{1/3}} \left(\frac{M}{R^3} \right)^{1/3} = \frac{\hbar}{m_e m_p^{1/3}} \left(\frac{(M)^{1/3}}{L_s M^{-1/3}} \right)$$

$$\Rightarrow v_e = \frac{\hbar}{m_e m_p^{1/3}} \frac{L_s}{R} M^{2/3}$$

At what M ; $\mathcal{G}_e = C$

Say at $M_{\max} \Rightarrow M_{\max}^{2/3} = \frac{C m e^{m_p/3}}{\hbar} \frac{k}{G}$

$$K = \frac{\hbar^2}{m e^{m_p/3}} \Rightarrow M_{\max}^{2/3} = \frac{C \cancel{m e^{m_p/3}}}{\cancel{k} G} \frac{\hbar^2}{\cancel{m e^{m_p/3}}} \\ = \frac{\hbar C}{G m_p^{4/3}}$$

$$\Rightarrow M_{\max} = \sqrt[3]{\frac{\hbar^3 C^3}{G^3 m_p^4}}$$

Mass of star (now)

< 7 Msun

> 7 Msun & < 18 Msun

> 18 Msun

Mass of star (at death)

< 1.4 Msun

> 1.4 Msun

> 1.4 Msun

Fate

White Dwarf

Neutron Star

Black Hole

