

Stars:

Radius of the Sun $\approx 7 \times 10^5$ km : Sun is a normal star.

Normal
Classical (1)
or
Main Sequence
star \rightarrow classical physics is mostly enough to explain the structure, stability and existence of Sun
Main Sequence in the H-R diagram.

(2) Quantum Stars: white dwarf, Neutron stars,
 \rightarrow Black Holes

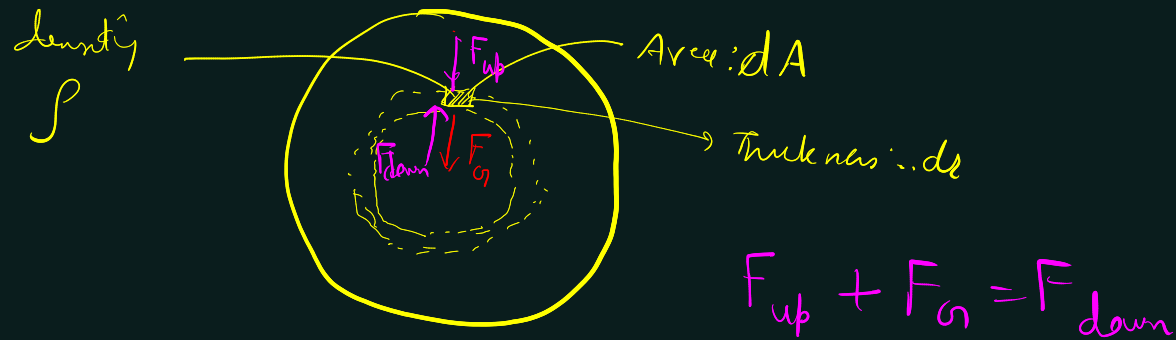
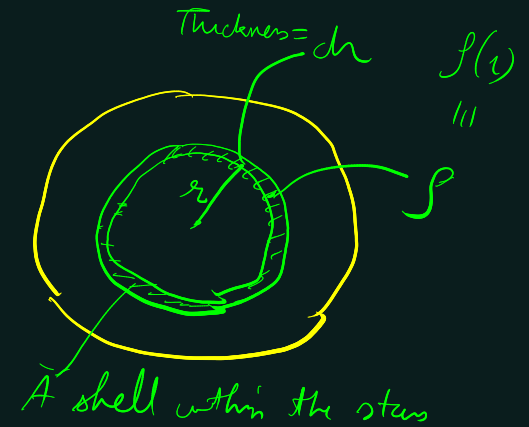
Astrophysics Books:

1. Astronomy: A Physical perspective, M. Kutner
2. Foundations of Astrophysics, B. Ryden
3. Physical Universe, F. Shu
4. An Introduction to Modern Astrophysics, Carroll & Ostlie

Stability of stars: Basic Equation

Mass of the shell: $dM = 4\pi r^2 \rho dr$

$$\Rightarrow \boxed{\frac{dM}{dr} = 4\pi r^2 \rho} \quad (1)$$



$$(P + dP) dA + \frac{GM}{r^2} dm = P dA$$

$$\Rightarrow (P + dP) dA + \frac{GM \rho dA dr}{r^2} = P dA$$

⇒

$$\boxed{\begin{aligned}\frac{dP}{dr} &= - \frac{GM}{r^2} \rho \\ \frac{dM}{dr} &= 4\pi r^2 \rho\end{aligned}} \quad \text{--- (2)}$$

Mass ↑ Pressure ↑ density ↑
 M, P, ρ
are unknowns

Student propose: Ideal gas law: $P \propto \rho T$ $\therefore P = \frac{\rho k_B T}{m_{\text{proton}}}$

But this introduce one more variable which is T

we need to consider some general way

Let's assume $P = K \rho^\gamma$

$K = \text{constant}$

A general power law
polytropic
relation
through index
 γ

Let's solve analytically:

Scaling (Grady) way

$$\begin{aligned} \frac{dM}{dr} &= 4\pi r^2 \rho \\ \frac{dp}{dr} &= -\frac{GM}{r^2} \rho \\ P &= K \rho^\gamma \end{aligned} \quad \rightarrow \quad \begin{aligned} \frac{M}{R} &= R^2 \frac{M}{R^3} = \frac{M}{R} \\ \frac{P}{R} &= \frac{M}{R^2} \frac{M}{R^3} \Rightarrow P = \frac{GM^2}{R^4} \\ P &= K \frac{M^\gamma}{R^{3\gamma}} \end{aligned}$$

$$\Rightarrow K \frac{M^\gamma}{R^{3\gamma}} = \frac{GM^2}{R^4} \Rightarrow M^{\gamma-2} = \frac{G}{K} R^{3\gamma-4}$$

for $\gamma = \frac{5}{3}$: $M^{\frac{5}{3}-2} = \frac{G}{K} R^{3 \cdot \frac{5}{3} - 4}$

$$\Rightarrow M^{-1/3} = \frac{G}{K} R$$

$$\Rightarrow \boxed{R \propto \frac{1}{M^{1/3}}}$$

$$\text{For } \gamma = \frac{4}{3}; \quad \Rightarrow \quad M^{\frac{4}{3}-2} = \frac{G}{K} R^{3 \cdot \frac{4}{3}-4}$$

$$\Rightarrow \quad M^{-2/3} = ? = \frac{G}{K}$$

Equation of state:

Boltzmann's ratio = $\frac{N_2}{N_1} = e^{-h\nu/k_B T}$

where $\nu = \nu_2 - \nu_1$

Kinetic Theory

Pressure: $P \propto n \text{ (K.E.)}$

Classical
Ideal gas
 $P \propto n k_B T$

Now what will happen for a quantum gas:

Let's see

Imagine Sun contracted to very small size and it is a plasma of e^- and p^+

Then what is K.E.?

Let's calculate:

Imagine we have N electrons packed closely in volume V

$$\text{then } (\Delta V) \text{ volume available to each electron} = \frac{V}{N} = \frac{1}{n}$$

$$\text{the space available to each } \Delta x \approx (\Delta V)^{1/3} \approx \frac{1}{n^{1/3}}$$

$$\Rightarrow \Delta x \Delta p \approx \hbar$$

Therefore roughly the momentum of such electrons is of the order $p \approx \frac{\hbar}{\Delta x}$

$$\Rightarrow \boxed{p_e \approx \hbar n^{1/3}}$$

$$\text{Kinetic energy: } K.E. = \frac{p_e^2}{2m_e} = \frac{\hbar^2 n^{2/3}}{2m_e}$$

$$\Rightarrow \text{Pressure} = P = n(K.E.) = \frac{\hbar^2 n^{5/3}}{2m_e}$$

$$\Rightarrow \rho = \frac{N m_e + N m_p}{\sqrt{\quad}} \approx n m_p \quad \text{became } m_e \ll m_p$$

$$\text{And } KE = \frac{p_e^2}{2m_e} + \frac{p_p^2}{2m_p} = \frac{\hbar^2 n^{2/3}}{2m_e} + \frac{\hbar^2 n^{2/3}}{2m_p} \approx \frac{\hbar^2 n^{2/3}}{2m_e}$$

$$\Rightarrow n = \rho / m_p$$

$$\Rightarrow P = \frac{\hbar^2}{2m_e m_p^{5/3}} \rho^{5/3} \Rightarrow P = K \rho^\gamma$$

$$\text{where } K = \frac{\hbar^2}{2m_e m_p^{5/3}} \text{ and } \gamma = \frac{5}{3}$$

stellar structure:

M
 R

$$M^{-1/3} = \frac{G}{K} R$$

$$\Rightarrow R = \frac{\hbar^2}{2G m_e m_p^{5/3}} M^{-1/3}$$

we can drop 2

only keeping physics

Find out R for $M = 2 \times 10^{30} \text{ kg}$

However $K.E. = \frac{p^2}{2m}$ is only valid for non-relativistic particles

what if the particles are relativistic

The $K.E. = \sqrt{p^2 c^2 + m_0^2 c^4} - m_0 c^2$
 \downarrow
rest mass energy

Relativistic case $pc \gg m_0 c^2 \Rightarrow K.E. \approx pc$

Non-relativistic case $pc \ll m_0 c^2 \Rightarrow K.E. = m_0 c^2 \left(1 + \frac{p^2}{m_0^2 c^2} \right)^{1/2} - m_0 c^2$

$$\Rightarrow K.E. \approx \cancel{m_0 c^2} + \frac{m_0 c^2 p^2}{2 m_0^2 c^2} - \cancel{m_0 c^2}$$

$$\Rightarrow K.E. \approx \frac{p^2}{2 m_0}$$

For relativistic case: $K.E. = pc$

$$\Rightarrow \text{Pressure: } P = n (K.E.) = n pc$$

$$\text{but } p_e = \hbar n^{1/3} \quad (\text{electrons}) \text{ as derived above}$$

$$\Rightarrow P = \hbar n^{4/3} = \frac{\hbar c}{m_p^{4/3}} \rho^{4/3} = K \rho^\gamma$$

$$\text{where } K = \frac{\hbar c}{m_p^{4/3}} \quad \text{and } \gamma = 4/3$$

From stellar structure equation: in this case: $M^{-2/3} = \frac{G}{K}$

$$\Rightarrow M = \left(\frac{K}{G} \right)^{3/2} = \left(\frac{\hbar c}{G m_p^{4/3}} \right)^{3/2}$$

$$\Rightarrow M = \sqrt{\frac{\hbar^3 c^3}{G^3 m_p^4}}$$

A constant mass for which
radius $R = 0$

$$\approx M_c = \sqrt{\frac{h^2 c^3}{G^3 m_p^4}}$$

$$h = 1.05 \times 10^{-34}$$

$$c = 3 \times 10^8$$

$$G = 6.67 \times 10^{-11}$$

$$m_p = 1.67 \times 10^{-24}$$

} S.I. units

Remember mass of Sun $= 2 \times 10^{30} \text{ Kg} = M_\odot$

$$M_c \approx 2 M_\odot \quad \} \text{ we get this but verify}$$

very close to actual

$$\underline{\underline{M_{ch} \approx 1.4 M_\odot}}$$

This is the Chandrasekhar's limit for w.D. stars:

Let's look at the speed of moving electron (degenerate electrons) in a white dwarf star.

$$\Delta x \Delta p \approx \hbar \quad \text{when} \quad v_e \approx \frac{p_e}{m_e} \approx \frac{\Delta p_e}{m_e} \approx \frac{\hbar}{\Delta x} \propto \frac{\hbar n^{1/3}}{m_e}$$

ΔV = Volume available to each electron: $\frac{V}{N} = \frac{1}{n}$

$$\Delta x \approx (\Delta V)^{1/3} \approx \frac{1}{n^{1/3}}$$

therefor $v_e = \frac{\hbar \rho^{1/3}}{m_e m_p^{1/3}} \approx v_e \propto \rho^{1/3}$

\downarrow

$\rho \propto v_e^3$

$$v_e = \frac{\hbar}{m_e m_p^{1/3}} \left(\frac{M}{R^3} \right)^{1/3} = \frac{\hbar}{m_e m_p^{1/3}} \left(\frac{(M)^{1/3}}{\frac{K}{G} M^{-1/3}} \right)$$

$$\Rightarrow v_e = \frac{\hbar}{m_e m_p^{1/3}} \frac{G}{K} M^{2/3}$$

At what M ; $G_e = C$

Say at M_{\max} $\Rightarrow M_{\max}^{2/3} = \frac{C m_e m_p^{1/3}}{\hbar} \frac{k}{G}$

$$k = \frac{\hbar^2}{m_e m_p^{5/3}} \Rightarrow M_{\max}^{2/3} = \frac{C m_e m_p^{1/3}}{\cancel{\hbar} G} \frac{\hbar^2}{\cancel{m_e m_p^{5/3}}}$$

$$= \frac{\hbar C}{G m_p^{4/3}}$$

$$\Rightarrow M_{\max} = \sqrt{\frac{\hbar^3 C^3}{G^3 m_p^4}}$$

Mass of star (now)

$< 7 \text{ Msun}$

$> 7 \text{ Msun} \text{ \& } < 18 \text{ Msun}$

$> 18 \text{ Msun}$

Mass of star (at death)

$< 1.4 \text{ Msun}$

$> 1.4 \text{ Msun}$

$> 1.4 \text{ Msun}$

Fate

White Dwarf

Neutron Star

Black Hole

